# The Kuramoto Model

Technical report for python numerical project

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The Kuramoto Model (Yoshki Kuramoto) is a mathematical model that describes the syncronization of coupled oscillators. This program, written by myself, Paradis Enzo, is based on this model to describe a set of 2D-coupled oscillators in various cases. In particular to reproduce data obtain in the case of the coupled arrays of Josephson junctions [1]. This program was edited for the last time the 10/23/2020. It was written in the language python 3 under ubuntu, so you just have to use python3 to execute it. There are 6 files. The first one is the main file, main.py, it is the only one you have to execute, you could have to comment or uncomment some lines. Instead of that you just have to edit one file, the settings file, settings.py were you can find all the settings. Then you have the kuramoto file, kuramoto.py, where you can find the class KuramotoModel which contains all the functions related with the model or the verifications. The data file, data.py, contains the class Data, which is used to manage data. The graphs file, graphs.py, with the class Graphs manages the displaying of the results. And the integrator file, integrator.py, where you can find the class Integration, which contains the three integrators: Euler, RK2, RK4. There is also the parameters directory which is important to save data. You can find all this files on github.

## I Functional requirement of the program

All the functions which are used for the computing or the displaying of the results are called in the main file. At first there are the functions wich create the initial values in respect of the parameters set in the settings file, then this values are stored in data (.dat) files in a directory named "parameters". With these data files, we don't have to repeat the calculations every time we want to test something. Then you have the functions that compute the results (the phase of the oscillators, the complex mean average, and the Shannon entropy), wich are stored in data files in the same directory. And finally there are the functions that display the graphs by using the module matplotlib.pyplot. In this section we will describe the functions that create data files and their datas. Firstly the initial datas are created throught the class Data wich is in the data file. Finally the computing of the other values is in the class KuramotoModel wich is in the kuramoto file.

## I.1 The initialisation of data

The initial datas are created by the function: data.init\_data(state). You can choose if you want to create new initial datas by changing the boolean value of the newData=False variable in the settings file.

Description	Input	Output
This function is using to create initial values, stored in data files in the directory parameters, according to the value of the argument state set by default to state="random". You can create data for random, chimera, inverse, or josephson states.	The argument state is a string. By default it takes the value "random" but you can give it these values:  • "random"  • "chimera"  • "inverse"  • "josephson"	This function will retrieve six data file in the parameters directory, computing in according to the state argument. The data files are:  • "omega.dat" • "theta0.dat" • "K.dat" • "eta.dat" • "alpha.dat" • "tau.dat"

Table 1: function data.init\_data()

Each state create six variables, wich are defined as follows:

- $\omega$  is a list of real of size N. It contains the natural oscillations of each oscillators.
- $\theta_0$  is a list of real of size N. It contains the intial phase of each oscillators. So at the  $\theta^i(t=t_0)$ .

- *K* is a matrix of real of size N×N. It contains the coupling coefficients between each oscillators depending on the nearest neighbour. The nearest neighbour of an oscillator are defined by M in the settings file. It is set as 30% of the totality of the oscillators, you can change it.
- $\alpha$  is a matrix of real of size N×N. It is the dephasing matrix of the coupling.
- $\tau$  is a matrix of real of size N×N. It is the delay matrix.
- $\eta$  is a matrix of real of size N×T. It is the representations of external noises for each oscillators at each time  $t \in ]0, T[$ .

They are different for each state, you can see their definitions in the tables 2 and 3. The function uniform() is from the module random, and provide random real numbers with uniform distribution, in the range given. The function randint() do the same things but for integers.

#### "random"

This state represent the case with random values defined by:

• 
$$\omega^i = uniform(0,3)$$

• 
$$\theta_0^i = uniform(0, \frac{2}{\pi})$$

• 
$$K_i^i = uniform(0, 1e10)$$

• 
$$\eta_i^i = uniform(0, 0.5)$$

• 
$$\alpha_j^i = uniform(0, \frac{2}{\pi})$$

• 
$$\tau_i^i = randint(0, N//2)$$

#### "chimera"

This state represent the case of a quantum chimera state[2] defined by:

• 
$$\omega^{i} = 0.2 + i * 0.4 * \sin(\frac{i^{2} * \pi}{(2 * N^{2})})$$

• 
$$\theta_0^i = uniform(0, \frac{2}{\pi})$$

• 
$$K_i^i = uniform(0, 1e10)$$

• 
$$\eta_i^i = uniform(0, 0.5)$$

• 
$$\alpha_i^i = 1.46$$

• 
$$\tau_i^i = randint(0, N//2)$$

Table 2: states

#### "inverse"

This state represent the case with the coupling matrix defined by the inverse of the distance between two oscillators, and the delay matrix is proportionnal to this distance.

• 
$$\omega^i = uniform(0,3)$$

• 
$$\theta_0^i = uniform(0, \frac{2}{\pi})$$

$$K_j^i = \left\{ \begin{array}{l} \frac{1}{|i-j|} \ if \ |i-j| \neq 0 \\ 1e20 \ otherwise \end{array} \right.$$

• 
$$\eta_i^i = uniform(0, 0.5)$$

• 
$$\alpha_i^i = 1.46$$

• 
$$\tau_i^i = |i - j|$$

## "josephson"

This state represent the case of the array of Josephson, data were drawn from this article [1]. Their datas are defined by:

• 
$$\omega^{i} = 0.2 + 0.4i \times sin(\frac{i^{2}\pi}{(2N^{2})})$$

• 
$$\theta_0^i = uniform(0, \frac{2}{\pi})$$

• 
$$K_j^i = \frac{Nr\omega^{i^2} 2 e/(\hbar r I_b) - \omega^i}{\sqrt{(L\omega^{i^2} - 1/C)^2 + \omega^{i^2} (R + rN)^2}}$$

• 
$$\eta_i^i = uniform(0, 0.5)$$

• 
$$cos(\alpha_j^i) = \frac{L\omega^{i^2} - 1/C}{\sqrt{(L\omega^{i^2} - 1/C)^2 + \omega^{i^2}(R + rN)^2}}$$

• 
$$\tau_i^i = randint(0, N//2)$$

Table 3: states

For each state choosen you have to define in the settings file the parameters Nr and Nc, to define the geometry of the system. Nr define the number of rows and Nc the number of columns, so N=Nr\*Nc is the number of oscillators that you have. For example if you choose N=12 in the geometry Nr=3, Nc=4, you will have this configuration:

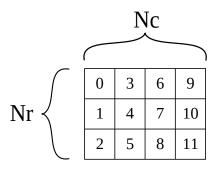


Figure 1: Configuration of the oscillators for Nr=3, Nc=4

The labels in the Figure are the labels of the oscillators. If you put a non-positive number you will have some issues, the program will not work. So be careful to respect the physics of the system.

#### The computing functions **I.2**

There are three computing functions that are imported from the class KuramotoModel, by the object kuramoto. You can choose if you want to compute new datas by changing the boolean value of the newComputing=False variable in the settings file.

<pre>kuramoto.integrate(f, theta0, tf=100, integrator="RK4")</pre>			
Description	Input	Output	
This function compute integrates the function f, with the initial	• f is the function to be integrated, such that $f : \mathbb{R}^n \to$	This function will retrieve two data files in the parameters direc-	

the function f, with the initial vector theta0 during a time tf, set by default to tf=100 with 1000 values(T), using the integrator, set by default to integrator="RK4"

- theta0 is the initial vector, such that  $\theta_0 \in \mathbb{R}^n$ . So it has to be a list of real of size N
- tf is the duration of the integration. Set to 100 by default. It is an integer
- integrator is the integrator that you want to choose. It can take only this string values: "Euler", "RK2", "RK4". Among the three integrators, RK4 is the one losing the least energy the longer the integration lasts, therefore it is the default value.

data files in the parameters directory.

- "theta.dat" : contains a matrix of real of size N×T. It is used like a list of vectors. Each vector contains the phase of each oscillators at one time, labeled in the same way as the Figure 1. So the matrix represent the evolution over time of this vector.
- "t.dat" : contains a list of real of size T, defined linspace(0, tf, T). linspace() is a function added by numpy.

Table 4: function kuramoto.integrate()

kuramoto.orders()		
Description	Input	Output
This function computes the orders parameters $R$ and $\Phi$ , of the $\theta^i$ over the time, by the relation : $R e^{i\Phi} = \frac{1}{N} \sum_{i=1}^{N} e^{i\theta^i}$ , for one time.	There are no arguments to this function because it takes only data files. It takes the theta file and the t file, see Table 4.	This function will retrieve two data files in the parameters directory.  • "R.dat": It contains a list of real of size T. R represents the complex mean of magnitudes of the $\theta^i$ over the time.
		• "phi.dat" : It contains a list of real of size T. phi represents the complex mean of angles of the $\theta^i$ over the time.

Table 5: function kuramoto.orders()

<pre>kuramoto.shannon_entropies()</pre>		
Description	Input	Output
This function computes the Shannon entropy over the time of $\theta$ . So for each vector of the list of vectors theta	There are no arguments to this function because it takes only data files. It takes the theta file and the t file, see Table 4.	This function will retrieve two data files in the parameters directory.  • "S.dat": It contains a list of real of size T. S represents the Shannon entropy over the time of each vector of theta

Table 6: function kuramoto.shannon entropies()

To summarize in the main file there are 4 functions that are called for the computing. The first one is for initializing data according to the state choosed by the user. The following three are here to compute the main datas by using the Kuramoto model. You can choose if you want to initialized new datas (1) or do a new computing (4, 5 & 6) by modifying the logical variables (newData, newComputing) in the settings file. In this file you can modify the state variable (2 & 3). You can also modify the geometry of the system by changing Nr and Nc, don't forget to redo the initial datas and the computing after that. Be careful not to set the number of rows or columns to a negative or zero number, it wouldn't make any sense. If you want to save or use other files than the predifined one you can modify it in the dictionnary FILE. You don't have to modify anything else that is not in the settings file. If you do it be careful, because you can break the program. For the rest of the main file there are 12 functions that are here to display the results. The last three are here to create gifs, they are the begining of the gold version. These displaying functions will be explained later.

## II Internal structure of the program

## II.1 Description of the physical model

In this program we had to describe 2D-coupled oscillators and tried to reproduce the results obtain in the article *Frequency locking in Josephson arrays: Connection with the Kuramoto model* [1]. So we will use the Kuramoto Model. Proposed by Yoshiki kuramoto this model is used to describe the syncronization of a large set of coupled oscillators. It is a mathematical model and the form that we have used is:

$$\dot{\theta}^{i}(t) = \omega_{i} + \frac{1}{N} \sum_{j=0}^{N-1} K_{ij} \sin(\theta^{j}(t - \tau_{ij}) - \theta^{i}(t) + \alpha_{ij}) + \eta_{i}(t)$$

where:

- $\dot{\theta}^i(t)$  represent the variation of the phase of the i-th oscillator over the time.
- $\omega_i$  represent the natural frequency of the i-th oscillator.
- N is the total number of oscillators in the system.
- $K_{ij}$  is the coupling matrix. It represent the coupling between each oscillators depending on the inverse of the physical distance between them. Sometimes you can consider only the nearest neighbour.
- $\theta^i$  is the phase of the i-th oscillator.
- $\tau_{ij}$  is the delay matrix. It is to take into account the phase propagation of the oscillators which have an influence on the i-th oscillator.
- $\alpha_{ij}$  is the dephasing matrix. It represent the dephasing due of the coupled oscillators.
- $\eta_i(t)$  represent the influence of the environement, the possible external noises.

As you can see this model is pretty complete. The formula is the sum of all the influences due to the environment and the coupled oscillators occuring on each oscillators at a given given time. That why it was chosen to build  $\theta$  like a list of vectors.

## II.2 Desciption of the used scientific computation algorithms

During the computation we have used 3 algorithm for the integration of the kuramoto model II.1, the Euler, RK2 and RK4 algorithm. But you can choose the one that you want to use. It is recommended to use the RK4 method because it loses less energy over time than the others, so it is more precise but if you want less computing time it is a good thing to try the other. You can see all these integrators in the class Integration in the integrator file.

We are looking to integrate  $\dot{\theta} = f(\theta)$ .

Let  $\theta \in \mathbb{R}^n$  and  $f : \mathbb{R}^n \to \mathbb{R}^n$ . We wil note  $\theta^i$  the i-th component of  $\theta$  and  $\theta_n$  the vector  $\theta$  at the time  $t_n$ ,  $\theta_n \equiv \theta(t_n)$ . So we have  $\theta_n^i$  the i-th component at a time  $t_n$  of  $\theta$ .

The Euler and Runge-Kutta methods are integrators of first-order differential equations and are based on the definition of the derivative:

$$\dot{\theta}(t_n) = \frac{\theta(t_n + h) - \theta(t_n)}{h} \simeq \frac{\theta_{n+1} - \theta_n}{\Delta t}$$

$$\iff \theta_{n+1} = \theta_n + \dot{\theta}(t_n)\Delta t = \theta_n + f(\theta_n)\Delta t$$
(II.1)

with  $\Delta t = \frac{T}{N_t}$  the step of the simulation ( $N_t \in \mathbb{N}$  the number of steps of the simulation and T the time of the simulation). We need  $\Delta t$  sufficiently small so  $N_t$  enought large and T sufficiently thin.

## II.2.1 Euler's method

The Euler algorithm is directly an application of the Equation II.1:

$$\theta_{n+1} = \theta_n + f(\theta_n) \Delta t \tag{II.2}$$

This method will diverge quicker than the two other but the computing time is the shortest.

#### II.2.2 RK2 method

The Runge-Kutta algorithm (RK2) is like the Euler's method but with a double use of the Equation II.1. There is a middle step to refine the estimation of f.

$$\theta_{n+1} = \theta_n + f(\theta_n + f(\theta_n) \frac{\Delta t}{2}) \Delta t$$
 (II.3)

This method will diverge quicker than the RK4 method and slower than the Euler's method. Its computing time will be longer than the one of the Euler's method but remains shorter than the one of RK4 method.

#### II.2.3 RK4 method

RK4, the fourth order Runge-kutta algorithm, is similar to the RK2 algorithm but with a refinement around the middle point.

$$\theta_{n+1} = \theta_n + (K_{1,n} + 2K_{2,n} + 2K_{3,n} + K_{4,n}) \frac{\Delta t}{6}$$
(II.4)

with

$$K_{1,n} = f(\theta_n)$$

$$K_{2,n} = f\left(\theta_n + K_{1,n} \frac{\Delta t}{2}\right)$$

$$K_{3,n} = f\left(\theta_n + K_{2,n} \frac{\Delta t}{2}\right)$$

$$K_{4,n} = f\left(\theta_n + K_{3,n} \Delta t\right)$$
(II.5)

Even if this is the slowest method of three orders of the Runge-Kutta algorithm, it is the most accurate one. You can find all this algorithm in the integrator file in the class Integration.

## II.3 List of the constitutive elements of the program

The Program includes 4 classes. There are 3 that are used in the main file and 1 in the class KuramotoModel, see Figure 3.

## II.3.1 The class Integration

This class doesn't have any attributes. It is just used to store the integrators. It can be used in an other program just by writting import integrator at the begining of you code, and create an object with the class. This class is using the module numpy. Here is a table of the class:

class Integration		
	Attributes	Methods
Ø		This class contains 3 methods, see <pre>subsection II.2:</pre> • euler(f, x0,t)
		• RK2(f, x0,t)
		• RK4(f, x0,t)

Table 7: class Integration

#### II.3.2 The class Data

This class is used to create the initial data. Here is a table of the class:

Attributes  Methods  This class contains 8 attributes, see subsection I.1:  M, type:int  N, type:real  K, type:ndarray, shape: (N,N)  alpha, type:ndarray, shape: (N,N)  eta, type:ndarray, shape: (N,1000)  omega, type:ndarray, shape: (N,N)  Methods  This class contains 5 methods:  init_data(state), see Table 1  chimera_states(), see Table 2  inverse_states(), see Table 3  random_states(), see Table 2  josephson_matrice(), see Table 3	class Data		
<pre>tion I.1:</pre>	Attributes	Methods	
<ul><li>tau,type:ndarray,shape:(N,N)</li><li>theta0,type:ndarray,shape:(N,N)</li></ul>	<pre>tion I.1:     M,type:int     N,type:real     K,type:ndarray,shape:(N,N)     alpha,type:ndarray,shape:(N,N)     eta,type:ndarray,shape:(N,1000)     omega,type:ndarray,shape:(N,N)     tau,type:ndarray,shape:(N,N)</pre>	<ul> <li>init_data(state), see Table 1</li> <li>chimera_states(), see Table 2</li> <li>inverse_states(), see Table 3</li> <li>random_states(), see Table 2</li> </ul>	

Table 8: class Data

You can add a new method in this class to add new configurations. But don't forget to add a new possible value to the state variable in the init\_data() function. As follows:

```
if state == "random":
    self.random_states()
elif state == "chimera":
    self.chimera_states()
elif state == "inverse":
    self.inverse_states()
elif state == "josephson":
    self.josephson_matrice()
elif state == "new method":
    self.new_method()
```

## II.3.3 The class Graphs

This class has no attributes. It is used to store all the displaying functions and the function to create gifs. This class contains 13 methods:

- kuramoto(t, show=True): plot  $i \mapsto \{\theta^i(t)\}_{i=0,\dots,N-1}$  at a time t. The argument show is used by the animation function, it allows to not display the graph but save it in the animation directory (if show==False).
- all\_kuramoto(): plot the N curves  $t \mapsto \theta^i(t)$
- anim\_kuramoto(): using kuramoto(t, show=False) it take all the "kuramoto\*.png" in the animation directory and create a gif with them.
- dens\_kuramoto(): plot  $(i,t) \mapsto \theta^i(t)$  using the numpy.contourf() method.
- dens\_kuramoto\_coord(t, show=True): plot  $(r,c) \mapsto \theta^{i_{r,d}}(t)$  at time t using the numpy.contourf() method.
- anim\_dens\_kuramoto\_coord(): using dens\_kuramoto\_coord(t, show=False) it create a gif.
- orders(): plot  $t \mapsto R(t)$  and  $t \mapsto \Phi(t)$ .
- shannon(t): plot  $i \mapsto S_i^{q,n}(t)$  at a time t.
- dens\_shannon(): plot  $(i,t) \mapsto S_i^{q,n}(t)$  using the numpy.contourf() method.
- dens\_shannon\_coord(t, show=True): plot  $(r,c) \mapsto S_i^{q,n}(t)$  at time t using the numpy.contourf() method.

- anim\_dens\_shannon\_coord(): using dens\_shannon\_coord(t, show=False) it create a gif.
- connectivity(K, kmin): it is a static method that return a list with all the edges connected in the coupling matrix K using a criterion kmin.
- graph\_connectivity(kmin=0): using the static method connectivity(K, kmin) it creates a graph of the connectivity of the oscillators with the criterion kmin.

You need to have the modules numpy, matplotlib.pyplot, networkx and PIL, installed on your computer to use this class. The module numpy is used to recover data from the data files, matplotlib.pyplot is used to plot all the graphs, networkx is used to create the connectivity graph, and PIL is used to create the gifs with all the graphs.

# II.4 Functional diagram of the program

The program is cut out in 6 modules: main, graphs, kuramoto, data, integrator, and settings. You can see how they interact with each other in the following figure: In these modules data, integrator graphs and kuramoto

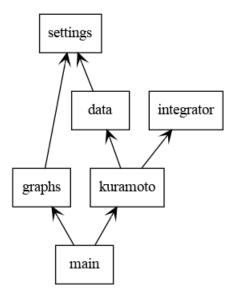


Figure 2: Diagram of the interactions between the modules.

contain a class. These classes are depending to each other like in the following diagram:

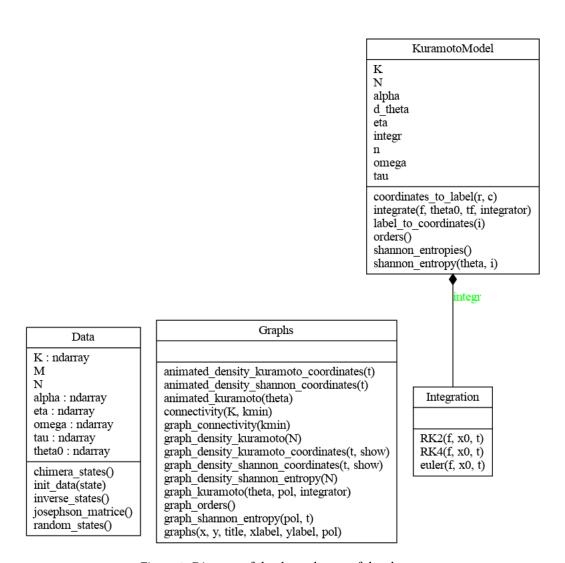


Figure 3: Diagram of the dependences of the classes.

# References

- [1] Steven H. Strogatz Kurt Wiesenfeld, Pere Colet. Frequency locking in josephson arrays: Connection with the kuramoto model. 57(2), February 1998.
- [2] David Viennot. Chaos et chimères quantiques. http://perso.utinam.cnrs.fr/ viennot/publi/chaos.pdf. En ligne; dernière visite : 31 octobre 2020.