

# The weak coupling limit of disordered copolymer models

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Joint work with Giambattista Giacomin (Université Paris Diderot)

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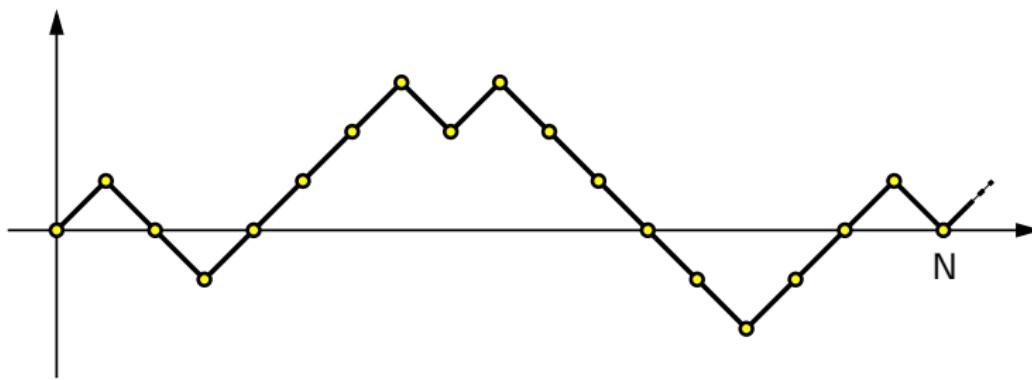
# Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

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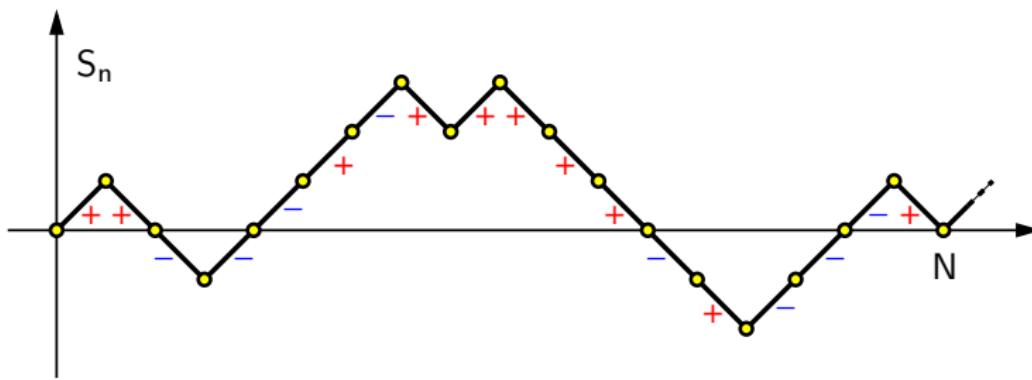
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# A random walk with a random potential



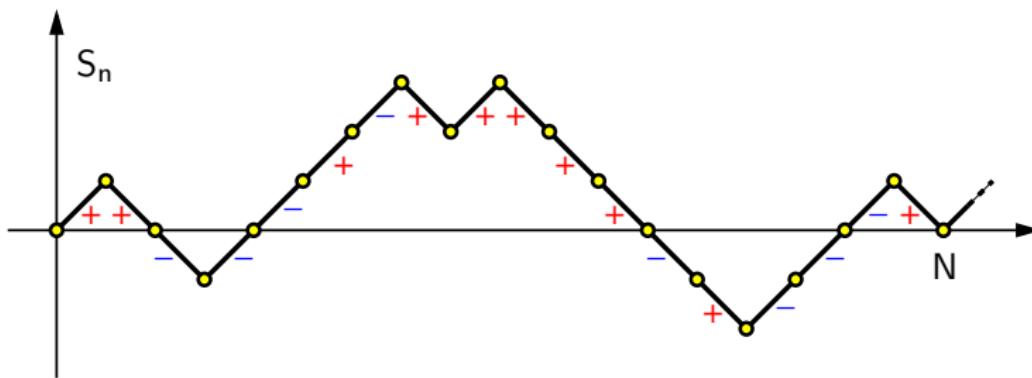
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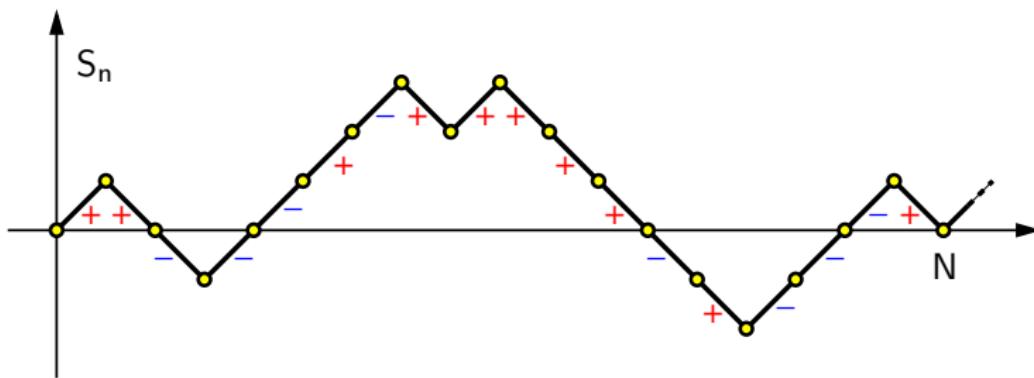
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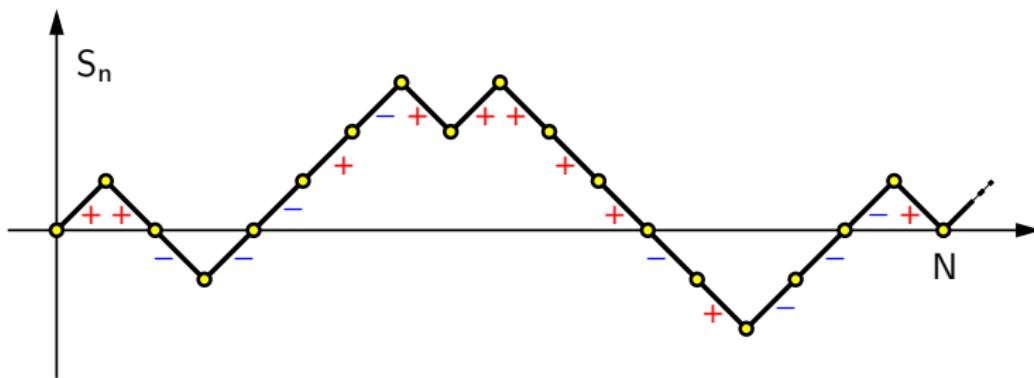


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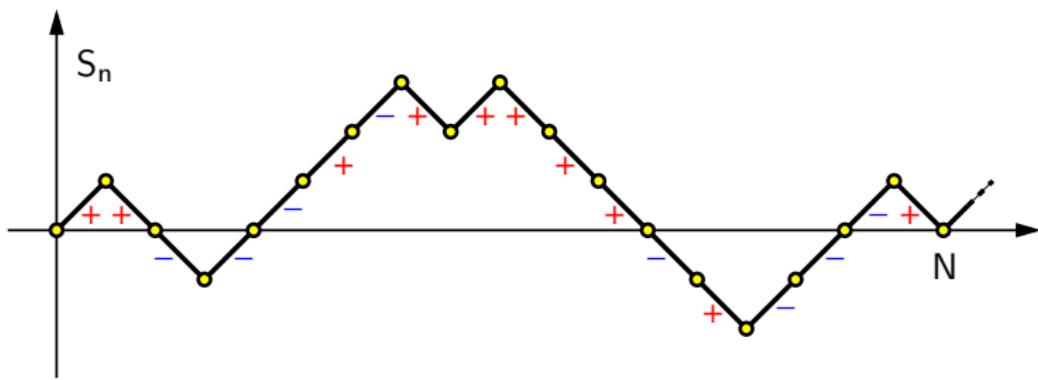


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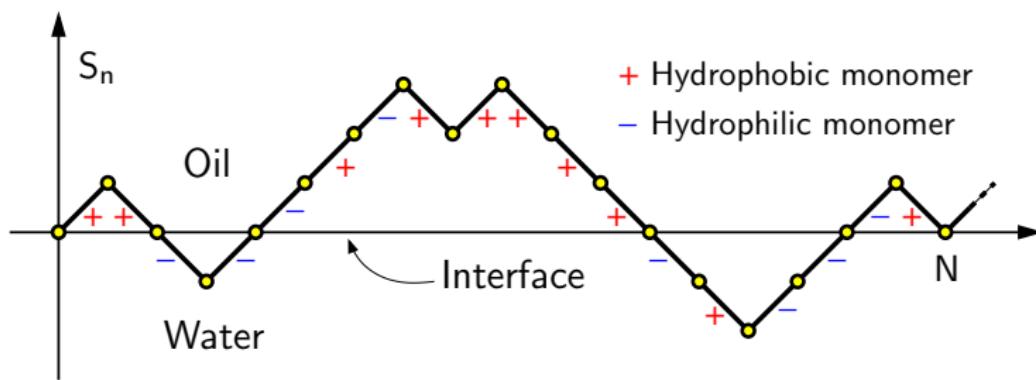
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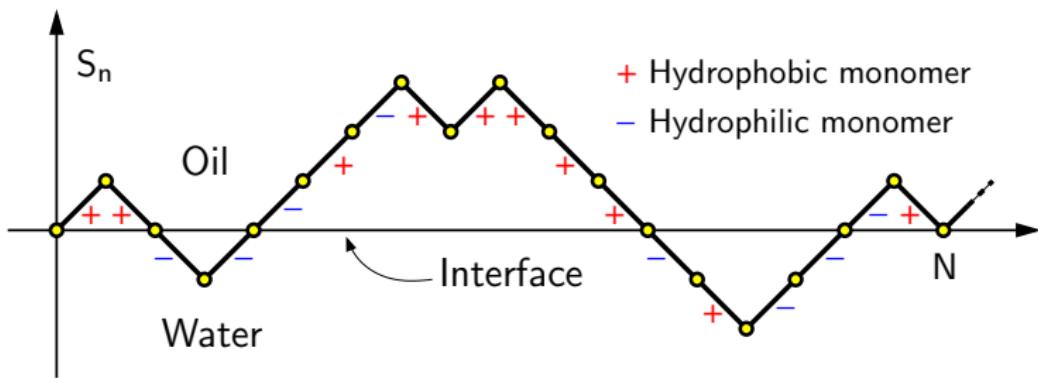
**Localization or Delocalization?**

# A polymer model interpretation



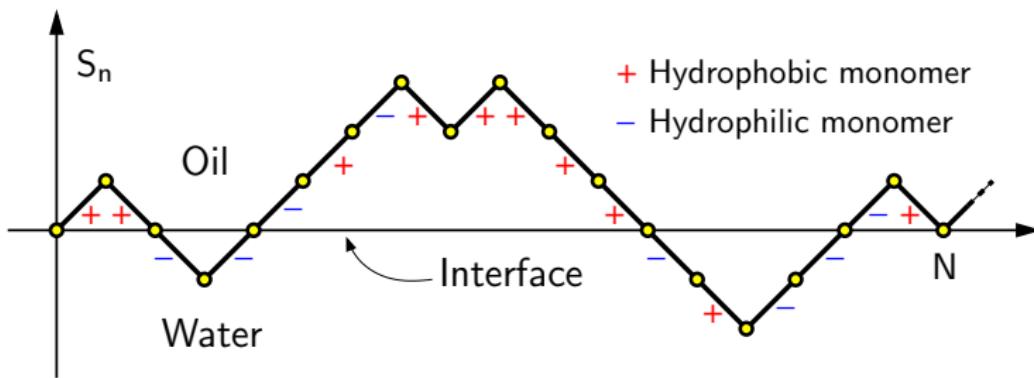
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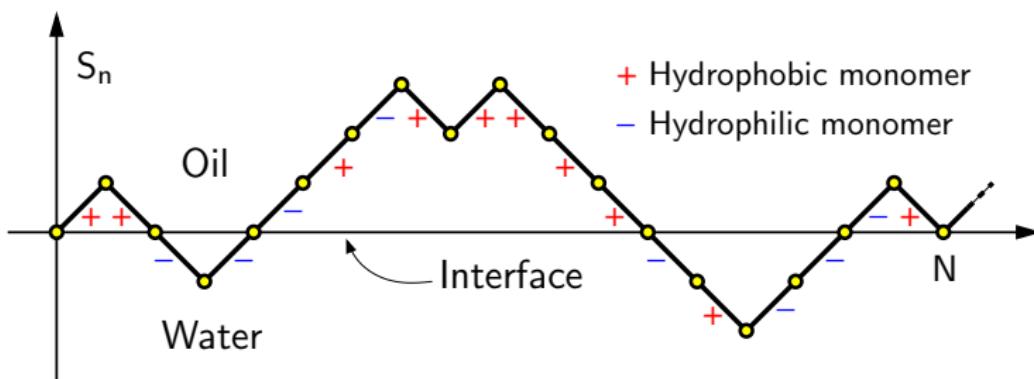
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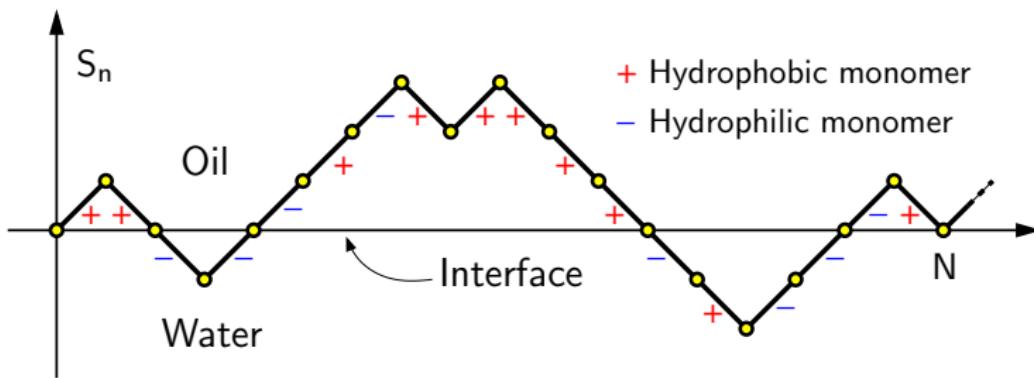
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- ▶ Several contributions in the physics literature too.

# The basic copolymer model

Definition of the model:  $\frac{d\mathbf{P}_{N,\omega}}{d\mathbf{P}}(S) := \frac{1}{Z_{N,\omega}} \exp(-H_{N,\omega}(S))$

$$-H_{N,\omega}(S) := \lambda \sum_{n=1}^N (\omega_n + h) \operatorname{sign}((S_{n-1}, S_n))$$

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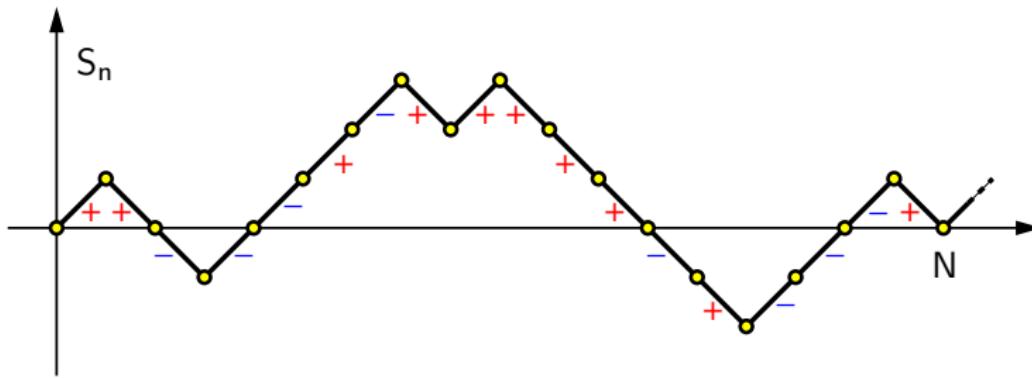
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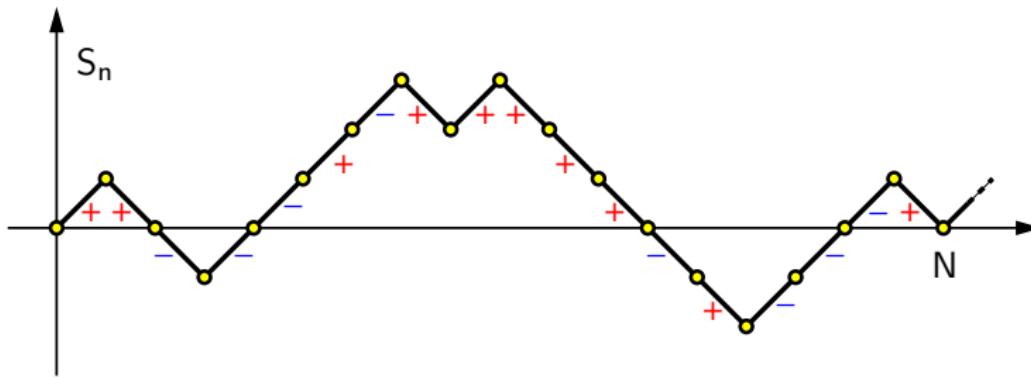
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We now define a continuum model.

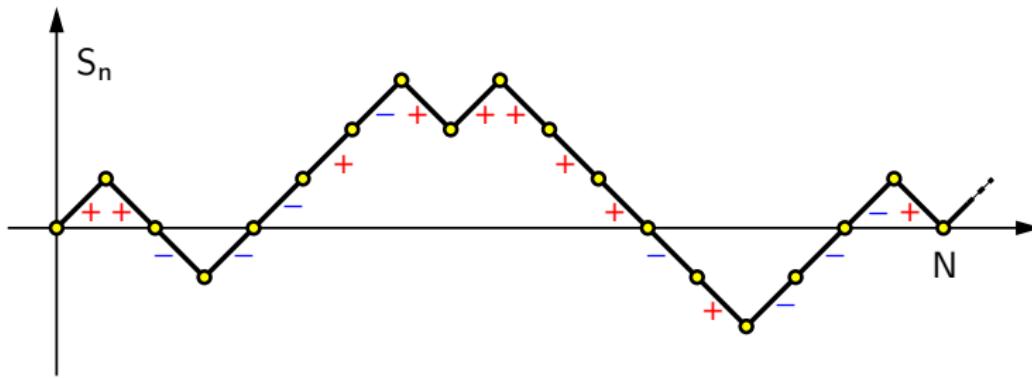
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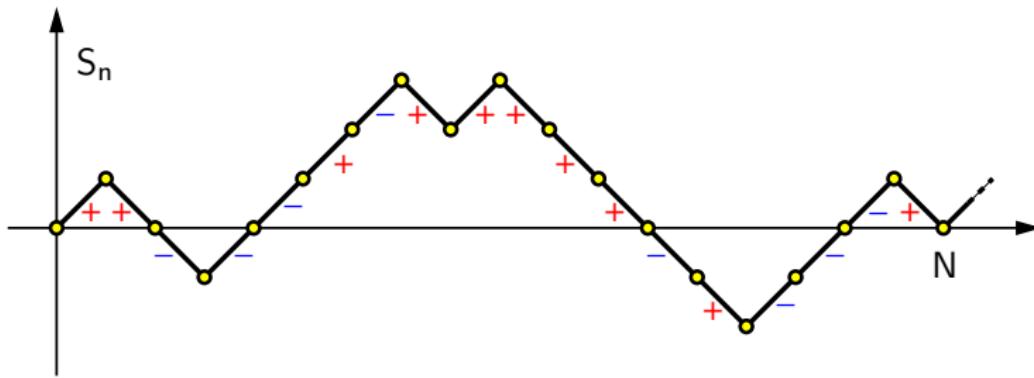
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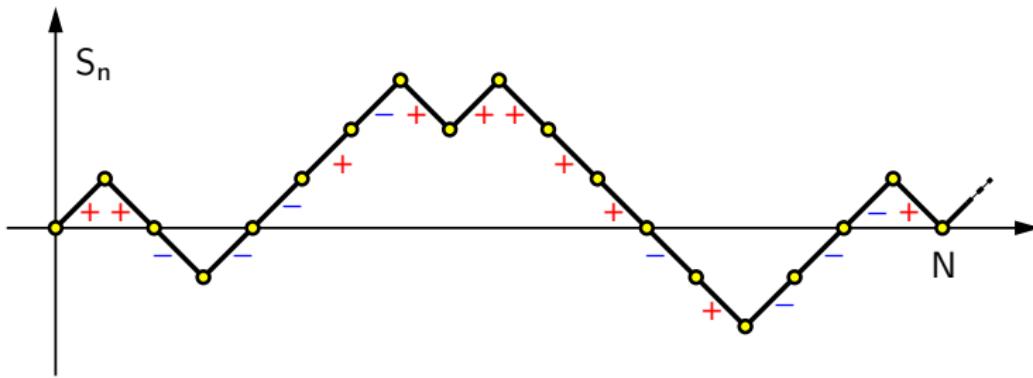


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Discrete model:  $- 2\lambda \sum_{n=1}^N (\textcolor{red}{\omega_n} + h) \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$

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(This definition does correspond to sharply different path behaviors)

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## Theorem

*The regions  $\mathcal{L}$  and  $\mathcal{D}$  are separated by a strictly increasing, continuous critical line  $h_c(\cdot)$ :*

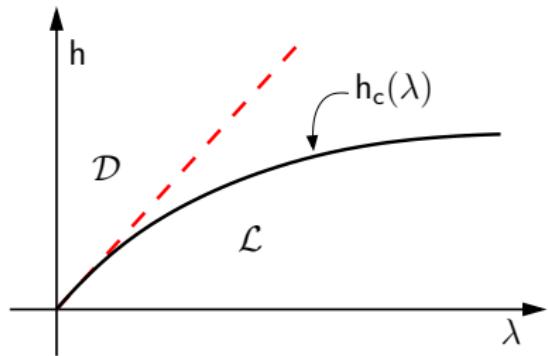
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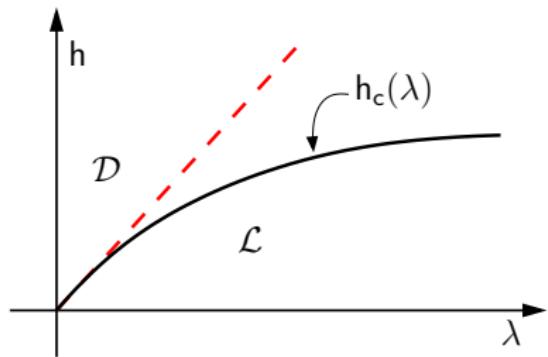


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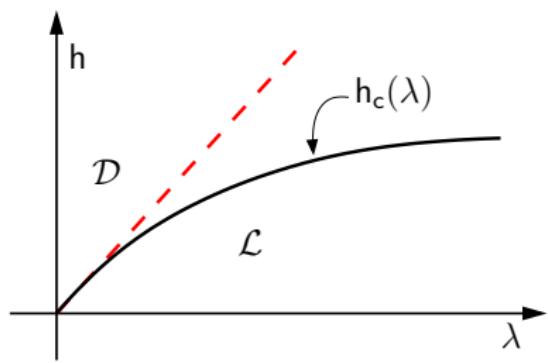
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with

$$\underline{h}'(0+) = \frac{2}{3}, \quad \bar{h}'(0+) = 1 - \epsilon.$$

# The phase diagram: continuum model

The continuum free energy  $\tilde{F}(\lambda, h)$  is defined analogously:

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Again  $\tilde{F}(\lambda, h) \geq 0$ . We then define  $\tilde{\mathcal{L}}$ ocalization and  $\tilde{\mathcal{D}}$ elocalization:

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For all  $\lambda, h \geq 0$

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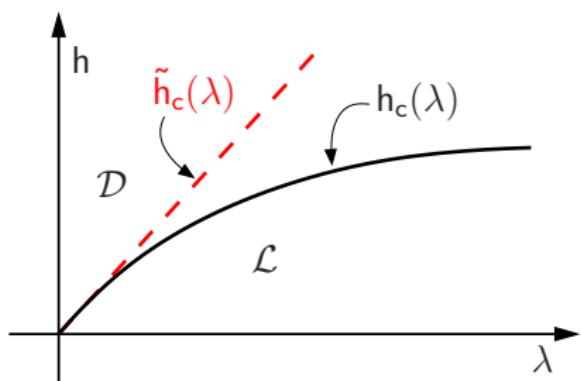
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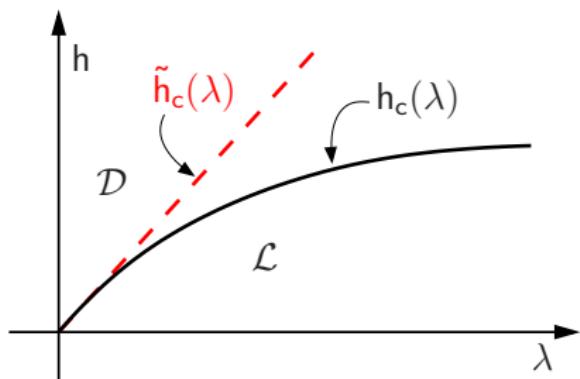
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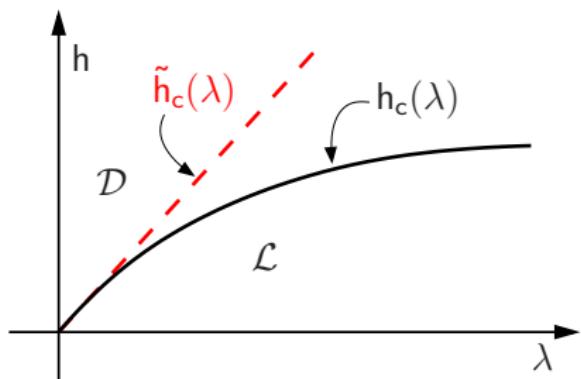
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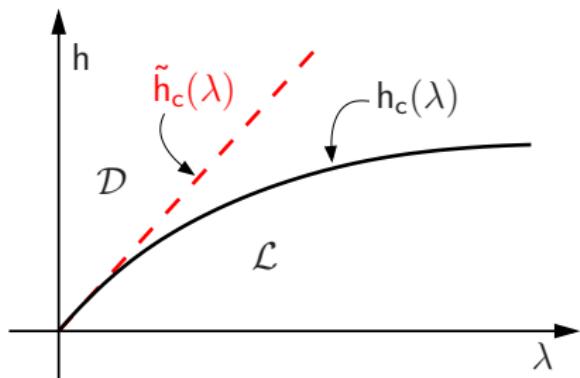
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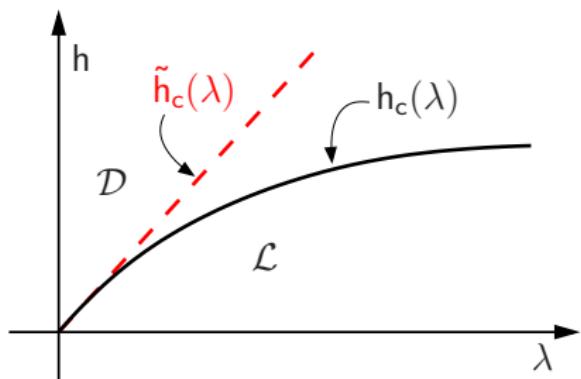
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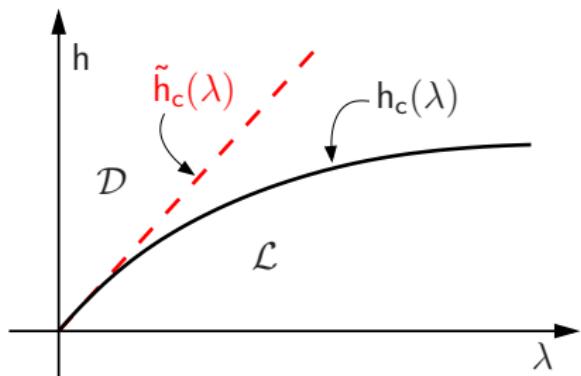
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... unfortunately for just **one** discrete model. Generalization?

# Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

# Beyond the simple random walk

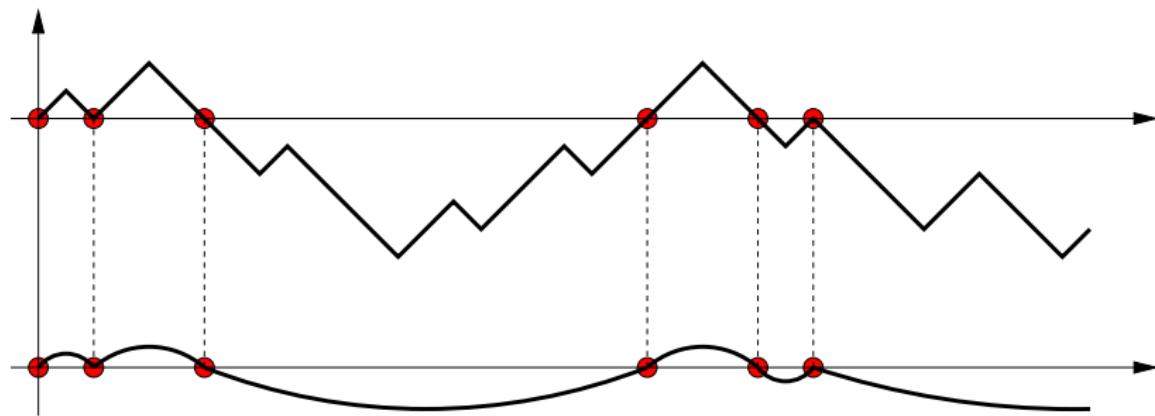
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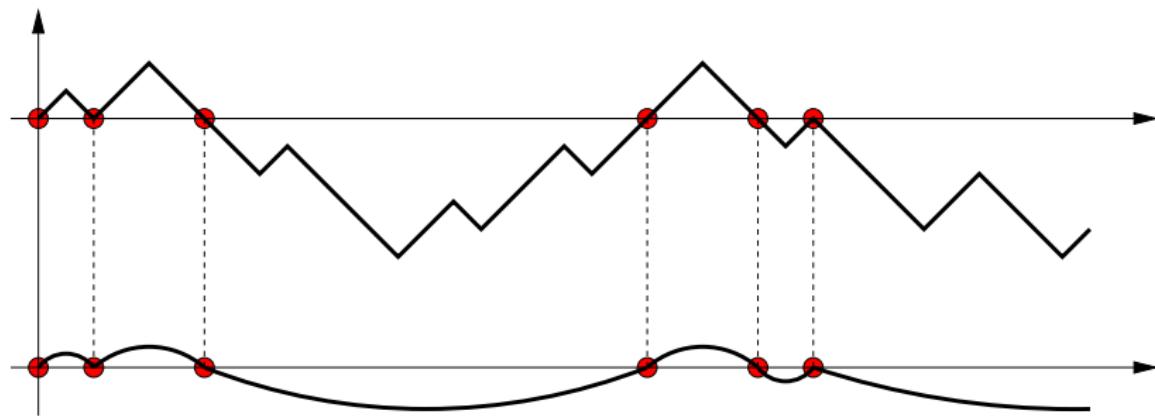
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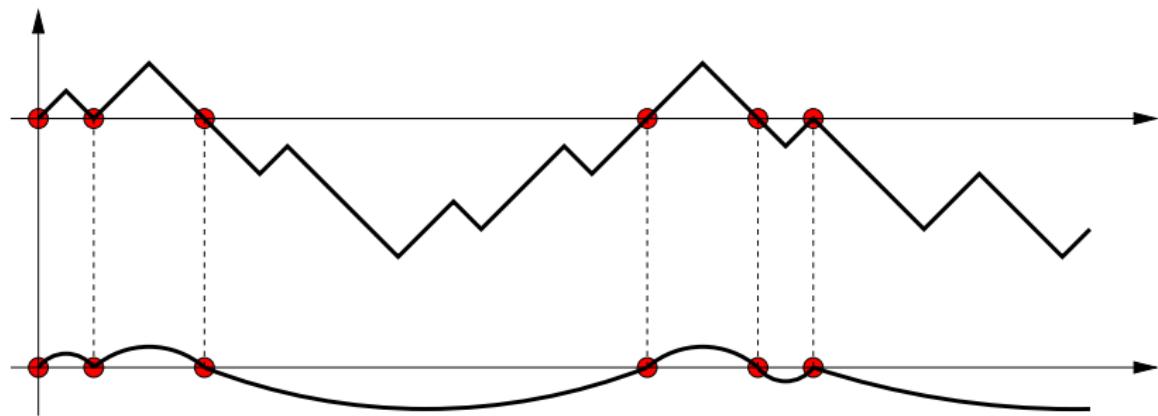


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- ▶ Excursions signs: fair coin tossing (independent of  $\{\tau_k\}_{k \geq 0}$ )

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Discrete Bessel-like process ( $c_\alpha = 1/2 - \alpha$ )

$$\mathbf{P}(S_{n+1} = x \pm 1 | S_n = x) = \frac{1}{2} \left( 1 \pm \frac{c_\alpha}{x} + o\left(\frac{1}{x}\right) \right) \text{ yields } (\star) \text{ asymp.}$$

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From  $\tilde{\tau}^\alpha$  we obtain  $(\tilde{\Delta}^\alpha = \{\tilde{\Delta}_t^\alpha\}_{t \geq 0}, \tilde{\mathbf{P}})$  (For  $\alpha = \frac{1}{2}$  we recover BM)

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$$\lim_{a \downarrow 0} \frac{F(a\lambda, ah)}{a^2} = \tilde{F}^\alpha(\lambda, h) \quad \lim_{\lambda \downarrow 0} \frac{h_c(\lambda)}{\lambda} = \tilde{m}^\alpha$$

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**Theorem**

For all  $\lambda, h > 0$  and  $\epsilon \in (0, 1)$  there exists  $a_0 > 0$  s.t. for all  $a < a_0$

$$\tilde{F}^\alpha \left( \frac{\lambda}{1+\epsilon}, \frac{h}{1-\epsilon} \right) \leq \frac{F(a\lambda, ah)}{a^2} \leq \tilde{F}^\alpha ((1+\epsilon)\lambda, (1-\epsilon)h)$$

# Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

# Strategy of the proof

Goal:  $\forall \lambda, h > 0, \epsilon \in (0, 1)$  one has for  $a \ll 1$

$$\frac{1}{a^2} F(a\lambda, ah) \leq \frac{1}{a^2} \tilde{F}((1 + \epsilon)a\lambda, (1 - \epsilon)ah)$$

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$$\limsup_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{E} \mathbf{E} \left[ \exp \left( -C(H_N(a\lambda, ah) - \tilde{H}_N(a\lambda, (1-\epsilon)ah)) \right) \right] \leq 0$$

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Recall that  $\Delta_n = \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$  and

$$-H_N^0(a\lambda, ah) = -2a\lambda \sum_{n=1}^N (\omega_n + ah)\Delta_n$$

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Recall that  $\Delta_n = \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$  and

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We need to show that  $\approx$  can be made  $\asymp$ .

# The proof

## Step 1: Coarse-graining of the renewal process.

Recall that  $\Delta_n = \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$  and

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Showing that  $H_N^0 \asymp H_N^1$  is delicate and very technical.

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## Step 2: From discrete charges to the white noise.

$H_N^2$  is obtained from  $H_N^1$  by replacing the charges  $\omega_n$  by i.i.d.  $N(0, 1)$  (discrete white noise).

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$H_N^3$  is obtained from  $H_N^2$  by replacing  $\Delta_n^{\eta, \delta}$  by an analogous coarse-grained version  $\tilde{\Delta}_t^{\eta, \delta}$  of the continuous-time process  $\tilde{\Delta}_t$ .

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$H_N^4 = \tilde{H}_N$  is obtained from  $H_N^3$  by replacing  $\tilde{\Delta}_t^{\eta, \delta}$  by the original (non coarse-grained) continuous-time process  $\tilde{\Delta}_t$ .

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This step is analogous to step 1.

Thank you.