

# **The Kuramoto Model**

## **Technical report for python numerical project**

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# I Functional requirement of the program

All the functions which are used for the computing or the displaying of the results are called in the main file. At first there are the functions which create the initial values in respect of the parameters set in the settings file, then these values are stocked in data (.dat) files in a directory named "parameters". With these data files, we don't have to repeat the calculations every time we want to test something. Then you have the functions that compute the results (the phase of the oscillators, the complex mean average, and the Shannon entropy), which are stocked in data files in the same directory. And finally there are the functions that display the graphs by using the module `matplotlib.pyplot`. In this section we will describe the functions that create the data files and their data. Firstly the initial data is created through the `class Data` which is in the data file. Finally the computing of the other values is in the `class KuramotoModel` which is in the kuramoto file.

## I.1 The initialisation of the data

The initial data is created by the function: `data.init_data(state)`. You can choose if you want to create new initial data by changing the boolean value of the `newData=False` variable in the settings file.

data.init_data(state="random")		
Description	Input	Output
This function is used to create initial values, stocked in data files in the directory parameters, according to the value of the argument state set by default to <code>state="random"</code> . You can create data for random, chimera, inverse, or josephson states.	<p>The argument state is a string. By default it takes the value "random" but you can give it these values:</p> <ul style="list-style-type: none"> <li>• "random"</li> <li>• "chimera"</li> <li>• "inverse"</li> <li>• "josephson"</li> </ul>	<p>This function will retrieve six data files in the parameters directory, computing in according to the state argument. The data files are:</p> <ul style="list-style-type: none"> <li>• "omega.dat"</li> <li>• "theta0.dat"</li> <li>• "K.dat"</li> <li>• "eta.dat"</li> <li>• "alpha.dat"</li> <li>• "tau.dat"</li> </ul>

Table 1: function data.init\_data()

Each state creates six variables, which are defined as follows:

- $\omega$  is a list of real of size N. It contains the natural oscillations of each oscillators.
- $\theta_0$  is a list of real of size N. It contains the initial phase of each oscillators. So at the  $\theta^i(t = t_0)$ .
- $K$  is a matrix of real of size  $N \times N$ . It contains the coupling coefficients between each oscillators depending on the nearest neighbour. The nearest neighbour of an oscillator are defined by M in the settings file. It is set as 30% of the totality of the oscillators, you can change it.
- $\alpha$  is a matrix of real of size  $N \times N$ . It is the dephasing matrix of the coupling.
- $\tau$  is a matrix of real of size  $N \times N$ . It is the delay matrix.
- $\eta$  is a matrix of real of size  $N \times T$ . It is the representations of external noises for each oscillators at each time  $t \in ]0, T[$ .

They are different for each state, you can see their definitions in the tables 2 and 3. The function `uniform()` is from the module `random`, and provides random real numbers with uniform distribution, in the range given. The function `randint()` does the same things but for integers.

"random"	"chimera"
<p>This state represent the case with random values defined by:</p> <ul style="list-style-type: none"> <li>• <math>\omega^i = \text{uniform}(0, 3)</math></li> <li>• <math>\theta_0^i = \text{uniform}(0, \frac{2}{\pi})</math></li> <li>• <math>K_j^i = \text{uniform}(0, 1e10)</math></li> <li>• <math>\eta_j^i = \text{uniform}(0, 0.5)</math></li> <li>• <math>\alpha_j^i = \text{uniform}(0, \frac{2}{\pi})</math></li> <li>• <math>\tau_j^i = \text{randint}(0, N // 2)</math></li> </ul>	<p>This state represent the case of a quantum chimera state[2] defined by:</p> <ul style="list-style-type: none"> <li>• <math>\omega^i = 0.2 + i * 0.4 * \sin(\frac{i^2 * \pi}{(2 * N^2)})</math></li> <li>• <math>\theta_0^i = \text{uniform}(0, \frac{2}{\pi})</math></li> <li>• <math>K_j^i = \text{uniform}(0, 1e10)</math></li> <li>• <math>\eta_j^i = \text{uniform}(0, 0.5)</math></li> <li>• <math>\alpha_j^i = 1.46</math></li> <li>• <math>\tau_j^i = \text{randint}(0, N // 2)</math></li> </ul>

Table 2: states

"inverse"	"josephson"
<p>This state represent the case with the coupling matrix defined by the inverse of the distance between two oscillators, and the delay matrix is proportionnal to this distance.</p> <ul style="list-style-type: none"> <li>• <math>\omega^i = \text{uniform}(0, 3)</math></li> <li>• <math>\theta_0^i = \text{uniform}(0, \frac{2}{\pi})</math></li> <li>• <math>K_j^i = \begin{cases} \frac{1}{ i-j } &amp; \text{if }  i-j  \neq 0 \\ 1e20 &amp; \text{otherwise} \end{cases}</math></li> <li>• <math>\eta_j^i = \text{uniform}(0, 0.5)</math></li> <li>• <math>\alpha_j^i = 1.46</math></li> <li>• <math>\tau_j^i =  i-j </math></li> </ul>	<p>This state represent the case of the array of Josephson, the datas were drawn from this article [1]. Their datas are defined by:</p> <ul style="list-style-type: none"> <li>• <math>\omega^i = 0.2 + 0.4i \times \sin(\frac{i^2 \pi}{(2N^2)})</math></li> <li>• <math>\theta_0^i = \text{uniform}(0, \frac{2}{\pi})</math></li> <li>• <math>K_j^i = \frac{Nr\omega^{i^2} 2e / (\hbar r I_b) - \omega^i}{\sqrt{(L\omega^{i^2} - 1/C)^2 + \omega^{i^2} (R + rN)^2}}</math></li> <li>• <math>\eta_j^i = \text{uniform}(0, 0.5)</math></li> <li>• <math>\cos(\alpha_j^i) = \frac{L\omega^{i^2} - 1/C}{\sqrt{(L\omega^{i^2} - 1/C)^2 + \omega^{i^2} (R + rN)^2}}</math></li> <li>• <math>\tau_j^i = \text{randint}(0, N // 2)</math></li> </ul>

Table 3: states

For each state choosen you have to define in the settings file the parameters Nr and Nc, to define the geometry of the system. Nr define the number of rows and Nc the number of columms, so  $N=Nr*Nc$  is the number of oscillators that you have. For example if you choose  $N=12$  in the geometry  $Nr=3$ ,  $Nc=4$ , you will have this configuration:

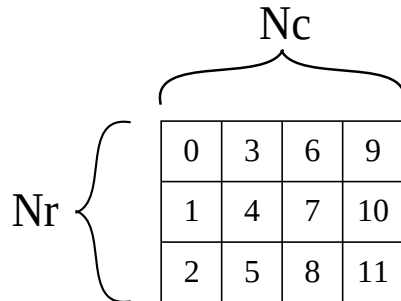


Figure 1: Configuration of the oscillators for  $N_r=3$ ,  $N_c=4$

The labels in the Figure are the labels of the oscillators. If you put a non-positive number you will have some issues, the program will not work. So be careful to respect the physics of the system.

## I.2 The computing functions

There are three computing functions that are imported from the `class KuramotoModel`, by the object `kuramoto`. You can choose if you want to compute new datas by changing the boolean value of the `newComputing=False` variable in the settings file.

kuramoto.integrate(f, theta0, tf=100, integrator="RK4")		
Description	Input	Output
<p>This function compute integrates the function <math>f</math>, with the initial vector <math>\theta_0</math> during a time <math>tf</math>, set by default to <math>tf=100</math> with 1000 values(<math>T</math>), using the integrator, set by default to <code>integrator="RK4"</code></p>	<ul style="list-style-type: none"> <li><math>f</math> is the function to be integrated, such that <math>f : \mathbb{R}^n \rightarrow \mathbb{R}^n</math></li> <li><math>\theta_0</math> is the initial vector, such that <math>\theta_0 \in \mathbb{R}^n</math>. So it has to be a list of real of size <math>N</math></li> <li><math>tf</math> is the duration of the integration. Set to 100 by default. It is an integer</li> <li>integrator is the integrator that you want to choose. It can take only this string values: "Euler", "RK2", "RK4". Among the three integrators, RK4 is the one losing the least energy the longer the integration lasts, therefore it is the default value.</li> </ul>	<p>This function will retrieve two data files in the parameters directory.</p> <ul style="list-style-type: none"> <li><code>"theta.dat"</code> : contains a matrix of real of size <math>N \times T</math>. It is used like a list of vectors. Each vector contains the phase of each oscillators at one time, labeled in the same way as the <a href="#">Figure 1</a>. So the matrix represent the evolution over time of this vector.</li> <li><code>"t.dat"</code> : contains a list of real of size <math>T</math>, defined by <code>linspace(0, tf, T)</code>. <code>linspace()</code> is a function added by numpy.</li> </ul>

Table 4: function kuramoto.integrate()

kuramoto.orders()		
Description	Input	Output
This function computes the orders parameters $R$ and $\Phi$ , of the $\theta^i$ over the time, by the relation : $R e^{i\Phi} = \frac{1}{N} \sum_{i=1}^N e^{i\theta^i}$ , for one time.	There are no arguments to this function because it takes only data files. It takes the theta file and the t file, see <a href="#">Table 4</a> .	This function will retrieve two data files in the parameters directory. <ul style="list-style-type: none"> <li>• <b>"R.dat"</b> : It contains a list of real of size T. R represents the complex mean of magnitudes of the <math>\theta^i</math> over the time.</li> <li>• <b>"phi.dat"</b> : It contains a list of real of size T. phi represents the complex mean of angles of the <math>\theta^i</math> over the time.</li> </ul>

Table 5: function kuramoto.orders()

kuramoto.shannon_entropies()		
Description	Input	Output
This function computes the Shannon entropy over the time of $\theta$ . So for each vector of the list of vectors theta	There are no arguments to this function because it takes only data files. It takes the theta file and the t file, see <a href="#">Table 4</a> .	This function will retrieve two data files in the parameters directory. <ul style="list-style-type: none"> <li>• <b>"S.dat"</b> : It contains a list of real of size T. S represents the Shannon entropy over the time of each vector of theta</li> </ul>

Table 6: function kuramoto.shannon\_entropies()

## References

- [1] Pere Colet Kurt Wiesenfeld and Steven H. Strogatz. Frequency locking in josephson arrays: Connection with the kuramoto model. 57(2), February 1998.
- [2] David Viennot. Chaos et chimères quantiques. <http://perso.utinam.cnrs.fr/viennot/publi/chaos.pdf>. En ligne; dernière visite : 31 octobre 2020.