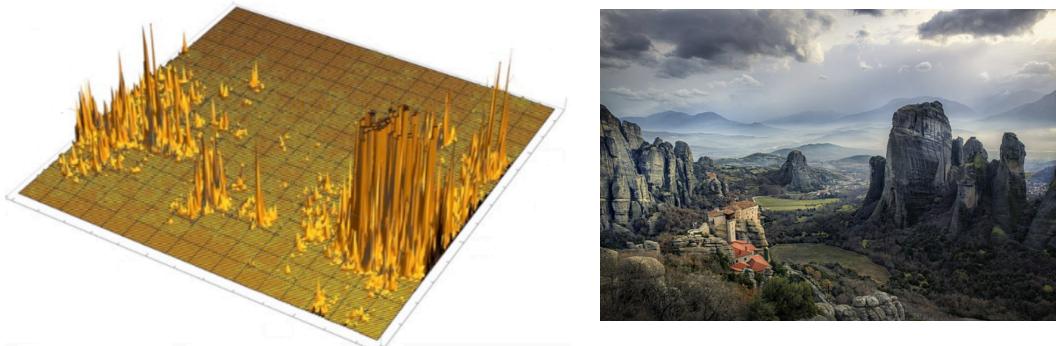


THE CRITICAL 2d STOCHASTIC HEAT FLOW AND RELATED MODELS

FRANCESCO CARAVENNA, RONGFENG SUN, AND NIKOS ZYGOURAS

ABSTRACT. In these lecture notes, we review recent progress in the study of the stochastic heat equation and its discrete analogue, the directed polymer model, in spatial dimension 2. It was discovered that a phase transition emerges on an intermediate disorder scale, with Edwards-Wilkinson (Gaussian) fluctuations in the sub-critical regime. In the critical window, a unique scaling limit has been identified and named the *critical 2d stochastic heat flow*. This gives a meaning to the solution of the stochastic heat equation in the critical dimension 2, which lies beyond existing solution theories for singular SPDEs. We outline the proof ideas, introduce the key ingredients, and discuss related literature on disordered systems and singular SPDEs. A list of open questions is also provided.



A picture of the critical 2d stochastic heat flow is shown on the left, and on the right, a resembling natural landscape located in central Greece and named Meteora (photo by Stathis Floros, CC BY-SA 4.0). The picture of the critical 2d stochastic heat flow has been obtained by simulating the partition function of the directed polymer model. The plateaus in the simulation are a result of a truncation of high peaks.

CONTENTS

1. INTRODUCTION	3
1.1. Stochastic heat equation	4
1.2. The KPZ equation	5
1.3. Directed polymer model	5
1.4. Phase transition on an intermediate disorder scale	6

Date: July 7, 2025.

2010 *Mathematics Subject Classification.* Primary: 82B44; Secondary: 35R60, 60H15, 82D60.

Key words and phrases. Coarse-Graining, Directed Polymer in Random Environment, Disordered Systems, KPZ Equation, Lindeberg Principle, Renormalization, Singular Stochastic Partial Differential Equation, Stochastic Heat Equation, Stochastic Heat Flow.

1.5. Outline	7
2. BASIC TOOLS AND PRELIMINARY RESULTS	8
2.1. Polynomial chaos expansion	8
2.2. A transition in the second moment	9
3. GAUSSIAN FLUCTUATIONS IN THE SUBCRITICAL REGIME	11
3.1. One point fluctuation	11
3.2. Random field fluctuation	14
4. THE 2d KPZ IN THE SUBCRITICAL REGIME	15
5. RELATED RESULTS OUTSIDE THE CRITICAL REGIME	18
5.1. Other marginally relevant directed polymer models	18
5.2. Stochastic Heat Equation	19
5.3. Alternative approach to subcritical 2d KPZ	19
5.4. Nonlinear stochastic heat equations	20
5.5. Anisotropic KPZ and related models	21
5.6. The quasi-critical regime	22
5.7. Higher dimensions	23
5.8. Space-time mollifications of noise	23
6. THE CRITICAL 2d SHF	24
6.1. Main result	24
6.2. Universality via a Lindeberg principle	26
6.3. Proof outline of Theorem 6.1	29
7. COARSE GRAINING AND LINDEBERG	30
7.1. Coarse graining	30
7.2. Applying Lindeberg to the coarse-grained model	36
8. SECOND MOMENT CALCULATIONS AND DICKMAN SUBORDINATOR	37
8.1. Renewal representation	37
8.2. Dickman subordinator	39
8.3. Asymptotics	41
8.4. Covariance kernel of the critical 2d SHF	44
9. HIGHER MOMENT BOUNDS	44
9.1. Moment bounds and proof strategy	45
9.2. Functional inequalities	51
9.3. Moment bounds for the coarse-grained model	52
9.4. Moments of the Critical 2d SHF	53
9.5. Related literature	54
10. FURTHER PROPERTIES OF CRITICAL 2d SHF	57
10.1. Shift invariance and scaling covariance	57
10.2. Flow property	57
10.3. Characterisation of the critical 2d SHF	59
10.4. Non-GMCness of fixed time marginals	59
10.5. Singularity and Regularity	61
11. DISCUSSIONS AND OPEN QUESTIONS	62
11.1. Disordered systems and singular SPDEs	62
11.2. More open questions	63
ACKNOWLEDGEMENTS	66
REFERENCES	66

1. INTRODUCTION

Our goal here is to review recent progress in making sense of the two-dimensional *stochastic heat equation* (SHE) [CSZ17b, CSZ19b, GQT21, CSZ23a, CSZ23b, Tsa24]

$$\partial_t u(t, x) = \frac{1}{2} \Delta u(t, x) + \beta \xi(t, x) u(t, x), \quad t > 0, x \in \mathbb{R}^2. \quad (\text{SHE})$$

This is a *singular stochastic partial differential equation* (SPDE) due to the presence of the term $\xi(t, x)u(t, x)$, where the potential (or disorder) $\xi(t, x)$ is a space-time white noise, i.e., a generalised Gaussian field with mean 0 and covariance $\mathbb{E}[\xi(t, x)\xi(s, y)] = \delta(t-s)\delta(x-y)$, and $u(t, x)$ is also expected to be a generalised function (i.e. distribution in the sense of Schwartz). The constant $\beta \geq 0$ tunes the strength of the interaction.

Via the Cole-Hopf transformation $h = \log u$, SHE is also related to the celebrated Kardar-Parisi-Zhang (KPZ) equation

$$\partial_t h(t, x) = \frac{1}{2} \Delta h(t, x) + \frac{1}{2} |\nabla h(t, x)|^2 + \beta \xi(t, x), \quad t > 0, x \in \mathbb{R}^d, \quad (\text{KPZ})$$

which is a model for random interface growth and has been studied extensively in $d = 1$ as a canonical example in the KPZ universality class [QS15, Cor12, Cor16, Zyg22].

In recent years, there have been fundamental breakthroughs in the study of singular SPDEs [Hai13, Hai14, GIP15, Kup16, GJ14], including the KPZ equation in $d = 1$. However, these theories only apply to sub-critical singular SPDEs, while dimension $d = 2$ is critical for SHE and KPZ. In a series of papers culminating in [CSZ23a], we were able to give a meaning to the solution of the 2d SHE on an *intermediate disorder scale* (i.e. by regularising the noise ξ and, simultaneously, *rescaling the disorder strength* $\beta \downarrow 0$ at a suitable rate, as the regularisation is removed). In particular, it was shown that a phase transition exists on this scale, with Edwards-Wilkinson (Gaussian) fluctuation in the sub-critical regime [CSZ17b]. For the 2d KPZ, the same Edwards-Wilkinson fluctuation in the subcritical regime was established in [CSZ20, Gu20]. Most interestingly, it was shown in [CSZ23a] that in a window around the critical point, we can make sense of the solution of the 2d SHE as a random measure-valued process, which is called the *critical 2d stochastic heat flow* (SHF). This is a rare example where a model in the critical dimension and at the phase transition point admits a non-Gaussian limit, in contrast to e.g. the Ising and Φ^4 models at the critical dimension $d = 4$ [ADC21].

Our motivation to make sense of the 2d SHE came from a different direction, which is the study of continuum limits of disordered systems (see discussions in Section 11.1), in particular, the *directed polymer model* (DPM). The DPM is one of the simplest and yet most challenging disordered system, which models a random walk interacting in equilibrium with a random environment (disorder). As the strength of the disorder (inverse temperature β) increases, the model undergoes a phase transition from diffusive behaviour for the random walk, with path delocalisation, to expected super-diffusive behaviour and path localisation (see [Com17, Zyg24] for more details). The DPM lies at the heart of two different areas of intense research in recent years. On the one hand, it is a canonical example in the KPZ universality class of random interface growth models, with the 0-temperature DPM in $d = 1$ being exactly the *last passage percolation* model. On the other hand, it provides a discretisation of the SHE (and the KPZ equation via the Cole-Hopf transformation), which are classical examples of singular SPDEs. It is this latter connection that we will exploit to make sense of the 2d SHE.

In the rest of the introduction, we will first explain heuristically why $d = 2$ is critical for SHE and what is the standard procedure to define the solution of a singular SPDE such as

the $2d$ SHE. We will then introduce the directed polymer model, its connection to SHE, and the intermediate disorder scale on which a phase transition exists and our analysis will be carried out. We will then conclude the introduction with an outline of the rest of the article.

1.1. STOCHASTIC HEAT EQUATION. We start with a heuristic scaling argument that shows why $d = 2$ is critical for the stochastic heat equation (SHE)

$$\partial_t u(t, x) = \frac{1}{2} \Delta u(t, x) + \beta \xi(t, x) u(t, x), \quad t > 0, \quad x \in \mathbb{R}^d. \quad (1.1)$$

In the spirit of renormalisation group theory [Kup16], we note that the space-time rescaled solution $\tilde{u}(t, x) := u(\varepsilon^2 t, \varepsilon x)$ formally solves

$$\partial_t \tilde{u} = \frac{1}{2} \Delta \tilde{u} + \beta \varepsilon^{1-\frac{d}{2}} \tilde{\xi} \tilde{u}, \quad \tilde{t} > 0, \quad \tilde{x} \in \mathbb{R}^d, \quad (1.2)$$

where $\tilde{\xi}(t, x) := \varepsilon^{1+\frac{d}{2}} \xi(\varepsilon^2 t, \varepsilon x)$ is a new space-time white noise with the same distribution as ξ . We now see that as we send $\varepsilon \downarrow 0$ and zoom into smaller and smaller space-time scales, the strength of the noise vanishes in dimensions $d < 2$ and diverges in dimensions $d > 2$. But in the critical dimension $d = 2$, the exponent $1 - \frac{d}{2}$ vanishes, and how scaling affects the effective strength of the noise depends on finer details of the model. In the language of singular SPDEs [Hai13, Hai14, GIP15], the regimes $d < 2$, $d = 2$ and $d > 2$ are called respectively *sub-critical*, *critical*, and *super-critical*.

The main difficulty in making sense of the solution of singular SPDEs such as (SHE) is the roughness on small scales of the solution (expected to be a generalised function), which makes terms such as $\xi(t, x)u(t, x)$ mathematically undefined. The standard procedure is to perform a regularisation, also called *ultraviolet cutoff*, to smoothen things below a small spatial scale ε , which leads to an approximate solution u^ε . We then check whether there exist suitable centering and scaling constants A_ε and B_ε , and suitable choice of the model parameter $\beta = \beta_\varepsilon$, such that $(u^\varepsilon - A_\varepsilon)/B_\varepsilon$ converges to a non-trivial limit as $\varepsilon \downarrow 0$. If such a non-trivial limit exists, then we can interpret it as the solution of (a renormalised version of) this singular SPDE.

There are several different ways of performing the ultraviolet cutoff. The most common approach is to mollify the space-time white noise ξ on the spatial scale ε . More precisely, let $j \in C_c^\infty(\mathbb{R}^2)$ be a smooth probability density function with compact support, and define the mollified noise $\xi^\varepsilon := j^\varepsilon * \xi$ by convolving ξ in space with $j^\varepsilon(x) := \varepsilon^{-2} j(\frac{x}{\varepsilon})$. This leads to the mollified SHE

$$\partial_t u^\varepsilon = \frac{1}{2} \Delta u^\varepsilon + \beta_\varepsilon u^\varepsilon \xi^\varepsilon, \quad u^\varepsilon(0, \cdot) \equiv 1, \quad (1.3)$$

which admits a classical Itô solution. Furthermore, the solution admits a Feynman-Kac representation (see [BC95, Section 3]) with

$$u^\varepsilon(t, x) = \mathbb{E} \left[e^{\int_0^t (\beta_\varepsilon \xi^\varepsilon(t-s, B_s) - \lambda_\varepsilon) ds} \middle| B_0 = x \right], \quad (1.4)$$

where $\mathbb{E}[\cdot]$ is w.r.t. a standard Brownian motion B in \mathbb{R}^2 , and $\lambda_\varepsilon := \frac{\beta_\varepsilon^2}{2} \text{Var}(\xi^\varepsilon(s, x)) = \frac{\beta_\varepsilon^2}{2\varepsilon^2} \|j\|_2^2$ ensures that $\mathbb{E}[u^\varepsilon(t, x)] = 1$.

Another approach to smoothen things on small scales is to discretize space and define an approximate solution on a grid with lattice spacing ε . This is the approach we will follow, where the approximate solution turns out to coincide with the partition functions of the directed polymer model. Yet another approach is to truncate the high frequency modes in the Fourier decomposition of $u(t, \cdot)$, see e.g. [CT24] and the references therein. It is expected

that modulo a change of parameters, the different procedures of ultraviolet cutoff will lead to the same limits, which are believed to be universal. This has been shown for the 2d SHE (and KPZ) in the subcritical regime in [CSZ17b], and in the critical window in [CSZ23a, Tsa24].

1.2. THE KPZ EQUATION. The KPZ equation (KPZ) in the critical dimension $d = 2$ is also singular, due to the presence of the term $|\nabla h|^2$. To make sense of its solution, we also need to perform an ultraviolet cutoff and then pass to the limit. In particular, if ξ^ε denotes the same spatially mollified noise as in (1.3), then we consider the mollified KPZ equation

$$\partial_t h^\varepsilon = \frac{1}{2} \Delta h^\varepsilon + \frac{1}{2} |\nabla h^\varepsilon|^2 + \beta_\varepsilon \xi^\varepsilon - C_\varepsilon, \quad h^\varepsilon(0, \cdot) \equiv 0, \quad (1.5)$$

where $C_\varepsilon := \beta_\varepsilon^2 \varepsilon^{-2} \|j\|_2^2$ ensures that h^ε is related to the solution of the mollified SHE (1.3) through the Cole-Hopf transformation $h^\varepsilon(t, x) = \log u^\varepsilon(t, x)$, see e.g. [CSZ20].

One could also consider the KPZ equation with different coefficients in front of the Laplacian and the non-linearity, that is,

$$\partial_t \tilde{h}^\varepsilon = \frac{\nu}{2} \Delta \tilde{h}^\varepsilon + \frac{\lambda}{2} |\nabla \tilde{h}^\varepsilon|^2 + D \xi^\varepsilon. \quad (1.6)$$

For example, we could set $\nu = D = 1$ and $\lambda = \beta_\varepsilon$ as in [CD20]. However, by scaling (see e.g. [CSZ20, Appendix A]), the three parameters ν, λ and D can be reduced to a single *effective coupling constant* $\beta^2 := \lambda^2 D / \nu^3$. Therefore it suffices to consider (1.5) with only the noise strength β_ε dependent on ε .

1.3. DIRECTED POLYMER MODEL. The directed polymer model (DPM) is defined as a random walk interacting with a random environment (disorder) [Com17, Zyg24]. To avoid complications caused by periodicity, instead of considering the simple symmetric random walk on \mathbb{Z}^2 as done in [CSZ23a], we consider an irreducible aperiodic random walk $S = (S_n)_{n \geq 0}$ whose increment $\xi := S_1 - S_0$ has mean 0, covariance matrix being the identity matrix, and $\mathbb{E}[e^{b|\xi|}] < \infty$ for some $b > 0$. Let P and E denote probability and expectation for S , and denote its transition probability kernels by

$$q_n(z) := P(S_n = z | S_0 = 0), \quad z \in \mathbb{Z}^2, n \in \mathbb{N}_0. \quad (1.7)$$

The disorder (random environment) is given by a family of i.i.d. random variables $\omega := (\omega(n, z))_{n \in \mathbb{N}, z \in \mathbb{Z}^2}$ with

$$\begin{aligned} \mathbb{E}[\omega] &= 0, & \mathbb{E}[\omega^2] &= 1, \\ \exists \beta_0 > 0 \text{ such that } \lambda(\beta) &:= \log \mathbb{E}[e^{\beta \omega}] < \infty & \forall \beta \in [0, \beta_0]. \end{aligned} \quad (1.8)$$

Probability and expectation for ω will be denoted by \mathbb{P} and \mathbb{E} .

Given $N \in \mathbb{N}$, $\beta > 0$, $z \in \mathbb{Z}^2$, and ω , the polymer with length $N \in \mathbb{N}$, starting point z , and disorder strength (inverse temperature) β in the random environment ω is defined by the Gibbs measure

$$dP_N^{\beta, \omega}(S | S_0 = z) := \frac{1}{Z_N^\beta(z)} e^{\sum_{n=1}^{N-1} \{\beta \omega(n, S_n) - \lambda(\beta)\}} dP(S | S_0 = z), \quad (1.9)$$

where

$$Z_N^\beta(z) = \mathbb{E}\left[e^{\sum_{n=1}^{N-1} \{\beta \omega(n, S_n) - \lambda(\beta)\}} \middle| S_0 = z\right] \quad (1.10)$$

is called *point-to-plane partition function*. Note that $\lambda(\beta)$ in the exponent ensures that $\mathbb{E}[Z_N^\beta(z)] = 1$ and, in fact, $(Z_N(z))_{N \in \mathbb{N}}$ is a martingale. It is also useful to consider the point-to-point partition functions: for $M \leq N \in \mathbb{N}_0 = \{0, 1, 2, \dots\}$ and $x, y \in \mathbb{Z}^2$, we define

$$Z_{M,N}^\beta(x, y) := \mathbb{E}\left[e^{\sum_{n=M+1}^{N-1}\{\beta\omega(n, S_n) - \lambda(\beta)\}} \mathbf{1}_{S_N=y} \mid S_M = x\right], \quad (1.11)$$

and denote $Z_{M,N}^\beta(x, \mathbf{1}) := \sum_{y \in \mathbb{Z}^2} Z_{M,N}^\beta(x, y)$ and $Z_{M,N}^\beta(\mathbf{1}, y) := \sum_{x \in \mathbb{Z}^2} Z_{M,N}^\beta(x, y)$.

Comparing (1.10) with (1.4), we note that the Feynman-Kac representation of the solution $u^\varepsilon(t, x)$ of the mollified SHE is the partition function of a Brownian directed polymer in a white-noise random environment, where the random walk $(S_i)_{0 \leq i \leq N}$ in (1.10) is replaced by the Brownian motion $(B_s)_{0 \leq s \leq t}$, and the i.i.d. space-time disorder $\omega(\cdot, \cdot)$ is replaced by the time-reversed mollified space-time white noise $\xi^\varepsilon(t - \cdot, \cdot)$. If we rescale space-time diffusively in the DPM, then $(S_{[Nt]}/\sqrt{N})_{t \geq 0}$ converges in law to $(B_t)_{t \geq 0}$, while $(\frac{1}{N}\omega([Nt], [\sqrt{N}x]))_{(t,x) \in \mathbb{R} \times \mathbb{R}^2}$ converges in law to the space-time white noise ξ . In particular, the family of partition functions

$$u^N(t, x) := Z_{0,tN}^{\beta_N}(\mathbf{1}, x\sqrt{N}), \quad t \in \frac{1}{N}\mathbb{N}_0, \quad x \in \frac{1}{\sqrt{N}}\mathbb{Z}^2, \quad (1.12)$$

gives an approximate solution to the $2d$ SHE through a discretisation of space and time. Compared with u^ε in (1.3), we see that N corresponds to ε^{-2} , and β_N corresponds to β_ε .

Instead of asking whether there are suitable choices of β_ε , A_ε and B_ε such that $(u^\varepsilon - A_\varepsilon)/B_\varepsilon$ has a non-trivial limit, we can consider instead whether there are suitable choices of β_N , A_N and B_N such that $(u^N - A_N)/B_N$ has a non-trivial limit. It turns out that the correct choice of β_N is exactly the scale on which the DPM undergoes a phase transition.

1.4. PHASE TRANSITION ON AN INTERMEDIATE DISORDER SCALE. It is known that (see [Com17, Zyg24] and the references therein) the DPM undergoes a phase transition in the sense that, there is a critical value $\beta_c \in [0, \infty)$ such that for $\beta < \beta_c$, the partition function $Z_N^\beta(0)$ defined in (1.10) converges a.s. to a positive random variable $Z_\infty(0)$, and the polymer measure $P_N^{\beta, \omega}$ converges to the Wiener measure under diffusive scaling of space-time and is hence delocalised in space [CY06]. On the other hand, when $\beta > \beta_c$, $Z_N^\beta(0)$ converges almost surely to 0, and the polymer measure $P_N^{\beta, \omega}$ experiences localisation in space [CSY03]. A recent breakthrough [JL24] has shown that when the disorder ω is bounded, there is also delocalisation at β_c , while localisation always manifests itself in a strong form.

It is known that $\beta_c = 0$ in dimensions $d = 1, 2$, and $\beta_c > 0$ in $d \geq 3$. However, it was discovered in [CSZ17b] that in $d = 2$, there is still a phase transition on an intermediate disorder scale. Recall that the expected replica overlap between two independent random walks with the same distribution as S is defined by

$$R_N := \sum_{n=1}^N \sum_{x \in \mathbb{Z}^2} q_n(x)^2 = \sum_{n=1}^N q_{2n}(0) \sim \frac{\log N}{4\pi} \quad \text{as } N \rightarrow \infty, \quad (1.13)$$

where the asymptotics follows from the local limit theorem, see (2.6) below.

It was shown in [CSZ17b] that if we choose β to be

$$\beta_N = \frac{\hat{\beta} + o(1)}{\sqrt{R_N}}, \quad \hat{\beta} \in (0, \infty), \quad (1.14)$$

then a phase transition occurs at the critical value $\hat{\beta}_c = 1$ (see Section 3 below for more details). We call $\hat{\beta} < 1$ the *subcritical regime*, and $\hat{\beta} > 1$ the *supercritical regime*. For the mollified SHE (1.3), the analogous choice is (see [CSZ17b, Theorem 2.15])

$$\beta_\varepsilon = (\hat{\beta} + o(1)) \sqrt{\frac{2\pi}{\log \varepsilon^{-1}}}, \quad \hat{\beta} \in (0, \infty). \quad (1.15)$$

This intermediate disorder scale turns out to be the correct scale on which u^N from (1.12) has non-trivial limits.

As will be discussed in more detail in Section 6, around the critical point $\hat{\beta} = 1$, there is in fact a finer window which we call the *critical window*, with β_N in this window satisfying

$$\sigma_N^2 := e^{\lambda(2\beta_N) - 2\lambda(\beta_N)} - 1 = \frac{1}{R_N} \left(1 + \frac{\vartheta + o(1)}{\log N} \right) \quad \text{for some fixed } \vartheta \in \mathbb{R}. \quad (1.16)$$

Note that for every $\vartheta \in \mathbb{R}$, we have $\beta_N \sim 1/\sqrt{R_N}$ because $\lambda(\beta) \sim \frac{1}{2}\beta^2$ as $\beta \downarrow 0$, see (1.8). So the parameter ϑ only appears in the second order asymptotics of β_N . However, different ϑ will lead to different scaling limits for the DPM partition functions regarded as a random field, which are the so-called *critical 2d stochastic heat flows*. It will become clear in Section 2.2 why it is more convenient to define the critical window in terms of $\sigma_N^2 = \sigma_N^2(\vartheta)$.

For the mollified SHE (1.3), the critical window corresponds to choosing

$$\beta_\varepsilon^2 = \frac{2\pi}{\log \frac{1}{\varepsilon}} \left(1 + \frac{\vartheta}{|\log \varepsilon|} \right). \quad (1.17)$$

1.5. OUTLINE. Assume that β_N is chosen as in (1.14) for some $\hat{\beta} > 0$, and in the case $\hat{\beta} = 1$, β_N is chosen in the critical window as in (1.16) for some $\vartheta \in \mathbb{R}$. Recall $u^N(t, x)$ from (1.12). We will discuss the following results in the rest of this article.

- Sections 3 and 4 will be devoted to results in the subcritical regime $\hat{\beta} < 1$, which include:
 - As $N \rightarrow \infty$, for each (t, x) , $u^N(t, x)$ defined in (1.12) converges to a log-normal limit when $\hat{\beta} < 1$ and converges to 0 when $\hat{\beta} \geq 1$. This shows that there is a phase transition with critical value $\hat{\beta}_c = 1$.
 - After suitable scaling, the centred random field $u^N(t, \cdot) - 1$ converges to a Gaussian limit, which is the solution of the additive SHE, also known as the Edwards-Wilkinson equation.
 - After suitable scaling, the centred random field $\log u^N(t, \cdot) - \mathbb{E}[\log u^N(t, \cdot)]$, which solves the 2d KPZ equation regularised through space-time discretisation, converges to the same Gaussian limit as $u^N(t, \cdot) - 1$.
- In Section 5, we discuss various results connected to the 2d SHE in the subcritical regime.
- Sections 6 and 7 will be devoted to the following key result (and its proof outline) in the critical window for some $\vartheta \in \mathbb{R}$:
 - Without centering and scaling, the random measure $u^N(t, \cdot)$ converges in law to a unique limit, which is called the *critical 2d stochastic heat flow* (SHF) with parameter ϑ , denoted by $\text{SHF}(\vartheta) = (\text{SHF}_{s,t}^\vartheta(dx, dy))_{0 \leq s \leq t < \infty}$.

- In Sections 8 and 9, we will explain some key technical ingredients in the proof, focusing on the second moment and higher moments calculations.
- In Section 10, we will discuss some properties of the critical 2d SHF, including:
 - An almost sure Chapman-Kolmogorov type property that justifies the name *flow*;
 - A recent axiomatic characterisation of the $\text{SHF}(\vartheta)$ by Tsai [Tsa24];
 - $\text{SHF}(\vartheta)$ cannot be obtained as the exponential of a generalised Gaussian field (i.e., it cannot be a Gaussian multiplicative chaos);
 - The marginal distribution of $\text{SHF}(\vartheta)$ at each time is almost surely singular w.r.t. the Lebesgue measure.
- In Section 11, we will discuss some open questions.

Before discussing the results above, in Section 2, we will first warm up by introducing the polynomial chaos expansion for the directed polymer partition function and performing a second moment calculation that identifies the intermediate disorder scale.

2. BASIC TOOLS AND PRELIMINARY RESULTS

2.1. POLYNOMIAL CHAOS EXPANSION. The starting point of our analysis is the *polynomial chaos expansion* of the partition function, based on the Mayer cluster expansion. More precisely, with $\beta = \beta_N$ allowed to depend on N , for each $(n, z) \in \mathbb{N} \times \mathbb{Z}^2$, we can write

$$e^{(\beta_N \omega(n, z) - \lambda(\beta_N)) \mathbf{1}_{S_n=z}} = 1 + \xi_N(n, z) \mathbf{1}_{S_n=z},$$

where $(\xi_N(n, z))_{(n, z) \in \mathbb{N} \times \mathbb{Z}^2}$ is an i.i.d. family defined by

$$\begin{aligned} \xi_N(n, z) &:= e^{\beta_N \omega(n, z) - \lambda(\beta_N)} - 1 \\ \text{with } \mathbb{E}[\xi_N(n, z)] &= 0, \quad \mathbb{V}\text{ar}[\xi_N(n, z)] = e^{\lambda(2\beta_N) - 2\lambda(\beta_N)} - 1 =: \sigma_N^2. \end{aligned} \tag{2.1}$$

We can then rewrite the partition function $Z_N^{\beta_N}(z)$ from (1.10) as follows:

$$\begin{aligned} Z_N^{\beta_N}(z) &= \mathbb{E} \left[e^{\sum_{n=1}^{N-1} (\beta_N \omega(n, S_n) - \lambda(\beta_N))} \right] \\ &= \mathbb{E} \left[\prod_{(n, x) \in \{1, \dots, N-1\} \times \mathbb{Z}^2} e^{(\beta_N \omega(n, x) - \lambda(\beta_N)) \mathbf{1}_{S_n=x}} \right] \\ &= \mathbb{E} \left[\prod_{(n, x) \in \{1, \dots, N-1\} \times \mathbb{Z}^2} (1 + \xi_N(n, x) \mathbf{1}_{S_n=x}) \right] \\ &= 1 + \sum_{r=1}^{\infty} \sum_{\substack{n_0=0 < n_1 < \dots < n_r < N \\ z_0 := z, z_1, \dots, z_r \in \mathbb{Z}^2}} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i). \end{aligned} \tag{2.2}$$

This is called a polynomial chaos expansion because it can be seen as a discrete analogue of a Wiener-Itô chaos expansion with respect to a space-time white noise on $[0, \infty) \times \mathbb{R}^2$. The terms of the expansion are L^2 -orthogonal, and each term is a multilinear polynomial in the i.i.d. random variables $(\xi_N(n, x))_{(n, x) \in \mathbb{N} \times \mathbb{Z}^2}$.

For the point-to-point partition function $Z_{M,L}^{\beta_N}(x,y)$ defined in (1.11), we have a similar expansion:

$$\begin{aligned} Z_{M,L}^{\beta_N}(x,y) &= q_{L-M}(y-x) \\ &\quad + \sum_{r=1}^{\infty} \sum_{\substack{z_0:=x, z_1, \dots, z_r \in \mathbb{Z}^2 \\ n_0 := M < n_1 < \dots < n_r < L}} \prod_{i=1}^r q_{n_i-n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i) \cdot q_{L-n_r}(y - z_r) \\ &= q_{L-M}(y-x) + \sum_{r=1}^{\infty} \sum_{\substack{z_0:=x, z_1, \dots, z_r \in \mathbb{Z}^2 \\ n_0 := M < n_1 < \dots < n_r < L}} \text{(Diagram)} , \end{aligned} \tag{2.3}$$

where in the graphical representation, the laces represent the transition kernels $q_{n_i-n_{i-1}}(z_i - z_{i-1})$, the solid dots represent the disorder variables $\xi_N(n_i, z_i)$, and the hollow circles at (M, x) and (L, y) means the absence of the disorder variables $\xi_N(M, x)$ and $\xi_N(L, y)$.

Similarly, the solution u^ε of the mollified SHE (1.3) admits a Wiener-Itô chaos expansion w.r.t. the underlying white noise ξ (see [CSZ20, (5.9)]), which can also be obtained by Picard iterations:

$$u^\varepsilon(t, x) \stackrel{\text{dist}}{=} 1 + \sum_{k \geq 1} \beta_\varepsilon^k \int \cdots \int_{\substack{z_1, \dots, z_k \in \mathbb{R}^2 \\ 0 < s_1 < \dots < s_k < t_\varepsilon}} \left(\int_{(\mathbb{R}^2)^k} \prod_{i=1}^k g_{s_i-s_{i-1}}(y_i - y_{i-1}) j(y_i - z_i) d\vec{y} \right) \prod_{i=1}^k \tilde{\xi}(s_i, z_i) ds_i dz_i , \tag{2.4}$$

where $s_0 := 0$, $y_0 := \varepsilon^{-1}x$, $g_s(\cdot)$ is the heat kernel, and $\tilde{\xi}$ is obtained from ξ via diffusive rescaling and has the same distribution.

2.2. A TRANSITION IN THE SECOND MOMENT. We now show that the second moment of the point-to-plane partition function $\mathbb{E}[Z_N^{\beta_N}(0)^2]$, see (1.10), undergoes a transition exactly at the critical point $\hat{\beta} = 1$ defined in (1.14). The fact that this L^2 transition point coincides with the actual phase transition point is unique to dimension $d = 2$ and far from trivial (see Section 3 for more details). We will also see in Corollary 2.2 that for β_N chosen as in (1.14), changing the time horizon from N to N^α effectively replaces $\hat{\beta}^2$ by $\alpha\hat{\beta}^2$.

We first recall the local central limit theorem. Let g_t be the heat kernel on \mathbb{R}^2 , i.e.,

$$g_t(x) := \frac{1}{2\pi t} e^{-\frac{|x|^2}{2t}}, \quad g_t(x, y) := g_t(y - x) . \tag{2.5}$$

Since the increments of the random walk $(S_n)_{n \geq 0}$ have an identity covariance matrix, by the local central limit theorem, see e.g. [LL10, Theorems 2.3.5 & 2.3.11], we have

$$q_n(x) = g_n(x) + O\left(\frac{1}{n^2}\right) = g_n(x) e^{O\left(\frac{1}{n}\right) + O\left(\frac{|x|^4}{n^3}\right)}, \quad n \in \mathbb{N}, x \in \mathbb{Z}^2 . \tag{2.6}$$

Lemma 2.1 (Transition in the second moment). *Let $\beta_N = \hat{\beta}/\sqrt{R_N}$ for some $\hat{\beta} \in (0, \infty)$. Then the point-to-plane partition function $Z_N := Z_N^{\beta_N}(0)$ defined in (1.10) satisfies*

$$\lim_{N \rightarrow \infty} \mathbb{E}[Z_N^2] = \begin{cases} \frac{1}{1 - \hat{\beta}^2} & \text{for } \hat{\beta} < 1, \\ \infty & \text{for } \hat{\beta} \geq 1. \end{cases} \quad (2.7)$$

We will show in Section 8 that $\mathbb{E}[Z_N^2] \sim C \log N$ at the critical point $\hat{\beta} = \hat{\beta}_c = 1$.

Proof. By the polynomial chaos expansion (2.2) for the point-to-plane partition function $Z_N := Z_N^{\beta_N}(0)$, which is an L^2 -orthogonal expansion, we have

$$\begin{aligned} \mathbb{E}[Z_N^2] &= 1 + \sum_{r=1}^{\infty} \sum_{\substack{n_0=0 < n_1 < \dots < n_r < N \\ z_1, \dots, z_r \in \mathbb{Z}^2}} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1})^2 \mathbb{E}[\xi_N(n_i, z_i)^2] \\ &= 1 + \sum_{r=1}^{\infty} \sum_{n_0=0 < n_1 < \dots < n_r < N} \sigma_N^{2r} \sum_{z_1, \dots, z_r \in \mathbb{Z}^2} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1})^2 \\ &= 1 + \sum_{r=1}^{\infty} \sum_{n_0=0 < n_1 < \dots < n_r < N} \sigma_N^{2r} \prod_{i=1}^r q_{2(n_i - n_{i-1})}(0). \end{aligned} \quad (2.8)$$

Note that since $\beta_N \rightarrow 0$, we have

$$\sigma_N^2 = e^{\lambda(2\beta_N) - 2\lambda(\beta_N)} - 1 \sim \beta_N^2 = \frac{\hat{\beta}^2}{R_N} \sim \frac{4\pi\hat{\beta}^2}{\log N} \quad \text{as } N \rightarrow \infty. \quad (2.9)$$

For each $r \in \mathbb{N}$, we then have the upper bound

$$\sum_{n_0=0 < n_1 < \dots < n_r < N} \sigma_N^{2r} \prod_{i=1}^r q_{2(n_i - n_{i-1})}(0) \leq \left(\sigma_N^2 \sum_{n=1}^N q_{2n}(0) \right)^r \sim \left(\frac{4\pi\hat{\beta}^2}{\log N} \sum_{n=1}^N \frac{1}{4\pi n} \right)^r = (\hat{\beta} + o(1))^{2r},$$

and the lower bound

$$\sum_{n_0=0 < n_1 < \dots < n_r < N} \sigma_N^{2r} \prod_{i=1}^r q_{2(n_i - n_{i-1})}(0) \geq \left(\sigma_N^2 \sum_{n=1}^{N/r} q_{2n}(0) \right)^r \sim \left(\frac{4\pi\hat{\beta}^2}{\log N} \sum_{n=1}^{N/r} \frac{1}{4\pi n} \right)^r \sim \hat{\beta}^{2r},$$

where we used that $R_N = \sum_{n=1}^N q_{2n}(0)$ and the local limit theorem (2.6). For each $r \in \mathbb{N}$, the upper and lower bounds match asymptotically. The conclusion (2.7) then follows readily. \square

The following corollary shows that changing the time horizon from N to N^α effectively changes $\hat{\beta}^2$ to $\alpha\hat{\beta}^2$. It also implies that in the subcritical regime $\hat{\beta} < 1$, the partition function $Z_N^{\beta_N}(0)$ can be approximated in L^2 by $Z_{N^\alpha}^{\beta_N}(0)$ with $\alpha \in (0, 1)$ close to 1. In particular, $Z_N^{\beta_N}(0)$ is essentially determined by disorder variables $\omega(n, x)$ (equivalently $\xi_N(n, x)$) with $n \ll N$. This fact will be an important feature that drives the phenomenology and the analysis for $\hat{\beta} < 1$. This feature will become clear in our sketch of the proof of Theorem 3.1.

Corollary 2.2 (Exponential time scale). *Let $\beta_N = \hat{\beta}/\sqrt{R_N}$ for some $\hat{\beta} \in (0, \infty)$, and let $\alpha \in (0, \infty)$. Then we have*

$$\lim_{N \rightarrow \infty} \mathbb{E}[(Z_{N^\alpha}^{\beta_N})^2] = \begin{cases} \frac{1}{1 - \alpha\hat{\beta}^2} & \text{for } \alpha\hat{\beta}^2 < 1, \\ \infty & \text{for } \alpha\hat{\beta}^2 \geq 1. \end{cases} \quad (2.10)$$

Furthermore, for $\hat{\beta} \in (0, 1)$ and $\alpha \in (0, 1]$, we have

$$\lim_{N \rightarrow \infty} \|Z_{N^\alpha}^{\beta_N} - Z_N^{\beta_N}\| := \lim_{N \rightarrow \infty} \mathbb{E}[(Z_{N^\alpha}^{\beta_N} - Z_N^{\beta_N})^2]^{1/2} = \frac{(1 - \alpha)\hat{\beta}^2}{(1 - \alpha\hat{\beta}^2)(1 - \hat{\beta}^2)} \xrightarrow{\alpha \uparrow 1} 0. \quad (2.11)$$

Proof. If we set $M := N^\alpha$, then using the asymptotics for R_N from (1.13), we can write

$$\beta_N = \frac{\hat{\beta}}{\sqrt{R_N}} = \frac{1}{\sqrt{R_M}} \cdot \frac{\hat{\beta}\sqrt{R_M}}{\sqrt{R_N}} = \frac{\hat{\beta}\sqrt{\alpha} + o(1)}{\sqrt{R_M}}.$$

The corollary then follows immediately from Lemma 2.1. Alternatively, we can also repeat the proof of Lemma 2.1 and note that the key calculation is now

$$\sigma_N^2 \sum_{n=1}^{N^\alpha} q_{2n}(0) = \frac{\hat{\beta}^2}{R_N} \cdot R_{N^\alpha} \sim \hat{\beta}^2 \frac{\log N^\alpha}{\log N} = \alpha\hat{\beta}^2,$$

which replaces the original parameter $\hat{\beta}^2$.

Note that for $\alpha < 1$, the polynomial chaos expansion for $Z_{N^\alpha}^{\beta_N}$, see (2.2), contains only a subset of terms in the polynomial chaos expansion for $Z_N^{\beta_N}$, which is an L^2 -orthogonal expansion. The claim in (2.11) then follows from (2.10). \square

3. GAUSSIAN FLUCTUATIONS IN THE SUBCRITICAL REGIME

In this section, we review results in the subcritical regime, that is, results (a)-(c) listed in Section 1.5. Throughout the rest of this section, we will consider β_N chosen as in (1.14) for some $\hat{\beta} \in (0, 1)$. Instead of considering $u^N(t, x)$, the solution of the space-time discretised SHE defined in (1.12), for simplicity, we will fix $t = 1$ and consider the time reversed field of point-to-plane partition functions $Z_N^{\beta_N}(x\sqrt{N}) = Z_{0,N}^{\beta_N}(x\sqrt{N}, \mathbf{1})$, $x \in \frac{1}{\sqrt{N}}\mathbb{Z}^2$.

3.1. ONE POINT FLUCTUATION. In [CSZ17b, Theorem 2.8], it was shown that on the intermediate disorder scale defined in (1.14), the DPM in dimension $d = 2$ undergoes a phase transition similar to the phase transition in $d \geq 3$.

Theorem 3.1 (Limit of individual partition function). *We have the following convergence result for the point-to-plane partition function $Z_N^{\beta_N}(0)$:*

$$Z_N^{\beta_N}(0) \xrightarrow[N \rightarrow \infty]{d} \mathbf{Z}_{\hat{\beta}} := \begin{cases} \exp\left(\sigma_{\hat{\beta}} W_1 - \frac{\sigma_{\hat{\beta}}^2}{2}\right) & \text{if } \hat{\beta} < 1 \\ 0 & \text{if } \hat{\beta} \geq 1 \end{cases}. \quad (3.1)$$

where W_1 is a standard Gaussian random variable and

$$\sigma_{\hat{\beta}}^2 := \log \frac{1}{1 - \hat{\beta}^2}. \quad (3.2)$$

Remark 3.2. In [CSZ17b, Theorem 2.12], it was shown that $\log Z_N^{\beta_N}(x_N)$ and $\log Z_N^{\beta_N}(y_N)$ converge to a pair of independent normal random variables in the subcritical regime $\hat{\beta} < 1$ if x_N/\sqrt{N} and y_N/\sqrt{N} converge to distinct $x, y \in \mathbb{R}^2$. A non-trivial bivariate normal limit arises only if $\|x_N - y_N\| = N^{\alpha+o(1)}$ for some $\alpha \in (0, 1/2)$. Similar multivariate normal limits hold for the joint distribution of k log-partition functions $\log Z_{N^{\gamma_i}}^{\beta_N}(x_N^i)$, $1 \leq i \leq k$, with non-trivial change of the covariance matrix as we vary the exponents $\gamma_i \in (0, 1)$ in the time horizon N^{γ_i} .

In what follows, we sketch several alternative proof strategies for Theorem 3.1, which will help shed light on the structure of the 2d DPM and SHE in the subcritical regime.

Proof strategy 1 [CSZ17b, Theorem 2.15]. We start with the polynomial chaos expansion (2.2) for the point-to-plane partition function

$$Z_N^{\beta_N}(0) = 1 + \sum_{r=1}^{\infty} \sum_{\substack{n_0=0 < n_1 < \dots < n_r < N \\ z_0:=0, z_1, \dots, z_r \in \mathbb{Z}^2}} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i), \quad (3.3)$$

where the parameter $\hat{\beta} < 1$ is contained in the variance of ξ_N , see (2.9). From the second moment calculations in Lemma 2.1 and Corollary 2.2, we see that the dominant contribution to the series in (3.3) comes from terms which of order $r = O(1)$ – there is an exponential decay of order $\hat{\beta}^r$ in the chaoses of order r – the times differences are of order $n_i - n_{i-1} =: N^{t_i}$ for $t_i \in (0, 1)$, and the spatial coordinates $\|z_i - z_{i-1}\| = O(\sqrt{n_i - n_{i-1}}) = O(N^{t_i/2})$, for $1 \leq i \leq r$. The contributions coming from $t_i \approx t_j$ for $i \neq j$ is negligible, and furthermore, the dominant contribution comes from $\|z_i\| = O(\sqrt{n_i})$, $1 \leq i \leq r$.

The key observation of [CSZ17b] is that, we should rewrite (3.3) by partitioning

$$(t_1, t_2, \dots, t_r) = (s_{1,1}, \dots, s_{1,r_1}; s_{2,1}, \dots, s_{2,r_2}; \dots; s_{k,1}, \dots, s_{k,r_k}),$$

where $s_{1,1}, s_{2,1}, \dots, s_{k,1}$ are the successive running maxima along the original sequence (t_1, \dots, t_r) , and each subsequence $(s_{1,1}, \dots, s_{i,r_i})$ is called a *dominated sequence* because the first element $s_i := s_{i,1}$ dominates the subsequent elements $s_{i,2}, \dots, s_{i,r_i}$. Recall that in (3.3), $N^{s_i} = n_{M+1} - n_M$ with $M = r_1 + \dots + r_{i-1}$, and

$$n_M = N^{t_1} + \dots + N^{t_M}.$$

The fact that $s_i = s_{i,1}$ is a running maxima of (t_1, \dots, t_r) means that $s_i > t_1, t_2, \dots, t_M$, and hence

$$n_{M+1} - n_M = N^{s_i} \gg N^{t_1} + \dots + N^{t_M} = n_M.$$

In particular, $n_{M+1} - n_M \approx n_{M+1}$. Therefore in (3.3), at the cost of a small error in L^2 , we can replace the factor $q_{n_{M+1} - n_M}(z_{M+1} - z_M)$ by $q_{n_{M+1}}(z_{M+1})$, since the dominant contribution to (3.3) comes from $\|z_M\| = O(\sqrt{n_M}) \ll \sqrt{n_{M+1}}$.

Now observe that the replacement $q_{n_{M+1} - n_M}(z_{M+1} - z_M) \rightsquigarrow q_{n_{M+1}}(z_{M+1})$ in (3.3) at each running maxima of the sequence $n_i - n_{i-1} = N^{t_i}$, $1 \leq i \leq r$, leads to the factorisation

$$Z_N^{\beta_N}(0) \approx 1 + \sum_{k=1}^{\infty} \sum_{\substack{0 < s_1 < \dots < s_k < 1 \\ N^{s_i} \in \mathbb{N}}} \prod_{i=1}^k \Xi_N^{\beta_N}(N^{s_i}), \quad (3.4)$$

where for $s \in (0, 1)$ with $n_1 := N^s \in \mathbb{N}$,

$$\Xi_N^{\beta_N}(n_1) := \sum_{r=1}^{\infty} \sum_{\substack{z_1, \dots, z_r \in \mathbb{Z}^2 \\ n_i - n_{i-1} < n_1}} q_{n_1}(z_1) \xi_N(n_1, z_1) \prod_{i=2}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i). \quad (3.5)$$

The key point is that for $s \in (0, 1)$, $\Xi_N^{\beta_N}(N^s)$ depends essentially only on $\xi_N(n, \cdot)$ with $n \in [N^s, N^{s+o(1)}]$. In particular, for $s < t$, $\Xi_N^{\beta_N}(N^s)$ and $\Xi_N^{\beta_N}(N^t)$ are essentially independent and approximate the increments of a time changed Brownian motion. More precisely, it turns out that

$$M_N^{\beta_N}(t) := \sum_{1 \leq n_1 \leq N^t} \Xi_N^{\beta_N}(n_1) \xrightarrow[N \rightarrow \infty]{d} \int_0^t \frac{\hat{\beta}}{\sqrt{1 - \hat{\beta}^2 s}} dW(s), \quad t \in [0, 1], \quad (3.6)$$

where W is a standard Brownian motion. Just to check that the variances match, one can use the same calculations as in Section 2.2 to show that $\text{Var}(M_N^{\beta_N}(t)) \rightarrow \int_0^t \frac{\hat{\beta}^2}{1 - \hat{\beta}^2 s} ds$. It follows that $\Xi_N^{\beta_N}(N^t) \approx \frac{\hat{\beta}}{\sqrt{1 - \hat{\beta}^2 t}} dW(t)$, which implies that the r.h.s. of (3.4) should converge to the Wick exponential : $\exp\left(\int_0^1 \frac{\hat{\beta}}{\sqrt{1 - \hat{\beta}^2 t}} dW(t)\right)$:, thus matching the conclusion of Theorem 3.1.

The actual proof in [CSZ17b] is fairly involved and gives more information. In particular, the terms in (3.5) corresponding to different $r \in \mathbb{N}$ can be shown to converge to increments of independent Brownian motions. In particular, the white noise dW in (3.6) is the sum of a sequence of independent white noises, which arise as the limits of random fields defined from degree r polynomials of $\xi_N(\cdot, \cdot)$, one for each $r \in \mathbb{N}$. The proof is based on the *Fourth Moment Theorem* for Gaussian limits, see [CSZ17b, Theorem 4.2] and the references therein. \square

Proof strategy 2 [CC22, Theorems 3.5 & 3.6]. Instead of proving that $Z_N^{\beta_N}(0)$ converges in distribution to a log-normal random variable, i.e. the exponential of a Gaussian, see (3.1), it is also possible to prove directly that $\log Z_N^{\beta_N}(0)$ converges to a Gaussian. To this end, we approximate $\log Z_N^{\beta_N}(0)$ in terms of a random variable X_N^{dom} obtained by restricting the sum in (3.3) to a *single dominated sequence* $n_1 \geq \max\{n_2 - n_1, \dots, n_k - n_{k-1}\}$:

$$X_N^{\text{dom}} := \sum_{r=1}^{\infty} \sum_{\substack{n_0 = 0 < n_1 < \dots < n_r < N \\ n_1 \geq \max\{n_2 - n_1, \dots, n_k - n_{k-1}\} \\ z_0 := 0, z_1, \dots, z_r \in \mathbb{Z}^2}} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i). \quad (3.7)$$

One can show that X_N^{dom} is asymptotically Gaussian, more precisely $\mathbb{E}[(X_N^{\text{dom}})^2] \rightarrow \sigma_{\hat{\beta}}^2$ and $X_N^{\text{dom}} \rightarrow N(0, \sigma_{\hat{\beta}}^2)$ in distribution, see [CC22, Theorems 3.6]. The proof is based on the classic Feller-Lindeberg Central Limit Theorems for sums of independent random variables (obtained by partitioning the sums in (3.7) in disjoint intervals). A strengthened version of Theorem 3.1 is then obtained by showing that $\log Z_N^{\beta_N}(0) \approx X_N^{\text{dom}} - \frac{1}{2}\mathbb{E}[(X_N^{\text{dom}})^2]$ in L^2 (not just in distribution), which is shown in [CC22, Theorems 3.5] by exploiting a decomposition in terms of dominated sequences similar to the one described in Proof strategy 1.

Proof strategy 3. In [CD24], yet another proof of Theorem 3.1 was presented. The main idea was to approximate $Z_N^{\beta_N}(0)$ by a product of independent random variables. More

precisely, $Z_N^{\beta_N}(0) \approx \prod_{l=1}^M Z_{N,k}^{\beta_N}$ where $Z_{N,k}^{\beta_N} := E\left[e^{\sum_{n=\tau_{k-1}+1}^{\tau_k} (\beta_N \omega(n, S_n) - \lambda(\beta_N))}\right]$, with $\tau_0 := 0$ and $\tau_k = \lceil N^{k/M} \rceil$ for $k = 1, \dots, M$. Note that (τ_k) encode the same exponential time scales as in Proof strategy 1, and $(Z_{N,k}^{\beta_N})_{1 \leq k \leq M}$ are independent random variables with mean 1 and variance of order $1/M$ (cf. Corollary 2.2). Applying a classic Lindeberg Central Limit Theorem to $\log Z_N^{\beta_N}(0) \approx \sum_{k=1}^N \log Z_{N,k}^{\beta_N}$ then gives the convergence of $\log Z_N^{\beta_N}(0)$ in distribution to a Gaussian random variable with mean $-\sigma_{\hat{\beta}}^2/2$ and variance $\sigma_{\hat{\beta}}^2$. This requires Taylor expanding $\log Z_{N,k}^{\beta_N} = \log(1 + (Z_{N,k}^{\beta_N} - 1))$ and bounding the $(2 + \varepsilon)$ moments of $Z_{N,k}^{\beta_N} - 1$. The latter is achieved by applying hypercontractivity estimates for polynomial chaos expansions from [MOO10], which bounds moments higher than 2 by the second moment of a modified polynomial chaos expansion. In our setting, this amounts to increasing $\hat{\beta}$ in (1.14) to some $\hat{\beta}' > \hat{\beta}$. Thanks to the fact that we are considering the subcritical regime $\hat{\beta} < 1$, we can choose $\hat{\beta}' < 1$ to be subcritical as well, which ensures that we have the desired control on the second moments of the partition functions.

3.2. RANDOM FIELD FLUCTUATION. In [CSZ17b, Theorem 2.13], it was shown that in the subcritical regime $\hat{\beta} < 1$, the fluctuation of the random field of partition functions $(Z_{tN}^{\beta_N}(x\sqrt{N}))_{t \in \frac{1}{N}[0,1], x \in \frac{1}{\sqrt{N}}\mathbb{Z}^2}$ (and its mollified SHE analogue u^ε from (1.3)) is asymptotically Gaussian and solves the Edwards-Wilkinson equation (or additive SHE).

Theorem 3.3 (Edwards-Wilkinson fluctuation for DPM). *For any test function $\phi \in C_c(\mathbb{R}^2)$ and any $t \in (0, 1]$, we have*

$$\frac{1}{\beta_N} \int_{\mathbb{R}^2} \phi(x) \left(Z_{tN}^{\beta_N}(\lfloor x\sqrt{N} \rfloor) - 1 \right) dx \xrightarrow[N \rightarrow \infty]{d} \int_{\mathbb{R}^2} \phi(x) v(t, x) dx, \quad (3.8)$$

where $v(t, x)$ is a generalised Gaussian field that solves the 2d Edwards-Wilkinson equation

$$\begin{aligned} \partial_t v(t, x) &= \frac{1}{2} \Delta v(t, x) + \sqrt{\frac{1}{1 - \hat{\beta}^2}} \xi(t, x), \\ v(0, x) &\equiv 1. \end{aligned} \quad (3.9)$$

Proof sketch. Without loss of generality, assume $t = 1$. Using the polynomial chaos expansion for $Z_N^{\beta_N}(z)$ from (2.2), we can rewrite the left-hand side of (3.8) as

$$\begin{aligned} \Phi_N &:= \frac{1}{\beta_N N} \sum_{r=1}^{\infty} \sum_{\substack{z_0, z_1, \dots, z_r \in \mathbb{Z}^2 \\ 0 = n_0 < n_1 < \dots < n_r < N}} \phi\left(\frac{z_0}{\sqrt{N}}\right) \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i) \\ &= \underbrace{\frac{1}{N} \sum_{\substack{z_1 \in \mathbb{Z}^2 \\ 0 < n_1 < N}} \left(\sum_{z_0 \in \mathbb{Z}^2} \phi\left(\frac{z_0}{\sqrt{N}}\right) q_{n_1}(z_1 - z_0) \right) \Xi_N^{\beta_N}(n_1, z_1)}_{q_{N,\phi}(n_1, z_1)}, \end{aligned} \quad (3.10)$$

where

$$\Xi_N^{\beta_N}(n_1, z_1) := \frac{\xi_N(n_1, z_1)}{\beta_N} \left(1 + \sum_{k=2}^{\infty} \sum_{\substack{z_2, \dots, z_k \in \mathbb{Z}^2 \\ n_1 < n_2 < \dots < n_k < N}} \prod_{i=2}^k q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i) \right)$$

is simply the product of $\xi_N(n_1, z_1)/\beta_N$ with the point-to-plane partition function $Z_N^{\beta_N}(n_1, z_1)$ starting at (n_1, z_1) and ending at time N .

Note that the sum in (3.10) is an L^2 -orthogonal decomposition, each $\Xi_N^{\beta_N}(n_1, z_1)$ has variance of order 1, and the averaging w.r.t. the kernel $q_{N,\phi}(n_1, z_1)$ ensures that the dominant contribution to Φ_N comes from (n_1, z_1) of the order (N, \sqrt{N}) . Now the key observation is that, by Corollary 2.2 and the calculations therein, each $\Xi_N^{\beta_N}(n_1, z_1)$ depends only on $\xi_N(n, z)$ with (n, z) in a sub-diffusive time-space window around (n_1, z_1) . Therefore Φ_N is an average of the field of random variables $\Xi_N^{\beta_N}(n_1, z_1)$ with very local dependence. It is then not surprising that such an average has a Gaussian limit as $N \rightarrow \infty$. This was first proved in [CSZ17b] using the Fourth Moment Theorem for polynomial chaos expansions. See also [CC22] for an alternative proof using more classic tools from the proof of CLT for triangular arrays. \square

4. THE 2d KPZ IN THE SUBCRITICAL REGIME

In this section we will discuss the fluctuation of the solution to the KPZ equation (1.5) in the subcritical regime $\hat{\beta} < 1$. The one-point statistics can be easily deduced from the log-normality of the SHE via the Cole-Hopf transformation $h^\varepsilon(t, x) = \log u^\varepsilon(t, x)$ as

$$h^\varepsilon(t, x) \xrightarrow[\varepsilon \downarrow 0]{d} \begin{cases} \sigma_{\hat{\beta}} W_1 - \frac{1}{2} \sigma_{\hat{\beta}}^2 & \text{if } \hat{\beta} < 1 \\ -\infty & \text{if } \hat{\beta} \geq 1 \end{cases} \quad \text{with } \sigma_{\hat{\beta}}^2 := \log \frac{1}{1-\hat{\beta}^2}, \quad W_1 \sim N(0, 1), \quad (4.1)$$

The random field fluctuations are more subtle as the operation of averaging does not commute with nonlinear functions such as $\log x$, i.e., $\int \phi(x) \log u^\varepsilon(t, x) dx \neq \log \int \phi(x) u^\varepsilon(t, x) dx$. Surprisingly, the limiting random field fluctuation for $\log u^\varepsilon$ turns out to be the same as the fluctuation for u^ε and its polymer partition function analogue in Theorem 3.3. The theorem we obtained in [CSZ20] (see also [Gu20]) is the following:

Theorem 4.1 (Edwards-Wilkinson fluctuation for subcritical 2d KPZ - [CSZ20, Gu20]). *Let h^ε be the solution of the mollified KPZ equation (1.5) with initial condition $h^\varepsilon(0, x) \equiv 0$ and $\beta_\varepsilon = \hat{\beta}\sqrt{2\pi/\log \varepsilon^{-1}}$, $\hat{\beta} \in (0, 1)$. Denote*

$$\mathfrak{h}^\varepsilon(t, x) := \frac{h^\varepsilon(t, x) - \mathbb{E}[h^\varepsilon(t, x)]}{\beta_\varepsilon} = \frac{\sqrt{\log \varepsilon^{-1}}}{\sqrt{2\pi} \hat{\beta}} (h^\varepsilon(t, x) - \mathbb{E}[h^\varepsilon(t, x)]), \quad (4.2)$$

where the centering satisfies $\mathbb{E}[h^\varepsilon(t, x)] = -\frac{1}{2} \sigma_{\hat{\beta}}^2 + o(1)$ as $\varepsilon \downarrow 0$, see (4.1).

For any $t > 0$ and $\phi \in C_c(\mathbb{R}^2)$, the following convergence in law holds:

$$\langle \mathfrak{h}^\varepsilon(t, \cdot), \phi(\cdot) \rangle = \int_{\mathbb{R}^2} \mathfrak{h}^\varepsilon(t, x) \phi(x) dx \xrightarrow[\varepsilon \downarrow 0]{d} \langle v(t, \cdot), \phi(\cdot) \rangle, \quad (4.3)$$

where $v(\cdot, \cdot)$ is the solution of the two-dimensional Edwards-Wilkinson equation in (3.9).

The analogue of Theorem 4.1 for the log partition functions of the directed polymer is also given in [CSZ20]. This will be the theorem whose proof we will outline. The proof of Theorem 4.1 follows exactly the same lines if instead of working with the polynomial chaos expansion of the partition function we work with the Wiener chaos expansion of u^ε as in (2.4). We refer to [CSZ20, Section 5] for details.

Theorem 4.2. Assume the same setting as in Theorem 3.3 for some $\hat{\beta} \in (0, 1)$. Apart from (1.8), suppose that the disorder ω also satisfies the concentration of measure property [†]:

$$\exists \gamma > 1, C_1, C_2 \in (0, \infty) : \text{for all } n \in \mathbb{N} \text{ and } f : \mathbb{R}^n \rightarrow \mathbb{R} \text{ convex and 1-Lipschitz}$$

$$\mathbb{P}\left(|f(\omega_1, \dots, \omega_N) - M_f| \geq t\right) \leq C_1 \exp\left(-\frac{t^\gamma}{C_2}\right), \quad (4.4)$$

where M_f denotes a median of $f(\omega_1, \dots, \omega_N)$. Denote[‡]

$$\mathfrak{h}_N(t, x) := \frac{\log Z_{tN}^{\beta_N}(x\sqrt{N}) - \mathbb{E}[\log Z_{tN}^{\beta_N}]}{\beta_N} = \frac{\sqrt{\log N}}{\sqrt{4\pi}\hat{\beta}} (\log Z_{tN}^{\beta_N}(x\sqrt{N}) - \mathbb{E}[\log Z_{tN}^{\beta_N}]), \quad (4.5)$$

For any $t > 0$ and $\phi \in C_c(\mathbb{R}^2)$, the following convergence in law holds with $v(\cdot, \cdot)$ as in (3.9):

$$\langle \mathfrak{h}_N(t, \cdot), \phi(\cdot) \rangle = \int_{\mathbb{R}^2} \mathfrak{h}_N(t, x) \phi(x) dx \xrightarrow[N \rightarrow \infty]{d} \langle v(t, \cdot), \phi(\cdot) \rangle, \quad (4.6)$$

Remark 4.3. We note that the constant 4π that appears in (4.5) is different than the constant 2π that appears in the corresponding statement in [CSZ20]. This discrepancy is due to the fact that in this review we work with aperiodic random walks.

Proof sketch for Theorem 4.2. The main idea is to “linearize” $\log Z_N^{\beta_N}(x)$ by Taylor expansion, but not around $\mathbb{E}[Z_N^{\beta_N}(x)] = 1$. Rather, we should Taylor expand $\log Z_N^{\beta_N}(x)$ around $Z_{N,\beta_N}^A(x)$, which depends only on disorder in a small time-space window around the starting point of the polymer, but with $Z_{N,\beta_N}^A(x) \approx Z_N^{\beta_N}(x)$ by Corollary 2.2. More precisely, we define the time-space window

$$A_N^x := \left\{(n, z) \in \mathbb{N} \times \mathbb{Z}^2 : n \leq N^{1-a_N}, |z - x| < N^{\frac{1}{2}-\frac{a_N}{4}}\right\}, \quad (4.7)$$

where

$$a_N = \frac{1}{(\log N)^{1-\gamma}} \quad \text{with} \quad \gamma \in (0, \gamma^*), \quad (4.8)$$

for some $\gamma^* > 0$ depending only on $\hat{\beta}$, which makes the time window $\ll N$. The precise choice of γ^* is more of a technical nature and we will not bother with it here; one can refer to [CSZ20] for details. The spatial window of the set A_N^x is slightly superdiffusive compared to the time window to ensure that the random walk starting from x at time 0 will stay inside A_N^x till the end of the time window with high probability.

We define now the partition function $Z_{N,\beta}^A(x)$ which only uses disorder in A_N^x , i.e.,

$$Z_{N,\beta_N}^A(x) := \mathbb{E}_x\left[e^{H_{A_N^x}^{\beta_N}}\right], \quad \text{where} \quad H_{A_N^x}^{\beta_N} := \sum_{(n,x) \in A_N^x} (\beta_N \omega_{n,x} - \lambda(\beta_N)) \mathbb{1}_{\{S_n=x\}}. \quad (4.9)$$

This allows us to decompose the original partition function $Z_N^{\beta_N}(x)$ as follows:

$$Z_N^{\beta_N}(x) = Z_{N,\beta_N}^A(x) + \hat{Z}_{N,\beta_N}^A(x), \quad (4.10)$$

[†]Condition (4.4) is satisfied if ω are bounded, Gaussian, or if they have a density $\exp(-V(\cdot) + U(\cdot))$, with V uniformly strictly convex and U bounded. We refer to [Led01] for more details.

[‡]The scaling constant $\sqrt{4\pi}$ here is different than the constant $\sqrt{\pi}$ in [CSZ20] due to the aperiodicity of the random walk we consider here.

where $\hat{Z}_{N,\beta_N}^A(x) \stackrel{L^2}{\approx} 0$ is the ‘‘remainder’’, although it will be the source of the limiting fluctuation in (4.6). We now perform the Taylor expansion

$$\log Z_N^{\beta_N}(x) = \log Z_{N,\beta_N}^A(x) + \log \left(1 + \frac{\hat{Z}_{N,\beta_N}^A(x)}{Z_{N,\beta_N}^A(x)} \right) = \log Z_{N,\beta_N}^A(x) + \frac{\hat{Z}_{N,\beta_N}^A(x)}{Z_{N,\beta_N}^A(x)} + O_N(x), \quad (4.11)$$

where $O_N(x)$ is the error term. This approximation is quantified via the following estimates.

Estimate 1. For suitable test functions $\phi(\cdot)$, we have

$$\sqrt{\log N} \cdot \frac{1}{N} \sum_{x \in \mathbb{Z}^2} (O_N(x) - \mathbb{E}[O_N(x)]) \phi\left(\frac{x}{\sqrt{N}}\right) \xrightarrow[N \rightarrow \infty]{L^2(\mathbb{P})} 0. \quad (4.12)$$

The proof of this estimate uses a simple Taylor expansion estimate, which says that, essentially, the error term $O_N(z)$ is bounded by $(\frac{\hat{Z}_{N,\beta_N}^A(x)}{Z_{N,\beta_N}^A(x)})^2$. In order to bound this error, we apply Hölder inequality to separate the numerator and denominator. This in turn requires bounds on the moments of the partition function of order slightly higher than 2, and bounds on arbitrarily negative moments. In [CSZ20], hyper-contractivity was used to bound moments of order higher than 2. Alternatively, one can also use more recent techniques from [LZ23, CZ23] to bound moments of all orders in the subcritical regime.

To bound the negative moments, we use Assumption (4.4) and concentration of measure estimates ([CSZ20, Proposition 3.1]):

Negative tails. For any $\hat{\beta} \in (0, 1)$, there exists $c = c(\hat{\beta}) \in (0, \infty)$ with the following property: for every $N \in \mathbb{N}$ and for every choice of $\Lambda \subseteq \{1, \dots, N\} \times \mathbb{Z}^2$, one has

$$\forall t \geq 0 : \quad \mathbb{P}(\log Z_{\Lambda, \beta_N} \leq -t) \leq c e^{-t^{\gamma}/c}, \quad (4.13)$$

where $\gamma > 1$ is the same exponent appearing in assumption (4.4).

Next, we need:

Estimate 2. For $Z_{N,\beta_N}^A(\cdot)$ defined as in (4.9) and suitable test function ϕ ,

$$\sqrt{\log N} \cdot \frac{1}{N} \sum_{x \in \mathbb{Z}^2} (\log Z_{N,\beta_N}^A(x) - \mathbb{E}[\log Z_{N,\beta_N}^A(x)]) \phi\left(\frac{x}{\sqrt{N}}\right) \xrightarrow[N \rightarrow \infty]{L^2(\mathbb{P})} 0. \quad (4.14)$$

The proof of this is a fairly simple $L^2(\mathbb{P})$ estimate and uses the fact that $Z_{N,\beta_N}^A(\cdot)$ has very local dependence on the disorder, i.e., $Z_{N,\beta_N}^A(x)$ and $Z_{N,\beta_N}^A(y)$ are independent if $|x - y| \geq 2N^{\frac{1}{2} - \frac{a_N}{4}} \ll \sqrt{N}$.

Estimates 1 and 2 imply that the fluctuations of the field $\log Z_N^{\beta_N}(\cdot)$ is governed by the fluctuations of the field $\frac{\hat{Z}_{N,\beta_N}^A(\cdot)}{Z_{N,\beta_N}^A(\cdot)}$. The crucial point here is that the numerator *approximately factorises* in a way that cancels the denominator, and what remains is a restricted polymer partition function for which we can apply a variant of Theorem 3.3. With reference to Figure 1 (B), we define the partition function

$$Z_{N,\beta_N}^B(x) := \mathbb{E}_x[e^{H_{B_N}^{\beta_N}}] \quad \text{where} \quad H_{B_N}^{\beta_N} := \sum_{(n,x) \in B_N} (\beta_N \omega_{n,x} - \lambda(\beta_N)) \mathbb{1}_{\{S_n=x\}}$$

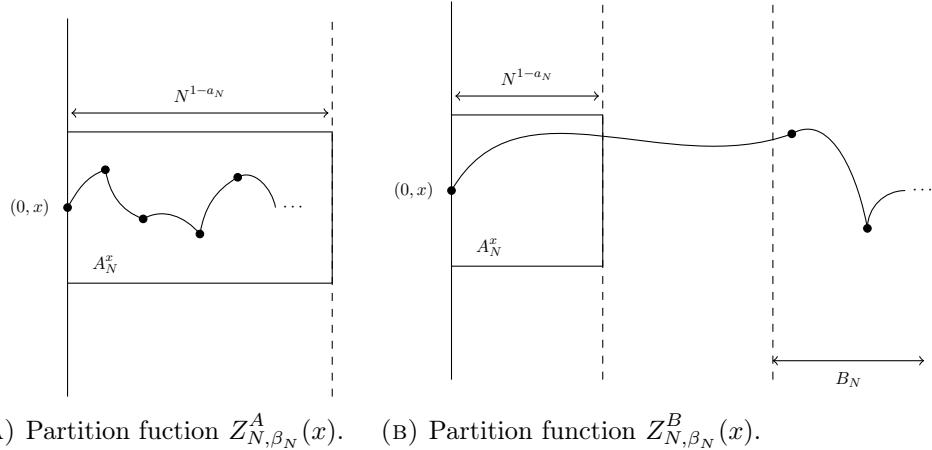


FIGURE 1. The above figures depict the chaos expansions of $Z_{N,\beta_N}^A(x)$ and $Z_{N,\beta_N}^B(x)$. The disorder used by $Z_{N,\beta_N}^A(x)$ is restricted to the set A_N^x , while the disorder used by $Z_{N,\beta_N}^B(x)$ is restricted to B_N .

where the set B_N is as in Figure 1. The crucial approximation that we establish in [CSZ20] is

$$\hat{Z}_{N,\beta_N}^A(x) \approx Z_{N,\beta_N}^A(x) (Z_{N,\beta_N}^B(x) - 1). \quad (4.15)$$

Here we use again the fact that the time-space window A_N^x is of microscopic scale compared to the diffusive scale (N, \sqrt{N}) . The quantitative estimate related to this approximation is

Estimate 3. For $Z_{N,\beta_N}^A(\cdot)$, $\hat{Z}_{N,\beta_N}^A(\cdot)$, $Z_{N,\beta_N}^B(\cdot)$ defined as above and for suitable test function ϕ , we have

$$\sqrt{\log N} \cdot \frac{1}{N} \sum_{x \in \mathbb{Z}^2} \left(\frac{\hat{Z}_{N,\beta_N}^A(x)}{Z_{N,\beta_N}^A(x)} - (Z_{N,\beta_N}^B(x) - 1) \right) \phi\left(\frac{x}{\sqrt{N}}\right) \xrightarrow[N \rightarrow \infty]{L^1(\mathbb{P})} 0.$$

The above estimates reduce the study of the fluctuation of the field $\log Z_N^{\beta_N}(\cdot)$ to those of the field $Z_{N,\beta_N}^B(\cdot)$, which is a restricted partition function and we can apply the same proof as for Theorem 3.3 to establish the Edward-Wilkinson limit. \square

5. RELATED RESULTS OUTSIDE THE CRITICAL REGIME

In this section we will briefly list some further related work. The focus will be on the critical dimension, but below the critical disorder strength. We will also discuss results in the quasi-critical regime and results in higher dimensions.

5.1. OTHER MARGINALLY RELEVANT DIRECTED POLYMER MODELS. Theorems 3.1 and 3.3 are special cases of more general results in [CSZ17b] concerning one-point and random field fluctuations of directed polymer models that are *marginally relevant*. Roughly speaking, a directed polymer model in the time-space domain $\mathbb{N} \times \mathbb{Z}^d$ is marginally relevant

if its replica overlap

$$R_N := \sum_{n=1}^N \sum_{x \in \mathbb{Z}^d} P(S_n = S'_n),$$

with S and S' being two independent random walks on \mathbb{Z}^d with law P , diverges as a slowly varying function in N . We will explain in more detail what is meant by the marginal relevance of a disordered systems and the connection with singular SPDE in Section 11.1.

As we have seen, for the 2d DPM, the divergence of the replica overlap R_N is logarithmic. Besides the 2d DPM, other examples include directed polymer in $\mathbb{N} \times \mathbb{Z}^d$ with $d = 1$, defined from a heavy-tailed random walk in the domain of attraction of the Cauchy distribution, i.e. $P(S_n > x) \sim x^{-1}$, or the disordered pinning model which can be regarded as a directed polymer on $\mathbb{N} \times \mathbb{Z}^d$ with $d = 0$, defined from a renewal process on $\mathbb{N} \times \{0\}$ (marginally relevant if the renewals are the times when a simple symmetric random walk on \mathbb{Z} returns to 0). It was shown in [CSZ17b] that these marginally relevant directed polymer models all exhibit the same asymptotic limit behavior as in Theorems 3.1 for the 2d DPM, which shows a surprising universality across different dimensions.

5.2. STOCHASTIC HEAT EQUATION. The analogue of Theorem 3.1 for the solution u^ε of the mollified SHE (1.3) was proved in [CSZ17b, Theorem 2.15]. More precisely,

Theorem 5.1 (1-point statistics for the mollified SHE). *If $\beta_\varepsilon = (\hat{\beta} + o(1))\sqrt{\frac{2\pi}{\log \frac{1}{\varepsilon}}}$ with $\hat{\beta} \in (0, 1)$, then for any $t > 0$ and $x \in \mathbb{R}^2$, we have*

$$u^\varepsilon(t, x) \xrightarrow[\varepsilon \rightarrow 0]{d} \begin{cases} \exp\left(\sigma_{\hat{\beta}} W_1 - \frac{\sigma_{\hat{\beta}}^2}{2}\right) & \text{if } \hat{\beta} < 1 \\ 0 & \text{if } \hat{\beta} \geq 1 \end{cases},$$

with W_1 a standard normal random variable and $\sigma_{\hat{\beta}}^2 = \log \frac{1}{1-\hat{\beta}^2}$.

Moreover, the analogue Theorem 3.3, i.e. Edwards-Wilkinson limit for the field of the solution $u^\varepsilon(t, \cdot)$ was also established in [CSZ17b]:

Theorem 5.2 (Edwards-Wilkinson fluctuation for the mollified SHE). *If $\beta_\varepsilon = (\hat{\beta} + o(1))\sqrt{\frac{2\pi}{\log \frac{1}{\varepsilon}}}$ with $\hat{\beta} \in (0, 1)$, then for any test function $\phi \in C_c(\mathbb{R}^2)$, we have*

$$\frac{1}{\beta_\varepsilon} \int_{\mathbb{R}^2} \phi(x) (u^\varepsilon(t, x) - 1) dx \xrightarrow[\varepsilon \rightarrow 0]{d} \int_{\mathbb{R}^2} \phi(x) v(t, x) dx,$$

where $v(t, x)$ is an in (3.9).

Theorem 2.17 in [CSZ17b] actually considered the fluctuation of the solution as a space-time field instead of as a field in space at a given time, as stated above. The proof of Theorem 5.2 is simpler. The joint limit of the solution at finitely many space-time points at sub-diffusive scales is identified in [CSZ17b, Theorem 2.15].

5.3. ALTERNATIVE APPROACH TO SUBCRITICAL 2d KPZ. In a subset of the subcritical regime for the 2d KPZ considered in Theorem 4.1, tightness of the fluctuation field was first established in [CD20], while [Gu20] gave an alternative proof using tools from stochastic analysis. The linearisation strategy is the same as in [CSZ20], but [Gu20] used

the Clark-Ocone formula (see [Nua06] for a reference and more background on Malliavin calculus) to express

$$\log u^\varepsilon(t, x) = \mathbb{E}[\log u^\varepsilon] + \int_0^t \int_{\mathbb{R}^2} \mathbb{E}[D_{s,y} \log u^\varepsilon(t, x) | \mathcal{F}_s] dW(s, y),$$

where $D_{s,y}$ is the Malliavin derivative and \mathcal{F}_s is the filtration generated by the white noise field $(W(r, x) : r < s)$ up to time s . Using Malliavin calculus, the above is equal to

$$\log u^\varepsilon(t, x) = \mathbb{E}[\log u^\varepsilon] + \int_0^t \int_{\mathbb{R}^2} \mathbb{E}\left[\frac{D_{s,y} u^\varepsilon(t, x)}{u^\varepsilon(t, x)} \mid \mathcal{F}_s\right] dW(s, y). \quad (5.1)$$

The linearisation is based on the same observation as in Section 4, namely, $u^\varepsilon(t, x)$ can be approximated by a quantity that only depends on the noise in a tiny neighborhood of (t, x) , and hence decouples from the numerator in the right-hand side of (5.1) as well as the σ -field \mathcal{F}_s . Once the linearisation is achieved, the Edwards-Wilkinson limit is obtained via the *second order Poincaré inequality* [Cha09, NPR09], which states that given a random variable X defined from a white noise W on a domain Λ , the total variation distance between the distribution of X and that of a normal random variable ζ with matching mean and variance can be bounded by

$$d_{TV}(X, \zeta) \leq C \mathbb{E}[\|DX\|_H^4]^{1/4} \mathbb{E}[\|D^2 X\|_{\text{op}}^4]^{1/4}, \quad (5.2)$$

where D is again the Malliavin derivative, H is the L^2 space of functions on Λ , and $\|\cdot\|_{\text{op}}$ is the operator norm of the Hessian operator $D^2 X$. The fact that the second order Poincaré inequality involves a fourth moment is why the result of [Gu20] is restricted to a subset of the subcritical regime.

We also mention that, spatial averages of the solution of the mollified 2d KPZ equation on the mesoscopic scale have been shown in [Tao24b] to also have Gaussian fluctuations. In [NN23], Theorem 4.1 has been generalised to the fluctuation field of $F(u^\varepsilon(t, x))$ with general initial condition $u^\varepsilon(0, \cdot)$ and for a suitable class of functions F that includes $F(z) = \log z$.

5.4. NONLINEAR STOCHASTIC HEAT EQUATIONS. A program to study nonlinear stochastic heat equations at the critical dimension two was initiated in [DG22] and studied further in [Tao24a, DG23, DG24]. In [DG22, Tao24a], they consider SPDEs of the form

$$\partial_t u^\varepsilon = \frac{1}{2} \Delta u^\varepsilon + \frac{\sqrt{2\pi}}{\sqrt{|\log \varepsilon|}} \sigma(u^\varepsilon) \xi^\varepsilon, \quad t > 0, x \in \mathbb{R}^2, \quad (5.3)$$

where $\sigma : [0, \infty) \rightarrow [0, \infty)$ is a globally Lipschitz function with Lipschitz constant $\text{Lip}(\sigma) < 1$, which is below the critical value $\hat{\beta}_c = 1$ of the linear SHE.

In [DG22], the analogue of Theorem 5.1 on the limit of one-point distribution was established. More precisely, if $u^\varepsilon(0, \cdot) \equiv a \geq 0$, then the solution $u^\varepsilon(t, x)$ of (5.3) at any $t > 0$ and $x \in \mathbb{R}^2$ converges in distribution to a random variable $\Xi_{a,Q}(Q)$ with $Q = 2$, which solves the following forward-backward SDE (FBSDE):

$$\begin{aligned} d\Xi_{a,Q}(q) &= J(Q - q, \Xi_{a,Q}(q)) dB(q), \quad q \in (0, Q], \\ \Xi_{a,Q}(0) &= a, \\ J(r, b) &= \frac{1}{\sqrt{2}} \mathbf{E}[\sigma^2(\Xi_{b,r}(r))]^{1/2}, \end{aligned} \quad (5.4)$$

where $(B(q))_{q \geq 0}$ is a one-dimensional Brownian motion, and a is the initial condition of the FBSDE $\Xi_{a,Q}(\cdot)$, with terminal time Q . When $\sigma(u) = cu$ with $c < 1$, $\Xi_{a,2}(2)$ is log-normal and this recovers Theorem 5.1.

The connection between (5.3) and the FBSDE (5.4) is that for $Q \in (0, 2]$, $u^\varepsilon(\varepsilon^{2-Q}, x) \sim \Xi_{a,Q}(Q)$. Furthermore, given $Q \in (0, 2]$, one has

$$u^\varepsilon(\varepsilon^{2-Q} - \varepsilon^{2-(Q-q)}, x) \approx \Xi_{a,Q}(q), \quad q \in [0, Q],$$

where $\varepsilon^{-(Q-q)}$ corresponds to the exponential time scale $N^{\frac{Q-q}{2}}$ in Corollary 2.2 with $N = \varepsilon^{-2}$. Such exponential time scales for the nonlinear SHE are motivated by the same type of second moment calculations as in Corollary 2.2. The heuristic justification for (5.4) is the following. Denote $T = \varepsilon^{2-Q}$ and $r = Q - q$. By second moment calculations, a small error is incurred if we turn off the noise ξ^ε (i.e., evolve according to the heat equation) in the time interval $(T - \varepsilon^{2-r}, T - \varepsilon^{2-r+\delta})$ for a very small $\delta > 0$. Smoothing by the heat flow on this time interval allows one to replace $u^\varepsilon(T - \varepsilon^{2-r+\delta}, \cdot)$ by a random constant distributed roughly as $\Xi_{a,Q}(q) = \Xi_{a,Q}(Q - r)$, which becomes the initial condition for the evolution of u^ε from time $T - \varepsilon^{2-r+\delta}$ onward. Consider the evolution of u^ε till a later time $T - \varepsilon^{2-r+\Delta}$ for some small $\Delta > \delta$. The time duration of this evolution is $T - \varepsilon^{2-r+\Delta} - (T - \varepsilon^{2-r+\delta}) \approx \varepsilon^{2-r}$, and therefore to compute the increment of the quadratic variation of $u^\varepsilon(\cdot, x)$ in this time interval, it suffices to consider (5.3) with initial condition $\Xi_{a,Q}(Q - r)$ and time duration ε^{2-r} , which leads to the SDE in (5.4) with diffusion coefficient $J(r, \Xi_{a,Q}(Q - r))$.

The analogue of Theorem 5.2 on the Edwards-Wilkinson limit for the non-linear SHE (5.3) was established in [Tao24a].

In [DG23, DG24], the results in [DG22] and [Tao24a] were extended to σ that do not necessarily satisfy the condition $\text{Lip}(\sigma) < 1$. For technical convenience, they considered a variant of (5.3) where smoothing is performed on $\sigma(u)\xi$ instead of on ξ , but the result is expected to be the same. The key idea of [DG23] is that if $\text{Lip}(\sigma) =: \sqrt{A} > 1$, then the result of [DG22] applies to u^ε up to a time scale of $\tilde{\varepsilon}^2 := \varepsilon^{2-\frac{2}{A}}$ with the corresponding spatial scale $\tilde{\varepsilon} = \varepsilon^{1-\frac{1}{A}}$. Heuristically, if one performs an averaging on this space-time scale, then $\tilde{\varepsilon}$ replaces ε as the basic spatial scale of smoothing, while $J(2/A, \cdot) =: \tilde{\sigma}(\cdot)$ replaces $\sigma(\cdot)$ as the renormalised diffusion function. If $\text{Lip}(\tilde{\sigma}) =: \sqrt{\tilde{A}} < \infty$, then one can apply [DG22] to this renormalised non-linear SHE up to a time scale of $\tilde{\varepsilon}^{2-\frac{2}{A}} = (\varepsilon^{(2-\frac{2}{A})})^{1-\frac{1}{A}} \gg \tilde{\varepsilon}^2$. This procedure can then be iterated to find the maximal time scale ε^{2-Q} at which one can identify the distributional limit of $u^\varepsilon(\varepsilon^{2-Q}, x)$. The corresponding Edwards-Wilkinson limit was obtained in [DG24].

5.5. ANISOTROPIC KPZ AND RELATED MODELS. The general 2-dimensional KPZ is formally written as

$$\partial_t h = \frac{1}{2} \Delta h + \langle \nabla h, Q \nabla h \rangle + \xi, \quad t > 0, x \in \mathbb{R}^2, \quad (5.5)$$

where $\langle \cdot, \cdot \rangle$ is the Euclidean inner product, ξ is the usual space-time white noise, and Q is a 2×2 matrix. When $Q = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, it is called the purely anisotropic KPZ (AKPZ). This model emerged in the physics literature in [Wol91]. It was shown in [CET23] that the correlation length of the interface grows like $\sqrt{t(\log t)^{1/2}}$.

Although the 2-dimensional anisotropic KPZ is also a critical singular SPDE, just as the 2-dimensional SHE and isotropic KPZ, the techniques are very different. We refer to the lecture notes by Cannizzaro and Toninelli [CT24] for more details, where their techniques apply to a family of critical and supercritical singular SPDEs, including the stochastic Navier Stokes equation with divergence-free driving noise, Burgers equation, diffusion in the curl of a 2d Gaussian Free Field, self-repelling Brownian polymers, etc.

The weak coupling limit of (5.5), with the same noise strength as in Theorem 4.1 for the isotropic KPZ, has also been studied [CES21, CET23]. More precisely, in [CES21, CET23], instead of mollifying the noise, all Fourier modes beyond order $1/\varepsilon$ are truncated, which is an alternative method of performing the so-called ultraviolet cutoff to smoothe things at spatial scale ε . The strength of non-linearity in the AKPZ (c.f. (1.6)) is then scaled down as $\lambda = \sqrt{\hat{\lambda}/\log \varepsilon^{-1}}$. The random field fluctuation of the AKPZ in this weak coupling regime is shown to converge to the solution of the following Edwards-Wilkinson equation:

$$\partial_t h = \frac{\nu_{\text{eff}}(\hat{\lambda})}{2} \Delta h + \sqrt{\nu_{\text{eff}}(\hat{\lambda})} \xi, \quad \text{with} \quad \nu_{\text{eff}}(\hat{\lambda}) = \sqrt{\frac{2}{\pi} \hat{\lambda}^2 + 1}. \quad (5.6)$$

We note that, in contrast to the isotropic 2d KPZ treated in Theorem 4.1, the coefficient of Δh is also modified. But there is no phase transition in $\hat{\lambda}$ in the sense that the strength of the limiting noise ξ is bounded for all $\hat{\lambda} > 0$. A key property needed in the analysis of the AKPZ is that it has an explicit Gaussian invariant measure (the Gaussian free field). This is a common feature shared by models in this class. Again, we refer to the lecture notes [CT24] for more details.

5.6. THE QUASI-CRITICAL REGIME. The fluctuations of the partition functions *as a random field* have very different features in the sub-critical and critical regimes:

- in the *sub-critical regime* $\hat{\beta} < 1$, fluctuations are *asymptotically Gaussian* after centering and *rescaling* by $\beta_N^{-1} \sim \hat{\beta}^{-1} \sqrt{\log N}$, see Theorem 3.3:

$$\sqrt{\log N} \int_{\mathbb{R}^2} \phi(x) \left(Z_{tN}^{\beta_N}([x\sqrt{N}]) - 1 \right) dx \xrightarrow[N \rightarrow \infty]{d} \frac{\hat{\beta}}{\sqrt{1 - \hat{\beta}^2}} \int_{\mathbb{R}^2} \phi(x) \hat{v}(t, x) dx, \quad (5.7)$$

where $\partial_t \hat{v}(t, x)$ solves a normalised version of the Edwards-Wilkinson equation (3.9), namely $\partial_t \hat{v}(t, x) = \frac{1}{2} \Delta \hat{v}(t, x) + \xi(t, x) \hat{v}(t, x)$ with $\partial_t \hat{v}(0, \cdot) = 1$ (we moved the scaling factor $(1 - \hat{\beta}^2)^{-1/2}$ into (5.7) for later comparison, see below);

- on the other hand, in the *critical regime* $\hat{\beta} = 1$, and actually in the whole critical window $\hat{\beta}^2 = 1 + O(\frac{1}{\log N})$, see (1.16), a *non-Gaussian* scaling limit emerges with *no need of centering and rescaling*, see Theorem 6.1:

$$\int_{\mathbb{R}^2} \phi(x) Z_{tN}^{\beta_N}([x\sqrt{N}]) dx \xrightarrow[N \rightarrow \infty]{d} \int_{\mathbb{R}^2} \varphi(x) \text{SHF}_{0,t}^\vartheta(dx, \mathbb{R}^2),$$

where $\text{SHF}_{s,t}^\vartheta(dx, dy)$ is the critical 2d stochastic heat flow, see Section 6.

To interpolate between these regimes, we can let $\hat{\beta} = \hat{\beta}(N) \uparrow 1$ from below the critical window. In view of (1.16), this amounts to taking $\beta = \beta_N$ in the *quasi-critical regime* defined by:

$$\sigma_{\beta_N}^2 = \frac{1}{R_N} \left(1 - \frac{\vartheta_N}{\log N} \right) \quad \text{with} \quad 1 \ll \vartheta_N \ll \log N. \quad (5.8)$$

It is proved in [CCR23] that throughout this quasi-critical regime, the partition function as a random field has *Gaussian fluctuations*, after centering and *rescaling by $\sqrt{\vartheta_N}$* , a factor which interpolates between $\sqrt{\log N}$ and 1:

$$\sqrt{\vartheta_N} \int_{\mathbb{R}^2} \phi(x) \left(Z_{tN}^{\beta_N}(\lfloor x\sqrt{N} \rfloor) - 1 \right) dx \xrightarrow[N \rightarrow \infty]{d} \int_{\mathbb{R}^2} \phi(x) \hat{v}(t, x) dx, \quad (5.9)$$

where $\hat{v}(t, x)$, as in (5.7), solves the Edwards-Wilkinson equation.[†] This shows that non-Gaussian behavior does not appear below the critical window. Note that setting formally $\hat{\beta}^2 = 1 - \frac{\vartheta_N}{\log N}$ in (5.7), as prescribed by (5.8), we would obtain (5.9).

The strategy to prove (5.9), inspired by [CC22], is to decompose the space-averaged partition function approximately into a sum of independent random variables and then apply a Central Limit Theorem under a Lyapunov condition, which requires estimating moments of the partition function of order higher than two. A key difficulty is the lack of hypercontractivity, which holds in the sub-critical regime $\hat{\beta} < 1$, but not in the quasi-critical regime (5.8) (nor in the critical regime $\hat{\beta} = 1$), since the main contribution now comes from terms of diverging order in the chaos expansion. To this end, novel moment estimates are derived in [CC22] by exploiting and extending the strategy in [CSZ23a, LZ23].

5.7. HIGHER DIMENSIONS. The Directed Polymer Model, the Stochastic Heat Equation and the KPZ equation have also been studied in dimensions $d \geq 3$. A difference in the higher dimensional setting is that, contrary to the two-dimensional case, the L^2 critical point (at which the second moment of the polymer partition function diverges) and the true critical point (where the polymer undergoes a localisation-delocalisation phase transition) do not coincide. This is a phenomenon known for some time in the directed polymer literature, see the references in [Com17, Zyg24]. The analogue of Theorems 3.3, 5.2, and 4.1 have been established in dimensions $d \geq 3$ below the L^2 critical point [MU18, CCM24, CCM20, CNN22, LZ22], which requires scaling of the noise strength in mollified SHE and KPZ by a factor of $\varepsilon^{\frac{d-2}{2}}$ (equivalent to no scaling in the inverse temperature β in the directed polymer model). Characterising the random field in $d \geq 3$ beyond the L^2 regime remains a challenging task. Some progress has been made in the directed polymer setting by Junk [Jun22, Jun23], and the equivalence of the fluctuation for the mollified SHE and the mollified KPZ has recently been established in [JN24], where the limiting field is believed to solve an analogue of the Edwards-Wilkinson equation driven by a stable noise field.

5.8. SPACE-TIME MOLLIFICATIONS OF NOISE. So far, we have been discussing spatial mollifications of the noise, which is convenient because the noise is still white in time and hence we can define the solution of the mollified SHE (1.3) in terms of Itô integral. However, it is also meaningful to consider a regularisation where the noise is mollified in both space and time. This approach was considered in [GRZ18] for the SHE in dimensions $d \geq 3$ and in [Kot24] in $d = 2$. Even though there is still Edwards-Wilkinson limit in the weak disorder regime, there are additional difficulties because the time mollification destroys the martingale structure of the process and the path measures are tilted by self-intersections local times. We refer to [GRZ18, Kot24] for details.

[†]To compare with [CCR23, Theorem 1.1], note that the RHS of (5.9) is a Gaussian random variable with variance $\sigma_{t,\phi}^2 := \int_{\mathbb{R}^2 \times \mathbb{R}^2} \phi(x) K_t(x, x') \phi(x') dx dx'$ with $K_t(x, x') := \int_0^t \frac{1}{2u} e^{-\frac{|x-x'|^2}{2u}} du$.

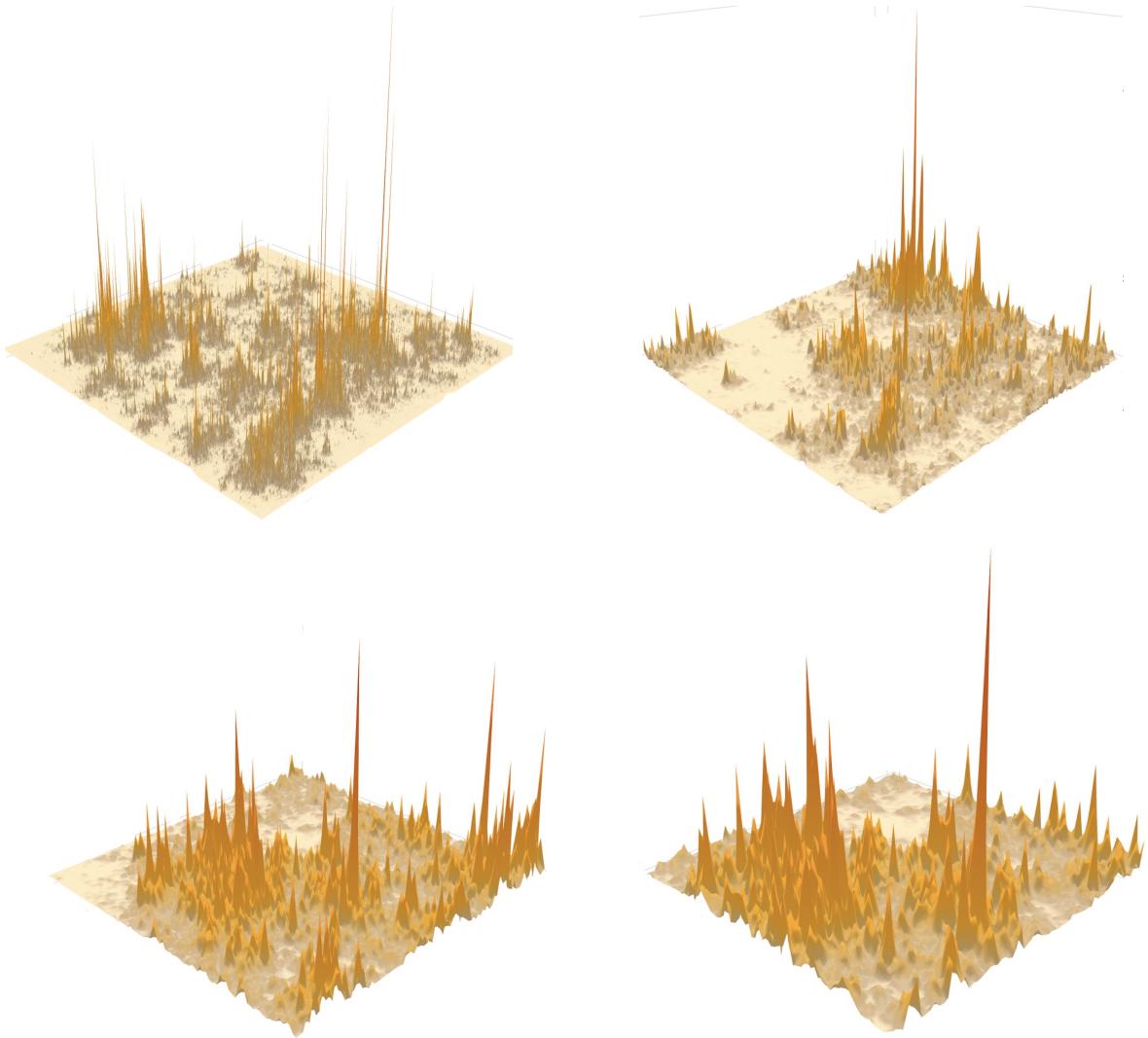


FIGURE 2. The pictures above illustrate the critical 2d stochastic heat flow, $\text{SHF}(\vartheta)$, and its scaling properties. The top-right picture and the two bottom pictures are zoomed-in snapshots of the upper-left picture, which is a simulation of $\text{SHF}(\vartheta)$ with $\vartheta = 0$. The singularity of the critical 2d SHF w.r.t. the Lebesgue measure can be seen from the upper-left picture. The zoomed-in snapshots demonstrate the scaling covariance of $\text{SHF}(\vartheta)$ stated in (10.1). In particular, the zoomed-in fields appear to be smoother than the original field, since (10.1) shows that zooming in has the effect of decreasing the disorder strength ϑ .

6. THE CRITICAL 2d SHF

6.1. MAIN RESULT. We assume throughout this section that β_N and σ_N^2 have been chosen in the critical window as in (1.16) for some $\vartheta \in \mathbb{R}$.

Recall the family of point-to-point partition functions $Z_{M,N}^{\beta_N}(x,y)$ from (1.11). To keep the notation simple, we will omit β_N from $Z_{M,N}^{\beta_N}$. We regard $Z_{M,N}(\cdot,\cdot)$ as a random measure

on $\mathbb{Z}^2 \times \mathbb{Z}^2$. Scaling space and time diffusively and also the mass of the random measure, we define the family of rescaled random measures $\mathcal{Z}_N := (\mathcal{Z}_{N;s,t}(dx, dy))_{0 \leq s \leq t}$ by

$$\iint_{\mathbb{R}^2 \times \mathbb{R}^2} \varphi(x)\psi(y)\mathcal{Z}_{N;s,t}(dx, dy) := \frac{1}{N} \sum_{x,y \in \mathbb{Z}^2} \varphi\left(\frac{x}{\sqrt{N}}\right)\psi\left(\frac{y}{\sqrt{N}}\right)Z_{[Ns],[Nt]}(x, y), \quad (6.1)$$

where $\varphi, \psi \in C_c(\mathbb{R}^2)$ are test functions with compact support. We regard $\mathcal{Z}_{N;s,t}(dx, dy)$ as a random variable taking values in the space of Radon measures $\mathcal{M}(\mathbb{R}^2 \times \mathbb{R}^2)$ on $\mathbb{R}^2 \times \mathbb{R}^2$, equipped with the vague topology.

The main result from [CSZ23a] is that in the critical window, the rescaled partition function $\mathcal{Z}_{N;s,t}(dx, dy)$, regarded as a stochastic process in the space $\mathcal{M}(\mathbb{R}^2 \times \mathbb{R}^2)$ indexed by $0 \leq s \leq t < \infty$, converges in law to a unique limit, that we denote by $\text{SHF}_{s,t}^\vartheta(dx, dy)$.

Theorem 6.1 (Critical 2d SHF). *Let β_N be chosen in the critical window (1.16) for some $\vartheta \in \mathbb{R}$. As $N \rightarrow \infty$, the family of random measures $\mathcal{Z}_N = (\mathcal{Z}_{N;s,t}(dx, dy))_{0 \leq s \leq t}$ converges in finite dimensional distribution to a unique limit*

$$\text{SHF}(\vartheta) := (\text{SHF}_{s,t}^\vartheta(dx, dy))_{0 \leq s \leq t < \infty},$$

called the critical 2d stochastic heat flow. The limit $\text{SHF}(\vartheta)$ does not depend on the law of the disorder ω except for the assumptions in (1.8). The first and second moments are

$$\mathbb{E}[\text{SHF}_{s,t}^\vartheta(dx, dy)] = g_{t-s}(y - x) dx dy, \quad (6.2)$$

$$\text{Cov}[\text{SHF}_{s,t}^\vartheta(dx, dy), \text{SHF}_{s,t}^\vartheta(dx', dy')] = K_{t-s}^\vartheta(x, x'; y, y') dx dy dx' dy', \quad (6.3)$$

with the kernel K_{t-s}^ϑ defined by

$$K_t^\vartheta(x, x'; y, y') := 4\pi g_{\frac{t}{2}}\left(\frac{y+y'}{2} - \frac{x+x'}{2}\right) \iint_{0 < a < b < t} g_{2a}(x' - x) G_\vartheta(b - a) g_{2(t-b)}(y' - y) da db, \quad (6.4)$$

where g is the heat kernel in \mathbb{R}^2 defined in (2.5), and G_ϑ is defined in (8.12)-(8.13).

Remark 6.2. The original definition of the Stochastic Heat Flow in [CSZ23a, Theorem 1.1], denoted by \mathcal{Z}^ϑ , was derived from directed polymers with simple symmetric random walk, which is periodic and has variance $\frac{1}{2}$ (in each component). The definition of $\text{SHF}(\vartheta)$ here is slightly different, because the random walk S is aperiodic and has unit variance, which matches the scaling limit of the solution of the mollified SHE. To match the two definitions, it suffices to rescale space, time, and the measure in order to match the first two moments: compare (6.2)-(6.4) with [CSZ23a, eqs. (1.13) and (3.56)]. This leads to the identification

$$\text{SHF}_{s,t}^\vartheta(dx, dy) \stackrel{d}{=} 4 \mathcal{Z}_{s,t}^\vartheta(d\frac{x}{\sqrt{2}}, d\frac{y}{\sqrt{2}}) \stackrel{d}{=} 2 \mathcal{Z}_{2s,2t}^{\vartheta-\log 2}(dx, dy), \quad (6.5)$$

where the second equality holds by scaling covariance, see [CSZ23a, Theorem 1.2] and Theorem 10.1 below, and $\nu(dx, dy) = \mu(d\frac{x}{\sqrt{2}}, d\frac{y}{\sqrt{2}})$ denotes the measure $\nu(A) := \mu(\frac{1}{\sqrt{2}}A)$ for $A \subseteq \mathbb{R}^2 \times \mathbb{R}^2$. Relation (6.5) corrects a misprint in [CSZ23a, Remark 1.5].

Remark 6.3. The covariance kernel $K_{t-s}^\vartheta(x, x'; y, y')$ was first identified in [BC98] (see also [CSZ19b]) and is logarithmically divergent near the diagonals $x = x'$ or $y = y'$. One can also express $K_{t-s}^\vartheta(x, x'; y, y')$ in terms of the Volterra function, see [CM23, CM24] and the references therein.

In what follows, we will sketch the proof Theorem 6.1 in [CSZ23a], which was carried out without any axiomatic characterisation of the limit SHF(ϑ). This difficulty was overcome by comparing with a family of coarse-grained models and then applying a Lindeberg principle. Recently, Tsai [Tsa24] gave an axiomatic characterisation of SHF(ϑ) and then applied it to show that the solution of the mollified 2d SHE in the critical window also converges to SHF(ϑ). We will discuss this in more detail in Section 10.3.

6.2. UNIVERSALITY VIA A LINDEBERG PRINCIPLE. To illustrate some of the proof ingredients for Theorem 6.1 in a simpler setting, we sketch here how one might apply a Lindeberg principle to show that the limit of the random measures \mathcal{Z}_N in Theorem 6.1, *if it exists*, does not depend on the law of the disorder variables $\omega = (\omega(n, z))_{(n,z) \in \mathbb{N} \times \mathbb{Z}^2}$.

Let us consider the averaged point-to-plane partition function

$$\mathcal{Z}_N(\varphi) := \frac{1}{N} \sum_{x \in \mathbb{Z}^2} \varphi\left(\frac{x}{\sqrt{N}}\right) Z_N^{\beta_N}(x), \quad \varphi \in C_c(\mathbb{R}^2), \quad (6.6)$$

which equals $\iint \varphi(x) \mathcal{Z}_{N;0,1}(dx, dy)$ in the notation introduced in (6.1).

By (2.2), $\mathcal{Z}_N(\varphi)$ admits the polynomial chaos expansion

$$\mathcal{Z}_N(\varphi) = \bar{\varphi}_N + \frac{1}{N} \sum_{r=1}^{\infty} \sum_{\substack{z_0, z_1, \dots, z_r \in \mathbb{Z}^2 \\ n_0 := 0 < n_1 < \dots < n_r < N}} \varphi\left(\frac{z_0}{\sqrt{N}}\right) \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i), \quad (6.7)$$

where

$$\bar{\varphi}_N := \frac{1}{N} \sum_{z_0 \in \mathbb{Z}^2} \varphi\left(\frac{z_0}{\sqrt{N}}\right). \quad (6.8)$$

Note that $\mathcal{Z}_N(\varphi)$ is a function of the i.i.d. random variables $\xi_N := (\xi_N(n, x))_{n,x \in \mathbb{N} \times \mathbb{Z}^2}$.

A Lindeberg principle is said to hold when the law of a function $\Psi(\zeta)$ of a family of random variables $\zeta = (\zeta_i)_{i \in \mathbb{T}}$ does not change much if ζ is replaced by another family of random variables $\eta = (\eta_i)_{i \in \mathbb{T}}$ with matching first few moments. Lindeberg principles provide powerful tools to prove universality, the simplest instance being the universality of the Gaussian distribution from the central limit theorem. The usual formulation such as in [Cha06] requires the family of random variables to be independent (or exchangeable), and Ψ needs to have bounded first three partial derivatives. This is not satisfied when Ψ is a multilinear polynomial, whose derivatives are unbounded. This case was treated in [Rot79, MOO10] (see also [CSZ17a]). Although the results in these references are not directly applicable to $\mathcal{Z}_N(\varphi)$ for reasons that we will explain shortly, we can still apply the general approach. We first recall the Lindeberg principle from [MOO10].

Lindeberg principle [MOO10]. Consider a function $\Psi(x)$ defined as a sum of multilinear polynomials in the variables $x = (x_i)_{i \in \mathbb{T}}$ by

$$\Psi(x) := \sum_{I \subset \mathbb{T}, |I| < \infty} \psi(I) \prod_{i \in I} x_i, \quad (6.9)$$

where ψ is a real-valued function defined on finite subsets of the index set \mathbb{T} .

Let $\zeta = (\zeta_i)_{i \in \mathbb{T}}$ and $\eta = (\eta_i)_{i \in \mathbb{T}}$ be two i.i.d. family with mean 0, variance 1, and $\mathbb{E}[|\zeta_i|^3] \vee \mathbb{E}[|\eta_i|^3] < \infty$. Then we note that

$$\mathbb{E}[\Psi(\zeta)] = \mathbb{E}[\Psi(\eta)] = \psi(\emptyset) \quad \text{and} \quad \mathbb{V}\text{ar}(\Psi(\zeta)) = \mathbb{V}\text{ar}(\Psi(\eta)) = \sum_{I \subset \mathbb{T}, 1 \leq |I| < \infty} \psi(I)^2.$$

Without loss of generality, we assume that Ψ has been normalised with $\text{Var}[\Psi(\zeta)^2] = 1$. To identify conditions that ensure $\Psi(\zeta)$ and $\Psi(\eta)$ are close in distribution, we consider

$$\mathbb{E}[f(\Psi(\eta))] - \mathbb{E}[f(\Psi(\zeta))] \quad \text{for } f \in C_b^3(\mathbb{R}),$$

i.e., bounded continuous test functions with bounded first three derivatives.

Assume for simplicity that \mathbb{T} is a finite set with elements labelled by $1, \dots, |\mathbb{T}|$. Then we replace ζ_i by η_i one variable at a time and for each $0 \leq k \leq |\mathbb{T}|$, define the k -th interpolating family by $\zeta^{(k)}$ with $\zeta_i^{(k)} = \eta_i$ for $1 \leq i \leq k$ and $\zeta_i^{(k)} = \zeta_i$ for $i > k$. Note that $\zeta^{(0)} = \zeta$ and $\zeta^{(|\mathbb{T}|)} = \eta$. We can then write the telescopic sum

$$\mathbb{E}[f(\Psi(\eta))] - \mathbb{E}[f(\Psi(\zeta))] = \sum_{k=1}^{|\mathbb{T}|} \mathbb{E}[f(\Psi(\zeta^{(k)})) - f(\Psi(\zeta^{(k-1)}))]. \quad (6.10)$$

Denote $g_{k,\zeta,\eta}(x) := f(\Psi(\eta_1, \dots, \eta_{k-1}, x, \zeta_{k+1}, \dots, \zeta_{|\mathbb{T}|}))$. By Taylor expansion,

$$\begin{aligned} f(\Psi(\zeta^{(k)})) - f(\Psi(\zeta^{(k-1)})) &= g_{k,\zeta,\eta}(\eta_k) - g_{k,\zeta,\eta}(\zeta_k) \\ &= g'_{k,\zeta,\eta}(0)(\eta_k - \zeta_k) + \frac{1}{2}g''_{k,\zeta,\eta}(0)(\eta_k^2 - \zeta_k^2) + R(\eta_k) - R(\zeta_k), \end{aligned}$$

where $R(x)$ is the remainder in the Taylor expansion of $g_{k,\zeta,\eta}(x)$ and satisfies

$$|R(x)| \leq \|g'''_{k,\zeta,\eta}\|_\infty |x|^3 \leq \|f'''\|_\infty |\partial_k \Psi(\zeta^{(k)})|^3 |x|^3.$$

Here $\partial_k \Psi$ denotes the partial derivative w.r.t. x_k and does not depend on x_k – note that because Ψ is multilinear, $\partial_k^2 \Psi = 0$. Substituting the above calculations into (6.10) and using the fact that (ζ_k, η_k) are independent of $(\zeta_i, \eta_i)_{i \neq k}$, we find that the first and second order terms in the Taylor expansion cancel out, and

$$\begin{aligned} |\mathbb{E}[f(\Psi(\eta))] - \mathbb{E}[f(\Psi(\zeta))]| &\leq \|f'''\|_\infty \sum_{k=1}^{|\mathbb{T}|} \mathbb{E}[|\partial_k \Psi(\zeta^{(k)})|^3] (\mathbb{E}[|\zeta_k|^3] + \mathbb{E}[|\eta_k|^3]) \\ &\leq \|f'''\|_\infty \max_k (\mathbb{E}[|\zeta_k|^3] + \mathbb{E}[|\eta_k|^3]) \sum_{k=1}^{|\mathbb{T}|} \mathbb{E}[|\partial_k \Psi(\zeta^{(k)})|^3]. \end{aligned} \quad (6.11)$$

This is the point where the analysis in [MOO10] is no longer adequate when considering the critical case.

Note that for each k , $\partial_k \Psi(\zeta)$ (we replaced $\zeta^{(k)}$ by ζ for simplicity) also admits a polynomial chaos expansion in the variables ζ . In [MOO10], $\mathbb{E}[|\partial_k \Psi(\zeta)|^3]$ is bounded using hypercontractivity (see also [CSZ20, Theorem B.1]), which bounds higher than 2 moments by the second moment with an additional weight c_3^{k-1} on each chaos component of order k . Therefore,

$$\mathbb{E}[|\partial_k \Psi(\zeta)|^3] \leq \left(\sum_{I \subset \mathbb{T}: k \in I} c_3^{2(|I|-1)} \psi(I)^2 \right)^{\frac{3}{2}} \leq c_3^{3\deg(\Psi)} \left(\sum_{I \subset \mathbb{T}: k \in I} \psi(I)^2 \right)^{\frac{3}{2}} = c_3^{3\deg(\Psi)} \text{Inf}_k(\Psi)^{\frac{3}{2}},$$

where $c_3 > 1$ depends only on the law of ζ , $\deg(\Psi) := \max\{|I| : \psi(I) > 0\}$ is the degree of Ψ , and

$$\text{Inf}_k(\Psi) := \mathbb{E}[(\partial_k \Psi(\zeta))^2] = \sum_{I \subset \mathbb{T}: k \in I} \psi(I)^2$$

is called the influence of ζ_k on $\Psi(\zeta)$ in [MOO10]. We can then bound the r.h.s. of (6.11) by

$$\begin{aligned} |\mathbb{E}[f(\Psi(\eta))] - \mathbb{E}[f(\Psi(\zeta))]| &\leq C_{f,\zeta,\eta} c_3^{3\deg(\Psi)} \left(\max_k \text{Inf}_k(\Psi) \right)^{\frac{1}{2}} \sum_{k=1}^{|\mathbb{T}|} \sum_{I \subset \mathbb{T}: k \in I} \psi(I)^2 \\ &= C_{f,\zeta,\eta} c_3^{3\deg(\Psi)} \left(\max_k \text{Inf}_k(\Psi) \right)^{\frac{1}{2}} \sum_{I \subset \mathbb{T}: |I| \geq 1} |I| \psi(I)^2 \\ &\leq C_{f,\zeta,\eta} c_3^{3\deg(\Psi)} \deg(\Psi) \left(\max_k \text{Inf}_k(\Psi) \right)^{\frac{1}{2}} \text{Var}(\Psi(\zeta)). \end{aligned} \quad (6.12)$$

When we consider a sequence Ψ_N with uniformly bounded degree and variance, the r.h.s. of (6.12) tends to 0 once we show that the maximum influence $\max_k \text{Inf}_k(\Psi_N) \rightarrow 0$ as $N \rightarrow \infty$.

The argument above fails when applied to the averaged partition function $\mathcal{Z}_N(\varphi)$ in the critical window, because we will see in Section 8 that in the polynomial chaos expansion of $\mathcal{Z}_N(\varphi)$ in (6.7), the dominant contribution comes from terms of degree $A \log N$ where $A > 0$ can be arbitrarily large. Therefore in (6.12), $c_3^{3\deg(\Psi)} \approx c_3^{3A \log N}$ can overwhelm $\max_k \text{Inf}_k(\Psi_N)$, and we will need sharper bounds than provided by hyper-contractivity.

Applying Lindeberg to $\mathcal{Z}_N(\varphi)$. Our starting point is (6.11), which was actually derived only using the assumption that for each $i \in \mathbb{T}$, η_k and ζ_k have matching first two moments. Recall the i.i.d. family $\xi_N = (\xi_N(n, z))_{(n,z) \in \mathbb{N} \times \mathbb{Z}^2}$ from (2.1), and denote

$$\Psi_N(\xi_N) := \mathcal{Z}_N(\varphi) = \bar{\varphi}_N + \frac{1}{N} \sum_{r=1}^{\infty} \sum_{\substack{z_0, z_1, \dots, z_r \in \mathbb{Z}^2 \\ n_0 := 0 < n_1 < \dots < n_r < N}} \varphi\left(\frac{z_0}{\sqrt{N}}\right) \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i).$$

To show that the distribution of $\Psi_N(\xi_N)$ hardly changes if we replace ξ_N by another family ζ_N with matching first two moments (and fourth moment of the same order), we apply the bound in (6.11), where the index set is now $\mathbb{T} = \mathbb{N} \times \mathbb{Z}^2$. For each $(n, z) \in \mathbb{T}$, we observe the factorisation

$$\partial_{(n,z)} \Psi_N(\xi_N) = \frac{1}{N} Z_{0,n}(\varphi_N, z) Z_{n,N}(z, \mathbf{1}), \quad (6.13)$$

where $\mathbf{1}(z) \equiv 1$, $\varphi_N(z) := \varphi(z/\sqrt{N})$ for $z \in \mathbb{Z}^2$, and for any $\psi : \mathbb{Z}^2 \rightarrow \mathbb{R}$ and $m < n$,

$$Z_{m,n}(\psi, z) = \sum_{x \in \mathbb{Z}^2} \psi(x) Z_{m,n}(x, z) \quad \text{and} \quad Z_{m,n}(z, \psi) = \sum_{y \in \mathbb{Z}^2} Z_{m,n}(z, y) \psi(y)$$

are the point-to-plane partition functions with the free endpoint weighted by the function ψ (recall the point-to-point partition function $Z_{m,n}(x, y)$ from (2.3)).

Note that $Z_{0,n}(\varphi_N, z)$ and $Z_{n,N}(z, \mathbf{1})$ are independent, while $Z_{0,n}(\varphi_N, z)$ has the same distribution as $Z_{0,n}(z, \varphi_N)$ and $Z_{n,N}(z, \mathbf{1})$ has the same distribution as $Z_{0,N-n}(z, \mathbf{1})$. Therefore we have

$$\begin{aligned} &|\mathbb{E}[f(\Psi_N(\xi_N))] - \mathbb{E}[f(\Psi_N(\zeta_N))]| \\ &\leq \|f'''\|_{\infty} \max_{(n,z)} (\mathbb{E}[|\xi_N(n, z)|^3] + \mathbb{E}[|\zeta_N(n, z)|^3]) \sum_{(n,z) \in \mathbb{N} \times \mathbb{Z}^2} \mathbb{E}[|\partial_{(n,z)} \Psi(\zeta_N^{(n,z)})|^3] \\ &\leq \frac{\|f'''\|_{\infty}}{N^3} \max_{(n,z)} (\mathbb{E}[|\xi_N(n, z)|^3] + \mathbb{E}[|\zeta_N(n, z)|^3]) \sum_{(n,z) \in \mathbb{N} \times \mathbb{Z}^2} \mathbb{E}[|Z_{0,n}(z, \varphi_N)|^3] \mathbb{E}[|Z_{0,N-n}(z, \mathbf{1})|^3], \end{aligned} \quad (6.14)$$

where $\zeta_N^{(n,z)}$ is a family that interpolates between ξ_N and ζ_N as in $\zeta^{(k)}$ in (6.10).

As will be discussed in Section 9, the key bound is that for any $\delta > 0$, we have

$$\mathbb{E}[Z_{0,N-n}(z, \mathbf{1})^4] \leq CN^\delta \quad \text{and} \quad \mathbb{E}[Z_{0,n}(z, \varphi_N)^4] \leq Ce^{-\frac{c|z|}{\sqrt{N}}} N^\delta, \quad (6.15)$$

where the exponential factor $e^{-\frac{c|z|}{\sqrt{N}}}$ comes from the assumption that $\varphi_N = \varphi(\cdot/\sqrt{N})$ is supported on a subset of \mathbb{Z}^2 within distance $C\sqrt{N}$ of the origin. Substituting these bounds into (6.14) and using the fact that

$$\mathbb{E}[|\xi_N(n, z)|^3] \leq \mathbb{E}[\xi_N(n, z)^4]^{3/4} \leq C\sigma_N^3 \leq \frac{C}{(\log N)^3}, \quad (6.16)$$

which follows easily from the definition of ξ_N in (2.1), this gives an upper bound of order $\frac{\sigma_N^3}{N^3} N^{2+2\delta}$, with the factor $\frac{1}{N^3}$ coming from the normalisation in (6.14). This upper bound tends to 0 as $N \rightarrow \infty$ if we choose $\delta < 1/2$.

Although this argument does not address whether $\mathcal{Z}_N(\varphi)$ converges in distribution to a unique limit, which is the heart of Theorem 6.1, it does bring out several key ingredients in the proof of Theorem 6.1, namely, the Lindeberg principle that led to (6.14), and the necessary moment bounds (6.15) for the partition function. We will see more complex variants of these when we sketch the proof of Theorem 6.1 in Section 7.

6.3. PROOF OUTLINE OF THEOREM 6.1. We now sketch the proof of Theorem 6.1 – the main result of [CSZ23a]. The key difficulty is the lack of a characterisation of the limit, i.e., the critical 2d stochastic heat flow (SHF). As will be discussed in Section 9, earlier results [CSZ19b, GQT21, Che24b] established the convergence of all positive integer moments of the averaged point-to-point partition functions

$$\mathcal{Z}_N^{\beta_N}(\varphi, \psi) := \frac{1}{N} \sum_{x \in \mathbb{Z}^2} \varphi\left(\frac{x}{\sqrt{N}}\right) Z_{0,N}^{\beta_N}(x, y) \psi\left(\frac{x}{\sqrt{N}}\right), \quad \varphi \in C_c(\mathbb{R}^2), \psi \in C_b(\mathbb{R}^2), \quad (6.17)$$

which equals $\iint \varphi(x) \psi(y) \mathcal{Z}_{N;0,1}(dx, dy)$ in the notation introduced in (6.1). However, as shown in [CSZ23b], the limiting moments $\lim_{N \rightarrow \infty} \mathbb{E}[\mathcal{Z}_N^{\beta_N}(\varphi, \psi)^k]$ diverge too fast in $k \in \mathbb{N}$ to uniquely characterise the distributional limit of $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$.

Our strategy is to show that the laws of $(\mathcal{Z}_N^{\beta_N}(\varphi, \psi))_{N \in \mathbb{N}}$ form a *Cauchy sequence*, i.e.,

$$\mathcal{Z}_M^{\beta_M}(\varphi, \psi) \text{ and } \mathcal{Z}_N^{\beta_N}(\varphi, \psi) \text{ are close in distribution for large } M, N \in \mathbb{N}. \quad (6.18)$$

We took inspiration from the work of Kozma [Koz07] on the convergence of the loop erased random walk on \mathbb{Z}^3 , which also lacked a characterisation of the scaling limit.

To establish (6.18), for each $\varepsilon \in (0, 1)$, we will define a *coarse-grained partition function* $\mathcal{Z}_\varepsilon^{(cg)}(\varphi, \psi | \Theta)$, which has a similar structure to $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ and admits a polynomial chaos expansion w.r.t. a family of random variables Θ . We will use $\mathcal{Z}_\varepsilon^{(cg)}$ as a bridge between $\mathcal{Z}_M^{\beta_M}$ and $\mathcal{Z}_N^{\beta_N}$. More precisely, we first perform a coarse-grained approximation for $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ on the time-space scale $(\varepsilon N, \sqrt{\varepsilon N})$ and show that

$$\mathcal{Z}_N^{\beta_N}(\varphi, \psi) \stackrel{L^2}{\approx} \mathcal{Z}_\varepsilon^{(cg)}(\varphi, \psi | \Theta_{N,\varepsilon}), \quad (6.19)$$

where $\Theta_{N,\varepsilon}$ is a family of weakly dependent coarse-grained disorder variables that depend on N and ε , and the approximation error can be made arbitrarily small in L^2 by choosing ε small and N large.

We then prove (6.18) by showing that

$$\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon}) \stackrel{\text{dist}}{\approx} \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{M,\varepsilon}), \quad (6.20)$$

which follows by applying a Lindeberg principle for sums of multilinear polynomials of a family of weakly dependent random variables Θ . Such a Lindeberg principle was developed in [CSZ23a, Appendix A].

Similar to the discussions in Section 6.2 where we tried to apply a Lindeberg principle directly to $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ as a function of the disorder variables $\xi_N(n, z)$, the key ingredients for the Lindeberg principle are uniform bounds on $\mathbb{E}[|\Theta_{N,\varepsilon}|^3]$, where $\Theta_{N,\varepsilon}$ turns out to be an averaged partition function itself, and similar moment bounds as in (6.15) for the point-to-plane versions of the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$.

Here is a summary of the key proof ingredients:

- A.** *Coarse-Graining*, which leads to *coarse-grained partition functions* with a similar structure to the original polymer partition functions;
- B.** *Time-Space Renewal Structure*, which gives a probabilistic representation for the second moment calculations and leads in the continuum limit to the so-called *Dickman subordinator*;
- C.** *Lindeberg Principle* for multilinear polynomials of *dependent* random variables, which controls the effect of changing Θ in the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta)$;
- D.** *Functional Inequalities for Green's Functions* of multiple random walks on \mathbb{Z}^2 , which yield sharp higher moment bounds for both the original partition functions and the coarse-grained partition functions.

Remark 6.4. *The proof steps outlined above can be applied to models with different microscopic details in the critical window, such as directed polymer models with different underlying random walks, or the solution of the mollified 2d SHE (1.3). As long as space-time is scaled to produce the same diffusion constant, they can all be compared to the same coarse-grained model $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta)$, just with different coarse-grained variables $\Theta_{N,\varepsilon}$. In particular, they should lead to the same limiting SHF \mathcal{Z}^ϑ in Theorem 6.1 if their mean are normalised to the same value, the disorder strength is chosen to produce matching values of $\lim_{N \rightarrow \infty} \mathbb{E}[\Theta_{N,\varepsilon}^2]$ for each $\varepsilon > 0$, and there are uniform bounds on $\mathbb{E}[|\Theta_{N,\varepsilon}|^3]$. See Section 7 for more details.*

7. COARSE GRAINING AND LINDEBERG

7.1. COARSE GRAINING. We now explain how to define the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$, see (6.19), and the family of coarse-grained disorder variables $\Theta_{N,\varepsilon}$. We will sketch the key ideas. The more precise details can be found in [CSZ23a, Section 4].

Given $\varepsilon \in (0, 1)$, we partition space-time into *mesoscopic* blocks indexed by $i \in \mathbb{N}$ and $a = (a_1, a_2) \in \mathbb{Z}^2$:

$$\mathcal{B}_{\varepsilon N}(i, a) := \underbrace{((i-1)\varepsilon N, i\varepsilon N]}_{\mathcal{T}_{\varepsilon N}(i)} \times \underbrace{((a - (1, 1))\sqrt{\varepsilon N}, a\sqrt{\varepsilon N}]}_{\mathcal{S}_{\varepsilon N}(a)} \cap \mathbb{Z}^3, \quad (7.1)$$

where

$$((a - (1, 1))\sqrt{\varepsilon N}, a\sqrt{\varepsilon N}] := ((a_1 - 1)\sqrt{\varepsilon N}, a_2\sqrt{\varepsilon N}] \times ((a_2 - 1)\sqrt{\varepsilon N}, a_2\sqrt{\varepsilon N}].$$

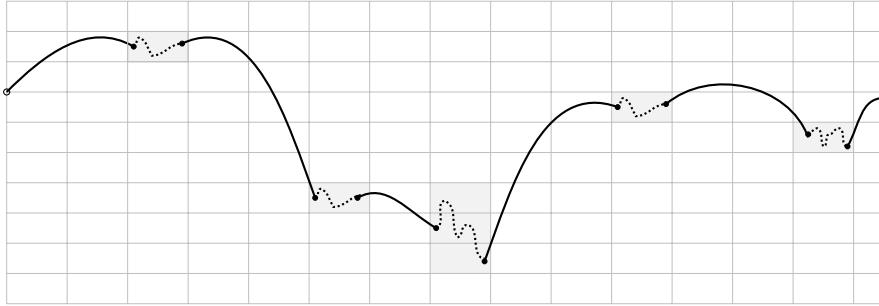


FIGURE 3. Space-time is partitioned into mesoscopic boxes $\mathcal{B}_{\varepsilon N}(i, a)$ as in (7.1). Given $\Gamma = ((n_1, z_1), \dots, (n_r, z_r))$, the sequence of space-time points associated with a given term in the chaos expansion in (7.2), the boxes containing the dotted lines are the ones that intersect Γ . The starting point $(d, x) \in \Gamma$ and the end point $(f, y) \in \Gamma$ of each dotted line are the points of entry and exit in Γ for that mesoscopic time interval $\mathcal{T}_{\varepsilon N}(i)$. Each dotted line contributes a factor $X_{d,f}(x, y)$ to the corresponding term in the chaos expansion, while each solid line contributes a random walk transition kernel.

By (6.17) and (2.3), we see that $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ admits a polynomial chaos expansion in the disorder variables ξ_N :

$$\begin{aligned} \mathcal{Z}_N^{\beta_N}(\varphi, \psi) &= q_{0,N}^N(\varphi, \psi) \\ &+ \frac{1}{N} \sum_{r=1}^{\infty} \sum_{\substack{z_0, \dots, z_{r+1} \in \mathbb{Z}^2 \\ n_0 := 0 < n_1 < \dots < n_r < N}} \varphi_N(z_0) \left\{ \prod_{j=1}^r q_{n_j - n_{j-1}}(z_j - z_{j-1}) \xi_N(n_j, z_j) \right\} q_{N-n_r}(z_{r+1} - z_r) \psi_N(z_{r+1}), \end{aligned} \quad (7.2)$$

where $\varphi_N(x) := \varphi(\frac{x}{\sqrt{N}})$, $\psi_N(x) := \psi(\frac{x}{\sqrt{N}})$, and

$$q_{m,n}^N(\varphi, \psi) := \frac{1}{N} \sum_{x,y \in \mathbb{Z}^2} \varphi\left(\frac{x}{\sqrt{N}}\right) q_{n-m}(y-x) \psi\left(\frac{y}{\sqrt{N}}\right).$$

Note that each term in (7.2) is associated with a sequence of space-time points

$$\Gamma := ((n_1, z_1), \dots, (n_r, z_r)),$$

which contribute the product of disorder variables $\prod_{i=1}^r \xi_N(n_i, z_i)$. We will group the terms in (7.2) according to the set of mesoscopic blocks $\mathcal{B}_{\varepsilon N}(i, a)$ visited by points in Γ . More precisely, we proceed as follows (see Figure 3):

- First, we decompose the sum in (7.2) according to the set of mesoscopic time intervals $\mathcal{T}_{\varepsilon N}(i)$ visited by points in $\Gamma = ((n_1, z_1), \dots, (n_r, z_r))$.
- For each time interval $\mathcal{T}_{\varepsilon N}(i)$ visited by points in Γ , let $(d, x) \in \Gamma$ be the point of entry into $\mathcal{T}_{\varepsilon N}(i)$, with $(d, x) \in \mathcal{B}_{\varepsilon N}(i, a)$ for some $a \in \mathbb{Z}^2$. Similarly, let $(f, y) \in \Gamma$ be the point of exit from $\mathcal{T}_{\varepsilon N}(i)$, with $(f, y) \in \mathcal{B}_{\varepsilon N}(i, a')$ for some $a' \in \mathbb{Z}^2$, which could be different from a (see Figure 3).

Given the above coarse-grained decomposition, we can rewrite

$$\begin{aligned} \mathcal{Z}_N^{\beta_N}(\varphi, \psi) &= q_{0,N}^N(\varphi, \psi) \\ &+ \frac{1}{N} \sum_{v,w \in \mathbb{Z}^2} \varphi_N(v) \sum_{k=1}^{\infty} \sum_{0 < i_1 < \dots < i_k \leq \frac{1}{\varepsilon}} \sum_{\substack{d_1 \leq f_1 \in \mathcal{T}_{\varepsilon N}(i_1), \dots, d_k \leq f_k \in \mathcal{T}_{\varepsilon N}(i_k) \\ x_1, y_1, \dots, x_k, y_k \in \mathbb{Z}^2}} \\ &q_{d_1}(x_1 - v) X_{d_1, f_1}(x_1, y_1) \left\{ \prod_{j=2}^k q_{d_j - f_{j-1}}(x_j - y_{j-1}) X_{d_j, f_j}(x_j, y_j) \right\} q_{N-f_k}(w - y_k) \psi_N(w), \end{aligned} \quad (7.3)$$

where, recalling the point-to-point partition function from (1.11),

$$X_{d,f}(x, y) := \begin{cases} \xi_N(d, x) \mathbb{1}_{\{x=y\}} & \text{if } d = f, \\ \xi_N(d, x) Z_{d,f}^{\beta_N}(x, y) \xi_N(f, y) & \text{if } d < f. \end{cases} \quad (7.4)$$

Ideally, we would like to use the local limit theorem (2.6) to replace the random walk transition kernels $q_{d_j - f_{j-1}}(x_j - y_{j-1})$ by a heat kernel

$$q_{d_j - f_{j-1}}(x_j - y_{j-1}) \rightsquigarrow g_{(i_j - i_{j-1})\varepsilon N}((a_j - a'_{j-1})\sqrt{\varepsilon N}) = \frac{1}{\varepsilon N} g_{i_j - i_{j-1}}(a_j - a'_{j-1}), \quad (7.5)$$

which depends only on the index of the mesoscopic blocks $\mathcal{B}_{\varepsilon N}(i_{j-1}, a'_{j-1}) \ni (f_{j-1}, y_{j-1})$ and $\mathcal{B}_{\varepsilon N}(i_j, a_j) \ni (d_j, x_j)$. Namely, approximate (7.3) by

$$\mathcal{Z}_{N,\varepsilon}^{(\text{mock-cg})}(\varphi, \psi | \Theta_{N,\varepsilon}) = g_1(\varphi, \psi) \quad (7.6)$$

$$+ \varepsilon \sum_{k=1}^{\lfloor \frac{1}{\varepsilon} \rfloor} \sum_{\substack{0 < i_1 < \dots < i_k \leq \lfloor \frac{1}{\varepsilon} \rfloor \\ a'_0, a'_1, \dots, a'_k, a_{k+1} \in \mathbb{Z}^2}} \varphi_\varepsilon(a'_0) \left(\prod_{j=1}^k g_{i_j - i_{j-1}}(a_j - a'_{j-1}) \Theta_{N,\varepsilon}(i_j, \vec{a}_j) \right) \cdot g_{\frac{1}{\varepsilon} - i_r}(a_{r+1} - a'_r) \psi_\varepsilon(a_{r+1}),$$

where $g_1(\varphi, \psi) := \int \varphi(x) g_1(y-x) \psi(y) dx dy$, and

$$\begin{aligned} \varphi_\varepsilon(a) &:= \frac{1}{\varepsilon N} \sum_{v \in \mathcal{S}_{\varepsilon N}(a) \cap \mathbb{Z}^2} \varphi_N(v), \quad \psi_\varepsilon(a) := \frac{1}{\varepsilon N} \sum_{w \in \mathcal{S}_{\varepsilon N}(a) \cap \mathbb{Z}^2} \psi_N(w) \\ \text{and} \quad \Theta_{N,\varepsilon}(i, \vec{a}) &:= \frac{1}{\varepsilon N} \sum_{\substack{(d,x) \in \mathcal{B}_{\varepsilon N}(i,a) \\ (f,y) \in \mathcal{B}_{\varepsilon N}(i,a') \\ \text{with } d \leq f}} X_{d,f}(x, y), \quad \text{for } \vec{a} := (a, a') \in (\mathbb{Z}^2)^2. \end{aligned} \quad (7.7)$$

However, the above approximation is not good enough (this is why we gave the name “mock-cg”) because to have a negligible error from the local limit approximation, we require that:

- (R1) The mesoscopic time intervals $\mathcal{T}_{\varepsilon N}(i_j)$, $j = 1, \dots, k$, are well separated, that is, $i_j - i_{j-1} > K_\varepsilon$ for some $K_\varepsilon \in \mathbb{N}$ that diverges fast enough as $\varepsilon \downarrow 0$, so that replacing any point in $\mathcal{B}_{\varepsilon N}(i_{j-1}, a'_{j-1})$, resp. $\mathcal{B}_{\varepsilon N}(i_j, a_j)$, by a single point induces a negligible error in the local limit theorem approximation. Furthermore, we need the cumulative error in the local limit theorem approximations to be small.
- (R2) The space-time coordinates of the transition kernels $g_{i_j - i_{j-1}}(a_j - a'_{j-1})$ obey diffusive scaling, that is, $|a_j - a'_{j-1}| < M_\varepsilon \sqrt{i_j - i_{j-1}}$ for some M_ε that diverges slow enough as $\varepsilon \downarrow 0$ ($M_\varepsilon = \log \log \frac{1}{\varepsilon}$ suffices), so that the local limit theorem approximation induces a

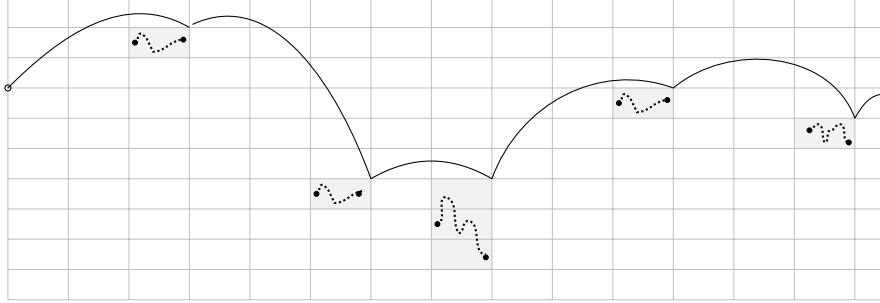


FIGURE 4. A simplified coarse-graining that defines $\mathcal{Z}_{N,\varepsilon}^{(\text{mock-cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$. A coarse-grained disorder variable $\Theta_{N,\varepsilon}(i, \vec{a})$ is defined for each visited mesoscopic time interval $\mathcal{T}_{\varepsilon N}(i)$. The random walk transition kernel connecting $\mathcal{T}_{\varepsilon N}(i)$ and $\mathcal{T}_{\varepsilon N}(j)$, with $i < j$, is replaced by a heat kernel connecting the corner of the mesoscopic spatial box of exit from $\mathcal{T}_{\varepsilon N}(i)$ to the corner of the box of entry in $\mathcal{T}_{\varepsilon N}(j)$, represented here by a solid line. Such replacements can lead to poor approximations if $j - i < K_\varepsilon$.

negligible error, while restricting the spatial coordinates to such an essentially diffusive window captures the dominant contribution to the partition function.

With regards to the first requirement (R1), a suitable choice is $K_\varepsilon := (\log \frac{1}{\varepsilon})^6$, which ensures that in the chaos expansion (7.2), configurations of $\Gamma = ((n_1, z_1), \dots, (n_r, z_r))$ that visit three mesoscopic time intervals $\mathcal{T}_{\varepsilon N}(i_j), \mathcal{T}_{\varepsilon N}(i_{j+1}), \mathcal{T}_{\varepsilon N}(i_{j+2})$ with $i_{j+1} - i_j < K_\varepsilon$ and $i_{j+2} - i_{j+1} < K_\varepsilon$ give a negligible contribution to the partition function. We can therefore restrict the chaos expansion for $\mathcal{Z}_N(\varphi, \psi)$ to such *no-triple configurations*. However, it can be shown by L^2 calculations that there will always be non-negligible contributions from configurations that contain $O(1)$ number of mesoscopic time indices i_j with $i_{j+1} - i_j < K_\varepsilon$. This is why $\mathcal{Z}_{N,\varepsilon}^{(\text{mock-cg})}(\varphi, \psi)$ is not a good approximation of $\mathcal{Z}_N(\varphi, \psi)$, although it captures some key features of the coarse-grained partition function that we need.

To obtain a good coarse-grained approximation of $\mathcal{Z}_N(\varphi, \psi)$, we proceed as follows. As discussed above, we may restrict the chaos expansion of $\mathcal{Z}_N(\varphi, \psi)$ to the *no-triple configurations*. Given a configuration $\Gamma = ((n_1, z_1), \dots, (n_r, z_r))$ in the chaos expansion (7.2), we call a visited mesoscopic time interval $\mathcal{T}_{\varepsilon N}(i_j)$ isolated if $i_j - i_{j-1}, i_{j+1} - i_j \geq K_\varepsilon$. If a_j (resp. a'_j) denotes the mesoscopic spatial box of entry (resp. exit) by Γ in the time interval $\mathcal{T}_{\varepsilon N}(i_j)$, then we define an associated coarse-grained disorder variable $\Theta_{N,\varepsilon}(i_j, \vec{a}_j)$ as in (7.7), with

$$\Theta_{N,\varepsilon}(i, \vec{a}) := \frac{1}{\varepsilon N} \sum_{\substack{(d,x) \in \mathcal{B}_{\varepsilon N}(i,a) \\ (f,y) \in \mathcal{B}_{\varepsilon N}(i,a') \\ \text{with } d \leq f}} X_{d,f}(x, y). \quad (7.8)$$

We also need to consider mesoscopic time intervals $\mathcal{T}_{\varepsilon N}(i_j)$ and $\mathcal{T}_{\varepsilon N}(i_{j+1})$ visited by Γ with $i_{j+1} - i_j < K_\varepsilon$, in which case, the no-triple configuration assumption ensures that mesoscopic time intervals $\mathcal{T}_{\varepsilon N}(i^*)$ with $i_j < i^* < i_{j+1}$, or $i^* \geq i_j - K_\varepsilon$, or $i^* \leq i_{j+1} + K_\varepsilon$ are not visited by the configuration Γ . If a_j (resp. a'_{j+1}) denotes the mesoscopic spatial box of entry (resp. exit) by Γ in $\mathcal{T}_{\varepsilon N}(i_j)$ (resp. from $\mathcal{T}_{\varepsilon N}(i_{j+1})$), then we define an associated

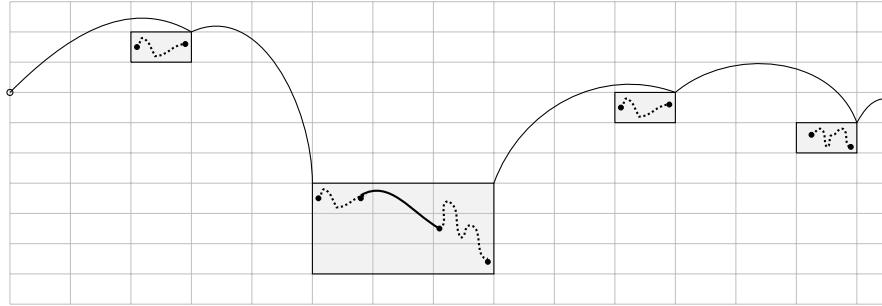


FIGURE 5. A more accurate coarse-graining approximation. Since the second and third visited boxes are separated by less than K_ε mesoscopic time intervals, they need to be grouped together as part of a single coarse-grained disorder variable $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$.

coarse-grained disorder variable $\Theta_{N,\varepsilon}(i_j, i_{j+1}, a_j, a'_{j+1})$ by

$$\begin{aligned} \Theta_{N,\varepsilon}(\vec{i}, \vec{a}) &:= \Theta_{N,\varepsilon}(i, i', a, a') \\ &:= \frac{1}{\varepsilon N} \sum_{\substack{(d,x) \in \mathcal{B}_{\varepsilon N}(i,a) \\ (f',y') \in \mathcal{B}_{\varepsilon N}(i',a')}} \sum_{\substack{b: |b-a| \leq M_\varepsilon \\ b': |b'-a'| \leq M_\varepsilon \\ \text{such that} \\ |b-b| \leq M_\varepsilon \sqrt{i'-i}}} \sum_{\substack{(f,y) \in \mathcal{B}_{\varepsilon N}(i,b) \\ (d',x') \in \mathcal{B}_{\varepsilon N}(i',b') \\ \text{such that} \\ d \leq f, d' \leq f'}} X_{d,f}(x,y) q_{d'-f}(x' - y) X_{d',f'}(x',y'), \end{aligned} \quad (7.9)$$

where $i < i'$, and we imposed a constraint on the mesoscopic spatial variables b, b' relative to a and a' . To unify notation, when $i = i'$, we will let $\Theta_{N,\varepsilon}(\vec{i}, \vec{a}) := \Theta_{N,\varepsilon}(i, \vec{a})$ as defined in (7.8). Schematically, the coarse-grained disorder variable can be represented as

$$\Theta_{N,\varepsilon}(\vec{i}, \vec{a}) := \begin{cases} \text{[Diagram of a single box with a wavy path]} & \text{if } i = i', \\ \text{[Diagram of three adjacent boxes with wavy paths]} & \text{if } i' - i \geq 1. \end{cases} \quad (7.10)$$

Comparing with (7.2), we can also interpret $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ as an averaged partition function, where the time window has length $(i' - i + 1)\varepsilon N$ and the spatial average is on the scale $\sqrt{\varepsilon N}$.

We are now ready to introduce a more accurate coarse-grained approximation of the partition function $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ than $\mathcal{Z}_{N,\varepsilon}^{(\text{mock-cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$ in (7.6).

Definition 7.1 (Coarse-grained partition function). Let $\varphi_\varepsilon, \psi_\varepsilon$ be defined in (7.7), and let $\Theta_{N,\varepsilon} := (\Theta_{N,\varepsilon}(\vec{i}, \vec{a}))_{\vec{i}=(i,i'), \vec{a}=(a,a')}$ be the family of coarse-grained disorder variables

defined in (7.8) and (7.9). Then we define the coarse-grained partition function

$$\begin{aligned} \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon}) &:= g_1(\varphi, \psi) + \varepsilon \sum_{k=1}^{(\log \frac{1}{\varepsilon})^2} \sum_{\mathbf{b}, \mathbf{c} \in \mathbb{Z}^2} \sum_{\substack{(\vec{i}_1, \dots, \vec{i}_k) \in \vec{\mathcal{A}}_\varepsilon^{(\text{no triple})} \\ (\vec{a}_1, \dots, \vec{a}_k) \in \vec{\mathcal{A}}_{\varepsilon; b, c}^{(\text{diff})}}} \varphi_\varepsilon(\mathbf{b}) g_{\vec{i}_1}(\mathbf{a}_1 - \mathbf{b}) \Theta_{N,\varepsilon}(\vec{i}_1, \vec{a}_1) \\ &\quad \times \left\{ \prod_{j=2}^k g_{\vec{i}_j - \vec{i}'_{j-1}}(\mathbf{a}_j - \mathbf{a}'_{j-1}) \Theta_{N,\varepsilon}(\vec{i}_j, \vec{a}_j) \right\} g_{\lfloor \frac{1}{\varepsilon} \rfloor - \vec{i}'_k}(\mathbf{c} - \mathbf{a}'_k) \psi_\varepsilon(\mathbf{c}). \end{aligned} \quad (7.11)$$

In (7.11), the indices of the mesoscopic time intervals $i_1 \leq i'_1 < i_2 \leq i'_2 < \dots < i_k \leq i'_k$ have been grouped into blocks $\vec{i}_j = (i_j, i'_j)$, each associated with a coarse-grained variable $\Theta_{N,\varepsilon}(\vec{i}_j, \vec{a}_j)$. The set $\vec{\mathcal{A}}_\varepsilon^{(\text{no triple})}$ ensures that for each j , $i_j - i'_{j-1} \geq K_\varepsilon$, $i'_j - i_j < K_\varepsilon$, and $i_{j+1} - i'_j \geq K_\varepsilon$, while the set $\vec{\mathcal{A}}_{\varepsilon; b, c}^{(\text{diff})}$ imposes diffusive constraints on the mesoscopic spatial indices $\vec{a}_j = (\mathbf{a}_j, \mathbf{a}'_j)$ such that $|\mathbf{a}'_j - \mathbf{a}_j| \leq M_\varepsilon \sqrt{i'_j - i_j + 1}$ and $|\mathbf{a}_j - \mathbf{a}'_{j-1}| \leq M_\varepsilon \sqrt{i_j - i'_{j-1}}$. For the precise definitions of $\vec{\mathcal{A}}_\varepsilon^{(\text{no triple})}$ and $\vec{\mathcal{A}}_{\varepsilon; b, c}^{(\text{diff})}$, see [CSZ23a, (4.4) and (4.6)].

Remark 7.2. Some remarks are in order:

1. In [CSZ23a], we imposed a further constraint in the definition of the coarse-grained variables $\Theta_{N,\varepsilon}$ so that each depends only on disorder in finitely many mesoscopic space-time blocks, which implies finite range dependence among the $\Theta_{N,\varepsilon}$'s. More precisely, in the definition of $X_{d,f}(x, y)$ in (7.4), we replace the point-to-point partition function $Z_{d,f}^{\beta_N}(x, y)$ by a variant $Z_{d,f}^{(\text{diff})}(x, y)$ that does not collect any disorder variables $\xi_N(n, z)$ outside a slightly super-diffusive window, see [CSZ23a, Section 4.3].
2. $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$ is a polynomial chaos expansion in the family of coarse-grained disorder variables $\Theta_{N,\varepsilon} = (\Theta_{N,\varepsilon}(\vec{i}, \vec{a}))$ and has the same structure as the chaos expansion for $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ in (7.2), which shows a degree of self-similarity. The time horizon N in $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ is replaced by $1/\varepsilon$, and the random walk transition kernels are replaced by heat kernels. Although the i.i.d. family of disorder variables ξ_N has been replaced by the dependent family $\Theta_{N,\varepsilon}$, the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$ is still critical since it approximates $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ with β_N in the critical window. See also (7.15).
3. If we change the family $\Theta_{N,\varepsilon}$ to $\Theta_{M,\varepsilon}$, then $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{M,\varepsilon})$ will approximate $\mathcal{Z}_M^{\beta_M}(\varphi, \psi)$ instead of $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$. It was shown in [CSZ23a, Theorem 4.7] that

$$\lim_{\varepsilon \downarrow 0} \limsup_{N \rightarrow \infty} \mathbb{E}[(\mathcal{Z}_N^{\beta_N}(\varphi, \psi) - \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon}))^2] = 0. \quad (7.12)$$

As we vary N , we only change the input $\Theta_{N,\varepsilon}$ for the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$. This paves the way to apply a Lindeberg principle to show that $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon}) \xrightarrow{\text{dist}} \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{M,\varepsilon})$.

4. In (7.11), the order of the chaos expansion in the variables $\Theta_{N,\varepsilon}$ is restricted to $r \leq (\log \frac{1}{\varepsilon})^2$. This induces a negligible error because it can be shown that the dominant contribution comes from terms of the order $\log \frac{1}{\varepsilon}$, similar to the fact that in the chaos

expansion for the averaged partition function $\mathcal{Z}_N^{\beta_N}(\varphi, \psi)$ in the critical window, the dominant contribution comes from terms of order $\log N$, see Section 8. On the other hand, restricting to $r \leq (\log \frac{1}{\varepsilon})^2$ allows us to control the cumulative error of replacing random walk transition kernels by heat kernels as in (7.5).

7.2. APPLYING LINDEBERG TO THE COARSE-GRAINED MODEL. We now sketch how to apply a Lindeberg principle to show that for each $\varepsilon > 0$,

$$\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon}) \xrightarrow{\text{dist}} \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{M,\varepsilon}) \quad \text{as } M, N \rightarrow \infty. \quad (7.13)$$

That is, the laws of $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$, $N \in \mathbb{N}$, form a Cauchy sequence.

The key difference from what was sketched in Section 6.2 is that,

$$\Psi_\varepsilon(\Theta) := \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta)$$

is a multilinear polynomial of the family of *dependant* random variables $\Theta = (\Theta(\vec{i}, \vec{a}))$, where we take $\Theta = \Theta_{N,\varepsilon}$ for $N \in \mathbb{N}$. Such a Lindeberg principle for multilinear polynomials of dependent random variables was formulated in [CSZ23a, Lemma A.4]. Instead of repeating the statement of this Lindeberg principle, we highlight here the key ingredients.

To establish (7.13), the first condition of the Lindeberg principle is that the covariances of $(\Theta_{N,\varepsilon}(\vec{i}, \vec{a}))$ converge to finite limits as $N \rightarrow \infty$ (see the bound in [CSZ23a, (A.11)]). It is clear from their definitions that distinct $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ are uncorrelated (but not independent). Therefore we only need to verify that for each $(\vec{i}, \vec{a}) = ((i, i'), (a, a'))$ with $1 \leq i \leq i' \leq \varepsilon^{-1}$ and $a, a' \in \mathbb{Z}^2$, the limit

$$\sigma_\varepsilon^2(\vec{i}, \vec{a}) := \lim_{N \rightarrow \infty} \mathbb{E}\left[\left(\Theta_{N,\varepsilon}(\vec{i}, \vec{a})\right)^2\right] \quad (7.14)$$

exists and is finite. We will sketch the proof of this in Section 8.3. Although not needed for the Lindeberg principle, we will also show in Section 8.3 that

$$\sigma_\varepsilon^2((i, i), (a, a)) \sim \frac{4\pi}{\log \frac{1}{\varepsilon}} \quad \text{as } \varepsilon \downarrow 0. \quad (7.15)$$

Remark 7.3. As will be explained later in (8.22), among all coarse-grained disorder variables, $\Theta_{N,\varepsilon}((i, i), (a, a))$ are the dominant ones as $\varepsilon \downarrow 0$. The precise asymptotics in (7.15) is a signature of the critical nature of the model, because it matches exactly the variance of the original disorder variable $\mathbb{E}[\xi_N(n, z)^2] = \sigma_N^2 \sim \frac{4\pi}{\log N}$ in the critical window (see (2.9)) if we identify the time horizon ε^{-1} with N . This shows a degree of self-similarity of the model under renormalisation.

The second condition of the Lindeberg principle in [CSZ23a, Lemma A.4] is that the input Θ of $\Psi_\varepsilon(\Theta)$ has a local dependency structure such that if $\Theta(\vec{i}, \vec{a}')$ belongs to the dependency neighbourhood of $\Theta(\vec{i}, \vec{a})$, and $\Theta(\vec{i}'', \vec{a}'')$ belongs to the dependency neighbourhood of $\{\Theta(\vec{i}, \vec{a}), \Theta(\vec{i}', \vec{a}')\}$, then no term in the chaos expansion of $\Psi_\varepsilon(\Theta)$ can contain at least 2 of the 3 factors among $\{\Theta(\vec{i}, \vec{a}), \Theta(\vec{i}', \vec{a}'), \Theta(\vec{i}'', \vec{a}'')\}$. In our setting, this is guaranteed by the facts that we only consider $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ with $i' - i < K_\varepsilon$, $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ and $\Theta_{N,\varepsilon}(\vec{j}, \vec{b})$ are independent when $[i, i'] \cap [j, j'] = \emptyset$, and each term in the chaos expansion of $\Psi_\varepsilon(\Theta)$ consists of well separated variables $\Theta(\vec{i}_i, \vec{a}_i)$ with $i_{i+1} - i'_i \geq K_\varepsilon$ thanks to the no-triple configuration condition in our coarse-graining. Details can be found in [CSZ23a], Lemma 9.3.

The last condition of the Lindeberg principle are fourth moment bounds on $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ and the coarse-grained point-to-plane partition function, analogous to (6.16) and (6.15). More

precisely, we need to show that

$$\limsup_{N \rightarrow \infty} \mathbb{E}[\Theta_{N,\varepsilon}(\vec{i}, \vec{a})^4] < \frac{C}{\log \frac{1}{\varepsilon}}, \quad (7.16)$$

where C is uniform in (\vec{i}, \vec{a}) and $\varepsilon > 0$ small. We will sketch the proof of this bound in Section 9.3. This bound is not sharp, but already more than enough. In fact, it suffices to have a bound of the order ε^{-c} for some $c > 0$ sufficiently small.

For $a_0 \in \mathbb{Z}^2$ and $m \in [1, 1/\varepsilon]$, the point-to-plane analogue of the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(cg)}(\varphi, \psi | \Theta)$ in (7.11) is defined by

$$\begin{aligned} \mathcal{Z}_{[0,m]}^{(cg)}(a_0, \psi_\varepsilon | \Theta_{N,\varepsilon}) := g_m(a_0, \psi_\varepsilon) + & \sum_{k=1}^{(\log \frac{1}{\varepsilon})^2} \sum_{c \in \mathbb{Z}^2} \sum_{\substack{(\vec{i}_1, \dots, \vec{i}_k) \in \vec{\mathcal{A}}_\varepsilon^{(\text{no triple})} \\ (\vec{a}_1, \dots, \vec{a}_k) \in \vec{\mathcal{A}}_{\varepsilon; a_0, c}^{(\text{diff})}}} g_{i_1}(a_1 - a_0) \Theta_{N,\varepsilon}(\vec{i}_1, \vec{a}_1) \\ & \times \left\{ \prod_{j=2}^k g_{i_j - i'_{j-1}}(a_j - a'_{j-1}) \Theta_{N,\varepsilon}(\vec{i}_j, \vec{a}_j) \right\} g_{m-i'_k}(c - a'_k) \psi_\varepsilon(c). \end{aligned}$$

Note that in contrast to (7.11), we do not multiply the sum by ε because there is no spatial averaging over the initial point a_0 . The plane-to-point partition function $\mathcal{Z}_{[0,m]}^{(cg)}(\varphi_\varepsilon, a_0 | \Theta_{N,\varepsilon})$ is defined similarly, and we can also shift the time window. We will then need the following analogue of (6.15): For any $\delta > 0$ and $\varphi \in C_c(\mathbb{R}^2)$, uniformly in $m \in [1, 1/\varepsilon]$, $a_0 \in \mathbb{Z}^2$, N large and ε small, we have

$$\mathbb{E}[\mathcal{Z}_{[0,m]}^{(cg)}(a_0, \mathbf{1} | \Theta_{N,\varepsilon})^4] \leq C\varepsilon^{-\delta} \quad \text{and} \quad \mathbb{E}[\mathcal{Z}_{[0,m]}^{(cg)}(a_0, \varphi_\varepsilon | \Theta_{N,\varepsilon})^4] \leq Ce^{-c\sqrt{\varepsilon}|a_0|}\varepsilon^{-\delta}. \quad (7.17)$$

The same reasoning as in (6.14) explains why the above bounds are sufficient for the application of the Lindeberg principle, and why the bound in (7.16) can be replaced by ε^{-c} for any $c \in (0, 1 - 2\delta)$. We will sketch the proof of (7.17) in Section 9.3.

8. SECOND MOMENT CALCULATIONS AND DICKMAN SUBORDINATOR

In this section, we explain how to compute the second moment of partition functions when β_N and σ_N^2 are chosen in the critical window defined in (1.16) for some $\vartheta \in \mathbb{R}$. This is based on a renewal representation that leads to the so-called Dickman subordinator in the scaling limit. We will then discuss how these calculations can be applied to the coarse-grained disorder variables $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ introduced in Section 7.

8.1. RENEWAL REPRESENTATION. We recall from (2.2) the polynomial chaos expansion of the point-to-plane partition function

$$Z_N^{\beta_N}(0) = 1 + \sum_{r=1}^{\infty} \sum_{\substack{n_0=0 < n_1 < \dots < n_r < N \\ z_0:=0, z_1, \dots, z_r \in \mathbb{Z}^2}} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1}) \xi_N(n_i, z_i).$$

Since the summands are L^2 -orthogonal and $\mathbb{E}[\xi_N(n_i, z_i)^2] = \sigma_N^2$ by (2.1), we have

$$\mathbb{E}[(Z_N^{\beta_N}(0))^2] = 1 + \sum_{r=1}^{\infty} \sum_{\substack{n_0=0 < n_1 < \dots < n_r < N \\ z_0:=0, z_1, \dots, z_r \in \mathbb{Z}^2}} \sigma_N^{2r} \prod_{i=1}^r q_{n_i - n_{i-1}}(z_i - z_{i-1})^2. \quad (8.1)$$

We now introduce a probabilistic representation of this sum by observing that

$$\frac{q_n(x)^2}{R_N} \mathbf{1}_{\{n \leq N\}}$$

defines a probability kernel on $\{1, \dots, N\} \times \mathbb{Z}^2$, which can be interpreted as the increment distribution of a time-space renewal process. More precisely, if $(T_i^{(N)}, X_i^{(N)})_{i \in \mathbb{N}}$ are i.i.d. random variables with distribution $\frac{q_n(x)^2}{R_N} \mathbf{1}_{\{n \leq N\}}$, and we denote their partial sums by

$$\tau_k^{(N)} := T_1^{(N)} + \dots + T_k^{(N)}, \quad S_k^{(N)} := X_1^{(N)} + \dots + X_k^{(N)}, \quad (8.2)$$

then $(\tau_k^{(N)})_{k \in \mathbb{N}_0}$ (with $\tau_0^{(N)} := 0$) is a renewal process with increment distribution $\frac{q_{2n}(0)}{R_N} \mathbf{1}_{\{n \leq N\}}$, and we have

$$\begin{aligned} \mathbb{E}[(Z_N^{\beta_N}(0))^2] &= 1 + \sum_{r=1}^{\infty} \sum_{\substack{n_0=0 < n_1 < \dots < n_r < N \\ z_0:=0, z_1, \dots, z_r \in \mathbb{Z}^2}} (\sigma_N^2 R_N)^r P((\tau_k^{(N)}, S_k^{(N)}) = (n_k, z_k) \text{ for } 1 \leq k \leq r) \\ &= \sum_{r=0}^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^r P(\tau_r^{(N)} < N). \end{aligned} \quad (8.3)$$

where we used (1.16) for σ_N^2 in the critical window and summed over the intermediate values.

We have a similar renewal representation for the second moment of the point-to-point partition function $Z_{n,m}^{\beta_N}(x,y)$, whose chaos expansion is given in (2.3). Actually, a more fundamental object in our analysis is

$$U_N(n, x) := \sigma_N^2 \operatorname{Var}(Z_{0,n}^{\beta_N}(0, x)) = \sigma_N^2 \operatorname{Var}(Z_{\ell,\ell+n}^{\beta_N}(z, x+z)),$$

for which we have

$$U_N(n, x) := \begin{cases} \mathbf{1}_{\{x=0\}} & \text{if } n=0, \\ \sigma_N^2 q_n(x)^2 + \sum_{k=1}^{\infty} \sum_{\substack{z_1, \dots, z_k \in \mathbb{Z}^2 \\ 0 < n_1 < \dots < n_k < n}} \prod_{j=1}^{k+1} \sigma_N^2 q_{n_j-n_{j-1}}(z_j - z_{j-1})^2 & \text{if } n \geq 1, \end{cases} \quad (8.4)$$

where $n_0 := 0$, $z_0 = 0$, $n_{k+1} := n$ and $z_{k+1} := x$. Similar to the graphical representation (2.3) of the chaos expansion of $Z_{n,m}^{\beta}(x,y)$, we can graphically represent the expansion for $U_N(m-n, y-x) = \sigma_N^2 \operatorname{Var}(Z_{n,m}^{\beta_N}(x,y))$ by

$$U_N(m-n, y-x) = \sum_{r \geq 1} \sum_{\substack{n < n_1 < \dots < n_r < m \\ z_1, \dots, z_r \in \mathbb{Z}^2}} \dots \quad \text{Diagram showing a sequence of points } (n_1, z_1), (n, x), (n_r, z_r), (m, y) \text{ connected by arcs. Each arc } (n_j, z_j) \text{ corresponds to a factor } q_{n_j-n_{j-1}}(z_j - z_{j-1}).$$

where each lace corresponds to one factor of $q_{n_j-n_{j-1}}(z_j - z_{j-1})$, each solid circle corresponds to one factor of σ_N^2 , while the empty circle at (n, x) means there is no factor σ_N^2 assigned to

(n, x) . A more compact graphical representation is

$$\begin{aligned} \sigma_N^2 U_N(m-n, y-x) &\equiv \text{Diagram showing two points } (n, x) \text{ and } (m, y) \text{ connected by a wavy line.} \quad \text{or} \\ &= \sum_{r \geq 1} \sum_{\substack{n_1, \dots, n_r \\ z_1, \dots, z_r}} \text{Diagram showing } (n, x) \text{ connected to } (n_1, z_1) \text{ which is connected to } (n_2, z_2) \dots \text{ which is connected to } (n_k, z_k) \text{ which is connected to } (m, y). \end{aligned} \quad (8.5)$$

where we included an extra factor of σ_N^2 to turn the empty circle at (n, x) into a solid circle.

Using the time-space renewal process introduced in (8.2), we can now write

$$U_N(n, x) = \sum_{r=1}^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^r P(\tau_r^{(N)} = n, S_r^{(N)} = x). \quad (8.6)$$

Summing over $x \in \mathbb{Z}^2$ then gives

$$U_N(n) := \sum_{x \in \mathbb{Z}^2} U_N(n, x) = \sum_{r=1}^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^r P(\tau_r^{(N)} = n). \quad (8.7)$$

To identify sharp asymptotics for $\mathbb{E}[(Z_N^{\beta_N}(0))^2]$ in (8.3) and for $U_N(n, x)$ and $U_N(n)$ as $N \rightarrow \infty$, the key observation is that $\frac{\tau_{[s \log N]}^{(N)}}{N}$ converges in distribution to a Lévy process called the *Dickman subordinator*.

8.2. DICKMAN SUBORDINATOR. We now introduce the *Dickman subordinator*, which is a truncated, 0-stable Lévy process $Y = (Y_s)_{s \geq 0}$ with Lévy measure

$$\nu(dt) := \frac{1}{t} \mathbf{1}_{(0,1)}(t) dt. \quad (8.8)$$

For $s \geq 0$, it has Laplace transform

$$\mathbb{E}[e^{\lambda Y_s}] = \exp \left\{ s \int_0^1 (e^{\lambda t} - 1) \frac{dt}{t} \right\}. \quad (8.9)$$

The density of the Y_s can be computed explicitly as (see [CSZ19a, Theorem 1.1])

$$f_s(t) := \frac{P(Y_s \in dt)}{dt} = \begin{cases} \frac{s t^{s-1} e^{-\gamma s}}{\Gamma(s+1)} & \text{for } t \in (0, 1], \\ \frac{s t^{s-1} e^{-\gamma s}}{\Gamma(s+1)} - s t^{s-1} \int_0^{t-1} \frac{f_s(a)}{(1+a)^s} da & \text{for } t \in (1, \infty). \end{cases} \quad (8.10)$$

The calculation in [CSZ19a] was partly based on the observation that conditional on all jumps of Y up to time s being smaller than a value $t \in (0, 1)$, Y_s/t has the same law as the original Y_s (see [CSZ19a, Prop. B.1]). More properties of the Dickman subordinator can be found in [GKLV24].

The name *Dickman subordinator* is motivated by the fact that

$$e^\gamma f_1(t) = \varrho(t),$$

where $\gamma = -\int_0^\infty (\log u) e^{-u} du$ is the Euler–Mascheroni constant and $\varrho(\cdot)$ is the so-called *Dickman function*. It was originally derived by Dickman [Dic30] in studying the distribution of the largest prime factor of a uniformly chosen number in $\{1, \dots, n\}$. It also arises from the distribution of the size of the longest cycle in a uniformly chosen random permutation of n elements. More instances can be found in [PW04] and the references therein.

The key observation in the asymptotic analysis of $\mathbb{E}[(Z_N^{\beta_N}(0))^2]$, $U_N(n, x)$ and $U_N(n)$ is the following convergence result for the time-space renewal process $(\tau^{(N)}, S^{(N)})$ defined in (8.2) (see [CSZ19a, Prop. 2.2] and [CSZ23a, Lemma 3.3]).

Lemma 8.1. *We have the following weak convergence in càdlàg space:*

$$\left(\frac{\tau_{[s \log N]}^{(N)}}{N}, \frac{S_{[s \log N]}^{(N)}}{\sqrt{N}} \right)_{s \geq 0} \xrightarrow[N \rightarrow \infty]{d} (Y_s, \frac{1}{\sqrt{2}} W_{Y_s})_{s \geq 0}, \quad (8.11)$$

where $(Y_s)_{s \geq 0}$ is the Dickman subordinator and $(W_s)_{s \geq 0}$ an independent standard Brownian motion on \mathbb{R}^2 .

Proof sketch. The finite-dimensional distribution convergence in (8.11) can be verified by showing the convergence of their Laplace transforms. We will illustrate this for $\tau^{(N)}$:

$$\begin{aligned} \mathbb{E}\left[e^{\frac{\lambda}{N} \tau_{s \log N}^{(N)}}\right] &= \mathbb{E}\left[e^{\frac{\lambda}{N} T_1^{(N)}}\right]^{s \log N} = \left(1 + \frac{1}{\log N} \sum_{n=1}^N (e^{\frac{\lambda}{N} n} - 1) \frac{1+o(1)}{n}\right)^{s \log N} \\ &\xrightarrow[N \rightarrow \infty]{} \exp\left(s \int_0^1 (e^{\lambda x} - 1) \frac{dx}{x}\right), \end{aligned}$$

where in the second equality we used the local central limit theorem to approximate $\frac{1}{R_N} q_{2n}(0) = \frac{1}{\log N} \frac{1}{n} (1+o(1))$ for n large. The full f.d.d. convergence can be found in the proof of [CSZ19a, Prop. 2.2], while process level tightness was proved in [CSZ23a, Lemma 3.3]. \square

In our analysis of the polymer partition functions in the critical window with parameter ϑ , defined in (1.16), we will need the following weighted Green's function of the Dickman subordinator:

$$G_\vartheta(t) := \int_0^\infty e^{\vartheta s} f_s(t) ds \quad t \in (0, \infty), \vartheta \in \mathbb{R}. \quad (8.12)$$

When $t \in (0, 1]$, by (8.10), we have the more explicit form

$$G_\vartheta(t) = \int_0^\infty \frac{e^{(\vartheta-\gamma)s} s t^{s-1}}{\Gamma(s+1)} ds, \quad t \in (0, 1], \vartheta \in \mathbb{R}. \quad (8.13)$$

We note that G_ϑ is related to the Volterra function

$$\nu(t) := \int_0^\infty \frac{t^s}{\Gamma(s+1)} ds.$$

See [Ape10] for more information on functions of this type.

We also record here the small t asymptotics of $G_\vartheta(t)$, which plays a crucial role in our analysis of the critical 2d SHF (see [CSZ19a, Prop. 1.6]):

Proposition 8.2. *For $t \in (0, 1]$, $G_\vartheta(t)$ is C^∞ and strictly positive. As $t \downarrow 0$, we have*

$$G_\vartheta(t) = \frac{1}{t(\log \frac{1}{t})^2} \left\{ 1 + \frac{2\vartheta}{\log \frac{1}{t}} + O\left(\frac{1}{(\log \frac{1}{t})^2}\right) \right\} \quad (8.14)$$

and

$$\int_0^t G_\vartheta(s) ds = \frac{1}{\log \frac{1}{t}} \left\{ 1 + \frac{\vartheta}{\log \frac{1}{t}} + O\left(\frac{1}{(\log \frac{1}{t})^2}\right) \right\}. \quad (8.15)$$

These asymptotics were derived in [CSZ19a] via a renewal framework. But they can also be derived from the asymptotic theory of Volterra functions [Ape10]. Note that G_ϑ is barely integrable near 0, which is the hallmark of a model at the critical dimension.

8.3. ASYMPTOTICS. We are now ready to state the asymptotics for $\mathbb{E}[(Z_N^{\beta_N}(0))^2]$, $U_N(n, x)$ and $U_N(n)$ in terms of the weighted Green's function $G_\vartheta(t)$ for the Dickman subordinator. We will also illustrate how to identify the asymptotics of the coarse-grained disorder variables $\Theta_{N,\varepsilon}(\vec{i}, \vec{s})$ introduced in Section 7.

Lemma 8.3. *Let β_N and σ_N^2 be chosen in the critical window defined in (1.16) for some $\vartheta \in \mathbb{R}$. Then we have*

$$\mathbb{E}[(Z_N^{\beta_N}(0))^2] = (\bar{G}_\vartheta + o(1)) \log N \quad \text{with} \quad \bar{G}_\vartheta := \int_0^1 G_\vartheta(t) dt. \quad (8.16)$$

Proof. By (8.3), we have

$$\begin{aligned} \mathbb{E}[(Z_N^{\beta_N}(0))^2] &= \sum_{r=0}^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^r P(\tau_r^{(N)} < N) \\ &= \log N \int_0^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^{\lfloor s \log N \rfloor} P\left(\frac{1}{N} \tau_{\lfloor s \log N \rfloor}^{(N)} < 1\right) ds \\ &= (1 + o(1)) \log N \int_0^{\infty} e^{\vartheta s} P(Y_s < 1) ds = \left(\int_0^1 G_\vartheta(t) dt + o(1)\right) \log N, \end{aligned}$$

where the last line follows from the convergence in Lemma 8.1 and the lower large deviation bounds for the renewal process $\tau^{(N)}$ stated in [CSZ19a, Lemma 6.1]. \square

We also record here the asymptotics for $U_N(n)$ and $U_N(n)$, which are of the local limit theorem type and play a crucial in the analysis in [CSZ23a]. These results were proved in [CSZ19a, Theorems 1.4, 2.3, 3.7], although the constants therein differ because the underlying random walk was the simple symmetric random walk on \mathbb{Z}^2 .

Proposition 8.4. *Let $\vartheta \in \mathbb{R}$ be as in Lemma 8.3. For any fixed $\delta > 0$ and uniformly in $\delta N \leq n \leq N$, as $N \rightarrow \infty$, we have*

$$U_N(n) = \frac{\log N}{N} G_\vartheta\left(\frac{n}{N}\right) (1 + o(1)), \quad (8.17)$$

$$U_N(n, x) = \frac{\log N}{N^2} G_\vartheta\left(\frac{n}{N}\right) g_{\frac{n}{2N}}\left(\frac{x}{\sqrt{N}}\right) (1 + o(1)). \quad (8.18)$$

Proof Sketch. Because the asymptotics are of the local limit theorem type, they do not follow directly from Lemma 8.1. Instead, using a renewal decomposition, (8.17) and (8.18) were deduced from their integrated versions in [CSZ19a, Sections 6.2 & 8.3].

We sketch here an alternative approach that is more transparent. First of all, with a bit more effort, the weak convergence in Lemma 8.3 can be upgraded to a local limit theorem:

$$P\left(\frac{1}{N} \tau_{\lfloor s \log N \rfloor}^{(N)} = \frac{n}{N}\right) = \frac{1}{N} f_s\left(\frac{n}{N}\right) (1 + o(1)), \quad (8.19)$$

$$P\left(\frac{1}{N} \tau_{\lfloor s \log N \rfloor}^{(N)} = \frac{n}{N}, \frac{1}{\sqrt{N}} S_{\lfloor s \log N \rfloor}^{(N)} = \frac{x}{\sqrt{N}}\right) = \frac{1}{N^2} f_s\left(\frac{n}{N}\right) g_{\frac{n}{2N}}\left(\frac{x}{\sqrt{N}}\right) (1 + o(1)), \quad (8.20)$$

where $f_s(\cdot)$ is the density of the Dickman subordinator in (8.10), $g_t(\cdot)$ is the heat kernel, and the $o(1)$ error term is uniform in $\delta N \leq n \leq N$ and $|x| \leq \frac{1}{\delta} \sqrt{N}$.

To see (8.17), note that by (8.7), we have

$$\begin{aligned} U_N(n) &= \sum_{r=1}^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^r P\left(\frac{1}{N} \tau_r^{(N)} = \frac{n}{N}\right) \\ &= \log N \int_0^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^{\lfloor s \log N \rfloor} P\left(\frac{1}{N} \tau_{\lfloor s \log N \rfloor}^{(N)} = \frac{n}{N}\right) ds \\ &= \frac{\log N}{N} \int_0^{\infty} (1 + o(1)) e^{\vartheta s} f_s\left(\frac{n}{N}\right) ds \\ &= (1 + o(1)) \frac{\log N}{N} G_{\vartheta}\left(\frac{n}{N}\right). \end{aligned}$$

Similarly, to see (8.18), we use (8.6) to write

$$\begin{aligned} U_N(n, x) &= \sum_{r=1}^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^r P\left(\frac{1}{N} \tau_r^{(N)} = \frac{n}{N}, \frac{1}{\sqrt{N}} S_r^{(N)} = \frac{x}{\sqrt{N}}\right) \\ &= \log N \int_0^{\infty} \left(1 + \frac{\vartheta + o(1)}{\log N}\right)^{\lfloor s \log N \rfloor} P\left(\frac{1}{N} \tau_{\lfloor s \log N \rfloor}^{(N)} = \frac{n}{N}, \frac{1}{\sqrt{N}} S_{\lfloor s \log N \rfloor}^{(N)} = \frac{x}{\sqrt{N}}\right) ds \\ &= \frac{\log N}{N^2} \int_0^{\infty} (1 + o(1)) e^{\vartheta s} f_s\left(\frac{n}{N}\right) g_{\frac{n}{2N}}\left(\frac{x}{\sqrt{N}}\right) ds \\ &= (1 + o(1)) \frac{\log N}{N^2} G_{\vartheta}\left(\frac{n}{N}\right) g_{\frac{n}{2N}}\left(\frac{x}{\sqrt{N}}\right). \end{aligned}$$

□

We now sketch the proof of (7.14), that is, the second moment of the coarse-grained disorder variables $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ converge as $N \rightarrow \infty$.

Proof Sketch for (7.14). For simplicity, we focus on $\Theta_{N,\varepsilon}(i, a)$ defined in (7.8) and (7.10), with $i = 1$ and $\vec{a} = (0, a)$ for some $a \in \mathbb{Z}^2$. By (7.8) and (7.4), we have

$$\Theta_{N,\varepsilon}(1, (0, a)) := \frac{1}{\varepsilon N} \sum_{(d,x) \in \mathcal{B}_{\varepsilon N}(1,0) \cap \mathcal{B}_{\varepsilon N}(1,a)} \xi_N(d, x) + \frac{1}{\varepsilon N} \sum_{\substack{(d,x) \in \mathcal{B}_{\varepsilon N}(1,0) \\ (f,y) \in \mathcal{B}_{\varepsilon N}(1,a) \\ \text{with } d < f}} \xi_N(d, x) Z_{d,f}^{\beta_N}(x, y) \xi_N(f, y).$$

Note that the first sum is 0 if $\mathcal{B}_{\varepsilon N}(1,0) \cap \mathcal{B}_{\varepsilon N}(1,a) = \emptyset$. The summands are L^2 -orthogonal, and

$$\mathbb{E}[\xi_N(d, x)^2 (Z_{d,f}^{\beta_N}(x, y))^2 \xi_N(f, y)^2] = \sigma_N^2 U_N(f - d, y - x).$$

Therefore,

$$\mathbb{E}[\Theta_{N,\varepsilon}(1, (0, a))^2] = \frac{1}{(\varepsilon N)^2} \sum_{(d,x) \in \mathcal{B}_{\varepsilon N}(1,0) \cap \mathcal{B}_{\varepsilon N}(1,a)} \sigma_N^2 + \frac{1}{(\varepsilon N)^2} \sum_{\substack{(d,x) \in \mathcal{B}_{\varepsilon N}(1,0) \\ (f,y) \in \mathcal{B}_{\varepsilon N}(1,a) \\ \text{with } d < f}} \sigma_N^2 U_N(f - d, y - x),$$

where the first term is bounded by $\sigma_N^2 \sim C/\log N$ and hence vanishes in the limit $N \rightarrow \infty$. For the second term, we interpret $(\varepsilon N)^{-2} \sum_{(d,x) \in \mathcal{B}_{\varepsilon N}(1,0)}$ as an expectation over a point (d, x)

chosen uniformly from $\mathcal{B}_{\varepsilon N}(1, 0)$. Using the same renewal representation of $U_N(f - d, y - x)$ as before, we find that

$$\lim_{N \rightarrow \infty} \frac{1}{(\varepsilon N)^2} \sum_{\substack{(d, x) \in \mathcal{B}_{\varepsilon N}(1, 0) \\ (f, y) \in \mathcal{B}_{\varepsilon N}(1, a) \\ \text{with } d < f}} \sigma_N^2 U_N(f - d, y - x) = 4\pi \int_0^\infty e^{\vartheta s} P_{\mathcal{B}_\varepsilon(1, 0)}((Y_s, V_s) \in \mathcal{B}_\varepsilon(1, a)) ds,$$

where 4π comes from $\sigma_N^2 \sim 4\pi/\log N$, and $P_{\mathcal{B}_\varepsilon(1, 0)}$ denotes the law of the continuum time-space renewal process $(Y_s, V_s) := (Y_s, \frac{1}{\sqrt{2}}W_{Y_s})$ in Lemma 8.1 with (Y_0, V_0) chosen uniformly from $\mathcal{B}_\varepsilon(1, 0)$, with $\mathcal{B}_\varepsilon(i, a) := ((i-1)\varepsilon, i\varepsilon) \times ((a-(1, 1))\sqrt{\varepsilon}, a\sqrt{\varepsilon})$, cf. (7.1). Therefore,

$$\sigma_\varepsilon^2(1, (0, a)) := \lim_{N \rightarrow \infty} \mathbb{E}[\Theta_{N, \varepsilon}(1, (0, a))^2] = 4\pi \int_0^\infty e^{\vartheta s} P_{\mathcal{B}_\varepsilon(1, 0)}((Y_s, V_s) \in \mathcal{B}_\varepsilon(1, a)) ds. \quad (8.21)$$

For more general coarse-grained disorder variables $\Theta_{N, \varepsilon}(\vec{i}, \vec{a})$, the calculations are similar. See [CSZ23a, Section 7.1] for more details. \square

Lastly, we sketch the proof of (7.15), which shows that the coarse-grained variable $\Theta_{N, \varepsilon}(i, (a, a))$ has an asymptotic variance of $4\pi/\log \varepsilon^{-1}$ as $N \rightarrow \infty$ and $\varepsilon \downarrow 0$, which is comparable to $\text{Var}(\xi_N(n, x)) = \sigma_N^2 \sim 4\pi/\log N$ for the original model if we identify ε^{-1} as the time horizon of the coarse-grained model.

Proof Sketch for (7.15). It was shown in [CSZ23a, Lemma 7.1] that whenever $i > 1$ or $a \neq 0$,

$$\sigma_\varepsilon^2((1, i), (0, a)) := \lim_{N \rightarrow \infty} \mathbb{V}\text{ar}(\Theta_{N, \varepsilon}((1, i), (0, a))) \leq \frac{Ce^{-c|a|^2/i}}{(\log \frac{1}{\varepsilon})^2}, \quad (8.22)$$

while $\sigma_\varepsilon^2(i, (a, a)) := \sigma_\varepsilon^2((i, i), (a, a))$ is of the order $1/\log \frac{1}{\varepsilon}$. Therefore among all the coarse-grained disorder variables, $\Theta_{N, \varepsilon}(i, (a, a))$ are dominant as $N \rightarrow \infty$ and $\varepsilon \downarrow 0$. Furthermore, as $\varepsilon \downarrow 0$,

$$\sigma_\varepsilon^2(1, (0, 0)) = (1 + o(1)) \sum_{a \in \mathbb{Z}^2} \sigma_\varepsilon^2(1, (0, a)) = (1 + o(1)) \lim_{N \rightarrow \infty} \mathbb{V}\text{ar} \left(\sum_{a \in \mathbb{Z}^2} \Theta_{N, \varepsilon}(1, (0, a)) \right),$$

where the $o(1)$ correction is due to (8.22). By (8.21), we have

$$\begin{aligned} \sigma_\varepsilon^2(1, (0, 0)) &= (1 + o(1)) 4\pi \int_0^\infty e^{\vartheta s} P_{\mathcal{B}_\varepsilon(1, 0)}(Y_s < \varepsilon) ds \\ &= \frac{4\pi + o(1)}{\varepsilon} \int_0^\varepsilon \left(\int_0^\infty e^{\vartheta s} P(Y_s < \varepsilon - u) ds \right) du \\ &= \frac{4\pi + o(1)}{\varepsilon} \int_0^\varepsilon \left(\int_0^{\varepsilon-u} G_\vartheta(t) dt \right) du \\ &= \frac{4\pi + o(1)}{\varepsilon} \int_0^\varepsilon \frac{1}{\log \frac{1}{\varepsilon-u}} du = \frac{4\pi + o(1)}{\log \frac{1}{\varepsilon}}, \end{aligned}$$

where we used that Y_0 is uniformly distributed in $[0, \varepsilon]$ under $P_{\mathcal{B}_\varepsilon(1, 0)}$ and $Y_0 = 0$ under P . We also used (8.15) in the last line. \square

8.4. COVARIANCE KERNEL OF THE CRITICAL $2d$ SHF. We sketch here the covariance kernel for the critical $2d$ stochastic heat flow $\text{SHF}(\vartheta) = (\text{SHF}_{s,t}^\vartheta(dx, dy))_{0 \leq s \leq t < \infty}$ given in Theorem 6.1.

Proof Sketch. The mean of $\text{SHF}_{s,t}^\vartheta(dx, dy)$ follows directly from the convergence in Theorem 6.1, the definition of $\mathcal{Z}_{N;s,t}$ in (6.1), and the local limit theorem in (2.6).

For illustration, we will identify the kernel

$$\tilde{K}_t^\vartheta(x, x') dx dx' := \mathbb{C}\text{ov}[\text{SHF}_{0,t}^\vartheta(dx, 1), \text{SHF}_{0,t}^\vartheta(dx', 1)].$$

By the convergence in Theorem 6.1, we see that

$$\tilde{K}_t^\vartheta(x, x') = \lim_{N \rightarrow \infty} \mathbb{C}\text{ov}[Z_{tN}^{\beta_N}(\sqrt{N}x), Z_{tN}^{\beta_N}(\sqrt{N}x')],$$

where $Z_{tN}^{\beta_N}(z)$ is the point-to-plane partition function with chaos expansion given in (2.2). By this chaos expansion, we note that

$$\begin{aligned} \mathbb{C}\text{ov}[Z_{tN}^{\beta_N}(\sqrt{N}x), Z_{tN}^{\beta_N}(\sqrt{N}x')] &= \sum_{z \in \mathbb{Z}^2, 1 \leq n \leq tN} q_n(z - \sqrt{N}x) q_n(z - \sqrt{N}x') \sigma_N^2 \mathbb{E}[Z_{tN-n}(0)^2] \\ &= \sum_{n=1}^{tN} q_{2n}(\sqrt{N}(x' - x)) \sigma_N^2 \mathbb{E}[Z_{tN-n}(0)^2], \end{aligned} \quad (8.23)$$

where σ_N^2 comes from $\mathbb{E}[\xi_N^2(n, z)]$, and we used the fact that the second moment of the point-to-plane partition function starting from (n, z) and terminal time tN is exactly $\mathbb{E}[Z_{tN-n}(0)^2]$.

Denote $u := n/N$. By the same calculations as in the proof of Lemma 8.3, we have

$$\mathbb{E}[Z_{tN-n}(0)^2] = \left(\int_0^{t-u} G_\vartheta(a) da + o(1) \right) \log N.$$

Substituting this into (8.23), using $\sigma_N^2 \sim 4\pi/\log N$, the local limit theorem (2.6), and a Riemann sum approximation, we obtain

$$\begin{aligned} \tilde{K}_t^\vartheta(x, x') &= (1 + o(1)) \sum_{n=1}^{tN} \frac{1}{4\pi n} e^{-\frac{N|x-x'|^2}{4n}} \frac{4\pi}{\log N} \cdot \log N \int_0^{t-u} G_\vartheta(a) da \\ &= 4\pi \iint_{0 < u < v < t} g_{2u}(x' - x) G_\vartheta(v - u) du dv. \end{aligned} \quad (8.24)$$

We refer to [CSZ23a, Prop. 3.6] for the complete derivation of K_{t-s}^ϑ , where the constants differ due to the periodicity of the simple symmetric random walk. \square

9. HIGHER MOMENT BOUNDS

In this section, we explain how to bound higher moments of the averaged partition function (cf. (6.1))

$$\mathcal{Z}_{N;0,t}^{\beta_N}(\varphi, \psi) = \frac{1}{N} \sum_{x,y \in \mathbb{Z}^2} \varphi\left(\frac{x}{\sqrt{N}}\right) \psi\left(\frac{y}{\sqrt{N}}\right) Z_{0,\lfloor Nt \rfloor}(x, y), \quad (9.1)$$

where β_N is chosen in the critical window defined in (1.16) for some $\vartheta \in \mathbb{R}$. As explained in Section 6.2 and in particular in (6.15), such moment bounds (4th moment will suffice) are needed to apply the Lindeberg principle. As explained in Section 6.3 and 7.2, to prove

Theorem 6.1 on the critical 2d stochastic heat flow (SHF), we will apply a Lindeberg principle for dependent random variables, which requires the moment bound (7.16) for the coarse-grained disorder variables $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ and the moment bound (7.17) for the coarse-grained partition function $\mathcal{Z}_\varepsilon^{(cg)}(\varphi, \psi | \Theta_{N,\varepsilon})$. We will focus mainly on the derivation of moment bounds for $\mathcal{Z}_{tN}^{\beta_N}(\varphi, \psi)$. We will then sketch how this implies the desired moment bound for $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ and how the same strategy can be used to derive moment bounds for $\mathcal{Z}_\varepsilon^{(cg)}(\varphi, \psi | \Theta_{N,\varepsilon})$. At the end of this section, we will also formulate moment bounds for the critical 2d SHF and discuss related results on the 2d Delta-Bose gas.

9.1. MOMENT BOUNDS AND PROOF STRATEGY. We now formulate the higher moment bounds for the averaged partition function $\mathcal{Z}_{N;0,t}^{\beta_N}(\varphi, \psi)$ defined in (9.1) for some parameter $t \in \mathbb{R}$ in the critical window (1.16). For consistency with [CSZ23a, Theorem 6.1] where such a bound was established, we will let $N' = tN$ and consider the partition averaged on the spatial scale $\sqrt{N'}$ instead of \sqrt{N} . Namely we consider

$$\mathcal{Z}_{N'}^{\beta_N}(\varphi, \psi) = \frac{1}{N'} \sum_{x,y \in \mathbb{Z}^2} \varphi_{N'}(x) Z_{0,N'}(x, y) \psi_{N'}(y), \quad (9.2)$$

where for $\phi : \mathbb{R}^2 \rightarrow \mathbb{R}$, $\phi_{N'} : \mathbb{Z}^2 \rightarrow \mathbb{R}$ is defined by setting $\phi_{N'}(z)$ to be the average of ϕ on the square of volume $1/N'$ centered at $z/\sqrt{N'}$.

Theorem 9.1 (Higher moments). Fix $p, q \in (1, \infty)$ with $\frac{1}{p} + \frac{1}{q} = 1$, any integer $h \geq 3$, and any weight function $w : \mathbb{R}^2 \rightarrow (0, \infty)$ such that $\log w$ is Lipschitz continuous. Then there exist $C = C(h), C' = C'(h) < \infty$ such that, uniformly in large $N' \leq N \in \mathbb{N}$ and locally integrable $\varphi, \psi : \mathbb{R}^2 \rightarrow \mathbb{R}$, we have

$$\left| \mathbb{E} \left[\left(\mathcal{Z}_{N'}^{\beta_N}(\varphi, \psi) - \mathbb{E}[\mathcal{Z}_{N'}^{\beta_N}(\varphi, \psi)] \right)^h \right] \right| \leq \frac{C}{\log(1 + \frac{N}{N'})} \frac{1}{(N')^h} \left\| \frac{\varphi_{N'}}{w_{N'}} \right\|_{\ell^p}^h \|\psi_{N'}\|_{\ell^\infty}^h \|w_{N'} \mathbf{1}_{B_{N'}}\|_{\ell^q}^h \quad (9.3)$$

$$\leq \frac{C'}{\log(1 + \frac{N}{N'})} \left\| \frac{\varphi}{w} \right\|_p^h \|\psi\|_\infty^h \|w \mathbf{1}_B\|_q^h, \quad (9.4)$$

where $\|\phi\|_{\ell^p} := (\sum_z |\phi(z)|^p)^{1/p}$, $B \subset \mathbb{R}^2$ is any ball (could be \mathbb{R}^2) that contains the support of ψ , and $B_{N'} := B\sqrt{N'}$.

Remark 9.2. Some remarks are in order:

- The bound in (9.4) follows from (9.3) via a simple Riemann sum approximation. For $\mathcal{Z}_{N;0,t}^{\beta_N}(\varphi, \psi)$ defined in (9.1) with averaging on the spatial scale \sqrt{N} , the same bound (9.4) holds with $t = N'/N$. This can be seen by applying (9.3) with φ and ψ therein replaced with $\varphi_t(x) := \varphi(\sqrt{tx})$ and $\psi_t(x) := \psi(\sqrt{tx})$, and using the weight function w_N instead of $w_{N'}$. But there is probably room for improvement since we expect that averaging on larger spatial scales will decrease the variance and the higher moments of the averaged partition function.
- Theorem 9.1 with $N' = \varepsilon N$ will be used later to deduce the moment bound (7.16) for the coarse-grained variables $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$.

- The weight function w in (9.3) and (9.4) allows us to include the case $\psi \equiv 1$, and to control the spatial decay when $\|\varphi\|_p < \infty$ and the support of ψ moves away to ∞ , which is needed for the bounds in (6.15) and (7.17).
- The bound (9.3) implies the moment bound for the point-to-plane partition function $Z_{0,N'}(z, \mathbf{1})$ in (6.15). Indeed, assume $z = 0$ and let $w(x) = e^{-\|x\|}$, $\psi \equiv 1$, and choose φ such that $\varphi_{N'}(y) = N' \cdot \mathbf{1}_{y=0}$ and $\mathcal{Z}_{N'}^{\beta_N}(\varphi, \mathbf{1}) = Z_{0,N'}(0, \mathbf{1})$. Then (9.3) gives

$$|\mathbb{E}[(Z_{0,N'}(0, \mathbf{1}) - 1)^h]| \leq C \|w_{N'}\|_{\ell^q}^h \leq C' N^{\frac{h}{q}}, \quad (9.5)$$

where the exponent can be made arbitrarily small by choosing q large.

- Theorem 9.1 can be strengthened by showing that the constants C and C' in (9.3) and (9.4) are proportional to pq , which allows one to send $p = p_N \downarrow 1$ and $q = q_N \uparrow \infty$. This was done in the subcritical regime in [LZ23, Theorem 1.5] and in the quasi-critical regime in [CCR23] (cf. Section 5.6). There is in fact a unified bound for all disorder strength β_N in or below the critical window, which says that the h -th centred moment of the averaged partition function $\mathcal{Z}_N^{\beta_N}(\varphi) := \mathcal{Z}_N^{\beta_N}(\varphi, \mathbf{1})$ can be controlled by its variance raised to the power $h/2$, provided φ is sufficiently regular and is supported on an appropriate spatial scale such that the variance remains bounded (we considered flat terminal condition $\psi = \mathbf{1}$ for simplicity): for some $C = C(h) < \infty$

$$|\mathbb{E}\left[\left(\mathcal{Z}_N^{\beta_N}(\varphi) - \mathbb{E}[\mathcal{Z}_N^{\beta_N}(\varphi)]\right)^h\right]| \leq C \operatorname{Var}[\mathcal{Z}_N^{\beta_N}(\varphi)]^{\frac{h}{2}}.$$

The details can be found in the forthcoming paper [CSZ25b]. Such a bound in the quasi-critical regime can be found in [CCR23, Proposition 2.2].

Proof Sketch. The proof of Theorem 9.1 is based on the analysis of collision diagrams of h independent random walks (see Figure 6). Let us first explain at the heuristic level before formalising it in terms of operators. For simplicity, we will assume $N' = N$.

Recalling from (7.2) the polynomial chaos expansion for $\mathcal{Z}_N(\varphi, \psi) := \mathcal{Z}_N^{\beta_N}(\varphi, \psi)$, we can write

$$\begin{aligned} |\mathbb{E}[(\mathcal{Z}_N(\varphi, \psi) - \mathbb{E}[\mathcal{Z}_N(\varphi, \psi)])^h]| &= \frac{1}{N^h} \times \\ &\mathbb{E}\left[\left(\sum_{k=1}^{\infty} \sum_{\substack{x_1, \dots, x_k \in \mathbb{Z}^2 \\ 0 < n_1 < \dots < n_k < N}} q_{n_1}(\varphi_N, x_1) \xi_N(n_1, x_1) \left\{ \prod_{j=2}^k q_{n_j-n_{j-1}}(x_j - x_{j-1}) \xi_N(n_j, x_j) \right\} q_{N-n_k}(x_k, \psi_N)\right)^h\right], \end{aligned} \quad (9.6)$$

where $\varphi_N(x) := \varphi(\frac{x}{\sqrt{N}})$, $\psi_N(x) := \psi(\frac{x}{\sqrt{N}})$,

$$q_m(\varphi_N, z) := \sum_{x \in \mathbb{Z}^2} \varphi_N(x) q_m(z - x) \quad \text{and} \quad q_m(z, \psi_N) := \sum_{y \in \mathbb{Z}^2} q_m(y - z) \psi_N(y). \quad (9.7)$$

We can expand the power inside the expectation in (9.6) as

$$\begin{aligned} & \sum_{k_1, \dots, k_h \geq 1} \sum_{\substack{(n_i^{(r)}, x_i^{(r)}) \in \mathbb{N} \times \mathbb{Z}^2 \\ 1 \leq i \leq k_r, 1 \leq r \leq h}} \prod_{r=1}^h q_{n_1^{(r)}}(\varphi_N, x_1^{(r)}) \left\{ \prod_{j=2}^{k_r} q_{n_j^{(r)} - n_{j-1}^{(r)}}(x_j^{(r)} - x_{j-1}^{(r)}) \right\} q_{N-n_{k_r}^{(r)}}(x_{k_r}^{(r)}, \psi_N) \\ & \quad \times \mathbb{E} \left[\prod_{r=1}^h \prod_{j=1}^{k_r} \xi_N(n_j^{(r)}, x_j^{(r)}) \right], \end{aligned} \quad (9.8)$$

which sums over h sequences of time-space points $\Gamma^{(r)} := (n_i^{(r)}, x_i^{(r)})_{1 \leq i \leq k_r}$ with $0 < n_1^{(r)} < \dots < n_{k_r}^{(r)} < N$, and consecutive points in each sequence are connected by the random walk transition kernel q . Since ξ_N is an i.i.d. family with zero mean, the factor $\mathbb{E}[\prod \xi_N(\cdot, \cdot)]$ is non-zero only when each $\xi_N(n, x)$ that appears in the product appears at least twice. It was shown in [CSZ23a] that the dominant contribution (as $N \rightarrow \infty$) to the expansion in (9.8) comes from configurations of $(\Gamma^{(r)})_{1 \leq r \leq h}$ where $\xi_N(n_j^{(r)}, x_j^{(r)})$ are matched in pairs, with the expectation of each matched pair contributing a factor of σ_N^2 as in (2.1). In other words, any $\xi_N(n, x)$ that is collected by one of the h sequences $\Gamma^{(r)}$ is collected by exactly 2 of the h sequences, and we say there are only pair collisions among $(\Gamma^{(r)})_{1 \leq r \leq h}$. We will focus on such configurations of $\Gamma^{(r)}$, $1 \leq r \leq h$, and keep track of the matchings between points in $\Gamma^{(r)}$. In particular, we will partition the sum in (9.8) according to streaks of consecutive matching between points in the same pair of sequences $\Gamma^{(i)}$ and $\Gamma^{(j)}$.

As illustrated in Figure 6, we will rewrite the expansion (9.8) of the h -th moment by first identifying the time intervals $[a_\ell, b_\ell]$, $1 \leq \ell \leq m$, where the streaks of consecutive matching occurs (the wiggle lines in Figure 6). Denote by $I_\ell = \{i_\ell, j_\ell\}$ the indices for the pair of $\Gamma^{(r)}$ for which matching occurs in the ℓ -th streak. Our decomposition ensures that $I_\ell \neq I_{\ell+1}$. Summing over m , a_ℓ , b_ℓ , I_ℓ , and the spatial locations of the solid and hollow circles at times a_ℓ and b_ℓ in Figure 6, we can then rewrite the expansion (9.8) using kernels:

$$\begin{aligned} & \frac{1}{N^h} \sum_{m=1}^{\infty} (\sigma_N^2)^m \sum_{\substack{I_1, \dots, I_m \subset \{1, \dots, h\} \\ |I_\ell|=2, I_\ell \neq I_{\ell+1}}} \sum_{\substack{0 < a_1 \leq b_1 < \dots < a_m \leq b_m < N \\ \mathbf{x}_1, \dots, \mathbf{x}_m, \mathbf{y}_1, \dots, \mathbf{y}_m \in (\mathbb{Z}^2)^h}} \mathbf{Q}_{a_1}^{*, I_1}(\varphi_N, \mathbf{x}_1) \mathbf{U}_N^{I_1}(b_1 - a_1, \mathbf{y}_1 - \mathbf{x}_1) \\ & \quad \times \prod_{\ell=2}^m \mathbf{Q}_{a_\ell - b_{\ell-1}}^{I_{\ell-1}, I_\ell}(\mathbf{y}_{\ell-1}, \mathbf{x}_\ell) \mathbf{U}_N^{I_\ell}(b_\ell - a_\ell, \mathbf{x}_\ell, \mathbf{y}_\ell) \mathbf{Q}_{N-b_m}^{I_m, *}(\mathbf{y}_m, \psi_N) + o(1), \end{aligned} \quad (9.9)$$

where $o(1)$ accounts for the contribution from configurations of $(\Gamma^{(r)})_{1 \leq r \leq h}$ that contain triple collisions, that is, there exists (n, x) that belongs to more than two different $\Gamma^{(r)}$.

The kernels $\mathbf{U}_N^{I_\ell}$, $\mathbf{Q}_{a_1}^{*, I_1}$, $\mathbf{Q}_{a_\ell - b_{\ell-1}}^{I_{\ell-1}, I_\ell}$, $\mathbf{Q}_{N-b_m}^{I_m, *}$ are defined as follows:

1. **Replica Evolution.** Recall U_N from (8.4). For $I = \{i < j\} \subset \{1, \dots, h\}$ and $\mathbf{x} = (x_a)_{1 \leq a \leq h}$, $\mathbf{y} = (y_a)_{1 \leq a \leq h} \in (\mathbb{Z}^2)^h$, we define

$$\mathbf{U}_N^I(n, \mathbf{x}, \mathbf{y}) := \mathbb{1}_{\{x_i=x_j\}} \cdot \left(\prod_{a \in \{1, \dots, h\} \setminus \{i, j\}} q_n(y_a - x_a) \right) U_N(n, y_i - x_i) \cdot \mathbb{1}_{\{y_i=y_j\}}. \quad (9.10)$$

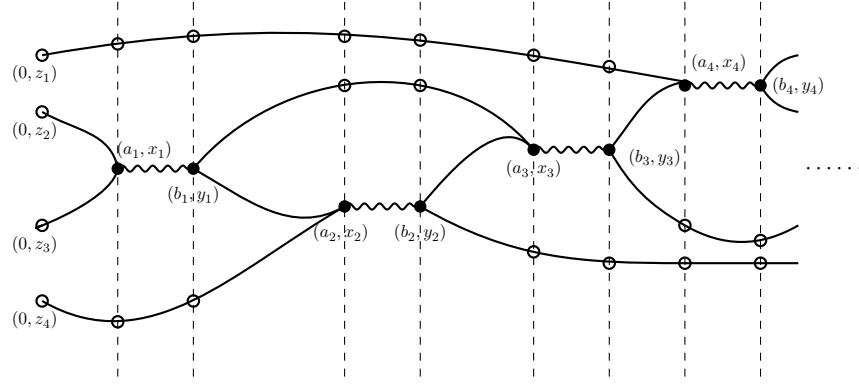


FIGURE 6. Illustration of the expansion in (9.8) for the fourth moment. Each wiggle line represents a streak of consecutive matchings of time-space points in the same pair of sequences $\Gamma^{(i)}$ and $\Gamma^{(j)}$, which is assigned weight $\sigma_N^2 U_N(b - a, y - x)$ (see (8.5)) if the streak starts at (a, x) and ends at (b, y) . Each solid curve between two circles is assigned a random walk transition kernel q , while each hollow circle at a point (t, y) arises from a sequence $\Gamma^{(r)}$ that is unmatched at time t and the Chapman-Kolmogorov decomposition of the associated transition kernel $q_{n_{j-1}^{(r)}, n_j^{(r)}}(x_{j-1}^{(r)}, x_j^{(r)}) = \sum_y q_{n_{j-1}^{(r)}, t}(x_{j-1}^{(r)}, y) q_{t, n_j^{(r)}}(y, x_j^{(r)})$ at the intermediate time t . In each time interval between consecutive dashed vertical lines, the product of the curve weights define a kernel which we call Replica Evolution (see (9.10)) if it contains a wiggle line, or Constrained Evolution (see (9.11)) if there is no wiggle line in this time interval.

2. Constrained Evolution. For $I = \{i < j\}$, $J = \{k < \ell\} \subset \{1, \dots, h\}$ and $\mathbf{x}, \mathbf{y} \in (\mathbb{Z}^2)^h$, we define

$$\mathbf{Q}_n^{I,J}(\mathbf{y}, \mathbf{x}) := \mathbf{1}_{\{y_i = y_j\}} \cdot \left(\prod_{a=1}^h q_n(x_a - y_a) \right) \cdot \mathbf{1}_{\{x_k = x_l\}}. \quad (9.11)$$

Removing the spatial constraint on either side, we define

$$\mathbf{Q}_n^{*,J}(\mathbf{y}, \mathbf{x}) := \prod_{a=1}^h q_n(x_a - y_a) \cdot \mathbf{1}_{\{x_k = x_l\}}, \quad \mathbf{Q}_n^{I,*}(\mathbf{y}, \mathbf{x}) := \mathbf{1}_{\{y_i = y_j\}} \prod_{a=1}^h q_n(x_a - y_a), \quad (9.12)$$

while

$$\mathbf{Q}_n^{*,J}(\varphi_N, \mathbf{x}) := \sum_{\mathbf{z} \in (\mathbb{Z}^2)^h} \left(\prod_{i=1}^n \varphi\left(\frac{z_i}{\sqrt{N}}\right) \right) \mathbf{Q}_n^{*,J}(\mathbf{z}, \mathbf{x}), \quad (9.13)$$

and similarly for $\mathbf{Q}_n^{I,*}(\mathbf{y}, \psi_N)$.

To decouple the summation range for the time increments $a_1, b_1 - a_1, a_2 - b_1, \dots$ in (9.9), we consider the Laplace transforms of the kernels above:

$$\begin{aligned} Q_{\lambda,N}^{I,J}(\mathbf{y}, \mathbf{z}) &:= \sum_{n=1}^{2N} e^{-\frac{\lambda}{N}n} Q_n^{I,J}(\mathbf{y}, \mathbf{z}), \quad \mathbf{y}, \mathbf{z} \in (\mathbb{Z}^2)^h, \\ U_{\lambda,N}^J(\mathbf{y}, \mathbf{z}) &:= \sum_{n=0}^{2N} e^{-\frac{\lambda}{N}n} U_N^J(n, \mathbf{y}, \mathbf{z}), \quad \mathbf{y}, \mathbf{z} \in (\mathbb{Z}^2)^h. \end{aligned} \quad (9.14)$$

Inserting the factor $e^\lambda e^{-\frac{\lambda}{N}\sum_{i=1}^{m-1}((a_i-b_{i-1})+(b_{i+1}-a_i))-\frac{\lambda}{N}(N-b_m)} = 1$ in (9.9) and enlarging the range of summation for $a_1, a_2 - b_1, \dots$ to $\{1, \dots, 2N\}$ and the range of summation for $b_i - a_i$ to $\{0, \dots, 2N\}$, we obtain the upper bound

$$\begin{aligned} \frac{1}{N^{h+1}} \sum_{m=1}^{\infty} (\sigma_N^2)^m \sum_{\substack{I_1, \dots, I_m \subset \{1, \dots, h\} \\ |I_\ell|=2, I_\ell \neq I_{\ell+1}}} &Q_{\lambda,N}^{*,I_1}(\varphi_N, \mathbf{x}_1) U_{\lambda,N}^{I_1}(b_1 - a_1, \mathbf{y}_1 - \mathbf{x}_1) \\ &\times \prod_{\ell=2}^m Q_{\lambda,N}^{I_{\ell-1}, I_\ell}(\mathbf{y}_{\ell-1}, \mathbf{x}_\ell) U_{\lambda,N}^{I_\ell}(\mathbf{x}_\ell, \mathbf{y}_\ell) \times Q_{\lambda,N}^{I_m, *}(y_m, \psi_N), \end{aligned} \quad (9.15)$$

where there is an additional factor N^{-1} because to introduce the last kernel $Q_{\lambda,N}^{I_m, *}$, we need to sum over $a_{m+1} - b_m \in \{1, \dots, 2N\}$, instead of $a_{m+1} = N$ as in (9.9). The additional sum over a_{m+1} is compensated by the averaging factor N^{-1} , and the fact this averaging over a_{m+1} is comparable to having fixed $a_{m+1} = N$ is justified by the following inequality (see [CSZ23a, (6.11)])

$$Q_{N-b_m}^{I_m, *}(y_m, z') \leq \frac{c}{N} \sum_{a_{m+1} \in \{N+1, \dots, 2N\}} Q_{a_{m+1}-b_m}^{I_m, *}(y_m, z').$$

Remark 9.3. Instead of averaging over the additional variable a_{m+1} , one could also take the maximum over b_m of the last random walk operator $Q_{N-b_m}^{I_m, *}(y_m, \psi_N)$ (which requires no assumption on ψ_N). This alternative approach, which has been developed in [CCR23], leads to a slightly sharper version of the bound (9.4). The key point is that the maximal random walk operator described above is bounded in L^q for $q > 1$, thanks to a variant of the Hardy-Littlewood maximal inequality, see [CCR23, Lemma 4.18].

The kernels $U_{\lambda,N}^I$ and $Q_{\lambda,N}^{*,I_1}, Q_{\lambda,N}^{I_{\ell-1}, I_\ell}, Q_{\lambda,N}^{I_m, *}$ define integral operators and we can bound (9.15) in terms of their operator norms. More precisely, the indicator constraints in (9.10) and (9.11) suggest that we should regard $U_{\lambda,N}^I$, with $I = \{i < j\}$, as an integral operator acting on functions defined on the following subset of $(\mathbb{Z}^2)^h$:

$$(\mathbb{Z}^2)_I^h := \{\mathbf{x} \in (\mathbb{Z}^2)^h : x_i = x_j\}. \quad (9.16)$$

Let $\ell^q((\mathbb{Z}^2)_I^h)$ be the space of functions $f : (\mathbb{Z}^2)_I^h \rightarrow \mathbb{R}$ with $\|f\|_q := (\sum_{x \in (\mathbb{Z}^2)_I^h} |f(x)|^q)^{1/q} < \infty$. For $q > 1$, we regard $U_{\lambda,N}^I$ as an operator from $\ell^q((\mathbb{Z}^2)_I^h)$ to $\ell^q((\mathbb{Z}^2)_I^h)$ with operator norm

$$\|U_{\lambda,N}^I\|_q := \sup_{\substack{f, g: (\mathbb{Z}^2)_I^h \rightarrow \mathbb{R} \\ \|f\|_p, \|g\|_q = 1}} \langle f, U_{\lambda,N}^I g \rangle,$$

where $\langle \cdot, \cdot \rangle$ is the inner product and $\frac{1}{p} + \frac{1}{q} = 1$.

For $I = \{i < j\}$ and $J = \{k < l\}$, we will regard $\mathbf{Q}_{\lambda,N}^{I,J}$ as an integral operator from $\ell^q((\mathbb{Z}^2)_J^h)$ to $\ell^q((\mathbb{Z}^2)_I^h)$ with operator norm $\|\mathbf{Q}_{\lambda,N}^{I,J}\|_q$ defined as before. Similarly, $\mathbf{Q}_{\lambda,N}^{*,I}$ will be an integral operator from $\ell^q((\mathbb{Z}^2)_I^h)$ to $\ell^q((\mathbb{Z}^2)_I^h)$, and $\mathbf{Q}_{\lambda,N}^{I,*}$ an integral operator from $\ell^q((\mathbb{Z}^2)_I^h)$ to $\ell^q((\mathbb{Z}^2)_I^h)$. Using the norms of these operators for a fixed $q > 1$, we can now bound (9.15) by

$$\begin{aligned} & \frac{Ce^\lambda}{N^{h+1}} \sum_{m=1}^{\infty} (\sigma_N^2)^m \sum_{I_1, \dots, I_m} \left\langle \varphi_N^{\otimes h}, \mathbf{Q}_{\lambda,N}^{*,I_1} \mathbf{U}_{\lambda,N}^{I_1} \mathbf{Q}_{\lambda,N}^{I_1, I_2} \dots \mathbf{U}_{\lambda,N}^{I_m} \mathbf{Q}_{\lambda,N}^{I_m, *} \psi_N^{\otimes h} \right\rangle \\ & \leq \frac{Ce^\lambda}{N^{h+1}} \sum_{m=1}^{\infty} (\sigma_N^2)^m \sum_{I_1, \dots, I_m} \|\varphi_N^{\otimes h}\|_p \|\mathbf{Q}_{\lambda,N}^{*,I_1}\|_q \|\mathbf{U}_{\lambda,N}^{I_1}\|_q \|\mathbf{Q}_{\lambda,N}^{I_1, I_2}\|_q \dots \|\mathbf{U}_{\lambda,N}^{I_m}\|_q \|\mathbf{Q}_{\lambda,N}^{I_m, *}\|_q \|\psi_N^{\otimes h}\|_q, \end{aligned} \quad (9.17)$$

where $\phi^{\otimes h}(\mathbf{x}) := \prod_{i=1}^h \phi(x_i)$. The key is to bound these operator norms, which is summarised in the following proposition.

Proposition 9.4. *Fix an integer $h \geq 2$ and $p, q > 1$ with $\frac{1}{p} + \frac{1}{q} = 1$. There exists $C = C(h) < \infty$ such that uniformly in $\lambda > 0$, N large, and $I \neq J \subset \{1, \dots, h\}$ with $|I| = |J| = 2$, we have*

$$\|\mathbf{Q}_{\lambda,N}^{I,J}\|_q \leq C pq; \quad (9.18)$$

$$\|\mathbf{Q}_{\lambda,N}^{*,I}\|_q \leq C N^{\frac{1}{q}}, \quad \|\mathbf{Q}_{\lambda,N}^{I,*}\|_q \leq C N^{\frac{1}{p}}; \quad (9.19)$$

$$\text{and} \quad \|\mathbf{U}_{\lambda,N}^I\|_q \leq \frac{C}{(\log \lambda) \sigma_N^2}. \quad (9.20)$$

Substituting these bounds into (9.17) leads to a geometric series that is convergent provided $\lambda > 0$ is chosen large enough. To arrive at the final form in Theorem 9.1, we need two modifications:

- Firstly, we need to modify (9.17) by introducing a weight function $w_N^{\otimes h}(\mathbf{x})$ with $w_n(x) = w(x/\sqrt{N})$. This is done by rewriting (9.15) in terms of the weighted operators

$$\widehat{\mathbf{Q}}_{\lambda,N}^{I,J}(\mathbf{y}, \mathbf{z}) := \mathbf{Q}_{\lambda,N}^{I,J}(\mathbf{y}, \mathbf{z}) \frac{w_N^{\otimes h}(\mathbf{y})}{w_N^{\otimes h}(\mathbf{z})} \quad \text{and} \quad \widehat{\mathbf{U}}_{\lambda,N}^J(\mathbf{y}, \mathbf{z}) := \mathbf{U}_N^J(\mathbf{y}, \mathbf{z}) \frac{w_N^{\otimes h}(\mathbf{y})}{w_N^{\otimes h}(\mathbf{z})}. \quad (9.21)$$

This allows us to replace the bound in (9.17) by

$$\begin{aligned} & \left| \mathbb{E} \left[\left(\mathcal{Z}_N^{\beta_N}(\varphi, \psi) - \mathbb{E}[\mathcal{Z}_N^{\beta_N}(\varphi, \psi)] \right)^h \right] \right| \\ & \leq \frac{Ce^\lambda}{N^{h+1}} \sum_{m \geq 1} (\sigma_N^2)^m \sum_{I_1, \dots, I_m} \left\langle \frac{\varphi_N^{\otimes h}}{w_N^{\otimes h}}, \widehat{\mathbf{Q}}_{\lambda,N}^{*,I_1} \widehat{\mathbf{U}}_{\lambda,N}^{I_1} \widehat{\mathbf{Q}}_{\lambda,N}^{I_1, I_2} \dots \widehat{\mathbf{U}}_{\lambda,N}^{I_m} \widehat{\mathbf{Q}}_{\lambda,N}^{I_m, *} w_N^{\otimes h} \psi_N^{\otimes h} \right\rangle \\ & \leq \frac{Ce^\lambda}{N^{h+1}} \sum_{m \geq 1} (\sigma_N^2)^m \sum_{I_1, \dots, I_m} \left\| \frac{\varphi_N^{\otimes h}}{w_N^{\otimes h}} \right\|_p \|\widehat{\mathbf{Q}}_{\lambda,N}^{*,I_1}\|_q \|\widehat{\mathbf{U}}_{\lambda,N}^{I_1}\|_q \|\widehat{\mathbf{Q}}_{\lambda,N}^{I_1, I_2}\|_q \dots \|\widehat{\mathbf{U}}_{\lambda,N}^{I_m}\|_q \|\widehat{\mathbf{Q}}_{\lambda,N}^{I_m, *}\|_q \|w_N^{\otimes h} \psi_N^{\otimes h}\|_q. \end{aligned}$$

The bounds in Proposition 9.4 remain valid for $\widehat{\mathbf{Q}}_{\lambda,N}$ and $\widehat{\mathbf{U}}_{\lambda,N}$, see [CSZ23a, Prop. 6.6].

- Secondly, to obtain the pre-factor $1/\log(1 + \frac{N}{N'})$ in Proposition 9.1 when $N' < N$, we need to replace the operator $\widehat{\mathbf{U}}_{\lambda,N}^I$ by a variant that takes into account the shorter

time horizon, even though β_N and σ_N^2 remain the same. For more details, see [CSZ23a, Prop. 6.6 and (6.24)]. \square

9.2. FUNCTIONAL INEQUALITIES. In this subsection, we sketch how to obtain the bounds in Proposition 9.4. We note that L^2 -variants of (9.18) first appeared in the work of Dell'Antonio, Figari and Teta [DFT94] (see the proof of Lemma 3.1 therein) and was then used in [GQT21, (5.2)] to bound the moments of the mollified SHE (1.3) in the critical window. Since it is non-trivial to adapt the Fourier transform techniques of [DFT94] to our discrete setting, we instead work directly in real space. Our method is robust enough to be applied to the coarse-grained model introduced in Section 7. Our extension from L^2 bounds to L^q bounds for arbitrary $q > 1$ is also important as discussed in (9.5).

The inequality (9.18) is reminiscent of the Hardy-Littlewood-Sobolev (HLS) inequality [LL01, Theorem 4.3]: for $f, g : \mathbb{R}^d \rightarrow \mathbb{R}$, $p, r > 1$ and $0 < \nu < d$ satisfying $\frac{1}{p} + \frac{1}{r} + \frac{\nu}{d} = 2$,

$$\left| \iint_{\mathbb{R}^d \times \mathbb{R}^d} \frac{f(x)g(y)}{|x-y|^\nu} dx dy \right| \leq C(d, \nu, p) \|f\|_p \|g\|_r. \quad (9.22)$$

Indeed, (9.18) is equivalent to showing that for all $f \in \ell^p((\mathbb{Z}^2)_I^h)$ and $g \in \ell^q((\mathbb{Z}^2)_J^h)$

$$\sum_{\mathbf{x} \in (\mathbb{Z}^2)_I^h, \mathbf{y} \in (\mathbb{Z}^2)_J^h} f(\mathbf{x}) Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) g(\mathbf{y}) \leq C p q \|f\|_p \|g\|_q, \quad (9.23)$$

where for $\mathbf{x} \in (\mathbb{Z}^2)_I^h$ and $\mathbf{y} \in (\mathbb{Z}^2)_J^h$, we have

$$Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) = \sum_{n=1}^{2N} e^{-\frac{\lambda}{N} n} \prod_{i=1}^h q_n(y_i - x_i) \leq \begin{cases} \frac{C}{(1 + |\mathbf{y} - \mathbf{x}|^2)^{h-1}} & \text{for } |\mathbf{x} - \mathbf{y}| \leq \sqrt{N}, \\ \frac{C}{N^{h-1}} e^{-\frac{|\mathbf{y} - \mathbf{x}|^2}{CN}} & \text{for } |\mathbf{x} - \mathbf{y}| > \sqrt{N}, \end{cases} \quad (9.24)$$

with a lower bound of the same order. In other words, for $|\mathbf{x} - \mathbf{y}| \leq \sqrt{N}$, $Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y})$ is comparable to the Green's function of a random walk in \mathbb{Z}^{2h} . Note that in (9.23), the variables \mathbf{x} and \mathbf{y} are being summed in a space of dimension $2(h-1)$ which corresponds to dimension d in (9.22), $r = q$, while at first glance, ν in (9.22) is comparable to the exponent $2(h-1)$ in $Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y})$. This will be exactly the borderline case where the HLS inequality fails because $\nu = d$. But what saves us is that when we restrict to $\mathbf{x} \in (\mathbb{Z}^2)_I^h$ and $\mathbf{y} \in (\mathbb{Z}^2)_J^h$, we have made the identifications $x_i = x_j$ and $y_k = y_l$, which makes $(1 + |\mathbf{y} - \mathbf{x}|^2)^{1-h}$ more regular because $|\mathbf{y} - \mathbf{x}|$ is greater than the Euclidean distance between $\mathbf{x} \in (\mathbb{Z}^2)_I^h$ and $\mathbf{y} \in (\mathbb{Z}^2)_J^h$ regarded as two points in $\mathbb{Z}^{2(h-1)}$. Indeed, $|\mathbf{x} - \mathbf{y}|^2$ contains terms such as $|x_i - y_i|^2 + |x_j - y_j|^2 \geq \frac{1}{3}(|x_i - y_i|^2 + |y_i - y_j|^2)$, which penalises the summation in y_i and y_j in (9.23) when they are far apart.

A naive attempt to bound (9.23) is to write $Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) = Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y})^{\frac{1}{p}} Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y})^{\frac{1}{q}}$ and then apply Hölder's inequality to bound

$$\langle f, Q_{\lambda, N}^{I, J} g \rangle \leq \left(\sum_{\mathbf{x} \in (\mathbb{Z}^2)_I^h, \mathbf{y} \in (\mathbb{Z}^2)_J^h} |f(\mathbf{x})|^p Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) \right)^{1/p} \left(\sum_{\mathbf{x} \in (\mathbb{Z}^2)_I^h, \mathbf{y} \in (\mathbb{Z}^2)_J^h} Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) |g(\mathbf{y})|^q \right)^{1/q}. \quad (9.25)$$

This will not work because it can be shown that

$$\sum_{\mathbf{y} \in (\mathbb{Z}^2)_J^h} Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) \geq C \log |x_i - x_j|. \quad (9.26)$$

The trick is to insert a factor $\frac{\gamma(\mathbf{x}, \mathbf{y})}{\gamma(\mathbf{x}, \mathbf{y})}$ and then apply Hölder such that the factor γ or $1/\gamma$ will remove the logarithmic singularity. More precisely, we choose a small $\alpha > 0$ and bound

$$\begin{aligned} \sum_{\mathbf{x} \in (\mathbb{Z}^2)_I^h, \mathbf{y} \in (\mathbb{Z}^2)_J^h} f(\mathbf{x}) Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) g(\mathbf{y}) &= \sum_{\mathbf{x}, \mathbf{y}} f(\mathbf{x}) Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) \frac{|x_k - x_\ell|^\alpha}{|y_i - y_j|^\alpha} \cdot \frac{|y_i - y_j|^\alpha}{|x_k - x_\ell|^\alpha} g(\mathbf{y}) \\ &\leq \left(\sum_{\mathbf{x}, \mathbf{y}} f(\mathbf{x})^p Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) \frac{|x_k - x_\ell|^{\alpha p}}{|y_i - y_j|^{\alpha p}} \right)^{\frac{1}{p}} \left(\sum_{\mathbf{x}, \mathbf{y}} Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) \frac{|y_i - y_j|^{\alpha q}}{|x_k - x_\ell|^{\alpha q}} g(\mathbf{y})^q \right)^{\frac{1}{q}}. \end{aligned} \quad (9.27)$$

A straightforward computation shows that

$$\sum_{\mathbf{y} \in (\mathbb{Z}^2)_J^h} Q_{\lambda, N}^{I, J}(\mathbf{x}, \mathbf{y}) \frac{|x_k - x_\ell|^{\alpha p}}{|y_i - y_j|^{\alpha p}} \leq C,$$

which gives (9.18) (see [CSZ23a, Section 6.2] for details). Keeping track of the constants more carefully and optimising over α will reveal the dependence on p and q as in (9.23) (see [LZ23, Prop. 3.3]).

The proof of (9.19) and (9.20) are similar. To see how the factor $(\sigma_N^2 \log \lambda)^{-1}$ arises in (9.20), we first recall that $U_{\lambda, N}^I$ differs from $Q_{\lambda, N}^{I, I}(\mathbf{x}, \mathbf{y})$ in that the transition kernels $q_n(y_i - x_i)q_n(y_j - x_j)$ in (9.11) are replaced by $U_N(n, y_i - x_i)$ in (9.10). Here we can apply the naive approach (9.25):

$$\sum_{\mathbf{x}, \mathbf{y} \in (\mathbb{Z}^2)_I^h} f(\mathbf{x}) U_{\lambda, N}^I(\mathbf{x}, \mathbf{y}) g(\mathbf{x}) \leq \left(\sum_{\mathbf{x}, \mathbf{y}} f(\mathbf{x})^p U_{\lambda, N}^I(\mathbf{x}, \mathbf{y}) \right)^{\frac{1}{p}} \left(\sum_{\mathbf{x}, \mathbf{y}} U_{\lambda, N}^I(\mathbf{x}, \mathbf{y}) g(\mathbf{x})^q \right)^{\frac{1}{q}},$$

where by (9.10), the definition of $U_N(n)$ in (8.7), and the asymptotics in (8.17) and (8.14),

$$\sum_{\mathbf{y} \in (\mathbb{Z}^2)_I^h} U_{\lambda, N}^I(\mathbf{x}, \mathbf{y}) = \sum_{n=1}^{2N} e^{-\frac{\lambda}{N} n} U_N(n) \leq C \log N \int_0^\infty e^{-\lambda t} G_\vartheta(t) dt = \frac{C}{\sigma_N^2 \log \lambda}.$$

9.3. MOMENT BOUNDS FOR THE COARSE-GRAINED MODEL. In this section, we explain how to obtain the moment bounds (7.16) and (7.17), which are crucial in the application of the Lindeberg principle to the coarse-grained model $\mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi|\Theta)$ as outlined in Section 7.2.

To deduce (7.16), the key observation is that the chaos expansion for $\Theta_{N, \varepsilon}(\vec{\mathbf{i}}, \vec{\mathbf{a}})$ in (7.9) and (7.10) is similar to the chaos expansion for the averaged partition function $\mathcal{Z}_N^{\beta_N}(\varphi, \mathbf{1})$ in (7.2). Indeed, if we follow the renewal interpretation of the sequence of time-space points $(n_1, z_1), \dots, (n_r, z_r)$ in the chaos expansion for $\mathcal{Z}_N^{\beta_N}(\varphi, \mathbf{1})$ as we did in Section 8.1, which is based on an expansion for $\mathbb{E}[\mathcal{Z}_N^{\beta_N}(\varphi, \mathbf{1})^2]$, then the first renewal point (n_1, z_1) is sampled according to the mass function

$$M(n_1, z_1) := \frac{1}{N^2} \left(\sum_{z_0 \in \mathbb{Z}^2} \varphi(\frac{z_0}{\sqrt{N}}) q_{n_1}(z_1 - z_0) \right)^2, \quad 0 < n_1 < N, z_1 \in \mathbb{Z}^2.$$

It is easy to see that the total mass $\sum_{0 < n_1 < N, z \in \mathbb{Z}^2} M(n_1, z_1)$ converges to a positive constant as $N \rightarrow \infty$, and if $\varphi(x) \geq 1$ for all $\|x\| \leq 1$, then uniformly in $0 < n_1 < N$ and $\|z_1\| \leq \sqrt{N}$, $M(n_1, z_1)$ is bounded from below by c/N^2 , i.e., a positive multiple of the uniform probability mass function on the set of such (n_1, z_1) .

If we apply the same renewal interpretation to the chaos expansion for $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ in (7.9), then we see that the first renewal point (n_1, z_1) is sampled uniformly from the time-space box with $0 < n_1 < \varepsilon N$ and $\|z_1\| \leq \sqrt{\varepsilon N}$, which can be compared to the chaos expansion for the averaged partition function $\mathcal{Z}_{\varepsilon N}^{\beta_N}(\varphi, \mathbf{1})$ defined in (9.2) with $\varphi(x) = 1$ for $\|x\| \leq 1$, time horizon $N' = \varepsilon N$ and spatial averaging on the scale $\sqrt{N'}$. This allows us to apply Theorem 9.1 with $N' = \varepsilon N$, $\varphi(x) = 1_{\{\|x\| \leq 1\}}$ and $\psi \equiv 1$ to deduce that for all N large,

$$\mathbb{E}[\Theta_{N,\varepsilon}(\vec{i}, \vec{a})^4] < \frac{C}{\log \frac{1}{\varepsilon}}. \quad (9.28)$$

See [CSZ23a, Section 7.2] for more details.

The moment bound (7.17) for the coarse-grained partition function follows from the following analogue of Theorem 9.1.

Theorem 9.5. *Let $\mathcal{Z}_{N,\varepsilon}(\varphi, \psi) := \mathcal{Z}_\varepsilon^{(\text{cg})}(\varphi, \psi | \Theta_{N,\varepsilon})$ be the coarse-grained partition function defined in (7.11). Further, assume that $\|\psi\|_\infty < \infty$ and ψ is supported on a ball B (possibly $B = \mathbb{R}^2$). Then for any $p, q \in (1, \infty)$ with $\frac{1}{p} + \frac{1}{q} = 1$ and any $w : \mathbb{R}^2 \rightarrow (0, \infty)$ such that $\log w$ is Lipschitz continuous, there exists $C \in (0, \infty)$ such that uniformly in $\varepsilon \in (0, 1)$,*

$$\limsup_{N \rightarrow \infty} \mathbb{E}\left[\left(\mathcal{Z}_{N,\varepsilon}(\varphi, \psi) - \mathbb{E}[\mathcal{Z}_{N,\varepsilon}(\varphi, \psi)]\right)^4\right] \leq C \varepsilon^{\frac{4}{p}} \left\| \frac{\varphi_{1/\varepsilon}}{w_{1/\varepsilon}} \right\|_{\ell^p(\mathbb{Z}^2)}^4 \|\psi\|_\infty^4 \|w \mathbf{1}_B\|_q^4, \quad (9.29)$$

where for $\phi : \mathbb{R}^2 \rightarrow \mathbb{R}$, $\phi_\varepsilon : \mathbb{Z}^2 \rightarrow \mathbb{R}$ is defined as above Theorem 9.1.

This is Theorem 8.1 in [CSZ23a], which is proved by adapting the proof for Theorem 9.1. Complications arise because the coarse-grained disorder variables $\Theta_{N,\varepsilon}(\vec{i}, \vec{a})$ is a family of dependent random variables. We refer to [CSZ23a, Section 8] for details.

9.4. MOMENTS OF THE CRITICAL 2d SHF. We already gave the second moment of the Critical 2d SHF in Theorem 6.1. Now we give a formula for higher integer moments of the SHF, which arises as the limit of the expansion in (9.9) and hence consists of an infinite series. In particular, the kernels $U_N^I, Q_{a_1}^{*, I_1}, Q_{a_\ell - b_{\ell-1}}^{I_{\ell-1}, I_\ell}, Q_{N-b_m}^{I_m, *}$ in (9.9) will be replaced by their continuum analogues as follows.

Similar to (9.16), for $I = \{i < j\} \subset \{1, \dots, h\}$, define

$$(\mathbb{R}^2)_I^h := \{\mathbf{x} = (x_1, \dots, x_h) \in (\mathbb{R}^2)^h : x_i = x_j\}. \quad (9.30)$$

1. Replica Evolution. The continuum analogue of $U_N^I(n, \mathbf{x}, \mathbf{y})$ in (9.10) is given by the kernel with density w.r.t. Lebesgue measure on $(\mathbb{R}^2)_I^h$

$$G_{\vartheta, t}^I(\mathbf{x}, \mathbf{y}) := \prod_{\ell \in \{1, \dots, h\} \setminus I} g_t(y_\ell - x_\ell) \cdot G_\vartheta(t) g_{\frac{t}{2}}(y_i - x_i), \quad \mathbf{x}, \mathbf{y} \in (\mathbb{R}^2)_I^h, \quad (9.31)$$

which is consistent with the definition of $U_N^I(n, \mathbf{x}, \mathbf{y})$ by (2.6) and (8.18). This kernel defines an integral operator from $L^p((\mathbb{R}^2)_I^h)$ to $L^p((\mathbb{R}^2)_I^h)$. Similar to $U_N(n, y_i - x_i)$ in (9.10), the factor $G_\vartheta(t) g_{\frac{t}{2}}(y_i - x_i)$ gives the weight of the wiggle line in Figure 6.

2. Constrained Evolution. For $I = \{i < j\}$ and $J = \{k < \ell\}$, the continuum analogue of $\mathbf{Q}_t^{I,J}$ in (9.11) is given by the kernel with density

$$\mathcal{Q}_t^{I,J}(\mathbf{y}, \mathbf{x}) := \prod_{a=1}^h g_t(x_a - y_a), \quad \mathbf{y} \in (\mathbb{R}^2)_J^h, \mathbf{x} \in (\mathbb{R}^2)_I^h. \quad (9.32)$$

This kernel defines an integral operator from $L^p((\mathbb{R}^2)_J^h)$ to $L^p((\mathbb{R}^2)_I^h)$ and corresponds to time strips in Figure 6 that contain no wiggle line. As in (9.12) and (9.13), we can define the boundary kernels $\mathcal{Q}_t^{I,*}(\mathbf{y}, \mathbf{x})$, $\mathcal{Q}_t^{*,J}(\mathbf{y}, \mathbf{x})$, $\mathcal{Q}_t^{*,J}(\varphi, \mathbf{x})$, and $\mathcal{Q}_t^{I,*}(\mathbf{y}, \psi)$.

We can now give the formula for higher integer moments of the critical 2d SHF, which was first derived in [GQT21]. The series representation for the 3rd moment was derived earlier in [CSZ19b].

Theorem 9.6. *Fix $h \in \mathbb{N}$ with $h \geq 3$. For $\varphi \in C_c(\mathbb{R}^2)$ and $\psi \in C_b(\mathbb{R}^2)$, the h -th moment of $\text{SHF}_{0,t}^\vartheta(\varphi, \psi)$ is finite and admits the expression*

$$\mathbb{E}[\text{SHF}_{0,t}^\vartheta(\varphi, \psi)^h] = \int_{(\mathbb{R}^2)^h} \cdots \int_{(\mathbb{R}^2)^h} \varphi^{\otimes h}(\mathbf{z}) \mathcal{K}_t^{(h)}(\mathbf{z}, \mathbf{w}) \psi^{\otimes h}(\mathbf{w}) d\mathbf{z} d\mathbf{w}, \quad (9.33)$$

where $\mathbf{x} = (x_1, \dots, x_h) \in (\mathbb{R}^2)^h$ and $\phi^{\otimes h}(\mathbf{x}) = \prod_{i=1}^h \phi(x_i)$, and

$$\begin{aligned} \mathcal{K}_t^{(h)}(\mathbf{z}, \mathbf{w}) := & \\ & 1 + \sum_{m=1}^{\infty} (4\pi)^m \sum_{\substack{I_1, \dots, I_m \subseteq \{1, \dots, h\} \\ |I_\ell|=2, I_\ell \neq I_{\ell+1}}} \int \cdots \int d\vec{a} d\vec{b} \int \cdots \int_{\substack{\mathbf{x}^{(\ell)}, \mathbf{y}^{(\ell)} \in (\mathbb{R}^2)_{I_\ell}^h \\ \text{for } \ell=1, \dots, m}} \prod_{\ell=1}^m d\mathbf{x}^{(\ell)} d\mathbf{y}^{(\ell)} \\ & \mathcal{Q}_{a_1}^{*,I_1}(\mathbf{z}, \mathbf{x}^{(1)}) G_{\vartheta, b_1 - a_1}^{I_1}(\mathbf{x}^{(1)}, \mathbf{y}^{(1)}) \left(\prod_{\ell=2}^m \mathcal{Q}_{a_\ell - b_{\ell-1}}^{I_{\ell-1}, I_\ell}(\mathbf{y}^{(\ell-1)}, \mathbf{x}^{(\ell)}) G_{\vartheta, b_\ell - a_\ell}^{I_\ell}(\mathbf{x}^{(\ell)}, \mathbf{y}^{(\ell)}) \right) \mathcal{Q}_{t-b_m}^{I_m, *}(\mathbf{y}^{(m)}, \mathbf{w}). \end{aligned} \quad (9.34)$$

The identity (9.34) can be proved by taking the limit $N \rightarrow \infty$ in (9.9), see [GQT21]. Similar bounds as in Theorem 9.1 also hold for $\text{SHF}_{0,t}^\vartheta(\varphi, \psi)$, which allows one to take more general φ and ψ in (9.33) as long as the terms in the r.h.s. of (9.4) are finite.

9.5. RELATED LITERATURE. The positive integer moments of the directed polymer partition function can be expressed in terms of the collision local times of independent random walks, which are also connected to the so-called Delta-Bose gas.

Exponential moments of collision local times. Recall the definition of the point-to-point partition function $Z_N^\beta(z, w) := Z_{0,N}^\beta(z, w)$ from (1.11). Let us consider mixed moments of the form $\mathbb{E}[Z_N^\beta(z_1, w_1) \cdots Z_N^\beta(z_h, w_h)]$. Assuming the disorder variables $\omega(n, z)$ are i.i.d.

standard normal, we can then compute directly

$$\begin{aligned} \mathbb{E}\left[\prod_{i=1}^h Z_N^\beta(z_i, w_i)\right] &= \mathbb{E}\mathbb{E}_{\mathbf{z}}^{\otimes h}\left[e^{\sum_{i=1}^h \sum_{n=1}^{N-1} (\beta\omega(n, S_n^{(i)}) - \lambda(\beta))} \prod_{i=1}^h \mathbb{1}_{\{S_N^{(i)} = w_i\}}\right] \\ &= \mathbb{E}_{\mathbf{z}}^{\otimes h}\left[\prod_{(n,x) \in \{1, \dots, N-1\} \times \mathbb{Z}^2} e^{\beta^2 \sum_{1 \leq i < j \leq h} \mathbb{1}_{\{S_n^{(i)} = S_n^{(j)} = x\}}} \prod_{i=1}^h \mathbb{1}_{\{S_N^{(i)} = w_i\}}\right] \\ &= \mathbb{E}_{\mathbf{z}}^{\otimes h}\left[e^{\beta^2 \sum_{1 \leq i < j \leq h} L_N^{\{i,j\}}} \prod_{i=1}^h \mathbb{1}_{\{S_N^{(i)} = w_i\}}\right], \end{aligned} \quad (9.35)$$

where $\mathbb{E}_{\mathbf{z}}^{\otimes h}$ denotes expectation w.r.t. h independent random walks $(S^{(i)})_{1 \leq i \leq h}$ starting from $\mathbf{z} = (z_1, \dots, z_h)$ respectively, and $L_N^{\{i,j\}} := \sum_{n=1}^{N-1} \mathbb{1}_{\{S_n^{(i)} = S_n^{(j)}\}}$ denotes the collision local time between $S^{(i)}$ and $S^{(j)}$.

When $z_i = 0$ for all $1 \leq i \leq h$, the classic Erdős-Taylor theorem [ET60] shows that each $L_N^{\{i,j\}}/R_N = L_N^{\{i,j\}}/\mathbb{E}[L_N^{\{i,j\}}]$ converges in distribution to an $\text{Exp}(1)$ random variable. When I is a set of pairs $\{i < j\}$ such that the graph with vertex set $\{1, \dots, h\}$ and edge set I contains no loops, it follows from [GS09] that $(L_N^{\{i,j\}}/R_N)_{\{i < j\} \subset I}$ converges jointly to a family of i.i.d. $\text{Exp}(1)$ random variables. Finally, [LZ24b] proved the joint convergence of the full family $(L_N^{\{i,j\}}/R_N)_{1 \leq i < j \leq h}$ to a family of i.i.d. $\text{Exp}(1)$ random variables. This implies that if we choose $\beta = \beta_N = (\hat{\beta}/R_N)^{1/2}$ for some $\hat{\beta} < 1$, i.e., the subcritical window defined in (1.14), then we can identify the limit of (9.35) as $N \rightarrow \infty$ in terms of the exponential moments of independent exponential random variables even if $z_i = z_j$ for some $i \neq j$.

However, when $\beta = \beta_N$ is critical with $\hat{\beta} = 1$, allowing $z_i = z_j$ for $i \neq j$ will lead to divergence in (9.35) as $N \rightarrow \infty$ because

$$\mathbb{E}[e^{\hat{\beta}(L_N^{\{i,j\}}/R_N)}] = \mathbb{E}[Z_N^{\beta_N}(0)^2] \sim C \log N,$$

where the asymptotics follows from Lemma 8.3. Nevertheless, as long as $(z_i/\sqrt{N})_{1 \leq i \leq h}$ converge to a vector of distinct points \mathbf{z}' as $N \rightarrow \infty$ and the same holds for $(w_i/\sqrt{N})_{1 \leq i \leq h} \rightarrow \mathbf{w}'$, we expect (9.35) to converge (after normalising by N^h) to the kernel $\mathcal{K}_1^{(h)}(\mathbf{z}', \mathbf{w}')$ in Theorem 9.6.

Delta-Bose gas and singular interacting diffusions. Formally, the continuum analogue of (9.35) is

$$\mathcal{S}_t^{(h)}(\mathbf{z}, \mathbf{w}) := \mathbb{E}_{\mathbf{z}}^{\otimes h}\left[e^{\beta^2 \sum_{1 \leq i < j \leq h} \int_0^t \delta(B_s^i - B_s^j) ds} \prod_{i=1}^h \delta(B_t^i - w_i)\right], \quad (9.36)$$

where $(B^i)_{1 \leq i \leq h}$ are i.i.d. standard Brownian motion on \mathbb{R}^2 , starting from $\mathbf{z} = (z_1, \dots, z_h)$ respectively, and $\delta(\cdot)$ is the delta function at the origin. This is the Feynman-Kac semigroup associated with the Schrödinger operator with delta potential on the diagonals

$$\frac{1}{2}\Delta + \beta^2 \sum_{1 \leq i < j \leq h} \delta(x_i - x_j), \quad (9.37)$$

known as the Delta-Bose gas. In other words, $\mathcal{S}_t^{(h)}(\mathbf{z}, \mathbf{w})$ solves (formally) the parabolic equation

$$\begin{aligned} \partial_t u(t, \mathbf{x}) &= \frac{1}{2} \Delta u(t, \mathbf{x}) + \beta^2 \sum_{1 \leq i < j \leq h} \delta(x_i - x_j) u(t, \mathbf{x}), \quad t > 0, \mathbf{x} \in (\mathbb{R}^2)^h, \\ u(0, \mathbf{x}) &= \prod_{i=1}^h \delta(x_i - w_i). \end{aligned} \tag{9.38}$$

Of course, the presence of the delta function makes (9.36)-(9.38) ill-defined, especially because Brownian motion in \mathbb{R}^2 does not hit points. The challenge is to make sense of the semigroup $\mathcal{S}_t^{(h)}(\mathbf{z}, \mathbf{w})$ and its associated operator. A natural approach is to replace the delta function by its mollified version $\delta_\varepsilon(x) = \frac{1}{\varepsilon^2} j(\frac{x}{\varepsilon})$ for some smooth probability density function j . This introduces a minimal spatial scale $\varepsilon > 0$ and has the same effect as discretising space by considering the discrete time-space kernel considered in (9.35). Similar to (1.15), the most interesting choice will be to choose

$$\beta^2 = \beta_\varepsilon^2 = \frac{2\pi}{\log \frac{1}{\varepsilon}} \left(1 + \frac{\vartheta}{|\log \varepsilon|} \right), \tag{9.39}$$

which lies in the critical window with $\hat{\beta} = 1$. This has been the choice considered in the literature, see e.g. [AFHeK⁺92, DFT94, ABD95, BC98, DR04, AGHeKH05]. In particular, [DFT94] defined self-adjoint extensions of the operator in (9.37) using an approach of quadratic forms and Γ -convergence, while [DR04] followed an approach using resolvents and L^2 Fourier analysis. The latter approach is closely related to the diagrammatic expansions used in [GQT21] and [CSZ23a], which is illustrated in Figure 6.

In light of the connection with the directed polymer and the stochastic heat equation, the semigroup $\mathcal{S}_t^{(h)}(\mathbf{z}, \mathbf{w})$ in (9.36) should in fact be a family of semigroups indexed by $\vartheta \in \mathbb{R}$, given by the kernels $\mathcal{K}_t^{(h)}(\mathbf{z}, \mathbf{w})$ in Theorem 9.6 for the h -th moment of the critical $2d$ SHF.

The formal semigroups $\mathcal{S}_t^{(h)}(\mathbf{z}, \mathbf{w})$ in (9.36) were defined rigorously in [Che24a] as the limit of approximate semigroups via mollification of the delta function and choosing β as in (9.39). This naturally leads to a family of $(\mathbb{R}^2)^h$ -valued time-inhomogeneous Markov processes $Y_t^{T,(h)} = (Y_{1,t}^{T,(h)}, \dots, Y_{h,t}^{T,(h)})$, such that given terminal time T , it has transition densities

$$g_{s,t}^{T,(h)}(\mathbf{x}, \mathbf{y}) := \mathcal{S}_{t-s}^{(h)}(\mathbf{x}, \mathbf{y}) \cdot \frac{\int \mathcal{S}_{T-t}^{(h)}(\mathbf{y}, \mathbf{z}) d\mathbf{z}}{\int \mathcal{S}_{T-s}^{(h)}(\mathbf{x}, \mathbf{z}) d\mathbf{z}}, \quad 0 \leq s < t \leq T, \mathbf{x}, \mathbf{y} \in (\mathbb{R}^2)^h.$$

When $h = 2$, [CM23] considered the difference $X_t^T := Y_{2,t}^{T,(2)} - Y_{1,t}^{T,(2)}$ between the two-components and showed that it solves a stochastic differential equation (SDE) with a singular drift toward the origin, which enables X_t^T to visit and accumulate a local time at the origin, defined as the limit (as $\varepsilon \downarrow 0$) of

$$L_t^\varepsilon := \frac{1}{2\varepsilon^2 (\log \frac{1}{\varepsilon})^2} |\{s \in [0, t] : |X_s^T| \leq \varepsilon\}|, \quad t \in [0, T].$$

This stands in contrast to a standard Brownian motion in \mathbb{R}^2 , which does not hit the origin.

The semigroups $\mathcal{S}_t^{(h)}(\mathbf{z}, \mathbf{w})$ were also constructed in [Che22, Che24b] via a stochastic representation in the spirit of (9.36), but in terms of a special diffusion process instead of

a Brownian motion. In particular, when $h = 2$, the analogue of the semigroup in (9.36) in terms of the relative motion $B_t^1 - B_t^2$ is given by

$$\begin{aligned} P_t^\vartheta f(z_0) &= \mathbb{E}^{\otimes h} \left[e^{\beta^2 \int_0^t \delta(B_s^1 - B_s^2) ds} f(B_t^1 - B_t^2) \middle| B_0^1 - B_0^2 = z_0 \right] \\ &:= \mathbb{E}_{z_0} \left[\frac{e^{\vartheta t} K_0(\sqrt{\varrho}|z_0|)}{K_0(\sqrt{\varrho}|Z_t|)} f(Z_t) \middle| Z_0 = z_0 \right], \end{aligned}$$

where K_0 is the Macdonald function or generalised Bessel function of second kind of order 0, the parameter ϱ is determined by ϑ in (9.39) and the mollifier j , and $(Z_t)_{t \geq 0}$ is an \mathbb{R}^2 -valued diffusion whose radial component $|Z_t|$ is a transformation of a Bessel process constructed by Donati-Martin and Yor [DMY06], with 0 being a point of instantaneous reflection for $|Z_t|$.

There have also been recent studies of delta-Bose gas in dimension $d \geq 3$ [Wan24].

10. FURTHER PROPERTIES OF CRITICAL 2d SHF

In this section, we list some further properties of the critical 2d SHF with parameter $\vartheta \in \mathbb{R}$, that we denote by $\text{SHF}(\vartheta) = (\text{SHF}_{s,t}^\vartheta(dx, dy))_{0 \leq s \leq t < \infty}$.

10.1. SHIFT INVARIANCE AND SCALING COVARIANCE. The critical 2d SHF satisfies the following shift invariance and scaling covariance properties.

Theorem 10.1 (Translation invariance and scaling covariance). *The critical 2d SHF is translation invariant in law:*

$$(\text{SHF}_{s+a,t+a}^\vartheta(d(x+b), d(y+b)))_{0 \leq s \leq t < \infty} \stackrel{\text{dist}}{=} (\text{SHF}_{s,t}^\vartheta(dx, dy))_{0 \leq s \leq t < \infty} \quad \forall a \geq 0, \quad \forall b \in \mathbb{R}^2,$$

and it satisfies the following scaling relation:

$$(\text{SHF}_{as,at}^\vartheta(d(\sqrt{a}x), d(\sqrt{a}y)))_{0 \leq s \leq t < \infty} \stackrel{\text{dist}}{=} (a \text{SHF}_{s,t}^{\vartheta+\log a}(dx, dy))_{0 \leq s \leq t < \infty} \quad \forall a > 0. \quad (10.1)$$

The shift invariance is inherited from the shift invariance of the directed polymer model and that of the underlying disorder ω . The scaling covariance relation can be checked by computing the covariance of the critical 2d SHF. It shows that zooming out in space-time diffusively has the effect of increasing the disorder strength ϑ .

10.2. FLOW PROPERTY. The heat equation induces the so-called heat flow in the sense that the solution is linear in the initial condition and can be written as a mixture of heat kernels that satisfy the Chapman-Kolmogorov equation $p_{s,u}(x, dy) = \int p_{s,t}(x, dz) p_{t,u}(z, dy)$. If the family of kernels $p_{s,t}(x, dy)_{s < t, x \in \mathbb{R}}$ is replaced by a random family $(p_{s,t}^\omega(x, dy))_{s < t, x \in \mathbb{R}}$ such that for all $s < t < u$ and $x \in \mathbb{R}$, almost surely p^ω satisfies the Chapman-Kolmogorov equation

$$p_{s,u}^\omega(x, dy) = \int p_{s,t}^\omega(x, dz) p_{t,u}^\omega(z, dy), \quad (10.2)$$

and p_{s_i, t_i}^ω are independent over disjoint time intervals $[s_i, t_i]$ and translation invariant in law under space-time shifts, then $(p_{s,t}^\omega(x, dy))_{s < t, x \in \mathbb{R}}$ is a *stochastic flow of kernels* introduced by Le Jan and Raimond [LJR04] and can be interpreted as the transition kernels of a random motion in an ‘i.i.d.’ space-time random environment (see e.g. [SSS14]). The critical 2d SHF $(\text{SHF}_{s,t}^\vartheta(dx, dy))_{s < t}$ satisfies a similar property, except that it is not meaningful to consider delta initial condition $\int \delta_x \text{SHF}_{s,t}^\vartheta(dx, dy)$ as in (10.2). Instead, we should consider $\text{SHF}_{s,t}^\vartheta(\mu, dy)$ for initial condition $\mu(dx)$ that is sufficiently regular.

In a forthcoming work [CSZ25a], we will show that the critical $2d$ SHF defines a linear measure-valued Markov process with a state space S , which consists of locally finite measures μ , such that for any $R > 0$,

$$\iint_{|x|,|y|\leq R} \log \frac{1}{|x-y|} \mu(dx) \mu(dy) < \infty,$$

and $\int e^{-|x|^2/2t} \mu(dx) < \infty$ for all $t > 0$. This is a natural state space for the critical $2d$ SHF, since

$$\text{Var}(\text{SHF}_{0,t}^\vartheta(\phi)) = \iint_{(\mathbb{R}^2)^2} \phi(x) \phi(x') \tilde{K}_t^\vartheta(x, x') dx dx',$$

where $\tilde{K}_t^\vartheta(x, x')$ is as in (8.24) and $\tilde{K}_t^\vartheta(x, x') \sim c \log \frac{1}{|x-x'|}$ as $|x - x'| \rightarrow 0$.

We will show in [CSZ25a] that the critical $2d$ SHF satisfies the following almost sure Chapman-Kolmogorov (flow) property: for all $s < t < u$ and initial measure $\mu \in S$, almost surely, $\text{SHF}_{s,t}^\vartheta(\mu, dz) \in S$ and

$$\text{SHF}_{s,u}^\vartheta(\mu, dy) = \int \text{SHF}_{s,t}^\vartheta(\mu, dz) \text{SHF}_{t,u}^\vartheta(z, dy) = \text{SHF}_{t,u}^\vartheta\left(\text{SHF}_{s,t}^\vartheta(\mu, dz), dy\right). \quad (10.3)$$

Clark and Mian [CM24] have given a different formulation of the flow property, which starts by defining a notion of ε -regularised convolution for measures. More precisely, let X, X' be Polish spaces and let μ_1, μ_2 be Borel measures on $X \times \mathbb{R}^2$ and $\mathbb{R}^2 \times X'$, respectively. Then the ε -regularised convolution of μ_1 and μ_2 is defined as a measure on $X \times \mathbb{R}^2 \times X'$ with

$$\mu_1 \circ_\varepsilon \mu_2(ds, dx, ds') := \iint_{x' \in \mathbb{R}^2} \mu_1(ds, dx) g_\varepsilon(x - x') \mu_2(dx', ds'),$$

where g_ε is the heat kernel at time ε . The following result was established in [CM24].

Theorem 10.2. *Let $\text{SHF}_{s,r}^\vartheta(dx, dy)$ and $\text{SHF}_{r,t}^\vartheta(dy, dz)$ be marginals of the critical $2d$ SHF at times $s < r < t$. Then there exists a random measure $\text{SHF}_{s,r,t}^\vartheta(dx, dy, dz)$ on $(\mathbb{R}^2)^3$ such that*

$$\text{SHF}_{s,r}^\vartheta(dx, dy) \circ_\varepsilon \text{SHF}_{r,t}^\vartheta(dy, dz) \xrightarrow[\varepsilon \downarrow 0]{L^2(\mathbb{P})} \text{SHF}_{s,r,t}^\vartheta(dx, dy, dz), \quad (10.4)$$

where the convergence is in $L^2(\mathbb{P})$. Moreover,

$$\text{SHF}_{s,t}^\vartheta(dx, dz) = \int_{y \in \mathbb{R}^2} \text{SHF}_{s,r,t}^\vartheta(dx, dy, dz).$$

This result also establishes an almost sure Chapman-Komogorov type of relation for the critical $2d$ SHF. The procedure can be iterated to show that for $0 < t_1 < \dots < t_n$, the critical $2d$ SHF satisfies

$$\text{SHF}_{t_1,t_2}^\vartheta(dx_1, dx_2) \circ_\varepsilon \dots \circ_\varepsilon \text{SHF}_{t_{n-1},t_n}^\vartheta(dx_{n-1}, dx_n) \xrightarrow[\varepsilon \downarrow 0]{L^2(\mathbb{P})} \text{SHF}_{t_1,\dots,t_n}^\vartheta(dx_1, \dots, dx_n).$$

10.3. CHARACTERISATION OF THE CRITICAL 2d SHF. Recently, Tsai [Tsa24] gave an axiomatic characterisation of the critical 2d SHF, which says that, a process $(Z_{s,t}(\cdot, \cdot))_{s \leq t}$ taking values in $M_+(\mathbb{R}^2 \times \mathbb{R}^2)$, the space of locally finite non-negative measures on $\mathbb{R}^2 \times \mathbb{R}^2$ equipped with the vague topology, must have the same law as the critical 2d SHF with parameter ϑ , denoted by $\text{SHF}(\vartheta)$, if $Z_{s,t}$ is continuous in s and t , satisfies the flow property formulated in (10.4), has independent increments, and has matching first four moments with $(\text{SHF}_{s,t}^\vartheta(\cdot, \cdot))_{s \leq t}$.

Theorem 10.3 (Characterisation of the critical 2d SHF). *Let $Z = (Z_{s,t}(\cdot, \cdot))_{s \leq t}$ be a stochastic process taking values in $M_+(\mathbb{R}^2 \times \mathbb{R}^2)$ such that:*

- (1) *for any $s < t < u$, the random measures $Z_{s,t}$, $Z_{t,u}$ and $Z_{s,u}$ satisfy the conclusions of Theorem 10.2;[†]*
- (2) *for any $s < t < u$, $Z_{s,t}$ and $Z_{t,u}$ are independent;*
- (3) *for any $s < t < u$, $1 \leq n \leq 4$, and $\phi_i, \psi_i \in L^2(\mathbb{R}^2)$ for $i = 1, \dots, 4$, the mixed moments $\mathbb{E}\left[\prod_{i=1}^n Z_{s,t}(\phi_i, \psi_i)\right]$ agree with that of $\text{SHF}(\vartheta)$, cf. (9.33), for some $\vartheta \in \mathbb{R}$.*

Then Z has the same law as the critical 2d stochastic heat flow $\text{SHF}(\vartheta)$. Furthermore, $(Z_{s,t}(\cdot, \cdot))_{s \leq t}$ admits a version that is almost surely continuous in $s \leq t$.

The characterization in Theorem 10.3 was proved in [Tsa24] via a Lindeberg principle, while the continuity was established by verifying the Kolmogorov moment criterion. If Z and \tilde{Z} are two processes satisfying the assumptions in Theorem 10.3 for the same ϑ , then by the Chapman-Kolmogorov property, $Z_{0,1}$ is a functional of $(Z_{(i-1)/n, i/n})_{1 \leq i \leq n}$, and the same holds for \tilde{Z} . The basic idea is that, in the reconstruction of $Z_{0,1}$ from $(Z_{(i-1)/n, i/n})_{1 \leq i \leq n}$, one can successively replace $Z_{(i-1)/n, i/n}$ by $\tilde{Z}_{(i-1)/n, i/n}$ and control the error of each replacement. If the cumulative error can be shown to tend to 0 as $n \rightarrow \infty$, then $Z_{0,1}$ and $\tilde{Z}_{0,1}$ must be equal in law, where we can take \tilde{Z} to be $\text{SHF}(\vartheta)$. Similar to the Lindeberg principle explained in Section 6.2, controlling the error of each replacement requires Taylor expansions and matching first two moments for $Z_{(i-1)/n, i/n}$ and $\tilde{Z}_{(i-1)/n, i/n}$, plus suitable bounds on higher moments (4-th moment should suffice). Of course $Z_{(i-1)/n, i/n}$ are much more complicated objects than \mathbb{R} -valued random variables, which requires more delicate analysis. We refer to [Tsa24] for further details.

Theorem 10.3 makes it much easier to prove convergence to $\text{SHF}(\vartheta)$ since one only needs to verify that every subsequential limit satisfies the axioms in Theorem 10.3. This was carried out in [Tsa24] for the solution u^ε of the 2d mollified SHE in the critical window.

We also mention that there have been some recent progress in formulating martingale problems for the SHF [Nak25, Che25], which might eventually lead to a well-posed martingale problem characterisation for the SHF.

10.4. NON-GMCNESS OF FIXED TIME MARGINALS. Let $(X(x))_{x \in \mathbb{R}^d}$ be a centered (generalised) gaussian field with correlation function $k(x, y)$ for $x, y \in \mathbb{R}^d$. The Gaussian Multiplicative Chaos (GMC) is a random measure on \mathbb{R}^d , which formally has the following

[†]In [Tsa24], the condition is formulated such that the Gaussian kernels g_ε in Theorem 10.2 can be replaced by any family of mollifiers that ensure L^2 convergence.

“density” with respect to a reference measure $\sigma(dx)$:

$$\mathcal{M}_\gamma(dx) := e^{\gamma X(x) - \frac{\gamma^2}{2} \mathbb{E}[X^2]} \sigma(dx). \quad (10.5)$$

Various assumptions can be imposed on the measure $\sigma(dx)$ (see [RV14]), but for simplicity, we will just assume that $\sigma(dx)$ is the Lebesgue measure. For any test function ϕ , we denote

$$\mathcal{M}_\gamma(\phi) := \int_{\mathbb{R}^d} \phi(x) \mathcal{M}_\gamma(dx).$$

The most interesting case is when the field X is log-correlated in the sense that $k(x, y) \sim c \log \frac{1}{|x-y|}$ as $|x-y| \rightarrow 0$. Since $k(x, x) = \infty$, X is a generalised Gaussian field and $\mathcal{M}_\gamma(dx)$ is a priori undefined. To rigorously define $\mathcal{M}_\gamma(dx)$, the standard procedure is to first replace X by a regularised field X_ε (e.g., via mollification) and then show that

$$\mathcal{M}_{\varepsilon, \gamma}(dx) := e^{\gamma X_\varepsilon(x) - \frac{\gamma^2}{2} \mathbb{E}[X_\varepsilon^2]} \sigma(dx), \quad (10.6)$$

has a limit as $\varepsilon \rightarrow 0$. Log-correlated Gaussian field X is of special interest because there is a phase transition in the parameter γ in (10.6) that is similar to the phase transition in the polymer partition function in Theorem 3.1. In particular, when $\sigma(dx)$ is the Lebesgue measure, $\mathcal{M}_{\varepsilon, \gamma}(dx)$ has a non-trivial limit if $\gamma < \gamma_c := \sqrt{2d}$, and $\mathcal{M}_{\varepsilon, \gamma}(dx)$ converges to the zero measure if $\gamma \geq \gamma_c$. We refer to [RV14] for more details.

It is natural to wonder whether the one-time marginal of $\text{SHF}(\vartheta)$ can be identified with a GMC measure on \mathbb{R}^2 . The answer turns out to be no, which suggests heuristically that the logarithm of $\text{SHF}(\vartheta)$ (if defined rigorously, it will give a meaning to the solution of the critical 2d KPZ equation with parameter ϑ), would be a Gaussian field.

Since GMC is defined from a Gaussian field, its law is entirely determined by its first two moments. If we denote by \mathcal{M}_t^ϑ the candidate GMC with the same first two moments as the critical 2d SHF SHF_t^ϑ at time t , then we can rule out the possibility of SHF_t^ϑ being a GMC by showing that its higher moments do not match that of \mathcal{M}_t^ϑ . This is what we proved in [CSZ23b].

Theorem 10.4 (The critical 2d SHF is not a GMC). *Let g_δ be the heat kernel at time $\delta > 0$. Let $\text{SHF}_t^\vartheta(\cdot)$ be the critical 2d SHF started from Lebesgue measure and evaluated at time t , and let $\mathcal{M}_t^\vartheta(\cdot)$ be the GMC measure on \mathbb{R}^2 with the same first two moments as $\text{SHF}_t^\vartheta(\cdot)$. For any $\vartheta \in \mathbb{R}$ and $t > 0$, there exists $\eta = \eta_{t, \vartheta} > 0$ such that for all $h \in \mathbb{N}$ with $h \geq 3$, we have*

$$\liminf_{\delta \downarrow 0} \frac{\mathbb{E}[\text{SHF}_t^\vartheta(g_\delta)^h]}{\mathbb{E}[\mathcal{M}_t^\vartheta(g_\delta)^h]} \geq 1 + \eta > 1. \quad (10.7)$$

To see heuristically why Theorem 9.1 might hold, we note that Wick’s theorem for the moments of gaussian variables implies that

$$\mathbb{E}[\mathcal{M}^\vartheta(g_\delta)^h] = \int_{(\mathbb{R}^2)^h} g_\delta(z_1) \cdots g_\delta(z_h) e^{\sum_{1 \leq i < j \leq h} k(z_i, z_j)} dz_1 \cdots dz_h \sim \mathbb{E}[\mathcal{M}^\vartheta(g_\delta)^2]^{\binom{h}{2}},$$

where each term $k(z_i, z_j)$ represents a two-body interaction among the h particles at locations z_1, \dots, z_h , and the last asymptotic uses the fact that

$$\mathbb{E}[\mathcal{M}^\vartheta(dx) \mathcal{M}^\vartheta(dy)] = \mathbb{E}[\text{SHF}_t^\vartheta(dx) \text{SHF}_t^\vartheta(dy)] =: K_t^\vartheta(x, y) dx dy$$

with $K_t^\vartheta(x, y) \sim c \log \frac{1}{|x-y|}$ as $|x - y| \rightarrow 0$. In other words, $\mathbb{E}[\mathcal{M}^\vartheta(g_\delta)^h]$ has the same asymptotics as if there are only independent two-body interactions. This is no longer the case for $\mathbb{E}[\text{SHF}_t^\vartheta(\varphi)^h]$, where the interactions among the h particles are illustrated in Figure 6. At the heart of the proof of Theorem 10.7 is the use of the Gaussian Correlation Inequality [Roy14, LaM17], which together with additional analysis on the collision local times of Brownian motions, shows that there are positive correlations among the two-body interactions.

10.5. SINGULARITY AND REGULARITY. In the forthcoming work [CSZ25b], we show that the marginal distribution of the critical 2d SHF at each time $t > 0$ is a random measure that is almost surely singular respect to Lebegues measure (see Figure 2).

Theorem 10.5 (Singularity of SHF). *Fix any $t > 0$ and $\vartheta \in \mathbb{R}$. Almost surely, $\text{SHF}_t^\vartheta(dx) := \int_{y \in \mathbb{R}^2} \text{SHF}_{0,t}^\vartheta(dy, dx)$ is singular with respect to the Lebesgue measure on \mathbb{R}^2 .*

However, $\text{SHF}_t^\vartheta(dx)$ barely fails to be a function.

Theorem 10.6 (Regularity of the SHF). *Fix any $t > 0$ and $\vartheta \in \mathbb{R}$. Almost surely, the critical 2d SHF $\text{SHF}_t^\vartheta(dx)$ belongs to $\mathcal{C}^{0-} := \bigcap_{\varepsilon > 0} \mathcal{C}^{-\varepsilon}$, where $\mathcal{C}^{-\varepsilon}$ is the negative Hölder space of order $-\varepsilon$.*

A consequence of Theorem 10.6 is that, almost surely, $\text{SHF}_t^\vartheta(dx)$ contains no atoms, since delta measures on \mathbb{R}^d belong to C^{-d} . Theorem 10.6 follows from moment estimates, see (9.5), which imply that for any $h \in \mathbb{N}$ and $\varepsilon > 0$, there exists a constant C , depending on $h, \varepsilon, t, \vartheta$, such that for all small δ we have

$$\mathbb{E}[\text{SHF}_t^\vartheta(\mathcal{U}_{B(x,\delta)})^h]^{1/h} \leq C \delta^{-\varepsilon} \quad \forall x \in \mathbb{R}^2. \quad (10.8)$$

Theorem 10.6 then follows from this moment bound and a general tightness criterion for negative Hölder spaces [FM17, Theorem 2.30].

Theorem 10.5 follows by applying the Lebesgue differentiation theorem and showing that, almost surely,

$$\lim_{\delta \downarrow 0} \text{SHF}_t^\vartheta(\mathcal{U}_{B(x,\delta)}) = 0 \quad \text{for Lebesgue a.e. } x \in \mathbb{R}^2. \quad (10.9)$$

where $\mathcal{U}_{B(x,\delta)}(\cdot) := \frac{1}{\pi \delta^2} \mathbb{1}_{B(x,\delta)}(\cdot)$ for the Euclidean ball $B(x, \delta) := \{y \in \mathbb{R}^2 : |y - x| < \delta\}$. This is accomplished by bounding the fractional moments of $\text{SHF}_t^\vartheta(\mathcal{U}_{B(x,\delta)})$, where we use the monotonicity of the fractional moment in ϑ to decrease ϑ and send $\vartheta = \vartheta(\delta) \rightarrow -\infty$ at a suitable rate as $\delta \downarrow 0$. More precisely, $\vartheta(\delta)$ will be chosen such that $\text{SHF}_t^\vartheta(\mathcal{U}_{B(x,\delta)})$ converges to a log-normal limit, similar to the point-to-plane partition function in the subcritical regime in Theorem 3.1.

Theorem 10.7. *Let $\vartheta(\delta) = 2 \log \delta$. For any $t > 0$ and $x \in \mathbb{R}^2$, the following convergence in distribution holds:*

$$\forall \varrho \in (0, \infty) : \quad \text{SHF}_t^\vartheta(\mathcal{U}_{B(x,\delta^\varrho)}) \xrightarrow[\delta \downarrow 0]{d} e^{\mathcal{N}(0, \sigma^2) - \frac{1}{2}\sigma^2} \quad \text{with } \sigma^2 = \log(1 + \varrho). \quad (10.10)$$

Remark 10.8. *The meta-theorem suggested by Theorem 10.7 is that, averaging the polymer partition functions (or averaging the critical 2d SHF) has the effect of reducing the disorder strength. If the averaging is on a suitably chosen spatial scale such that the mean and*

variance remain of a constant order, then the averaged partition function behaves like the point-to-plane partition function of a 2d polymer in the subcritical regime with a log-normal limit. Theorem 3.1 is one special case. We refer to [CSZ25b] for details.

11. DISCUSSIONS AND OPEN QUESTIONS

11.1. DISORDERED SYSTEMS AND SINGULAR SPDEs. Our motivation for studying the 2d SHE came from the study of continuum limits of disordered systems, with the DPM being one particular example.

Recalling its definition from Section 1.3, we can regard the DPM as a disorder perturbation of the underlying random walk S . In dimensions $d \geq 3$, this perturbation is called *irrelevant* in the language of renormalisation group theory because the critical disorder strength (inverse temperature) $\beta_c(d) > 0$, and it has been shown that [CY06] for $\beta < \beta_c$, the polymer measure converges to the same limiting Wiener measure as the underlying random walk S . On the other hand, in $d = 1$ and 2 , $\beta_c(d) = 0$ and hence at any $\beta > 0$, the polymer measure experiences path localisation [CSY03] and the path is expected to be super-diffusive. Therefore the disorder perturbation is *relevant* in $d = 1$ and 2 . Dimension $d = 2$ turns out to be critical, and only finer details of the model determine whether disorder perturbation is relevant (called *marginal relevance*) or irrelevant (called *marginal irrelevance*). In general, a disorder perturbation of an underlying pure model (without disorder) is called *relevant* if any fixed disorder strength $\beta > 0$, no matter how small, changes the large scale behaviour of the model (in particular, its scaling limit), and it is called *irrelevant* if small $\beta > 0$ does not change the large scale behaviour and the scaling limit. We refer to [Gia11] for more discussions on disordered systems and the *Harris criterion* [Har74] on when disorder perturbation is predicted by physicists to be relevant/irrelevant.

The connection between disordered systems and singular SPDEs is that, we can regard such SPDEs as a disorder perturbation of the deterministic PDE without the noise term. As we saw in (1.2), in dimensions $d < 2$, the effective strength of the noise tends to zero as we zoom into smaller and smaller space-time scales, which makes the disorder perturbation of the heat equation an irrelevant perturbation, while in dimensions $d > 2$, the disorder perturbation is a relevant perturbation, and $d = 2$ is marginal. If we consider instead large scale behaviour by sending $\varepsilon \uparrow \infty$ in (1.2), then the disorder perturbation is relevant in $d < 2$, irrelevant in $d > 2$, and marginal in $d = 2$, which corresponds exactly to the DPM. Therefore the notion of disorder relevance (resp. irrelevance) for disordered systems, which is concerned with large scale behaviour, corresponds to the notion of the singular SPDE being subcritical (resp. supercritical), which is concerned with small scale behaviour. Marginality for disordered systems corresponds to criticality for singular SPDEs.

If a disorder perturbation of a pure model is relevant, then heuristically, the effective strength of disorder will diverge as we zoom out in space-time (equivalently, take the continuum limit by sending the lattice spacing to 0). It should then be possible to send the disorder strength to 0 at a suitable rate as the lattice spacing tends to 0, such that we obtain a continuum limit that has non-trivial disorder dependence. For disordered systems that are disorder relevant (not marginal), this was first carried for the one-dimensional directed polymer in [AKQ14b], which led to the continuum directed polymer model [AKQ14a]. Subsequently, it was shown in [CSZ17a] that similar results should hold for more general disorder relevant binary-valued spin systems, including long range DPM [CSZ17a], the disordered pinning model [CSZ16], and random field perturbation of the critical 2d Ising

model [CSZ17a, BS22]. See [CSZ17a] for a more detailed discussion, and [LMS24] for an extension to non-binary valued spin systems.

It is our attempt to investigate the disordered continuum limit of the DPM in the critical dimension $d = 2$ that led to the results discussed in this article. Following a similar approach, the scaling limit of disordered pinning model with critical tail exponent $\alpha = \frac{1}{2}$ was recently obtained in [WY24], which is also connected to singular stochastic Volterra equations.

Note that disordered systems in general, such as the random field Ising model, have no time dimension in the noise and hence do not correspond to any singular SPDE. In light of our results for the $2d$ DPM, a natural question is:

Q. Are there marginally relevant disordered systems without a time dimension, for which we can obtain a non-trivial disordered continuum limit?

11.2. MORE OPEN QUESTIONS. We close by presenting some open questions and possible directions for future research.

1. **Universality.** It would be interesting to show that the critical $2d$ SHF arises as the universal scaling limit of models beyond the directed polymer model and the mollified SHE. Potential candidates include the non-linear SHE considered in [DG22, Tao24a, DG23, DG24], although even identifying the correct critical point appears challenging. It will also be interesting to explore connections to other statistical mechanics models.
2. **Structure of the critical $2d$ SHF at a given time.** We have established in [CSZ25b] that the critical $2d$ SHF at any deterministic time is almost surely singular with respect to the Lebesgue measure. This raises the question: is the measure supported on a set of fractal Hausdorff dimension? It is plausible that this dimension is marginally below 2. Another unresolved question is whether the critical $2d$ SHF has infinite speed of propagation, that is, starting with an initial measure with compact support, at any later time, the SHF will almost surely assign positive mass to any open ball.
3. **Structure of the maxima.** The critical $2d$ SHF is a non-negative measure-valued process, which is a log-correlated field at each time. Over the past decade there has been significant interests in studying log-correlated fields or their exponential (Gaussian Multiplicative Chaos), and in particular the structure of their maxima. Some examples are [RV14, BDZ16, DRSV17, CGRV19, Bis20, CFLW21], though the list is much longer. A natural question is whether a similarly detailed picture can be achieved for the peaks of the critical $2d$ SHF. **The study of maxima is already very interesting in the sub-critical regime, which is more closely connected to Gaussian Multiplicative Chaos (see [CNZ25] for some recent progress).**
4. **Relation to GMC.** We have shown in [CSZ23b] that at any given time, the critical $2d$ SHF cannot be realised as the (Wick) exponential of a generalised Gaussian field, and hence is not a Gaussian Multiplicative Chaos (GMC). However, it remains open whether the critical $2d$ SHF could be the exponential of a perturbation of a Gaussian field. Alternatively, one might ask whether the law of the critical $2d$ SHF at each time is absolutely continuous with respect to the law of a GMC.

Given the form of the exponential weight in (1.4), it is natural to regard it as the exponential of a Gaussian field indexed by paths in the Wiener space, which could lead to a GMC on path space. This was carried out in [BM22, BLM22] in dimension $d \geq 3$ in the subcritical regime. In dimension 2, Clark and Mian [CM24] constructed a continuum polymer measure that is the path space extension of the critical $2d$ SHF.

Although this polymer measure cannot be realised as a GMC with respect to the underlying Wiener measure, an interesting open question raised in [CM24] is whether the continuum polymer measure satisfies the ‘*conditional GMC*’ structure similar to what is known on the continuum diamond hierarchical lattice [Cla23]. More precisely, this means that for any $\vartheta > \vartheta'$, the continuum polymer measure with parameter ϑ can be realised as a GMC with the reference measure being the continuum polymer measure with parameter ϑ' .

5. **On the moments.** It follows from the results in [GQT21] that all moments of $\text{SHF}_{s,t}^\vartheta(\phi, \psi)$, with $\phi, \psi \in L^2(\mathbb{R}^2)$, are finite[†]. The upper bound on the growth of the n -th moment is $\exp(e^{cn^2})$ (see [GQT21, (8.26)]), while existing lower bounds are of the order $\exp(cn^2)$ [CSZ23b]. However, predictions in the physics literature [Raj99] suggest that the growth should be $\exp(e^n)$. Although the rapid growth of moments does not allow one to uniquely determine the distribution, it would be valuable to identify the correct order of growth and understand the underlying mechanism. Such asymptotics could shed light on the tail statistics of $\text{SHF}_{s,t}^\vartheta(\phi, \psi)$ (for suitable ϕ, ψ). On the other hand, it will also be interesting to investigate whether $\text{SHF}_{s,t}^\vartheta(\phi, \psi)$ possesses negative moments. Recently, moment asymptotics have also been investigated in other regimes, such as in the quasi-critical window [CN25] or over shrinking balls [LZ24a]. Extending and sharpening these results will also be interesting.
6. **Interpolating nature and going beyond the critical window.** The critical $2d$ SHF sits precisely at the boundary between the weak and strong disorder phases ($\hat{\beta} < 1$ vs $\hat{\beta} > 1$ in (1.14)). The parameter ϑ in $\text{SHF}_{s,t}^\vartheta$ provides an interpolation between the weak disorder phase (as $\vartheta \rightarrow -\infty$) and the strong disorder phase (as $\vartheta \rightarrow +\infty$). An interesting question is whether we can obtain information on the strong disorder phase by taking $\vartheta \rightarrow +\infty$. A first step is to determine whether the random measure $\text{SHF}_{s,t}^\vartheta$ converges locally to 0 as $\vartheta \uparrow \infty$. Any progress in understanding potential scaling limits in the strong disorder phase $\hat{\beta} > 1$, or the very strong disorder phase $\beta_N \equiv \beta > 0$ in (1.14), would be very interesting.
7. **Construct the Critical $2d$ Polymer Measure.** The critical $2d$ SHF arises as the scaling limit of the directed polymer partition functions on the intermediate disorder scale. From a statistical physics point of view, it will be very interesting to construct a corresponding continuum polymer model and investigate its path properties and phase transition. The continuum polymer measure constructed in [CM24] is an infinite measure whose law is translation invariant in space. An open question is to construct a continuum polymer probability measure started at a single point, similar to its discrete analogue defined in (1.9). The main difficulty is that, the point-to-plane partition function, which serves as a normalising constant in the discrete polymer measure in (1.9), converges to 0 in the critical window as shown in Theorem 3.1.
8. **Black noise.** The critical $2d$ SHF determines a family of σ -fields $\mathcal{F}_{s,t} := \sigma(\text{SHF}_{u,v}^\vartheta : s \leq u < v \leq t)$ on the underlying probability space. The almost sure Chapman-Kolmogorov property established by Clark and Mian [CM24] and stated in Theorem 10.2 already implies that this family of σ -fields is a *noise* in the sense of Tsirelson, which is a

[†]The extension to $L^p(\mathbb{R}^2)$ spaces with $p \in (1, \infty)$ was achieved in [CSZ23a].

continuum generalisation of the notion of a family of i.i.d. random variables.[†] It would be very interesting to show that the noise generated by the critical 2d SHF is in fact a (one-dimensional) *black noise*[‡] (we refrain from giving here a precise definition of black noise and refer to the survey by Tsirelson [Tsi04a]). Examples of one-dimensional black noise include the Brownian web [Tsi04b] and the directed landscape [HP24], and examples of two-dimensional black noise, which is much more challenging to prove, include the continuum limit of critical planar percolation [SS11] and the Brownian web [EF16].

9. **Renormalisation of the 1-point statistics.** Theorem 3.1 establishes that, in the critical window of the intermediate disorder scale, the point-to-plane partition function $Z_N^{\beta_N}(0)$ of the directed polymer converges in distribution to 0. This aligns with the fact that the critical 2d SHF is almost surely singular with respect to the Lebesgue measure. A natural question is to determine the rate at which $\log Z_N^{\beta_N}(0)$ (i.e., the one-point statistic of the discretized 2d KPZ equation in the critical window) converges to $-\infty$. For critical GMC, which is connected to the directed polymer on trees, such a rate of convergence to 0 has been identified [RV14].

Based on extrapolation from the subcritical regime, we conjecture that in the critical window,

$$\mathbb{E}[\log Z_N^{\beta_N}(0)] \sim -\frac{1}{2} \log \log N,$$

with a variance of the same order $\log \log N$. It will be very interesting to identify the limiting law of $\log Z_N^{\beta_N}(0)$ after proper centering and scaling, which would provide a partial interpolation between the Gaussian distribution in the weak disorder phase $\hat{\beta} < 1$ and the 2d analogue of the Tracy-Widom distribution in the very strong disorder phase $\beta_n \equiv \beta > 0$ in (1.14).

10. **Construct the Critical 2d KPZ.** As shown in Theorem 4.1, in the subcritical regime $\hat{\beta} < 1$, the 2d KPZ has Gaussian fluctuations and solves the 2d Edwards-Wilkinson equation. It remains open to define the solution of the 2d KPZ in the critical window since we cannot perform the Cole-Hopf transformation on the critical 2d SHF, which is a random measure and not a function. The challenge is to study the field of log partition functions $\log Z_N^{\beta_N}(z)$, not only its mean and variance at each $z \in \mathbb{Z}^2$, but also its covariance at different $z \in \mathbb{Z}^2$. It would also be interesting to compare with physics works [FT94] where a dynamic renormalisation approach suggests that a non-Gaussian limit (referred to as ‘fixed point’ in the physics language) exists if the parameter $\vartheta = \vartheta_N$ in the critical window (1.16) is chosen to diverge at the rate of $\log \log N$.
11. **Step into critical singular SPDEs.** The theory of singular SPDEs has undergone revolutionary developments thanks to the frameworks of regularity structures [Hai14], paracontrolled calculus [GIP15], renormalisation group approach [Kup16, Duc22], energy solutions [GJ14], etc. However, these theories are restricted to subcritical singular SPDEs. For SHE and KPZ, this means dimension $d < d_c = 2$. The critical 2d SHF provides a rare example of a model in the critical dimension and at the critical

[†]The family $(\mathcal{F}_{s,t})_{s < t}$ being a noise means that, for $s < t < u$, the σ -fields $\mathcal{F}_{s,t}$ and $\mathcal{F}_{t,u}$ are independent and $\mathcal{F}_{s,u}$ is generated by $\mathcal{F}_{s,t}$ and $\mathcal{F}_{t,u}$; moreover, there is a group $(\vartheta_h)_{h \in \mathbb{R}}$ of measure preserving maps (time shifts) such that $\vartheta_h(A) \in \mathcal{F}_{s+h,t+h}$ for $A \in \mathcal{F}_{s,t}$.

[‡]After the completion of this article, Gu and Tsai [GT25] proved that the SHF is indeed a one-dimensional black noise.

point, which has a non-Gaussian scaling limit. A natural question is whether the critical $2d$ SHF can help shed some light on other critical singular SPDEs, such as dynamics for the ϕ^4 model [MW17, AK20, BG20, GH21] and the Yang-Mills model [CCHS22, CCHS24, CS23] at the critical dimension 4. Recently, there have much progress in the analysis of critical and supercritical singular SPDEs that admit an explicit Gaussian stationary measure, see the lecture notes [CT24] and the references therein. However, these models do not exhibit a phase transition as we see in the $2d$ SHE. It will be interesting to find other critical or supercritical singular SPDEs that exhibit a phase transition.

ACKNOWLEDGEMENTS

The material of these notes was motivated by and presented in mini-courses in the Summer School “*Statistical Mechanics and Stochastic PDEs*” at Cetraro, Italy, under Fondazione CIME, in the “*Spring School on Critical Singular Stochastic PDEs*” at the Beijing International Center for Mathematical Research, and in the “*Intensive Lecture Series*” at Seoul National University. We thank the organisers of these meetings – CIME foundation, Weijun Xu, Insuk Seo and Sung-Soo Byun – for providing the opportunity to present these lectures. We are also grateful to Stathis Floros for his kind permission to use his picture of Meteora.

F.C. is supported by INdAM/GNAMPA. R.S. is supported by NUS grant A-8001448-00-00 and NSFC grant 12271475.

REFERENCES

- [ABD95] S. Albeverio, Z. Brzeźniak, and L. Dąbrowski. Fundamental solution of the heat and Schrödinger equations with point interaction. *J. Funct. Anal.*, 130(1):220–254, 1995.
- [ADC21] M. Aizenman and H. Duminil-Copin. Marginal triviality of the scaling limits of critical 4D Ising and ϕ_4^4 models. *Ann. of Math. (2)*, 194(1):163–235, 2021.
- [AFHeK⁺92] S. Albeverio, J. E. Fenstad, R. Høegh Krohn, W. Karwowski, and T. Lindstrøm. Schrödinger operators with potentials supported by null sets. In *Ideas and methods in quantum and statistical physics (Oslo, 1988)*, pages 63–95. Cambridge Univ. Press, Cambridge, 1992.
- [AGHeKH05] S. Albeverio, F. Gesztesy, R. Høegh Krohn, and H. Holden. *Solvable models in quantum mechanics*. AMS Chelsea Publishing, Providence, RI, second edition, 2005. With an appendix by Pavel Exner.
- [AK20] S. Albeverio and S. Kusuoka. The invariant measure and the flow associated to the Φ_3^4 -quantum field model. *Ann. Sc. Norm. Super. Pisa Cl. Sci. (5)*, 20(4):1359–1427, 2020.
- [AKQ14a] T. Alberts, K. Khanin, and J. Quastel. The continuum directed random polymer. *J. Stat. Phys.*, 154(1-2):305–326, 2014.
- [AKQ14b] T. Alberts, K. Khanin, and J. Quastel. The intermediate disorder regime for directed polymers in dimension $1 + 1$. *Ann. Probab.*, 42(3):1212–1256, 2014.
- [Ape10] A. Apelblat. *Integral transforms and Volterra functions*. Nova Science Pub Inc., 2010.
- [BC95] L. Bertini and N. Cancrini. The stochastic heat equation: Feynman-Kac formula and intermittence. *J. Statist. Phys.*, 78(5-6):1377–1401, 1995.
- [BC98] L. Bertini and N. Cancrini. The two-dimensional stochastic heat equation: renormalizing a multiplicative noise. *J. Phys. A*, 31(2):615–622, 1998.
- [BDZ16] M. Bramson, J. Ding, and O. Zeitouni. Convergence in law of the maximum of the two-dimensional discrete Gaussian free field. *Comm. Pure Appl. Math.*, 69(1):62–123, 2016.
- [BG20] N. Barashkov and M. Gubinelli. A variational method for Φ_3^4 . *Duke Math. J.*, 169(17):3339–3415, 2020.
- [Bis20] M. Biskup. Extrema of the two-dimensional discrete Gaussian free field. In *Random graphs, phase transitions, and the Gaussian free field*, volume 304 of *Springer Proc. Math. Stat.*, pages 163–407. Springer, Cham, [2020] ©2020.

- [BLM22] R. Bazaes, I. Lammers, and C. Mukherjee. Subcritical gaussian multiplicative chaos in the wiener space: construction, moments and volume decay, arXiv:2211.08996, 2022.
- [BM22] Y. Bröker and C. Mukherjee. Geometry of the Gaussian multiplicative chaos in the Wiener space, arXiv:2008.04290, 2022.
- [BS22] A. Bowditch and R. Sun. The two-dimensional continuum random field Ising model. *Ann. Probab.*, 50(2):419–454, 2022.
- [CC22] F. Caravenna and F. Cottini. Gaussian limits for subcritical chaos. *Electron. J. Probab.*, 27:Paper No. 81, 35, 2022.
- [CCHS22] A. Chandra, I. Chevyrev, M. Hairer, and H. Shen. Langevin dynamic for the 2D Yang-Mills measure. *Publ. Math. Inst. Hautes Études Sci.*, 136:1–147, 2022.
- [CCHS24] A. Chandra, I. Chevyrev, M. Hairer, and H. Shen. Stochastic quantisation of Yang-Mills-Higgs in 3D. *Invent. Math.*, 237(2):541–696, 2024.
- [CCM20] F. Comets, C. Cosco, and C. Mukherjee. Renormalizing the Kardar-Parisi-Zhang equation in $d \geq 3$ in weak disorder. *J. Stat. Phys.*, 179(3):713–728, 2020.
- [CCM24] F. Comets, C. Cosco, and C. Mukherjee. Space-time fluctuation of the Kardar-Parisi-Zhang equation in $d \geq 3$ and the Gaussian free field. *Ann. Inst. Henri Poincaré Probab. Stat.*, 60(1):82–112, 2024.
- [CCR23] F. Caravenna, F. Cottini, and M. Rossi. Quasi-critical fluctuations for 2d directed polymers, arXiv:2307.02453, 2023.
- [CD20] S. Chatterjee and A. Dunlap. Constructing a solution of the $(2+1)$ -dimensional KPZ equation. *Ann. Probab.*, 48(2):1014–1055, 2020.
- [CD24] C. Cosco and A. Donadini. On the central limit theorem for the log-partition function of 2d directed polymers, arXiv:2402.14647, 2024.
- [CES21] G. Cannizzaro, D. Erhard, and P. Schönauer. 2D anisotropic KPZ at stationarity: scaling, tightness and nontriviality. *Ann. Probab.*, 49(1):122–156, 2021.
- [CET23] G. Cannizzaro, D. Erhard, and F. Toninelli. The stationary AKPZ equation: logarithmic superdiffusivity. *Comm. Pure Appl. Math.*, 76(11):3044–3103, 2023.
- [CFLW21] T. Claeys, B. Fahs, G. Lambert, and C. Webb. How much can the eigenvalues of a random Hermitian matrix fluctuate? *Duke Math. J.*, 170(9):2085–2235, 2021.
- [CGRV19] L. Chevillard, C. Garban, R. Rhodes, and V. Vargas. On a skewed and multifractal unidimensional random field, as a probabilistic representation of Kolmogorov’s views on turbulence. *Ann. Henri Poincaré*, 20(11):3693–3741, 2019.
- [Cha06] S. Chatterjee. A generalization of the Lindeberg principle. *Ann. Probab.*, 34(6):2061–2076, 2006.
- [Cha09] S. Chatterjee. Fluctuations of eigenvalues and second order Poincaré inequalities. *Probab. Theory Related Fields*, 143(1-2):1–40, 2009.
- [Che22] Y.-T. Chen. Two-dimensional delta-bose gas: skew-product relative motion and exact non-Gaussianity for the stochastic heat equation, arXiv:2207.06331, 2022.
- [Che24a] Y.-T. Chen. Delta-Bose gas from the viewpoint of the two-dimensional stochastic heat equation. *Ann. Probab.*, 52(1):127–187, 2024.
- [Che24b] Y.-T. Chen. Stochastic motions of the two-dimensional many-body delta-bose gas, arXiv:2401.17243, 2024.
- [Che25] Y.-T. Chen. Martingale problem of the two-dimensional stochastic heat equation at criticality, arXiv:2504.21791, 2025.
- [Cla23] J. T. Clark. The conditional Gaussian multiplicative chaos structure underlying a critical continuum random polymer model on a diamond fractal. *Ann. Inst. Henri Poincaré Probab. Stat.*, 59(3):1203–1222, 2023.
- [CM23] J. Clark and B. Mian. On planar Brownian motion singularly tilted through a point potential, arXiv:2306.14849, 2023.
- [CM24] J. Clark and B. Mian. Continuum polymer measures corresponding to the critical 2d stochastic heat flow, arXiv:2409.01510, 2024.
- [CN25] C. Cosco and S. Nakajima. High moments of 2d directed polymers up to quasi-criticality, arXiv:2503.17232, 2025.
- [CNN22] C. Cosco, S. Nakajima, and M. Nakashima. Law of large numbers and fluctuations in the sub-critical and L^2 regions for SHE and KPZ equation in dimension $d \geq 3$. *Stochastic Process. Appl.*, 151:127–173, 2022.

- [CNZ25] C. Cosco, S. Nakajima, and O. Zeitouni. The maximum of the two dimensional gaussian directed polymer in the subcritical regime, arXiv:2503.17236, 2025.
- [Com17] F. Comets. *Directed polymers in random environments*, volume 2175 of *Lecture Notes in Mathematics*. Springer, Cham, 2017. Lecture notes from the 46th Probability Summer School held in Saint-Flour, 2016.
- [Cor12] I. Corwin. The Kardar-Parisi-Zhang equation and universality class. *Random Matrices Theory Appl.*, 1(1):1130001, 76, 2012.
- [Cor16] I. Corwin. Kardar-Parisi-Zhang universality. *Notices Amer. Math. Soc.*, 63(3):230–239, 2016.
- [CS23] I. Chevyrev and H. Shen. Invariant measure and universality of the 2D Yang-Mills Langevin dynamic, arXiv:2302.12160, 2023.
- [CSY03] F. Comets, T. Shiga, and N. Yoshida. Directed polymers in a random environment: path localization and strong disorder. *Bernoulli*, 9(4):705–723, 2003.
- [CSZ16] F. Caravenna, R. Sun, and N. Zygouras. The continuum disordered pinning model. *Probab. Theory Related Fields*, 164(1-2):17–59, 2016.
- [CSZ17a] F. Caravenna, R. Sun, and N. Zygouras. Polynomial chaos and scaling limits of disordered systems. *J. Eur. Math. Soc. (JEMS)*, 19(1):1–65, 2017.
- [CSZ17b] F. Caravenna, R. Sun, and N. Zygouras. Universality in marginally relevant disordered systems. *Ann. Appl. Probab.*, 27(5):3050–3112, 2017.
- [CSZ19a] F. Caravenna, R. Sun, and N. Zygouras. The Dickman subordinator, renewal theorems, and disordered systems. *Electron. J. Probab.*, 24:Paper No. 101, 40, 2019.
- [CSZ19b] F. Caravenna, R. Sun, and N. Zygouras. On the moments of the $(2+1)$ -dimensional directed polymer and stochastic heat equation in the critical window. *Comm. Math. Phys.*, 372(2):385–440, 2019.
- [CSZ20] F. Caravenna, R. Sun, and N. Zygouras. The two-dimensional KPZ equation in the entire subcritical regime. *Ann. Probab.*, 48(3):1086–1127, 2020.
- [CSZ23a] F. Caravenna, R. Sun, and N. Zygouras. The critical 2d stochastic heat flow. *Invent. Math.*, 233(1):325–460, 2023.
- [CSZ23b] F. Caravenna, R. Sun, and N. Zygouras. The critical 2d stochastic heat flow is not a Gaussian multiplicative chaos. *Ann. Probab.*, 51(6):2265–2300, 2023.
- [CSZ25a] F. Caravenna, R. Sun, and N. Zygouras. The flow property of the critical 2d stochastic heat flow, 2025+.
- [CSZ25b] F. Caravenna, R. Sun, and N. Zygouras. Singularity and regularity of the critical 2D stochastic heat flow, arXiv:2504.06128, 2025.
- [CT24] G. Cannizzaro and F. Toninelli. Lecture notes on stationary critical and super-critical SPDEs, arXiv:2403.15006, 2024.
- [CY06] F. Comets and N. Yoshida. Directed polymers in random environment are diffusive at weak disorder. *Ann. Probab.*, 34(5):1746–1770, 2006.
- [CZ23] C. Cosco and O. Zeitouni. Moments of partition functions of 2d Gaussian polymers in the weak disorder regime—I. *Comm. Math. Phys.*, 403(1):417–450, 2023.
- [DFT94] G. F. Dell’Antonio, R. Figari, and A. Teta. Hamiltonians for systems of N particles interacting through point interactions. *Ann. Inst. H. Poincaré Phys. Théor.*, 60(3):253–290, 1994.
- [DG22] A. Dunlap and Y. Gu. A forward-backward SDE from the 2D nonlinear stochastic heat equation. *Ann. Probab.*, 50(3):1204–1253, 2022.
- [DG23] A. Dunlap and C. Graham. The 2d nonlinear stochastic heat equation: pointwise statistics and the decoupling function, arXiv:2308.11850, 2023.
- [DG24] A. Dunlap and C. Graham. Edwards-wilkinson fluctuations in subcritical 2d stochastic heat equations, arXiv:2405.09520, 2024.
- [Dic30] K. Dickman. On the frequency of numbers containing prime factors of a certain relative magnitude. *Ark. Math. Astr. Fys.*, 22:1–14, 1930.
- [DMY06] C. Donati-Martin and M. Yor. Some explicit Krein representations of certain subordinators, including the gamma process. *Publ. Res. Inst. Math. Sci.*, 42(4):879–895, 2006.
- [DR04] J. Dimock and S. G. Rajeev. Multi-particle Schrödinger operators with point interactions in the plane. *J. Phys. A*, 37(39):9157–9173, 2004.
- [DRSV17] B. Duplantier, R. Rhodes, S. Sheffield, and V. Vargas. Log-correlated Gaussian fields: an overview. In *Geometry, analysis and probability*, volume 310 of *Progr. Math.*, pages 191–216. Birkhäuser/Springer, Cham, 2017.

- [Duc22] P. Duch. Flow equation approach to singular stochastic PDEs, arXiv:2109.11380, 2022.
- [EF16] T. Ellis and O. N. Feldheim. The Brownian web is a two-dimensional black noise. *Ann. Inst. Henri Poincaré Probab. Stat.*, 52(1):162–172, 2016.
- [ET60] P. Erdős and S. J. Taylor. Some problems concerning the structure of random walk paths. *Acta Math. Acad. Sci. Hungar.*, 11:137–162. (unbound insert), 1960.
- [FM17] M. Furlan and J.-C. Mourrat. A tightness criterion for random fields, with application to the Ising model. *Electron. J. Probab.*, 22:Paper No. 97, 29, 2017.
- [FT94] E. Frey and U. C. Täuber. Two-loop renormalization-group analysis of the Burgers-Kardar-Parisi-Zhang equation. *Phys. Rev. E (3)*, 50(2):1024–1045, 1994.
- [GH21] M. Gubinelli and M. Hofmanová. A PDE construction of the Euclidean ϕ_3^4 quantum field theory. *Comm. Math. Phys.*, 384(1):1–75, 2021.
- [Gia11] G. Giacomin. *Disorder and critical phenomena through basic probability models*, volume 2025 of *Lecture Notes in Mathematics*. Springer, Heidelberg, 2011. Lecture notes from the 40th Probability Summer School held in Saint-Flour, 2010, École d’Été de Probabilités de Saint-Flour. [Saint-Flour Probability Summer School].
- [GIP15] M. Gubinelli, P. Imkeller, and N. Perkowski. Paracontrolled distributions and singular PDEs. *Forum Math. Pi*, 3:e6, 75, 2015.
- [GJ14] P. Gonçalves and M. Jara. Nonlinear fluctuations of weakly asymmetric interacting particle systems. *Arch. Ration. Mech. Anal.*, 212(2):597–644, 2014.
- [GKLV24] N. Gupta, A. Kumar, N. Leonenko, and J. Vaz. Generalized fractional derivatives generated by Dickman subordinator and related stochastic processes. *Fract. Calc. Appl. Anal.*, 27(4):1527–1563, 2024.
- [GQT21] Y. Gu, J. Quastel, and L.-C. Tsai. Moments of the 2D SHE at criticality. *Probab. Math. Phys.*, 2(1):179–219, 2021.
- [GRZ18] Y. Gu, L. Ryzhik, and O. Zeitouni. The Edwards-Wilkinson limit of the random heat equation in dimensions three and higher. *Comm. Math. Phys.*, 363(2):351–388, 2018.
- [GS09] J. Gärtner and R. Sun. A quenched limit theorem for the local time of random walks on \mathbb{Z}^2 . *Stochastic Process. Appl.*, 119(4):1198–1215, 2009.
- [GT25] Y. Gu and L.-C. Tsai. Stochastic heat flow is a black noise, arXiv:2506.16484, 2025.
- [Gu20] Y. Gu. Gaussian fluctuations from the 2D KPZ equation. *Stoch. Partial Differ. Equ. Anal. Comput.*, 8(1):150–185, 2020.
- [Hai13] M. Hairer. Solving the KPZ equation. *Ann. of Math. (2)*, 178(2):559–664, 2013.
- [Hai14] M. Hairer. A theory of regularity structures. *Invent. Math.*, 198(2):269–504, 2014.
- [Har74] A.B. Harris. Effect of random defects on the critical behaviour of ising models. *Journal of Physics C: Solid State Physics*, 7(9):1671, 1974.
- [HP24] Z. Himwich and S. Parekh. The directed landscape is a black noise, arXiv:2404.16801, 2024.
- [JL24] S. Junk and H. Lacoin. Strong disorder and very strong disorder are equivalent for directed polymers, arXiv:2402.02562, 2024.
- [JN24] S. Junk and S. Nakajima. Equivalence of fluctuations of discretized SHE and KPZ equations in the subcritical weak disorder regime, arXiv:2410.06801, 2024.
- [Jun22] S. Junk. Fluctuations of partition functions of directed polymers in weak disorder beyond the L^2 -phase, arXiv:2202.02907, 2022. To appear in *Ann. Probab.*
- [Jun23] S. Junk. Local limit theorem for directed polymers beyond the L^2 -phase, arXiv:2307.05097, 2023.
- [Kot24] S. Kotitsas. The heat equation with time-correlated random potential in $d = 2$: Edwards-Wilkinson fluctuations, arXiv:2405.01519, 2024.
- [Koz07] G. Kozma. The scaling limit of loop-erased random walk in three dimensions. *Acta Math.*, 199(1):29–152, 2007.
- [Kup16] A. Kupiainen. Renormalization group and stochastic PDEs. *Ann. Henri Poincaré*, 17(3):497–535, 2016.
- [LaM17] R. Latał a and D. Matlak. Royen’s proof of the Gaussian correlation inequality. In *Geometric aspects of functional analysis*, volume 2169 of *Lecture Notes in Math.*, pages 265–275. Springer, Cham, 2017.
- [Led01] M. Ledoux. *The concentration of measure phenomenon*, volume 89 of *Mathematical Surveys and Monographs*. American Mathematical Society, Providence, RI, 2001.

- [LJR04] Y. Le Jan and O. Raimond. Flows, coalescence and noise. *Ann. Probab.*, 32(2):1247–1315, 2004.
- [LL01] E. H. Lieb and M. Loss. *Analysis*, volume 14 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, second edition, 2001.
- [LL10] G. F. Lawler and V. Limic. *Random walk: a modern introduction*, volume 123 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 2010.
- [LMS24] Liuyan Li, V. Margarint, and R. Sun. Scaling limits of disorder relevant non-binary spin systems, arXiv:2410.00579, 2024.
- [LZ22] D. Lygkonis and N. Zygouras. Edwards-Wilkinson fluctuations for the directed polymer in the full L^2 -regime for dimensions $d \geq 3$. *Ann. Inst. Henri Poincaré Probab. Stat.*, 58(1):65–104, 2022.
- [LZ23] D. Lygkonis and N. Zygouras. Moments of the 2D directed polymer in the subcritical regime and a generalisation of the Erdős-Taylor theorem. *Comm. Math. Phys.*, 401(3):2483–2520, 2023.
- [LZ24a] Z. Liu and N. Zygouras. On the moments of the mass of shrinking balls under the critical $2d$ stochastic heat flow, arXiv:2410.14601, 2024.
- [LZ24b] D. Lygkonis and N. Zygouras. A multivariate extension of the Erdős-Taylor theorem. *Probab. Theory Related Fields*, 189(1-2):179–227, 2024.
- [MOO10] E. Mossel, R. O'Donnell, and K. Oleszkiewicz. Noise stability of functions with low influences: invariance and optimality. *Ann. of Math. (2)*, 171(1):295–341, 2010.
- [MU18] J. Magnen and J. Unterberger. The scaling limit of the KPZ equation in space dimension 3 and higher. *J. Stat. Phys.*, 171(4):543–598, 2018.
- [MW17] J.-C. Mourrat and H. Weber. Global well-posedness of the dynamic Φ^4 model in the plane. *Ann. Probab.*, 45(4):2398–2476, 2017.
- [Nak25] M. Nakashima. Martingale measure associated with the critical $2d$ stochastic heat flow, arXiv:2503.20171, 2025.
- [NN23] S. Nakajima and M. Nakashima. Fluctuations of two-dimensional stochastic heat equation and KPZ equation in subcritical regime for general initial conditions. *Electron. J. Probab.*, 28:Paper No. 1, 38, 2023.
- [NPR09] I. Nourdin, G. Peccati, and G. Reinert. Second order Poincaré inequalities and CLTs on Wiener space. *J. Funct. Anal.*, 257(2):593–609, 2009.
- [Nua06] D. Nualart. *The Malliavin calculus and related topics*. Probability and its Applications (New York). Springer-Verlag, Berlin, second edition, 2006.
- [PW04] M. D. Penrose and A. R. Wade. Random minimal directed spanning trees and Dickman-type distributions. *Adv. in Appl. Probab.*, 36(3):691–714, 2004.
- [QS15] J. Quastel and H. Spohn. The one-dimensional KPZ equation and its universality class. *J. Stat. Phys.*, 160(4):965–984, 2015.
- [Raj99] S.G. Rajeev. A condensation of interacting Bosons in two dimensional space, arXiv:hep-th/9905120, 1999.
- [Rot79] V. I. Rotar'. Limit theorems for polylinear forms. *J. Multivariate Anal.*, 9(4):511–530, 1979.
- [Roy14] T. Royen. A simple proof of the Gaussian correlation conjecture extended to some multivariate gamma distributions. *Far East J. Theor. Stat.*, 48(2):139–145, 2014.
- [RV14] R. Rhodes and V. Vargas. Gaussian multiplicative chaos and applications: a review. *Probab. Surv.*, 11:315–392, 2014.
- [SS11] O. Schramm and S. Smirnov. On the scaling limits of planar percolation. *Ann. Probab.*, 39(5):1768–1814, 2011. With an appendix by Christophe Garban.
- [SSS14] E. Schertzer, R. Sun, and J. M. Swart. Stochastic flows in the Brownian web and net. *Mem. Amer. Math. Soc.*, 227(1065):vi+160, 2014.
- [Tao24a] R. Tao. Gaussian fluctuations of a nonlinear stochastic heat equation in dimension two. *Stoch. Partial Differ. Equ. Anal. Comput.*, 12(1):220–246, 2024.
- [Tao24b] R. Tao. Mesoscopic averaging of the two-dimensional KPZ equation. *J. Stat. Phys.*, 191(1):Paper No. 6, 28, 2024.
- [Tsa24] L.-C. Tsai. Stochastic heat flow by moments, arXiv:2410.14657, 2024.
- [Tsi04a] B. Tsirelson. Nonclassical stochastic flows and continuous products. *Probab. Surv.*, 1:173–298, 2004.

- [Tsi04b] B. Tsirelson. Scaling limit, noise, stability. In *Lectures on probability theory and statistics*, volume 1840 of *Lecture Notes in Math.*, pages 1–106. Springer, Berlin, 2004.
- [Wan24] T.-C. Wang. Asymptotics and the sub-limit at L^2 -criticality of higher moments for the SHE in dimension $d \geq 3$, arXiv:2410.06426, 2024.
- [Wol91] E. Wolf. Kinetic roughening of vicinal surfaces. *Phys. Rev. Lett.*, 67:1783–1786, 1991.
- [WY24] R. Wei and J. Yu. The critical disordered pinning measure, arXiv:2402.17642, 2024.
- [Zyg22] N. Zygouras. Some algebraic structures in KPZ universality. *Probab. Surv.*, 19:590–700, 2022.
- [Zyg24] N. Zygouras. Directed polymers in a random environment: a review of the phase transitions. *Stochastic Process. Appl.*, 177:Paper No. 104431, 34, 2024.

DIPARTIMENTO DI MATEMATICA E APPLICAZIONI, UNIVERSITÀ DEGLI STUDI DI MILANO-BICOCCA, VIA COZZI 55, 20125 MILANO, ITALY

Email address: francesco.caravenna@unimib.it

DEPARTMENT OF MATHEMATICS, NATIONAL UNIVERSITY OF SINGAPORE, 10 LOWER KENT RIDGE ROAD, 119076 SINGAPORE

Email address: matsr@nus.edu.sg

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WARWICK, COVENTRY CV4 7AL, UK

Email address: N.Zygouras@warwick.ac.uk