

On the 2d KPZ and Stochastic Heat Equation

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Overview

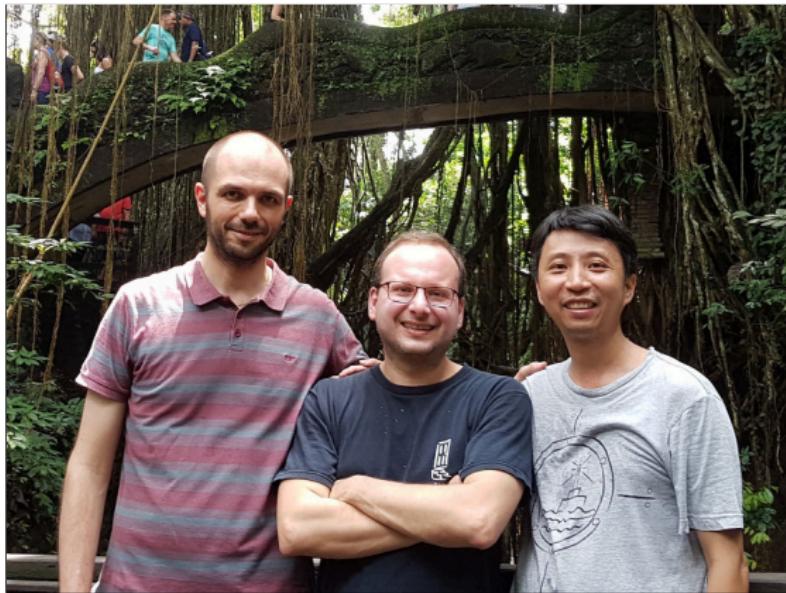
I will talk about two stochastic PDEs on \mathbb{R}^d (mainly $d = 2$)

- ▶ Kardar-Parisi-Zhang Equation (KPZ)
- ▶ Multiplicative Stochastic Heat Equation (SHE)

- ▶ Very interesting, yet ill-defined equations !
- ▶ GOAL: find a “reasonable solution”

Stochastic Analysis \rightsquigarrow Statistical Mechanics

Collaborators



Nikos Zygouras (Warwick) and Rongfeng Sun (NUS)

References

- ▶ [CSZ 17]
Universality in marginally relevant disordered systems
Ann. Appl. Probab. 2017
 - ▶ [CSZ 18a]
On the moments of the (2+1)-dimensional directed polymer and Stochastic Heat Equation in the critical window
to appear in Commun. Math. Phys.
 - ▶ [CSZ 18b]
The two-dimensional KPZ equation in the entire subcritical regime
arXiv, Dec 2018
- ($d = 2$) [Bertini Cancrini 98]
[Chatterjee Dunlap 18] [Gu 18] [Gu Quastel Tsai 19]
- ($d \geq 3$) [Magnen Unterberger 18] [Gu Ryzhik Zeitouni 18]
[Dunlap Gu Ryzhik Zeitouni 19] [Comets Cosco Mukherjee 18 19a 19b]

White noise

What is (space-time) white noise $\xi = \xi(t, x)$ on \mathbb{R}^{1+d} ?

- ▶ A centered Gaussian field with

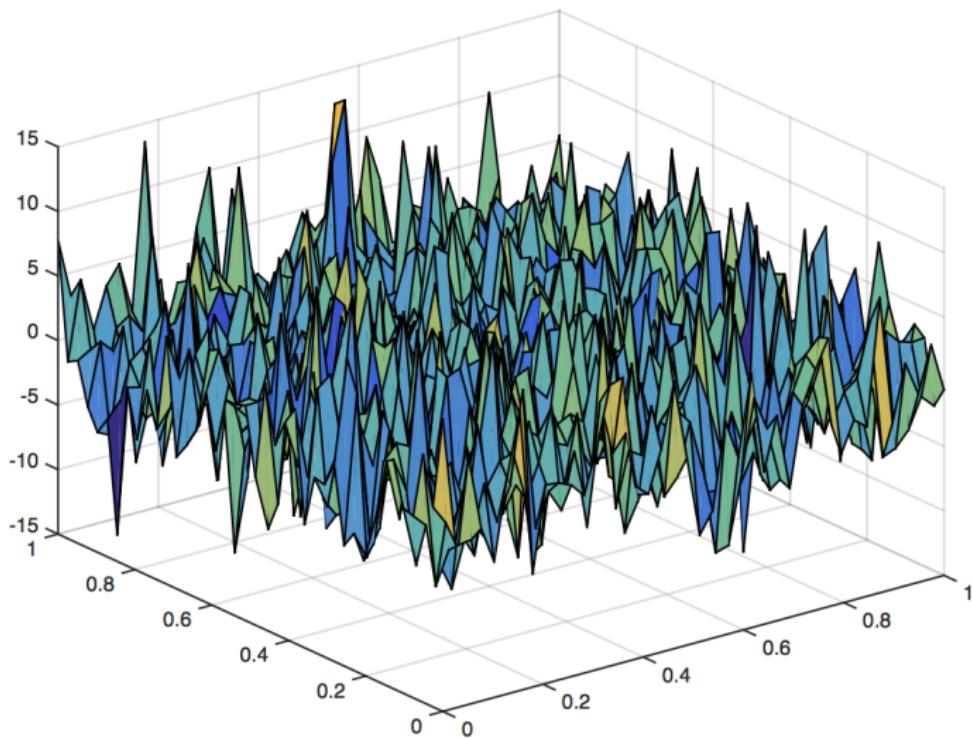
$$\text{Cov}[\xi(t, x), \xi(t', x')] = \delta(t - t') \delta(x - x')$$

- ▶ A random (Schwartz) distribution such that

$$\langle \phi, \xi \rangle = \int_{\mathbb{R}^{1+d}} \phi(t, x) \xi(t, x) dt dx \sim N(0, \|\phi\|_{L^2}^2)$$

- ▶ The scaling limit of i.i.d. $N(0, 1)$ random variables indexed by \mathbb{Z}^{1+d}

White noise



The simplest stochastic PDE

Additive Stochastic Heat Equation $(t > 0, x \in \mathbb{R}^d)$

$$\partial_t v = \frac{1}{2} \Delta_x v + c \xi \quad (\text{EW})$$

also known as Edwards-Wilkinson equation $(c \in \mathbb{R})$

- ▶ Well-posed (pathwise) in any dimension d
- ▶ Explicit solution $v = c g * \xi$

$$v(t, x) = c \int_0^t \int_{\mathbb{R}^2} g_{t-s}(x-y) \xi(s, y) ds dy \quad g_t(x) = \frac{e^{-\frac{|x|^2}{2t}}}{(2\pi t)^{d/2}}$$

v is a function for $d = 1$, a genuine distribution for $d \geq 2$

- ▶ v is Gaussian process with explicit covariance

The KPZ equation

KPZ

[Kardar Parisi Zhang PRL'86]

$$\partial_t h = \frac{1}{2} \Delta_x h + \frac{1}{2} |\nabla_x h|^2 + \beta \xi \quad (\text{KPZ})$$

Model for random interface growth

$h = h(t, x)$ = interface height at time $t \geq 0$, space $x \in \mathbb{R}^d$

$\xi = \xi(t, x)$ = space-time white noise $\beta > 0$ noise strength

$|\nabla_x h|^2$ ill-defined $(\nabla_x h(t, \cdot))$ exp. distribution

For smooth ξ we can linearize KPZ by Cole-Hopf transformation

$$u(t, x) := e^{h(t, x)}$$

The multiplicative Stochastic Heat Equation (SHE)

SHE

 $(t > 0, \ x \in \mathbb{R}^d)$

$$\partial_t u = \frac{1}{2} \Delta_x u + \beta u \xi \quad (\text{SHE})$$

Product $u \xi$ ill-defined (non smooth function \times distribution)

- ▶ $(d = 1)$ Well-posed by **Ito integration** [Walsh 80's]
- ▶ $(d = 1)$ SHE and KPZ are now well-understood in a **robust sense**
Regularity Structures (Hairer) or **Paracontrolled Distributions** (GIP)
- ▶ $(d \geq 2)$ No general theory. How to “find a solution”?

(KPZ) and (SHE) in dimensions $d \geq 2$

Mollify the white noise $\xi(t, x)$ in space on scale $\varepsilon > 0$

$$\xi^\varepsilon(t, \cdot) := \xi(t, \cdot) * \varrho_\varepsilon$$

$$\varrho_\varepsilon(x) := \varepsilon^{-d} \varrho(\varepsilon^{-1}x) \quad \text{probability density } \varrho \in C_c^\infty(\mathbb{R}^d)$$

- ▶ Replace ξ by ξ^ε \rightsquigarrow solutions $h^\varepsilon(t, x)$, $u^\varepsilon(t, x)$ well-defined (**Ito**)
- ▶ Do $h^\varepsilon(t, x)$ and $u^\varepsilon(t, x)$ have a limit as $\varepsilon \downarrow 0$?
- ▶ Disorder strength $\beta = \beta_\varepsilon$ needs to be renormalized

Mollified equations

Mollified SHE

$$\begin{cases} \partial_t u^\varepsilon = \frac{1}{2} \Delta u^\varepsilon + \beta_\varepsilon u^\varepsilon \xi^\varepsilon \\ u^\varepsilon(0, \cdot) \equiv 1 \end{cases} \quad (\varepsilon\text{-SHE})$$

Set $h^\varepsilon(t, x) := \log u^\varepsilon(t, x)$ (Cole-Hopf). By Ito's formula

Mollified KPZ

$$\begin{cases} \partial_t h^\varepsilon = \frac{1}{2} \Delta h^\varepsilon + \frac{1}{2} |\nabla h^\varepsilon|^2 + \beta_\varepsilon \xi^\varepsilon - C_\varepsilon \\ h^\varepsilon(0, \cdot) \equiv 0 \end{cases} \quad (\varepsilon\text{-KPZ})$$

where $C_\varepsilon = \beta_\varepsilon^2 \varepsilon^{-d} \|\varrho\|_{L^2}^2$

Key problem

Can we choose $\beta_\varepsilon \in (0, \infty)$ so that
 $u^\varepsilon(t, x)$ and $h^\varepsilon(t, x)$ have a limit as $\varepsilon \downarrow 0$?
 YES (...)

$$\beta_\varepsilon = \begin{cases} \hat{\beta} \text{ (fixed)} & d = 1 \\ \hat{\beta} \frac{1}{\sqrt{\log \varepsilon^{-1}}} & d = 2 \\ \hat{\beta} \varepsilon^{\frac{d-2}{2}} & d \geq 3 \end{cases} \quad \hat{\beta} \in (0, \infty)$$

Note that $\beta_\varepsilon \rightarrow 0$ for $d = 2$ and $d \geq 3$

Main result I. Phase transition and LLN

Space dimension $d = 2$

$$\beta_\varepsilon = \frac{\hat{\beta}}{\sqrt{\log \varepsilon^{-1}}} \quad \hat{\beta} \in (0, \infty)$$

Theorem (SHE marginal laws)

[CSZ 17]

Phase transition with critical value $\hat{\beta}_c = \sqrt{2\pi}$ ("weak to strong disorder")

- ▶ Fix $t > 0$, $x \in \mathbb{R}^2$:

$$u^\varepsilon(t, x) \xrightarrow[\varepsilon \downarrow 0]{d} \begin{cases} \exp(\sigma_{\hat{\beta}} Z - \frac{1}{2} \sigma_{\hat{\beta}}^2) & \text{if } \hat{\beta} < \sqrt{2\pi} \\ 0 & \text{if } \hat{\beta} \geq \sqrt{2\pi} \end{cases}$$

where

$$Z \sim N(0, 1) \quad \sigma_{\hat{\beta}}^2 := \log \frac{2\pi}{2\pi - \hat{\beta}^2}$$

Sub-critical regime $\hat{\beta} < \sqrt{2\pi}$

► Fix $\hat{\beta} < \sqrt{2\pi}$. For distinct $x_1, \dots, x_n \in \mathbb{R}^2$

$u^\varepsilon(t, x_i)$ become asymptotically independent as $\varepsilon \downarrow 0$

Look at $u^\varepsilon(t, \cdot)$ as a random distribution on \mathbb{R}^2 $(\mathbb{E}[u^\varepsilon(t, x)] \equiv 1)$

Corollary (Law of large numbers)

$$u^\varepsilon(t, \cdot) \xrightarrow[\varepsilon \downarrow 0]{d} 1 \quad \text{as a distribution}$$

$$\forall \phi \in C_c(\mathbb{R}^2) : \quad \int_{\mathbb{R}^2} u^\varepsilon(t, x) \phi(x) dx \xrightarrow[\varepsilon \downarrow 0]{d} \int_{\mathbb{R}^2} \phi(x) dx$$

Analogous results for the KPZ solution $h^\varepsilon(t, x) = \log u^\varepsilon(t, x)$

Main result II. Fluctuations for SHE

Recall that $\beta_\varepsilon = \frac{\hat{\beta}}{\sqrt{\log \varepsilon^{-1}}}$ sub-critical $\hat{\beta} \in (0, \sqrt{2\pi})$

Rescaled SHE solution $\mathcal{U}^\varepsilon(t, x) := \frac{1}{\beta_\varepsilon}(u^\varepsilon(t, x) - 1)$

Theorem (EW fluctuations for SHE)

[CSZ 17]

$$\mathcal{U}^\varepsilon(t, \cdot) \xrightarrow[\varepsilon \downarrow 0]{d} v(t, \cdot) \quad \text{as a distribution}$$

v = solution of Edwards-Wilkinson equation

$$\partial_t v = \frac{1}{2} \Delta_x v + c \xi \quad \text{where} \quad c = \sqrt{\frac{2\pi}{2\pi - \hat{\beta}^2}} > 1 \quad (\text{EW})$$

$$\partial_t \mathcal{U}^\varepsilon = \frac{1}{2} \Delta_x \mathcal{U}^\varepsilon + (1 + \beta_\varepsilon \mathcal{U}^\varepsilon) \xi^\varepsilon$$

From SHE to KPZ?

Key tools

1. SHE solution $u^\varepsilon(t, x)$ has explicit Wiener-Chaos expansion
2. 4th Moment Theorem for Gaussianity [Nourdin, Peccati, Reinert]

Problem: Wiener Chaos not available for KPZ sol. $h^\varepsilon(t, x) = \log u^\varepsilon(t, x)$

Naive idea: Since $u^\varepsilon(t, x) \rightarrow 1$ (as a distribution), by Taylor expansion

$$h^\varepsilon(t, x) \approx (u^\varepsilon(t, x) - 1) ?$$

NO, because $u^\varepsilon(t, x)$ is **not close to 1 pointwise**

We can correct Taylor expansion and deduce fluctuations of h^ε from u^ε

Main result III. Sub-critical fluctuations for KPZ

Recall that $\beta_\varepsilon = \frac{\hat{\beta}}{\sqrt{\log \varepsilon^{-1}}}$ sub-critical $\hat{\beta} \in (0, \sqrt{2\pi})$

Rescaled KPZ solution $\mathcal{H}^\varepsilon(t, x) := \frac{1}{\beta_\varepsilon} (h^\varepsilon(t, x) - \mathbb{E}[h^\varepsilon])$

Theorem (EW fluctuations for KPZ)

[CSZ 18b]

$$\mathcal{H}^\varepsilon(t, \cdot) \xrightarrow[\varepsilon \downarrow 0]{d} v(t, \cdot) \quad \text{as a distribution}$$

v = solution of Edwards-Wilkinson equation

$$\partial_t v = \frac{1}{2} \Delta_x v + c \xi \quad \text{where} \quad c = \sqrt{\frac{2\pi}{2\pi - \hat{\beta}^2}} > 1 \quad (\text{EW})$$

Same as for SHE (but different centering)

Alternative proof by [Gu 18] (only for small $\hat{\beta}$)

A variation on KPZ

Last result was motivated by a paper of Chatterjee and Dunlap [CD 18] who considered a variation of KPZ

$$\partial_t \tilde{h}^\varepsilon = \frac{1}{2} \Delta_x \tilde{h}^\varepsilon + \frac{1}{2} \beta_\varepsilon |\nabla_x \tilde{h}^\varepsilon|^2 + \xi^\varepsilon$$

The same $\beta_\varepsilon = \frac{\hat{\beta}}{\sqrt{\log \varepsilon^{-1}}}$ now multiplies the **non-linearity** instead of ξ^ε

Theorem

[Chatterjee Dunlap 18]

For $\hat{\beta}$ sufficiently small, the centered solution $\tilde{h}^\varepsilon(t, \cdot) - \mathbb{E}[\tilde{h}^\varepsilon]$ admits subsequential limits in law as $\varepsilon \downarrow 0$ (as a random distribution on \mathbb{R}^2)

Any limit is **not** the solution of (EW) with $c = 1$
 (what one would get simply removing the non-linearity)

Relation with our results

Recall “our” KPZ : $\partial_t h^\varepsilon = \frac{1}{2} \Delta_x h^\varepsilon + \frac{1}{2} |\nabla_x h^\varepsilon|^2 + \beta_\varepsilon \xi^\varepsilon - C_\varepsilon$

Scaling relation

$$\tilde{h}^\varepsilon(t, x) - \mathbb{E}[\tilde{h}^\varepsilon] = \frac{1}{\beta_\varepsilon} (h^\varepsilon(t, x) - \mathbb{E}[h^\varepsilon]) = \mathcal{H}^\varepsilon(t, x)$$

Theorem

[CSZ 18b]

For every sub-critical $\hat{\beta} < 1$, the centered solution $\tilde{h}^\varepsilon(t, \cdot) - \mathbb{E}[\tilde{h}^\varepsilon]$ admits a unique limit in law as $\varepsilon \downarrow 0$ (as a random distribution on \mathbb{R}^2)

The limit is the solution of (EW) with $c = \frac{1}{\sqrt{1-\hat{\beta}^2}} > 1$

Grazie!

The critical regime

What about the critical point $\hat{\beta} = \sqrt{2\pi}$?

More generally, **critical window** [Bertini Cancrini 98]

$$\beta_\varepsilon = \sqrt{\frac{2\pi}{\log \varepsilon^{-1}} \left(1 + \frac{\vartheta}{\log \varepsilon^{-1}} \right)} \quad \text{with } \vartheta \in \mathbb{R}$$

Nothing is known for KPZ $h^\varepsilon(t, x)$, some progress for SHE $u^\varepsilon(t, x)$

Key conjecture

$u^\varepsilon(t, \cdot)$ has a limit $\mathcal{U}(t, \cdot)$ for $\varepsilon \downarrow 0$, as a random distribution on \mathbb{R}^2

$$\langle u^\varepsilon(t, \cdot), \phi \rangle := \int_{\mathbb{R}^2} u^\varepsilon(t, x) \phi(x) dx \xrightarrow[\varepsilon \downarrow 0]{d} \int_{\mathbb{R}^2} \mathcal{U}(t, x) \phi(x) dx$$

(actually a random measure, since $u^\varepsilon \geq 0$)

Second moment in the critical window

What is known

[Bertini Cancrini 98]

Tightness via second moment bounds

$$\mathbb{E}[\langle u^\varepsilon(t, \cdot), \phi \rangle] \equiv \langle 1, \phi \rangle \quad \sup_{\varepsilon > 0} \mathbb{E}[\langle u^\varepsilon(t, \cdot), \phi \rangle^2] < \infty$$

More precisely $\mathbb{E}[\langle u^\varepsilon(t, \cdot), \phi \rangle^2] \xrightarrow{\varepsilon \downarrow 0} \langle \phi, K\phi \rangle < \infty$

Explicit kernel $K(x, x') \sim C \log \frac{1}{|x - x'|}$ as $|x - x'| \rightarrow 0$

Corollary

\exists subsequential limits $\langle u^{\varepsilon_k}(t, \cdot), \phi \rangle \xrightarrow[k \rightarrow \infty]{d} \langle \mathcal{U}, \phi \rangle$

Can the limit be trivial $\mathcal{U}(t, \cdot) \equiv 1$?

Main result III. Third moment in the critical window

We compute the sharp asymptotics of third moment

Theorem

[CSZ 18a]

$$\lim_{\varepsilon \downarrow 0} \mathbb{E}[\langle u^\varepsilon(t, \cdot), \phi \rangle^3] = C(\phi) < \infty$$

Corollary

Any subsequential limit $u^{\varepsilon_k}(t, \cdot) \xrightarrow{d} \mathcal{U}(t, \cdot)$ has covariance $\mathbf{K}(x, x')$
 $\rightsquigarrow \mathcal{U}(t, \cdot) \not\equiv 1$ is non-degenerate!

Very recently, [Gu Quastel Tsai 19] improved our result showing that
 all moments converge to a finite limit (not only 3rd moment)