

Assignment 3

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a) First, let's consider the reference frames \mathcal{F}_a and \mathcal{F}_b such that $\mathbf{C}_{ba} = \underline{\mathcal{F}}_{\rightarrow b} \cdot \underline{\mathcal{F}}_{\rightarrow a}^\top$. Moreover, let's consider the intermediate reference frames \mathcal{F}_k and \mathcal{F}_l such that:

$$\begin{aligned} \mathbf{C}_{ka} &= \underline{\mathcal{F}}_{\rightarrow k} \cdot \underline{\mathcal{F}}_{\rightarrow a}^\top = \mathbf{C}_2(\alpha), & \underline{a}^2 &= \underline{k}^2 \\ \mathbf{C}_{lk} &= \underline{\mathcal{F}}_{\rightarrow l} \cdot \underline{\mathcal{F}}_{\rightarrow k}^\top = \mathbf{C}_1(\beta), & \underline{k}^1 &= \underline{l}^1 \\ \mathbf{C}_{bl} &= \underline{\mathcal{F}}_{\rightarrow b} \cdot \underline{\mathcal{F}}_{\rightarrow l}^\top = \mathbf{C}_3(\gamma), & \underline{l}^3 &= \underline{b}^3 \end{aligned}$$

Now, from the definition of the Euler angles, we can find the Angular Velocities between these successive reference frames. In particular, we recall that the Angular Velocity physical vector between two reference frames, say \mathcal{F}_x and \mathcal{F}_y , is resolved identically in frames \mathcal{F}_x and \mathcal{F}_y . Plus, it is simply the multiplication between the angular time derivative and the unit physical vector (resolved in frame \mathcal{F}_x or \mathcal{F}_y) about which the rotation occurs.

In our case, \mathcal{F}_k rotates relative to \mathcal{F}_a about the physical vector \underline{a}^2 , at an angular time rate of change of $\dot{\alpha}$. Since \underline{a}^2 resolved in frame \mathcal{F}_a is simply $\mathbf{1}_2$, we have:

$$\omega_a^{ka} = \omega_k^{ka} = \dot{\alpha} \mathbf{1}_2 = \begin{bmatrix} 0 & \dot{\alpha} & 0 \end{bmatrix}^\top$$

Similarly for ω_k^{lk} and ω_l^{bl} :

$$\begin{aligned} \omega_k^{lk} &= \omega_l^{lk} = \dot{\beta} \mathbf{1}_1 = \begin{bmatrix} \dot{\beta} & 0 & 0 \end{bmatrix}^\top \\ \omega_l^{bl} &= \omega_b^{bl} = \dot{\gamma} \mathbf{1}_3 = \begin{bmatrix} 0 & 0 & \dot{\gamma} \end{bmatrix}^\top \end{aligned}$$

Using the facts that the Angular Velocity physical vectors add, it yields:

$$\underline{\omega}^{ba} = \underline{\omega}^{bl} + \underline{\omega}^{lk} + \underline{\omega}^{ka}$$

Therefore:

$$\begin{aligned}
 \underline{\mathcal{F}}_b^\top \omega_b^{ba} &= \underline{\mathcal{F}}_b^\top \omega_b^{bl} + \underline{\mathcal{F}}_l^\top \omega_l^{lk} + \underline{\mathcal{F}}_k^\top \omega_k^{ka} \\
 &= \underline{\mathcal{F}}_b^\top \omega_b^{bl} + \underline{\mathcal{F}}_b^\top \mathbf{C}_{bl} \omega_l^{lk} + \underline{\mathcal{F}}_b^\top \mathbf{C}_{bl} \mathbf{C}_{lk} \omega_k^{ka} \\
 &= \underline{\mathcal{F}}_b^\top \left(\omega_b^{bl} + \mathbf{C}_3(\gamma) \omega_l^{lk} + \mathbf{C}_3(\gamma) \mathbf{C}_1(\beta) \omega_k^{ka} \right) \\
 &= \underline{\mathcal{F}}_b^\top \underbrace{\left(\dot{\gamma} \mathbf{1}_3 + \mathbf{C}_3(\gamma) \dot{\beta} \mathbf{1}_1 + \mathbf{C}_3(\gamma) \mathbf{C}_1(\beta) \dot{\alpha} \mathbf{1}_2 \right)}_{\omega_b^{ba}}
 \end{aligned}$$

And finally:

$$\omega_b^{ba} = \underbrace{\begin{bmatrix} \mathbf{C}_3(\gamma) \mathbf{C}_1(\beta) \mathbf{1}_2 & \mathbf{C}_3(\gamma) \mathbf{1}_1 & \mathbf{1}_3 \end{bmatrix}}_{\mathbf{S}_b^{ba}(\gamma, \beta)} \underbrace{\begin{bmatrix} \dot{\alpha} \\ \dot{\beta} \\ \dot{\gamma} \end{bmatrix}}_{\dot{\boldsymbol{\theta}}}$$

Componentwise:

$$\begin{aligned}
 \mathbf{C}_3(\gamma) \mathbf{C}_1(\beta) \mathbf{1}_2 &= \begin{bmatrix} \cos(\gamma) & \sin(\gamma) & 0 \\ -\sin(\gamma) & \cos(\gamma) & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\beta) & \sin(\beta) \\ 0 & -\sin(\beta) & \cos(\beta) \end{bmatrix} \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix} \\
 &= \begin{bmatrix} \cos(\gamma) & \sin(\gamma) & 0 \\ -\sin(\gamma) & \cos(\gamma) & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ \cos(\beta) \\ -\sin(\beta) \end{bmatrix} \\
 &= \begin{bmatrix} \sin(\gamma) \cos(\beta) \\ \cos(\gamma) \cos(\beta) \\ -\sin(\beta) \end{bmatrix}
 \end{aligned}$$

$$\mathbf{C}_3(\gamma) \mathbf{1}_1 = \begin{bmatrix} \cos(\gamma) & \sin(\gamma) & 0 \\ -\sin(\gamma) & \cos(\gamma) & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} = \begin{bmatrix} \cos(\gamma) \\ -\sin(\gamma) \\ 0 \end{bmatrix}$$

$$\mathbf{S}_b^{ba}(\gamma, \beta) = \begin{bmatrix} \sin(\gamma) \cos(\beta) & \cos(\gamma) & 0 \\ \cos(\gamma) \cos(\beta) & -\sin(\gamma) & 0 \\ -\sin(\beta) & 0 & 1 \end{bmatrix}$$

b) $\mathbf{S}_b^{ba}(\gamma, \beta)$ is singular if and only if $\det(\mathbf{S}_b^{ba}(\gamma, \beta)) = 0$.

$$\det(\mathbf{S}_b^{ba}(\gamma, \beta)) = -\sin^2(\gamma) \cos(\beta) - \cos^2(\gamma) \cos(\beta) = -\cos(\beta)$$

Therefore, $\mathbf{S}_b^{ba}(\gamma, \beta)$ is singular for $\beta = \beta_k^* = \frac{\pi}{2} + k\pi$, $k \in \mathbb{Z}$.

Since $\mathbf{S}_b^{ba}(\gamma, \beta_k^*)$ is not invertible, the equation $\omega_b^{ba} = \mathbf{S}_b^{ba}(\gamma, \beta_k^*) \dot{\boldsymbol{\theta}}$ cannot be solved for $\dot{\boldsymbol{\theta}}$.

c) For $\beta = \beta_k^*$, we have:

$$\mathbf{S}_b^{ba}(\gamma, \beta_k^*) = \begin{bmatrix} 0 & \cos(\gamma) & 0 \\ 0 & -\sin(\gamma) & 0 \\ (-1)^{k+1} & 0 & 1 \end{bmatrix}$$

We want to find a 3×1 column matrix $\mathbf{n} = [n_1 \ n_2 \ n_3]^\top$ such that:

$$\mathbf{S}_b^{ba}(\gamma, \beta_k^*)\mathbf{n} = \mathbf{0} \Leftrightarrow \begin{bmatrix} 0 & \cos(\gamma) & 0 \\ 0 & -\sin(\gamma) & 0 \\ (-1)^{k+1} & 0 & 1 \end{bmatrix} \begin{bmatrix} n_1 \\ n_2 \\ n_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}, \quad \forall \gamma \in \mathbb{R}$$

By inspection, $n_1 = c$, $n_2 = 0$ and $n_3 = (-1)^k c$, $c \in \mathbb{R}$. Thus, the solutions for \mathbf{n} create a one-dimensional subspace of \mathbb{R}^3 which is spanned by $[1 \ 0 \ (-1)^k]^\top$. This subspace is the null space of $\mathbf{S}_b^{ba}(\gamma, \beta_k^*)$, also written as $\mathcal{N}(\mathbf{S}_b^{ba}(\gamma, \beta_k^*))$.

d) For $\dot{\boldsymbol{\theta}} \in \mathcal{N}(\mathbf{S}_b^{ba}(\gamma, \beta_k^*))$, $\boldsymbol{\omega}_b^{ba} = \mathbf{S}_b^{ba}(\gamma, \beta_k^*)\dot{\boldsymbol{\theta}} = \mathbf{0}$. Physically, this means that the reference frames \mathcal{F}_a and \mathcal{F}_b are not rotating relative to each other. This comes from the fact that the Euler angle rates cancel themselves.

In particular, for $k = 0$, we have $\beta_0^* = \frac{\pi}{2}$ and $\underline{a}^2 = -\underline{b}^3$. If additionally $\dot{\boldsymbol{\theta}}$ is in the null space of $\mathbf{S}_b^{ba}(\gamma, \beta_k^*)$, it yields $\dot{\alpha} = \dot{\gamma}$ and the rotation at the rate $\dot{\alpha}$ about \underline{a}^2 is cancelled by the rotation at the rate $\dot{\gamma} = \dot{\alpha}$ about $\underline{b}^3 = -\underline{a}^2$. Lastly, this cancellation will persist in time since $\dot{\beta} = 0$ ($\beta(t) = \beta_0^* = \frac{\pi}{2} \quad \forall t$). The same reasoning can be done for $k \neq 0$.

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$$\dot{\mathbf{a}} = \frac{1}{2} \left[\mathbf{a}^\times - \cot\left(\frac{\phi}{2}\right) \mathbf{a}^\times \mathbf{a}^\times \right] \boldsymbol{\omega}_b^{ba} \quad (1)$$

$$\dot{\phi} = \mathbf{a}^\top \boldsymbol{\omega}_b^{ba} \quad (2)$$

a) From the definition of the quaternion:

$$\boldsymbol{\epsilon} = \mathbf{a} \sin\left(\frac{\phi}{2}\right) \quad (3)$$

$$\eta = \cos\left(\frac{\phi}{2}\right) \quad (4)$$

Taking the time derivative of (3) and (4):

$$\dot{\boldsymbol{\epsilon}} = \dot{\mathbf{a}} \sin\left(\frac{\phi}{2}\right) + \mathbf{a} \frac{\dot{\phi}}{2} \cos\left(\frac{\phi}{2}\right) \quad (5)$$

$$\dot{\eta} = -\frac{\dot{\phi}}{2} \sin\left(\frac{\phi}{2}\right) \quad (6)$$

Substituting (1) and (2) into (5) and (6):

$$\begin{aligned} \dot{\epsilon} &= \left(\frac{1}{2} \left[\mathbf{a}^\times - \cot\left(\frac{\phi}{2}\right) \mathbf{a}^\times \mathbf{a}^\times \right] \omega_b^{ba}\right) \sin\left(\frac{\phi}{2}\right) + \frac{1}{2} \mathbf{a} \mathbf{a}^\top \omega_b^{ba} \cos\left(\frac{\phi}{2}\right) \\ &= \frac{1}{2} \left\{ \underbrace{-\cos\left(\frac{\phi}{2}\right) \mathbf{a}^\times \mathbf{a}^\times}_{\eta} + \underbrace{\mathbf{a} \mathbf{a}^\top \cos\left(\frac{\phi}{2}\right)}_{\eta} + \left[\underbrace{\mathbf{a} \sin\left(\frac{\phi}{2}\right)}_{\epsilon} \right]^\times \right\} \omega_b^{ba} \\ &= \frac{1}{2} \left\{ \eta \left[\mathbf{a} \mathbf{a}^\top - (\mathbf{a} \mathbf{a}^\top - \mathbf{1}) \right] + \epsilon^\times \right\} \omega_b^{ba} \\ &= \frac{1}{2} (\eta \mathbf{1} + \epsilon^\times) \omega_b^{ba} \quad \square \end{aligned}$$

$$\begin{aligned} \dot{\eta} &= -\frac{1}{2} \sin\left(\frac{\phi}{2}\right) \mathbf{a}^\top \omega_b^{ba} \\ &= -\frac{1}{2} \left[\underbrace{\sin\left(\frac{\phi}{2}\right) \mathbf{a}}_{\epsilon} \right]^\top \omega_b^{ba} \\ &= -\frac{1}{2} \epsilon^\top \omega_b^{ba} \quad \square \end{aligned}$$

Where we have used the following identities:

$$\begin{aligned} \mathbf{a}^\top \mathbf{a} &= 1 \\ \mathbf{a}^\times \mathbf{a}^\times &= \mathbf{a} \mathbf{a}^\top - \mathbf{a}^\top \mathbf{a} \mathbf{1} = \mathbf{a} \mathbf{a}^\top - \mathbf{1} \\ \cot(x) \sin(x) &= \cos(x) \end{aligned}$$

b)

$$\begin{aligned}
 \dot{\epsilon} &= \frac{1}{2}(\eta \mathbf{1} + \epsilon^\times) \omega_b^{ba} \\
 2(\eta \mathbf{1} + \epsilon^\times)^{-1} \dot{\epsilon} &= \omega_b^{ba} \\
 2(\eta \mathbf{1} + \eta^{-1} \epsilon \epsilon^\top - \epsilon^\times) \dot{\epsilon} &= \omega_b^{ba} \\
 2(\eta \mathbf{1} - \epsilon^\times) \dot{\epsilon} + 2\eta^{-1} \epsilon \epsilon^\top \dot{\epsilon} &= \omega_b^{ba} \\
 2(\eta \mathbf{1} - \epsilon^\times) \dot{\epsilon} + 2\eta^{-1} \epsilon \epsilon^\top \dot{\epsilon} - 2\eta^{-1} \underbrace{\epsilon (\epsilon^\top \dot{\epsilon} + \eta \dot{\eta})}_{=0} &= \omega_b^{ba} \\
 2(\eta \mathbf{1} - \epsilon^\times) \dot{\epsilon} + 2\eta^{-1} \epsilon \epsilon^\top \dot{\epsilon} - 2\eta^{-1} \epsilon \epsilon^\top \dot{\epsilon} - 2\eta^{-1} \eta \dot{\eta} \epsilon &= \omega_b^{ba} \\
 2(\eta \mathbf{1} - \epsilon^\times) \dot{\epsilon} - 2\dot{\eta} \epsilon &= \omega_b^{ba} \\
 2 \begin{bmatrix} (\eta \mathbf{1} - \epsilon^\times) & -\epsilon \end{bmatrix} \begin{bmatrix} \dot{\epsilon} \\ \dot{\eta} \end{bmatrix} &= \omega_b^{ba} \quad \square
 \end{aligned}$$

c) The equations for the quaternion rate doesn't contain any trigonometric function, making it easier to compute. More importantly, it can be shown the the quaternion rate does not suffer from any kinematic singularity, i.e. the equation

$$\begin{bmatrix} \dot{\epsilon} \\ \dot{\eta} \end{bmatrix} = \mathbf{\Gamma}_b^{ba}(\epsilon, \eta) \omega_b^{ba}$$

can always be solved for ω_b^{ba} . On the other hand, the equation for the rates of axis/angle parameters contains a kinematic singularity, which can be a problem.

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a) First, using the fact that angular velocity physical vectors add:

$$\underline{\omega}^{ba} = \underline{\omega}^{bq} + \underline{\omega}^{qa} \quad (7)$$

Moreover, from the definition of the problem:

$$\begin{aligned}
 \mathbf{C}_{qa} = \mathbf{C}_3(\alpha) &= \begin{bmatrix} \cos(\alpha) & \sin(\alpha) & 0 \\ -\sin(\alpha) & \cos(\alpha) & 0 \\ 0 & 0 & 1 \end{bmatrix} \\
 \mathbf{C}_{bq} = \mathbf{C}_1(\beta) &= \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\beta) & \sin(\beta) \\ 0 & -\sin(\beta) & \cos(\beta) \end{bmatrix}
 \end{aligned}$$

Resolving the angular velocity physical vectors in the different reference frames:

$$\underline{\omega}^{ba} = \mathcal{F}_b^\top \omega_b^{ba} \quad (8)$$

$$\underline{\omega}^{bq} = \underline{\mathcal{F}}_b^T \omega_b^{bq} \quad (9)$$

$$\underline{\omega}^{qa} = \underline{\mathcal{F}}_q^T \omega_q^{qa} = \underline{\mathcal{F}}_b^T \mathbf{C}_{bq} \omega_q^{qa} \quad (10)$$

Substituting (8),(9) and (10) into (7), and factoring $\underline{\mathcal{F}}_b^T$:

$$\underline{\omega}^{ba} = \underline{\mathcal{F}}_b^T \omega_b^{ba} = \underline{\mathcal{F}}_b^T \underbrace{\left(\omega_b^{bq} + \mathbf{C}_{bq} \omega_q^{qa} \right)}_{\omega_b^{ba}} \quad (11)$$

By inspection:

$$\omega_a^{qa} = \omega_q^{qa} = \dot{\alpha} \mathbf{1}_3 \quad (12)$$

$$\omega_b^{bq} = \omega_q^{bq} = \dot{\beta} \mathbf{1}_1 \quad (13)$$

Lastly, substituting (12) and (13) into (11):

$$\begin{aligned} \underline{\omega}^{ba} &= \underline{\mathcal{F}}_b^T \left(\dot{\beta} \mathbf{1}_1 + \mathbf{C}_{bq} (\dot{\alpha} \mathbf{1}_3) \right) \\ &= \underline{\mathcal{F}}_b^T \left(\begin{bmatrix} \dot{\beta} \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\beta) & \sin(\beta) \\ 0 & -\sin(\beta) & \cos(\beta) \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ \dot{\alpha} \end{bmatrix} \right) \\ &= \underline{\mathcal{F}}_b^T \underbrace{\begin{bmatrix} \dot{\beta} \\ \sin(\beta) \dot{\alpha} \\ \cos(\beta) \dot{\alpha} \end{bmatrix}}_{\omega_b^{ba}} \\ &= \begin{bmatrix} \dot{\beta} & \sin(\beta) \dot{\alpha} & \cos(\beta) \dot{\alpha} \end{bmatrix} \underline{\mathcal{F}}_b \quad \square \end{aligned}$$

b)

$$\begin{aligned} r_{\rightarrow}^{pw \cdot a} &= r_{\rightarrow}^{pz \cdot a} + r_{\rightarrow}^{zw \cdot a} \\ &= r_{\rightarrow}^{zw \cdot a} + r_{\rightarrow}^{pz \cdot b} + \underline{\omega}^{ba} \times r_{\rightarrow}^{pz} \\ &= \underline{\mathcal{F}}_a^T \dot{\mathbf{r}}_a^{zw} + \underline{\mathcal{F}}_b^T \dot{\mathbf{r}}_b^{pz} + \underline{\mathcal{F}}_b^T \omega_b^{ba \times} \mathbf{r}_b^{pz} \\ &= \underline{\mathcal{F}}_b^T \mathbf{C}_{ba} \dot{\mathbf{r}}_a^{zw} + \underline{\mathcal{F}}_b^T \dot{\mathbf{r}}_b^{pz} + \underline{\mathcal{F}}_b^T \omega_b^{ba \times} \mathbf{r}_b^{pz} \\ &= \underline{\mathcal{F}}_b^T \underbrace{\left(\mathbf{C}_{ba} \dot{\mathbf{r}}_a^{zw} + \dot{\mathbf{r}}_b^{pz} + \omega_b^{ba \times} \mathbf{r}_b^{pz} \right)}_{\mathbf{v}_b^{pw/a}} \end{aligned}$$

Where we have used the Transport Theorem.

By inspection and from a):

$$\begin{aligned}
\mathbf{C}_{ba} &= \mathbf{C}_3(\alpha)\mathbf{C}_1(\beta) \\
&= \begin{bmatrix} \cos(\alpha) & \sin(\alpha) & 0 \\ -\sin(\alpha) & \cos(\alpha) & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\beta) & \sin(\beta) \\ 0 & -\sin(\beta) & \cos(\beta) \end{bmatrix} \\
&= \begin{bmatrix} \cos(\alpha) & \sin(\alpha)\cos(\beta) & \sin(\alpha)\sin(\beta) \\ -\sin(\alpha) & \cos(\alpha)\cos(\beta) & \cos(\alpha)\sin(\beta) \\ 0 & -\sin(\beta) & \cos(\beta) \end{bmatrix} \\
\\
\mathbf{r}_a^{zw} &= \begin{bmatrix} 0 \\ 0 \\ -l^{zw} \end{bmatrix}, \quad \dot{\mathbf{r}}_a^{zw} = \begin{bmatrix} 0 \\ 0 \\ -\dot{l}^{zw} \end{bmatrix} \\
\mathbf{r}_b^{pz} &= \begin{bmatrix} 0 \\ l^{pz} \\ 0 \end{bmatrix}, \quad \dot{\mathbf{r}}_b^{pz} = \begin{bmatrix} 0 \\ \dot{l}^{pz} \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \quad (p \text{ fixed}) \\
\\
\omega_b^{ba \times} \mathbf{r}_b^{pz} &= \begin{bmatrix} \dot{\beta} \\ \sin(\beta)\dot{\alpha} \\ \cos(\beta)\dot{\alpha} \end{bmatrix}^\times \begin{bmatrix} 0 \\ l^{pz} \\ 0 \end{bmatrix} \\
&= \begin{bmatrix} 0 & -\cos(\beta)\dot{\alpha} & \sin(\beta)\dot{\alpha} \\ \cos(\beta)\dot{\alpha} & 0 & -\dot{\beta} \\ -\sin(\beta)\dot{\alpha} & \dot{\beta} & 0 \end{bmatrix} \begin{bmatrix} 0 \\ l^{pz} \\ 0 \end{bmatrix} \\
&= \begin{bmatrix} -l^{pz}\cos(\beta)\dot{\alpha} \\ 0 \\ l^{pz}\dot{\beta} \end{bmatrix} \\
\\
\mathbf{C}_{ba}\dot{\mathbf{r}}_a^{zw} &= \begin{bmatrix} \cos(\alpha) & \sin(\alpha)\cos(\beta) & \sin(\alpha)\sin(\beta) \\ -\sin(\alpha) & \cos(\alpha)\cos(\beta) & \cos(\alpha)\sin(\beta) \\ 0 & -\sin(\beta) & \cos(\beta) \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ -\dot{l}^{zw} \end{bmatrix} \\
&= \begin{bmatrix} -\dot{l}^{zw}\sin(\alpha)\sin(\beta) \\ -\dot{l}^{zw}\cos(\alpha)\sin(\beta) \\ -\dot{l}^{zw}\cos(\beta) \end{bmatrix}
\end{aligned}$$

And finally, substituting these identities into the expression for $\mathbf{v}_b^{pw/a}$:

$$\mathbf{v}_b^{pw/a} = \begin{bmatrix} -\dot{l}^{zw}\sin(\alpha)\sin(\beta) - l^{pz}\cos(\beta)\dot{\alpha} \\ -\dot{l}^{zw}\cos(\alpha)\sin(\beta) \\ -\dot{l}^{zw}\cos(\beta) + l^{pz}\dot{\beta} \end{bmatrix}$$

(As far as I understand the problem, l^{zw} and β are not constant in time but l^{pz} is.)