

Assignment 2

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1

From the definition of \underline{u} and \underline{v} , we have:

$$\mathbf{u}_a = \begin{bmatrix} 4 \\ u_{a2} \\ -2 \end{bmatrix}, \quad \mathbf{v}_a = \begin{bmatrix} 1 \\ 3 \\ 5 \end{bmatrix}$$

a) Using the definition of the dot product:

$$\underline{u} \cdot \underline{v} = \mathbf{u}_a^T \mathbf{v}_a = 4 + 3u_{a2} - 10 = -6 + 3u_{a2} \stackrel{!}{=} 12$$

Solving for u_{a2} :

$$u_{a2} = \frac{12 + 6}{3} = 6$$

b) Using the definition of the cross product:

$$\underline{u} \times \underline{v} = \underline{\mathcal{F}}_a^T \mathbf{u}_a^\times \mathbf{v}_a = \underline{\mathcal{F}}_a^T \begin{bmatrix} 0 & 2 & u_{a2} \\ -2 & 0 & -4 \\ -u_{a2} & 4 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 3 \\ 5 \end{bmatrix} = \underline{\mathcal{F}}_a^T \begin{bmatrix} 6 + 5u_{a2} \\ -22 \\ -u_{a2} + 12 \end{bmatrix} \stackrel{!}{=} \underline{\mathcal{F}}_a^T \begin{bmatrix} 16 \\ -22 \\ 10 \end{bmatrix}$$

Solving for u_{a2} , we have the following two equations:

$$u_{a2} = \frac{16 - 6}{5} = 2$$

$$u_{a2} = -(10 - 12) = 2$$

Therefore, $u_{a2} = 2$ is the unique solution.

2

$$\begin{aligned}
 \det \boldsymbol{\omega}_a^\times &= \det \begin{bmatrix} 0 & -\omega_{a3} & \omega_{a2} \\ \omega_{a3} & 0 & -\omega_{a1} \\ -\omega_{a2} & \omega_{a1} & 0 \end{bmatrix} = \begin{bmatrix} 0 \\ \omega_{a3} \\ -\omega_{a2} \end{bmatrix}^\top \begin{bmatrix} -\omega_{a3} \\ 0 \\ \omega_{a1} \end{bmatrix}^\times \begin{bmatrix} \omega_{a2} \\ -\omega_{a1} \\ 0 \end{bmatrix} \\
 &= \begin{bmatrix} 0 & \omega_{a3} & -\omega_{a2} \end{bmatrix} \begin{bmatrix} 0 & -\omega_{a1} & 0 \\ \omega_{a1} & 0 & \omega_{a3} \\ 0 & -\omega_{a3} & 0 \end{bmatrix} \begin{bmatrix} \omega_{a2} \\ -\omega_{a1} \\ 0 \end{bmatrix} \\
 &= \begin{bmatrix} 0 & \omega_{a3} & -\omega_{a2} \end{bmatrix} \begin{bmatrix} \omega_{a1}^2 \\ \omega_{a1}\omega_{a2} \\ \omega_{a3}\omega_{a1} \end{bmatrix} = \omega_{a3}\omega_{a1}\omega_{a2} - \omega_{a2}\omega_{a3}\omega_{a1} = 0
 \end{aligned}$$

Therefore, $\boldsymbol{\omega}_a^\times$ is not invertible.

3

a) By inspection, we can see that:

$$r_{b1} = r_{a1} = 1, \quad r_{b2} = r_{a3} = 4, \quad r_{b3} = -r_{a2} = -3$$

Therefore, $\mathbf{r}_b = \begin{bmatrix} 1 & 4 & -3 \end{bmatrix}^\top$

b) \mathcal{F}_b is obtained by rotating \mathcal{F}_a about \underline{a}^1 by 90° .

$$\begin{aligned}
 \mathbf{C}_{ba} &= \underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_a^\top = \begin{bmatrix} \underline{b}^1 \\ \underline{b}^2 \\ \underline{b}^3 \end{bmatrix} \cdot \begin{bmatrix} \underline{a}^1 & \underline{a}^2 & \underline{a}^3 \end{bmatrix} = \begin{bmatrix} \underline{b}^1 \cdot \underline{a}^1 & \underline{b}^1 \cdot \underline{a}^2 & \underline{b}^1 \cdot \underline{a}^3 \\ \underline{b}^2 \cdot \underline{a}^1 & \underline{b}^2 \cdot \underline{a}^2 & \underline{b}^2 \cdot \underline{a}^3 \\ \underline{b}^3 \cdot \underline{a}^1 & \underline{b}^3 \cdot \underline{a}^2 & \underline{b}^3 \cdot \underline{a}^3 \end{bmatrix} \\
 &= \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix}
 \end{aligned}$$

where the dot products were found by looking at Figure 1.

c)

$$\mathbf{r}_b = \mathbf{C}_{ba}\mathbf{r}_a = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & -1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 3 \\ 4 \end{bmatrix} = \begin{bmatrix} 1 \\ 4 \\ -3 \end{bmatrix}$$

This corresponds to the solution found in a).

4

$$\begin{aligned}
 \det \mathbf{Q} &= \det \begin{bmatrix} 0 & \frac{1}{2} & -\frac{\sqrt{3}}{2} \\ 1 & 0 & 0 \\ 0 & \frac{\sqrt{3}}{2} & \frac{1}{2} \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}^T \begin{bmatrix} \frac{1}{2} \\ 0 \\ \frac{\sqrt{3}}{2} \end{bmatrix} \times \begin{bmatrix} -\frac{\sqrt{3}}{2} \\ 0 \\ \frac{1}{2} \end{bmatrix} \\
 &= \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & -\frac{\sqrt{3}}{2} & 0 \\ \frac{\sqrt{3}}{2} & 0 & -\frac{1}{2} \\ 0 & \frac{1}{2} & 0 \end{bmatrix} \begin{bmatrix} -\frac{\sqrt{3}}{2} \\ 0 \\ \frac{1}{2} \end{bmatrix} \\
 &= \begin{bmatrix} 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ -1 \\ 0 \end{bmatrix} = -1
 \end{aligned}$$

In conclusion, \mathbf{Q} is not a valid direction cosine matrix because its determinant is different than +1.

5

From the definition of the cross product:

$$\underline{u} \times \underline{v} = \underline{\mathcal{F}}_a^T \mathbf{u}_a^\times \mathbf{v}_a = \underline{\mathcal{F}}_b^T \mathbf{u}_b^\times \mathbf{v}_b$$

Multiplying on the left by $\underline{\mathcal{F}}_b$ (dot product), we get:

$$\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_a^T \mathbf{u}_a^\times \mathbf{v}_a = \underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_b^T \mathbf{u}_b^\times \mathbf{v}_b \quad (1)$$

Recall:

$$\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_a^T = \mathbf{C}_{ba} \quad (2)$$

$$\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_b^T = \mathbf{1} \quad (3)$$

Substituting (2) and (3) into (1), we obtain:

$$\mathbf{C}_{ba} \mathbf{u}_a^\times \mathbf{v}_a = \mathbf{u}_b^\times \mathbf{v}_b \quad (4)$$

Now, from the definition of the DCM:

$$\mathbf{v}_a = \mathbf{C}_{ab} \mathbf{v}_b = \mathbf{C}_{ba}^T \mathbf{v}_b \quad (5)$$

$$\mathbf{u}_b = \mathbf{C}_{ba} \mathbf{u}_a \quad (6)$$

Substituting (5) and (6) into (4), we obtain:

$$\mathbf{C}_{ba} \mathbf{u}_a^\times \mathbf{C}_{ba}^T \mathbf{v}_b = (\mathbf{C}_{ba} \mathbf{u}_a)^\times \mathbf{v}_b \quad (7)$$

Since (7) must be valid for any \underline{v} (hence any \mathbf{v}_b):

$$\mathbf{C}_{ba} \mathbf{u}_a^\times \mathbf{C}_{ba}^T = (\mathbf{C}_{ba} \mathbf{u}_a)^\times \quad \square$$

6

a) We have:

$$\underline{T} = \underline{\mathcal{F}}_a^\top \mathbf{T}_a \underline{\mathcal{F}}_a = \underline{\mathcal{F}}_b^\top \mathbf{T}_b \underline{\mathcal{F}}_b \quad (8)$$

Taking the dot product on the left with $\underline{\mathcal{F}}_b$ and on the right with $\underline{\mathcal{F}}_b^\top$:

$$\underbrace{\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_a^\top}_{\mathbf{C}_{ba}} \mathbf{T}_a \underbrace{\underline{\mathcal{F}}_a \cdot \underline{\mathcal{F}}_b^\top}_{\mathbf{C}_{ab}} = \underbrace{\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_b^\top}_{\mathbf{I}} \mathbf{T}_b \underbrace{\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_b^\top}_{\mathbf{I}}$$

Where we used the definition of the direction cosine matrix. Rewriting both sides:

$$\mathbf{C}_{ba} \mathbf{T}_a \mathbf{C}_{ab} = \mathbf{T}_b \quad \square$$

b) Using (8) and $\underline{u} = \underline{\mathcal{F}}_a^\top \mathbf{u}_a$, we find:

$$\underline{T} \cdot \underline{u} = \underline{\mathcal{F}}_b^\top \mathbf{T}_b \underbrace{\underline{\mathcal{F}}_b \cdot \underline{\mathcal{F}}_a^\top}_{\mathbf{C}_{ba}} \mathbf{u}_a = \underline{\mathcal{F}}_b^\top \mathbf{T}_b \mathbf{C}_{ba} \mathbf{u}_a \quad \square$$

7

a)

$$\begin{aligned} (\lambda - 1)(\lambda^2 + \lambda(1 - \text{tr} \mathbf{C}_{qp}) + 1) &= (\lambda - 1)\lambda^2 + (\lambda - 1)(\lambda - \lambda \text{tr} \mathbf{C}_{qp}) + (\lambda - 1) \\ &= \lambda^3 - \lambda^2 + \lambda^2 - \lambda^2 \text{tr} \mathbf{C}_{qp} - \lambda + \lambda \text{tr} \mathbf{C}_{qp} + \lambda - 1 \\ &= \lambda^3 - \lambda^2 \text{tr} \mathbf{C}_{qp} + \lambda \text{tr} \mathbf{C}_{qp} - 1 \\ &= 0 \quad \square \end{aligned}$$

One obvious solution of this equation is $\lambda_1 = +1$, which is therefore also an eigenvalue of \mathbf{C}_{qp} .

b) Recall that the trace is a linear operator, i.e. $\forall \alpha, \beta \in \mathbb{C}, \forall \mathbf{A}, \mathbf{B} \in \mathbb{C}^{n \times n}, n \in \mathbb{N}$:

$$\text{tr}(\alpha \mathbf{A} + \beta \mathbf{B}) = \alpha \text{tr} \mathbf{A} + \beta \text{tr} \mathbf{B}$$

Therefore, we can express the trace of \mathbf{C}_{qp} as follows:

$$\text{tr} \mathbf{C}_{qp} = \cos \phi \text{tr}(\mathbf{1}) + (1 - \cos \phi) \text{tr}(\mathbf{a} \mathbf{a}^\top) - \sin \phi \text{tr}(\mathbf{a}^\times) \quad (9)$$

We can now compute the trace of each element.

- Trivially:

$$\text{tr}(\mathbf{1}) = 3 \quad (10)$$

- Let $\mathbf{a} = \begin{bmatrix} a_1 & a_2 & a_3 \end{bmatrix}^\top$, therefore:

$$\mathbf{a}\mathbf{a}^\top = \begin{bmatrix} a_1 \\ a_2 \\ a_3 \end{bmatrix} \begin{bmatrix} a_1 & a_2 & a_3 \end{bmatrix} = \begin{bmatrix} a_1^2 & a_1a_2 & a_1a_3 \\ a_2a_1 & a_2^2 & a_2a_3 \\ a_3a_1 & a_3a_2 & a_3^2 \end{bmatrix}$$

By inspection:

$$\text{tr}(\mathbf{a}\mathbf{a}^\top) = a_1^2 + a_2^2 + a_3^2 = \|\underline{a}\|_2^2 = 1 \quad (11)$$

- Taking the same generic notation for \mathbf{a} :

$$\mathbf{a}^\times = \begin{bmatrix} 0 & -a_3 & a_2 \\ a_3 & 0 & -a_1 \\ -a_2 & a_1 & 0 \end{bmatrix}$$

By inspection:

$$\text{tr}(\mathbf{a}^\times) = 0 \quad (12)$$

Substituting (10), (11) and (12) into (9):

$$\text{tr}\mathbf{C}_{qp} = 3 \cos \phi + (1 - \cos \phi) = 1 + 2 \cos \phi \quad \square$$

Using this result and the fact that $\mathbf{C}_{qp} \in \text{SO}(3)$, we have the two following conditions:

$$\det \mathbf{C}_{qp} = \lambda_1 \lambda_2 \lambda_3 = 1 \quad (13)$$

$$\text{tr} \mathbf{C}_{qp} = \lambda_1 + \lambda_2 + \lambda_3 = 1 + 2 \cos \phi \quad (14)$$

Since we have already found that $\lambda_1 = 1$, (13) and (14) become:

$$\lambda_2 \lambda_3 = 1 \quad (15)$$

$$\lambda_2 + \lambda_3 = 2 \cos \phi \quad (16)$$

From (16), we find:

$$\lambda_3 = 2 \cos \phi - \lambda_2 \quad (17)$$

Substituting (17) into (15):

$$\lambda_2(2 \cos \phi - \lambda_2) = 1 \quad (18)$$

We can now rearrange (18) as follows:

$$\lambda_2^2 - 2 \cos \phi \lambda_2 + 1 = 0 \quad (19)$$

This is simply a quadratic equation that we can solve as follows:

$$\begin{aligned} \lambda_2 &= \cos \phi \pm \sqrt{\cos^2 \phi - 1} = \cos \phi \pm \sqrt{(-1)(1 - \cos^2 \phi)} \\ &= \cos \phi \pm \sqrt{-1} \sqrt{\sin^2 \phi} = \cos \phi \pm i \sin \phi = e^{\pm i \phi} \end{aligned}$$

Substituting this solution into (15):

$$e^{\pm i \phi} \lambda_3 = 1 \quad \Rightarrow \quad \lambda_3 = \frac{1}{e^{\pm i \phi}} = e^{\mp i \phi}$$

In other words, λ_2 is the complex conjugate of λ_3 and vice versa. It doesn't matter which one has the positive sign in the exponent.

8

First, let's derive some useful properties:

$$\begin{aligned}
 \mathbf{a}^\times \mathbf{a}^\times &= \mathbf{a} \mathbf{a}^\top - \mathbf{a}^\top \mathbf{a} \mathbf{1} = \mathbf{a} \mathbf{a}^\top - \mathbf{1} \\
 (\mathbf{a} \mathbf{a}^\top - \mathbf{1})^2 &= \overbrace{\mathbf{a} \mathbf{a}^\top \mathbf{a} \mathbf{a}^\top}^{=1} - 2\mathbf{a} \mathbf{a}^\top + \mathbf{1} = \mathbf{1} - \mathbf{a} \mathbf{a}^\top \\
 (\mathbf{a} \mathbf{a}^\top - \mathbf{1})^3 &= (\mathbf{a} \mathbf{a}^\top - \mathbf{1})^2 (\mathbf{a} \mathbf{a}^\top - \mathbf{1}) = (\mathbf{1} - \mathbf{a} \mathbf{a}^\top) (\mathbf{a} \mathbf{a}^\top - \mathbf{1}) = \mathbf{a} \mathbf{a}^\top - \mathbf{1} \\
 &\vdots \\
 (\mathbf{a} \mathbf{a}^\top - \mathbf{1})^n &= (-1)^n (\mathbf{1} - \mathbf{a} \mathbf{a}^\top), \quad n \in \mathbb{N} \setminus \{0\}
 \end{aligned}$$

Moreover:

$$\mathbf{a}^\times (\mathbf{1} - \mathbf{a} \mathbf{a}^\top) = \mathbf{a}^\times - \overbrace{\mathbf{a}^\times \mathbf{a} \mathbf{a}^\top}^{=0} = \mathbf{a}^\times$$

Now, let's expand the expression for $e^{-\phi \mathbf{a}^\times}$ using the definition of the matrix exponential and the former properties:

$$\begin{aligned}
 e^{-\phi \mathbf{a}^\times} &= \sum_{k=0}^{\infty} \frac{(-\phi)^k}{k!} (\mathbf{a}^\times)^k \\
 &= \sum_{l=0}^{\infty} \frac{\overbrace{(-1)^{2l}}^{=1} \phi^{2l}}{(2l)!} (\mathbf{a}^\times)^{2l} + \sum_{m=0}^{\infty} \frac{\overbrace{(-1)^{2m+1}}^{=-1} \phi^{2m+1}}{(2m+1)!} (\mathbf{a}^\times)^{2m+1} \\
 &= \sum_{l=0}^{\infty} \frac{\phi^{2l}}{(2l)!} (\overbrace{\mathbf{a}^\times \mathbf{a}^\times}^{=\mathbf{a} \mathbf{a}^\top - \mathbf{1}})^l - \sum_{m=0}^{\infty} \frac{\phi^{2m+1}}{(2m+1)!} \mathbf{a}^\times (\overbrace{\mathbf{a}^\times \mathbf{a}^\times}^{=\mathbf{a} \mathbf{a}^\top - \mathbf{1}})^m \\
 &= \mathbf{1} + \sum_{l=1}^{\infty} \frac{\phi^{2l}}{(2l)!} \overbrace{(\mathbf{a} \mathbf{a}^\top - \mathbf{1})^l}^{=(-1)^l (\mathbf{1} - \mathbf{a} \mathbf{a}^\top)} - \left(\phi \mathbf{a}^\times + \sum_{m=1}^{\infty} \frac{\phi^{2m+1}}{(2m+1)!} \mathbf{a}^\times \overbrace{(\mathbf{a} \mathbf{a}^\top - \mathbf{1})^m}^{=(-1)^m (\mathbf{1} - \mathbf{a} \mathbf{a}^\top)} \right) \\
 &= \mathbf{1} + \sum_{l=1}^{\infty} \frac{\overbrace{(-1)^l \phi^{2l}}^{=\cos \phi - 1}}{(2l)!} (\mathbf{1} - \mathbf{a} \mathbf{a}^\top) - \left(\phi \mathbf{a}^\times + \sum_{m=1}^{\infty} \frac{(-1)^m \phi^{2m+1}}{(2m+1)!} \overbrace{\mathbf{a}^\times (\mathbf{1} - \mathbf{a} \mathbf{a}^\top)}^{=\mathbf{a}^\times} \right) \\
 &= \mathbf{1} + (\cos \phi - 1) (\mathbf{1} - \mathbf{a} \mathbf{a}^\top) - \sum_{m=0}^{\infty} \frac{\overbrace{(-1)^m \phi^{2m+1}}{=\sin \phi}}{(2m+1)!} \mathbf{a}^\times \\
 &= \cos \phi \mathbf{1} + (1 - \cos \phi) \mathbf{a} \mathbf{a}^\top - \sin \phi \mathbf{a}^\times \quad \square
 \end{aligned}$$