

Chapter 4. Electric Fields in Matter

4 Electric Fields in Matter

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4.3 The Electric Displacement

4.3.1 Gauss's Law in the Presence of Dielectrics

The effect of polarization (\mathbf{P}) is to produce accumulations of bound charge:

$$\rho_b \equiv -\nabla \cdot \mathbf{P} \quad : \text{within the dielectric} \quad \sigma_b \equiv \mathbf{P} \cdot \hat{\mathbf{n}} \quad : \text{on the surface}$$


If any free charge exists within a dielectric, therefore, the total charge density can be written as

$$\rho = \rho_b + \rho_f$$

→ Gauss's law reads

$$\epsilon_0 \nabla \cdot \mathbf{E} = \rho = \rho_b + \rho_f = -\nabla \cdot \mathbf{P} + \rho_f \longrightarrow \nabla \cdot (\epsilon_0 \mathbf{E} + \mathbf{P}) = \rho_f$$

$$\mathbf{D} \equiv \epsilon_0 \mathbf{E} + \mathbf{P} \quad (C / m^2) \quad \rightarrow \text{Electric displacement}$$


$$\begin{aligned} \nabla \cdot \mathbf{D} &= \rho_f \\ \oint \mathbf{D} \cdot d\mathbf{a} &= Q_{fenc} \end{aligned}$$

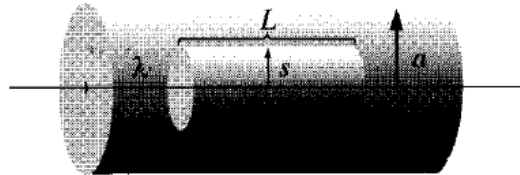
Similar to Gauss's law

Q_{fenc} denotes the total free charge enclosed in the volume.

**In a typical problem, we know ρ_f , but we do not (initially) know ρ_b .
Whenever the requisite symmetry is present, we can immediately calculate D by the Gauss's law method, because it makes reference only to free charges.**

Gauss's Law: $\nabla \cdot \mathbf{D} = \rho_f$ $\oint \mathbf{D} \cdot d\mathbf{a} = Q_{f\text{enc}}$

Example 4.4 **A long straight wire, carrying uniform line charge λ , is surrounded by rubber insulation out to a radius a . Find the electric displacement \mathbf{D} and the field \mathbf{E} .**



→ Drawing a cylindrical Gaussian surface, of radius s and length L ,

$$D(2\pi sL) = \lambda L \quad \longrightarrow \quad \mathbf{D} = \frac{\lambda}{2\pi s} \hat{\mathbf{s}}.$$

: both within the insulation and outside it.

→ *Inside* the rubber the electric field cannot be determined, since we do not know P .

→ *Outside* it, $P = 0$: $\mathbf{E} = \frac{1}{\epsilon_0} \mathbf{D} = \frac{\lambda}{2\pi \epsilon_0 s} \hat{\mathbf{s}}, \quad \text{for } s > a.$

Gauss's Law: $\nabla \cdot \mathbf{D} = \rho_f \quad \oint \mathbf{D} \cdot d\mathbf{a} = Q_{f_{\text{enc}}}$

Problem 4.15 A thick **spherical shell** (inner radius **a**, outer radius **b**) is made of dielectric material with a "frozen-in" polarization **P(r)**. (No free charge inside)

Find the electric field in all three regions by **two different methods**:

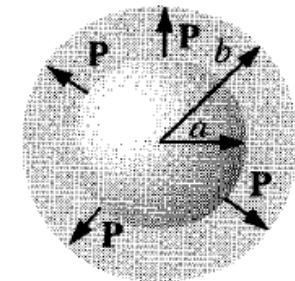
$$\mathbf{P}(\mathbf{r}) = \frac{k}{r} \hat{\mathbf{r}}$$

(a) Identify all the bound charge, and use Gauss's law to calculate the field it produces.

$$\rho_b = -\nabla \cdot \mathbf{P} = -\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{k}{r} \right) = -\frac{k}{r^2};$$

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = \begin{cases} +\mathbf{P} \cdot \hat{\mathbf{r}} = k/b & (\text{at } r = b), \\ -\mathbf{P} \cdot \hat{\mathbf{r}} = -k/a & (\text{at } r = a). \end{cases}$$

$$\mathbf{P}(\mathbf{r}) = \frac{k}{r} \hat{\mathbf{r}}$$



For $r < a$, $Q_{\text{enc}} = 0$, so Gauss's law $\Rightarrow \mathbf{E} = 0$.

$$\text{For } r > b, \quad Q_{\text{tot}} = \oint_S \sigma_b da + \int_V \rho_b d\tau = \oint_S \mathbf{P} \cdot d\mathbf{a} - \int_V \nabla \cdot \mathbf{P} d\tau$$

(divergence theorem)

$$\oint_S \mathbf{P} \cdot d\mathbf{a} = \int_V \nabla \cdot \mathbf{P} d\tau, \text{ so } Q_{\text{enc}} = 0 \quad \text{Gauss's law} \Rightarrow \mathbf{E} = 0$$

$$\text{For } a < r < b, \quad Q_{\text{enc}} = \left(\frac{-k}{a} \right) (4\pi a^2) + \int_a^r \left(\frac{-k}{\bar{r}^2} \right) 4\pi \bar{r}^2 d\bar{r} = -4\pi k a - 4\pi k (r - a) = -4\pi k r; \text{ so}$$

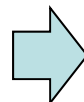
$$\mathbf{E} = -(k/\epsilon_0 r) \hat{\mathbf{r}}.$$

(b) Use the equations of electric displacement **D: $\mathbf{D} \equiv \epsilon_0 \mathbf{E} + \mathbf{P}$**

$$\oint \mathbf{D} \cdot d\mathbf{a} = Q_{f_{\text{enc}}} = 0 \Rightarrow \mathbf{D} = 0 \text{ everywhere. } \mathbf{E} = (-1/\epsilon_0) \mathbf{P}, \text{ so}$$

$$\mathbf{E} = 0 \text{ (for } r < a \text{ and } r > b);$$

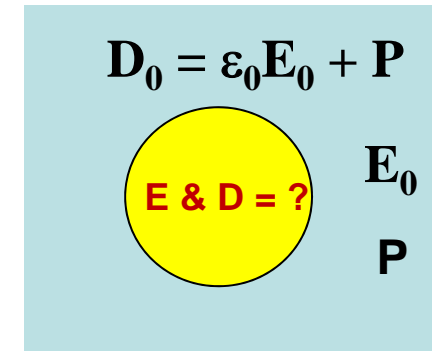
$$\mathbf{E} = -(k/\epsilon_0 r) \hat{\mathbf{r}} \text{ (for } a < r < b)$$



*Much simple to use **D**.*

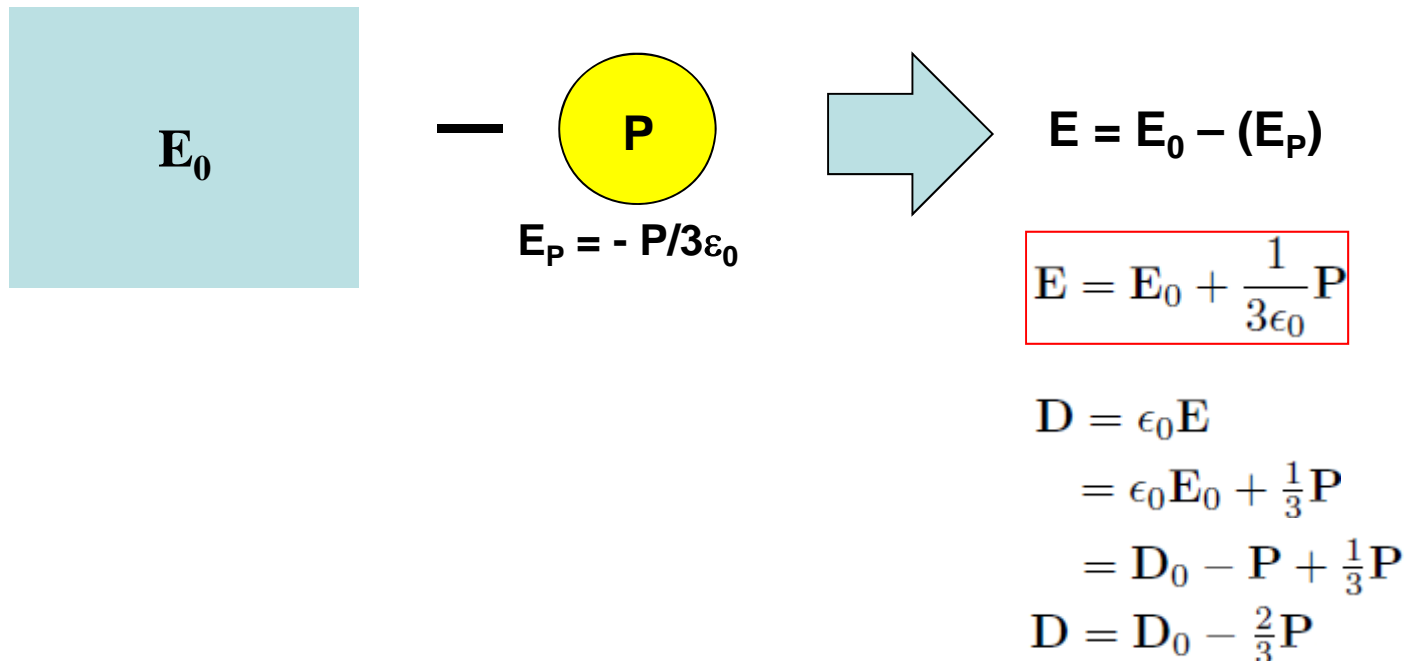
Gauss's Law: $\nabla \cdot \mathbf{D} = \rho_f$ $\oint \mathbf{D} \cdot d\mathbf{a} = Q_{f\text{enc}}$

Problem 4.16 Suppose the field inside a large piece of dielectric is \mathbf{E}_0 , so that the electric displacement is: $\mathbf{D}_0 = \epsilon_0 \mathbf{E}_0 + \mathbf{P}$. Now a small spherical cavity is hollowed out of the material. Find the field at the center of the cavity in terms of \mathbf{E}_0 and \mathbf{P} . Also find the displacement at the center of the cavity in terms of \mathbf{D}_0 and \mathbf{P} .



Assume the cavities are small enough that \mathbf{P} , \mathbf{E}_0 , and \mathbf{D}_0 are essentially uniform.

Hint → Carving out a cavity is the same as superimposing an object of the same shape but opposite polarization.



4.3.2 A Deceptive Parallel: Misleading in comparison between E and D

$$\nabla \cdot \mathbf{D} = \rho_f \quad \xleftarrow{\text{Let's compare it with Gauss's law}} \quad \epsilon_0 \nabla \cdot \mathbf{E} = \rho$$

It looks just like Gauss's law,

- Only the *total* charge density ρ is replaced by the *free* charge density ρ_f , \mathbf{D} is substituted for $\epsilon_0 \mathbf{E}$.
- \mathbf{D} is "just like" \mathbf{E} (apart from the factor ϵ_0), except that its source is ρ_f instead of ρ .

Therefore, one may conclude that

"to solve problems involving dielectrics, you just forget all about the bound charge, calculate the field as you ordinarily would, only call the answer \mathbf{D} instead of \mathbf{E} ."

→ This conclusion in similarity between \mathbf{E} and \mathbf{D} is false!

There is **no "Coulomb's law" for \mathbf{D}** : $\mathbf{D}(\mathbf{r}) \neq \frac{1}{4\pi} \int \frac{\hat{\mathbf{r}}}{r^2} \rho_f(\mathbf{r}') d\tau'$

For the divergence alone is insufficient to determine a vector field; you need to know the curl as well.

In the case of electrostatic fields, **the curl of \mathbf{E} is always zero**. $\nabla \times \mathbf{E} = 0$

→ But **the curl of \mathbf{D} is not always zero**. $\nabla \times \mathbf{D} = \epsilon_0(\nabla \times \mathbf{E}) + (\nabla \times \mathbf{P}) = \nabla \times \mathbf{P}$

→ Because $\nabla \times \mathbf{D} \neq 0$, \mathbf{D} cannot be expressed as the gradient of a scalar.
→ **there is no "potential" for \mathbf{D} .**

Advice: When you are asked to compute the electric displacement, **first look for symmetry**.

→ If the problem exhibits symmetry, then you can **get \mathbf{D} directly by the usual Gauss's law methods**.

→ If symmetry is absent, you'll have to think of another approach.
you **must not assume that \mathbf{D} is determined exclusively by the free charge**.

4.3.3 Boundary Conditions

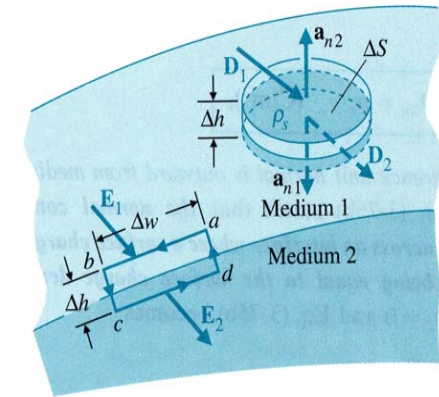
$$\nabla \cdot \mathbf{D} = \rho_f \rightarrow \oint_S \mathbf{D} \cdot d\mathbf{a} = (\mathbf{D}_1 \cdot \mathbf{a}_{n2} + \mathbf{D}_2 \cdot \mathbf{a}_{n1}) \Delta S = \mathbf{a}_{n2} \cdot (\mathbf{D}_1 - \mathbf{D}_2) \Delta S = (D_{1n} - D_{2n}) \Delta S$$

$$\rightarrow \int_V \rho_f dV = (\rho_f \Delta h) \Delta S$$

$$\Rightarrow D_{1n} - D_{2n} = \sigma_f \quad \sigma_f = \lim_{\Delta h \rightarrow 0} \rho_f \Delta h$$

$$\nabla \times \mathbf{D} = \nabla \times \mathbf{P} \rightarrow \oint_C \mathbf{D} \cdot d\mathbf{l} = (D_1 \cdot \mathbf{l}_2 + D_2 \cdot \mathbf{l}_1) = \mathbf{l}_2 \cdot (\mathbf{D}_1 - \mathbf{D}_2) = (D_{1t} - D_{2t})$$

$$\Rightarrow D_{1t} - D_{2t} = P_{1t} - P_{2t}$$



Note for E:

$$E_{1n} - E_{2n} = \frac{\sigma_{enclosed}}{\epsilon_0} = \frac{\sigma_f + \sigma_b}{\epsilon_0}$$

$$E_{1t} - E_{2t} = 0$$

Summary of Boundary conditions

$$\mathbf{E}_{above}^{\parallel} - \mathbf{E}_{below}^{\parallel} = 0 \quad \mathbf{D}_{above}^{\parallel} - \mathbf{D}_{below}^{\parallel} = \mathbf{P}_{above}^{\parallel} - \mathbf{P}_{below}^{\parallel}$$

$$E_{above}^{\perp} - E_{below}^{\perp} = \frac{1}{\epsilon_0} \sigma$$

$$D_{above}^{\perp} - D_{below}^{\perp} = \sigma_f$$

$$(\sigma = \sigma_f + \sigma_b)$$