

# The weak coupling limit of disordered copolymer models

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Joint work with Giambattista Giacomin (Université Paris Diderot)

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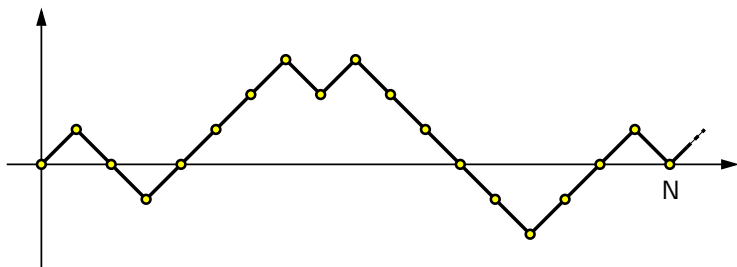
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1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

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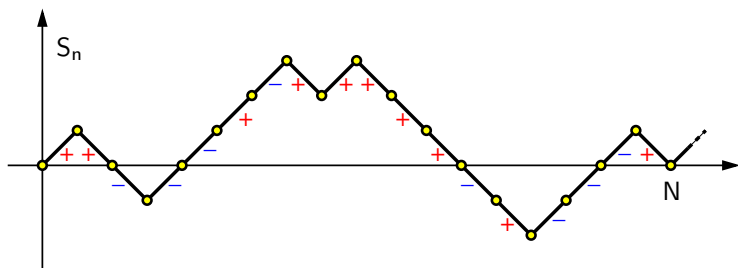
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# A random walk with a random potential



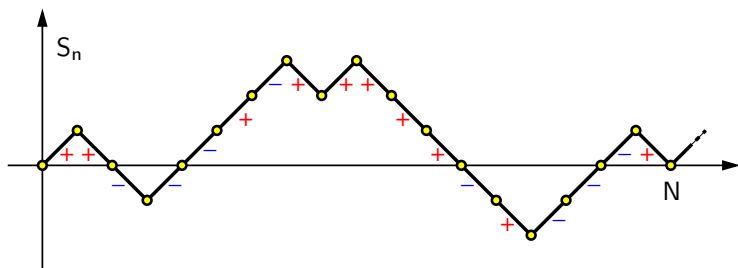
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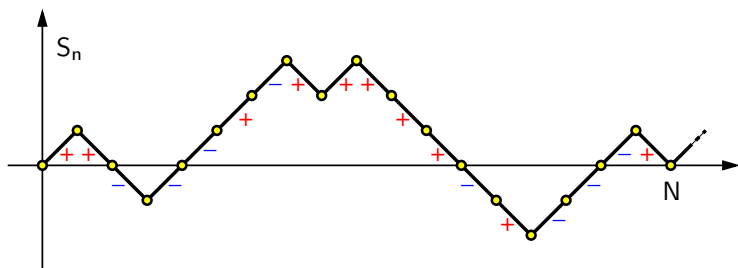
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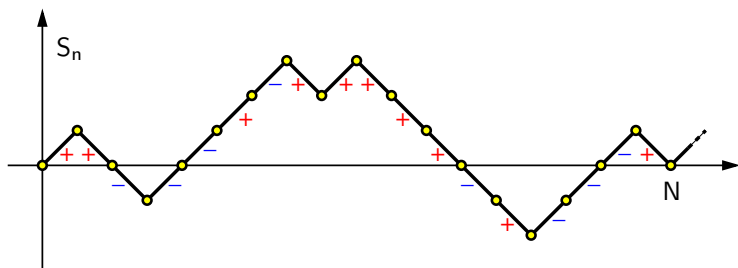


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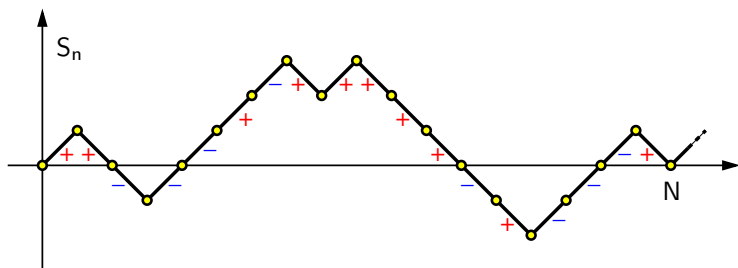
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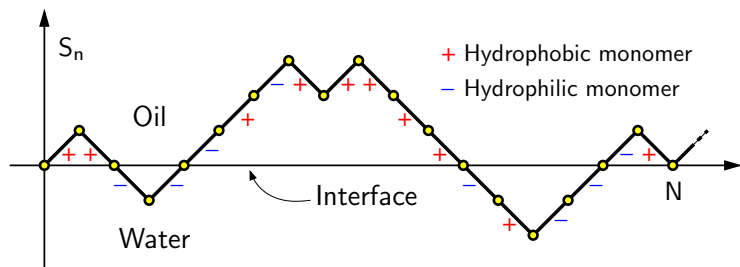
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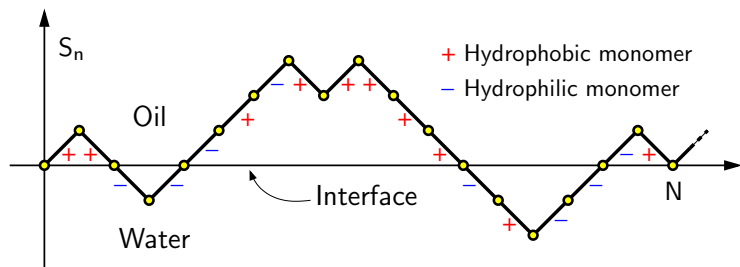
**Localization** or **Delocalization**?

# A polymer model interpretation



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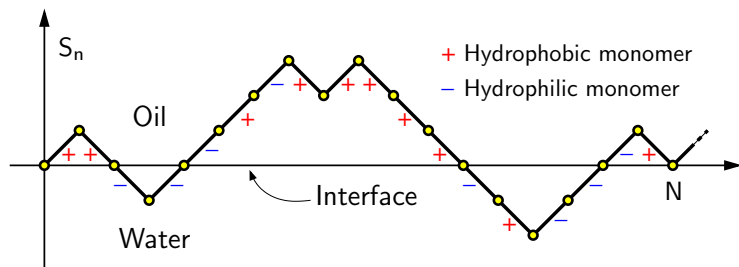
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- ▶ M. Biskup, G. Giacomin, T. Bodineau, F.L. Toninelli, F.C., M. Gubinelli, N. Pétrélis, L. Zambotti, B. de Tilière, H. Lacoin, D. Cheliotis, ...

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# The basic copolymer model

Definition of the model:  $\frac{d\mathbf{P}_{N,\omega}}{d\mathbf{P}}(S) := \frac{1}{Z_{N,\omega}} \exp(-H_{N,\omega}(S))$

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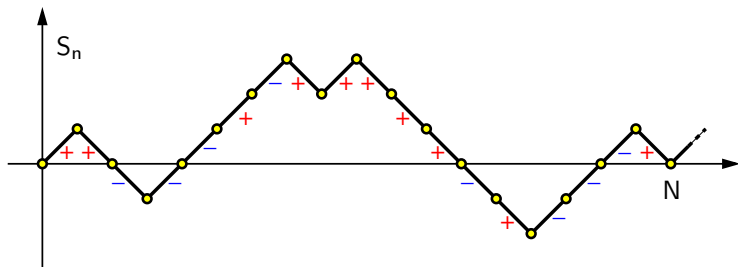
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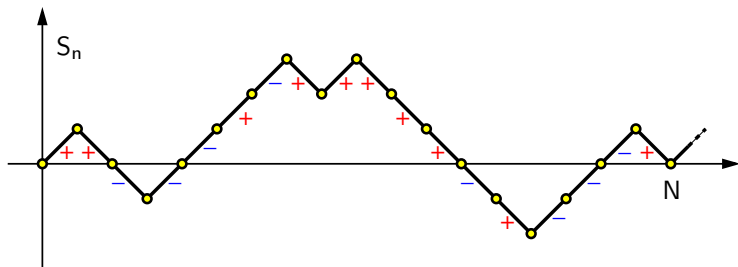
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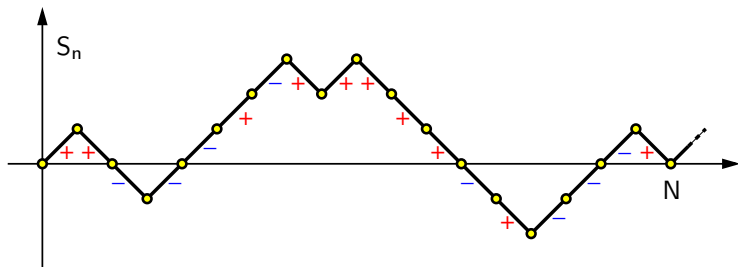
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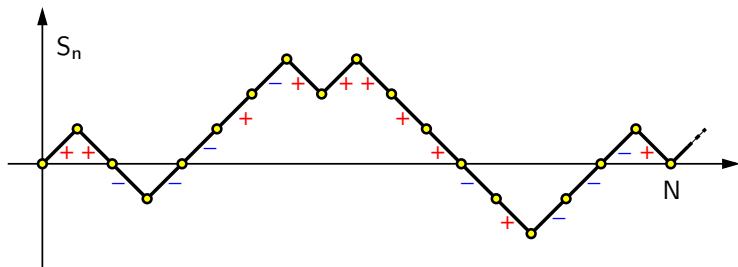
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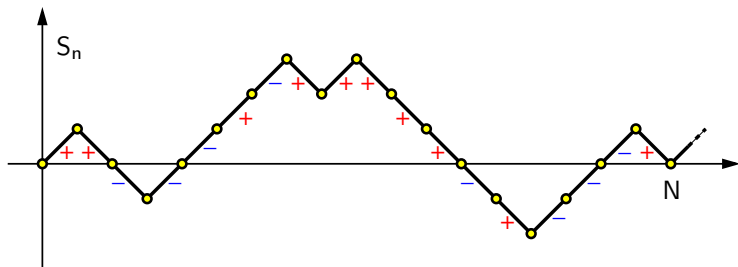


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(This definition does correspond to sharply different path behaviors)

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## Theorem

The regions  $\mathcal{L}$  and  $\mathcal{D}$  are separated by a strictly increasing, continuous *critical line*  $h_c(\cdot)$ :

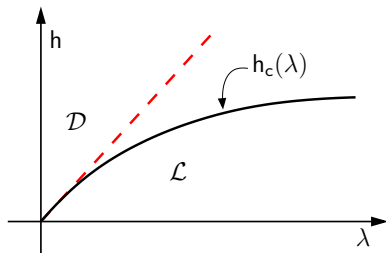
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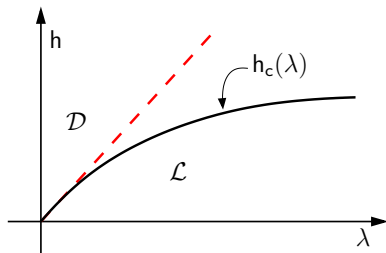


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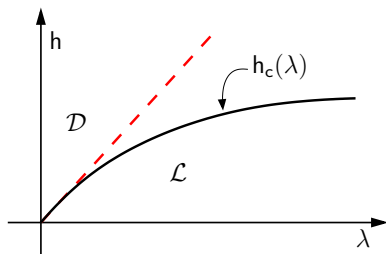
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$$\underline{h}'(0+) = \frac{2}{3}, \quad \bar{h}'(0+) = 1 - \epsilon.$$

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The **continuum free energy**  $\tilde{F}(\lambda, h)$  is defined analogously:

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By Brownian scaling  $\tilde{F}(a\lambda, ah) = a^2 \tilde{F}(\lambda, h)$  for all  $a, \lambda, h \geq 0$ .

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Again  $\tilde{F}(\lambda, h) \geq 0$ . We then define  **$\tilde{\mathcal{L}}$ ocalization** and  **$\tilde{\mathcal{D}}$ elocalization**:

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By Brownian scaling  $\tilde{F}(a\lambda, ah) = a^2 \tilde{F}(\lambda, h)$  for all  $a, \lambda, h \geq 0$ .

Therefore  $\tilde{h}_c(\cdot)$  is a straight line:  $\tilde{h}_c(\lambda) = \tilde{m} \lambda$ .

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## Theorem ([BdH 97])

For all  $\lambda, h \geq 0$

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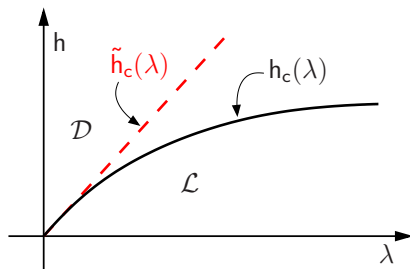
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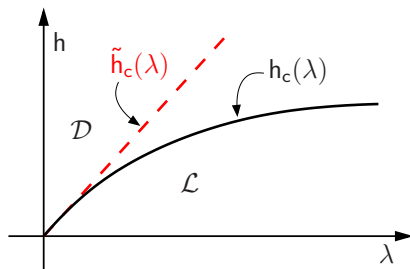
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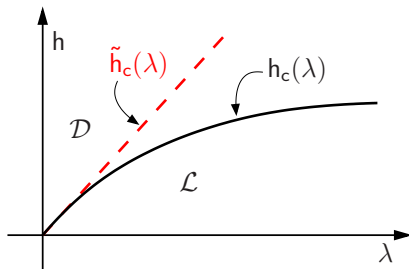
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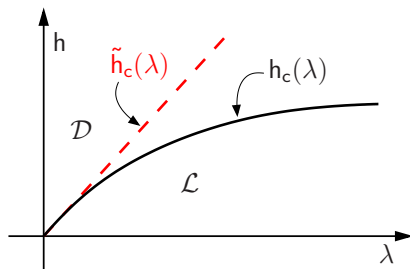
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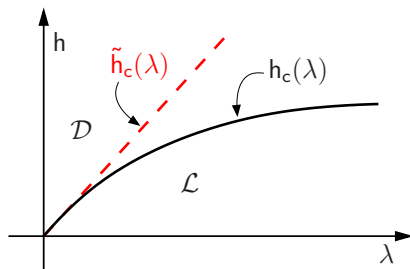
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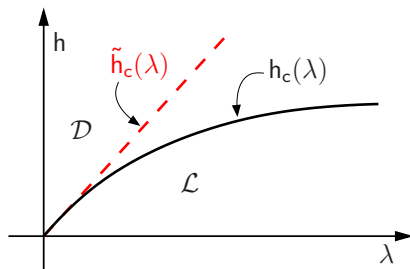
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... unfortunately for just **one** discrete model. Generalization?

# Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

# Beyond the simple random walk

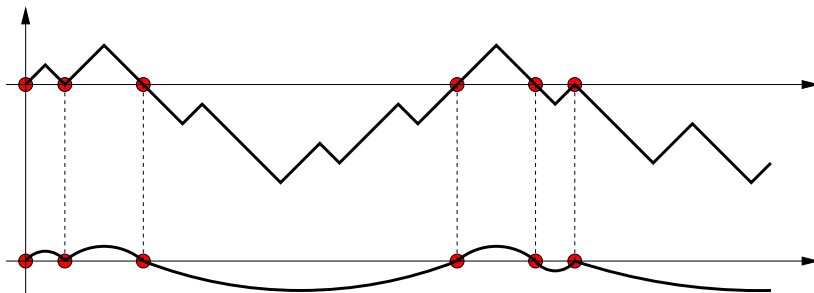
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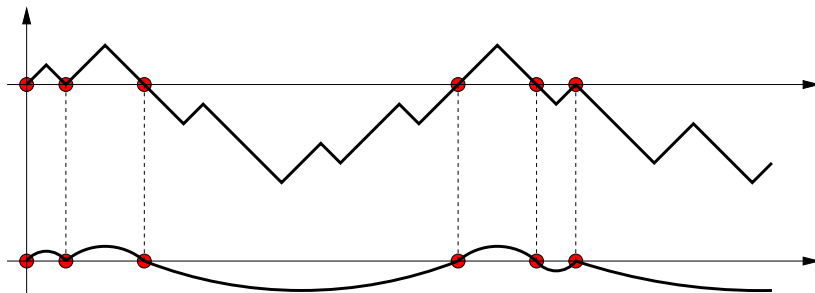
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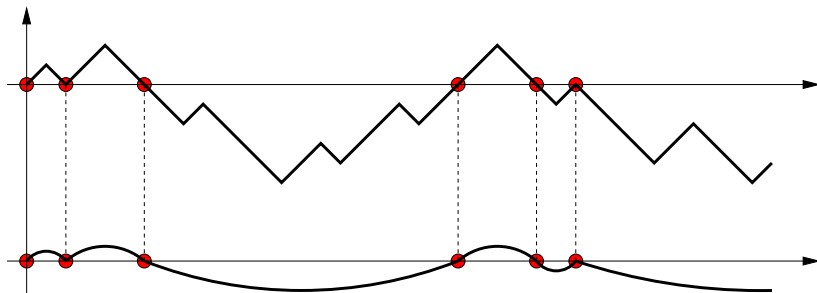
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- ▶ **Excursions signs**: fair coin tossing (independent of  $\{\tau_k\}_{k \geq 0}$ )

# Generalized discrete $\alpha$ -copolymer models

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Discrete Bessel-like process ( $c_\alpha = 1/2 - \alpha$ )

$$\mathbf{P}(S_{n+1} = x \pm 1 | S_n = x) = \frac{1}{2} \left( 1 \pm \frac{c_\alpha}{x} + o\left(\frac{1}{x}\right) \right) \text{ yields } (\star) \text{ asymp.}$$

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Note that  $F(\cdot, \cdot)$  and  $h_c(\cdot)$  **do depend** on the choice of  $\mathbf{P}$  and  $\mathbb{P}$

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From  $\tilde{\tau}^\alpha$  we obtain  $(\tilde{\Delta}^\alpha = \{\tilde{\Delta}_t^\alpha\}_{t \geq 0}, \tilde{\mathbf{P}})$  (For  $\alpha = \frac{1}{2}$  we recover BM)

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- Non-trivial, highly technical proof (also for  $\alpha = \frac{1}{2}$ )



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$\tilde{F}^\alpha(\lambda, h)$  exists and is self-averaging

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# The continuum $\alpha$ -copolymer model: free energy

The continuum charges  $(\{d\beta_s\}_{s \geq 0}, \tilde{\mathbb{P}})$  are always white noise

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- ▶ Continuity theory of Gaussian processes plays essential role

# The continuum $\alpha$ -copolymer model: scaling limit

Scaling properties of  $\beta_s$  and  $\tilde{\Delta}_s^\alpha \longrightarrow \tilde{F}^\alpha(a\lambda, ah) = a^2 \tilde{F}^\alpha(\lambda, h)$ .

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$$\lim_{a \downarrow 0} \frac{F(a\lambda, ah)}{a^2} = \tilde{F}^\alpha(\lambda, h) \qquad \lim_{\lambda \downarrow 0} \frac{h_c(\lambda)}{\lambda} = \tilde{m}^\alpha$$

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Theorem

For all  $\lambda, h > 0$  and  $\epsilon \in (0, 1)$  there exists  $a_0 > 0$  s.t. for all  $a < a_0$

$$\tilde{F}^\alpha\left(\frac{\lambda}{1+\epsilon}, \frac{h}{1-\epsilon}\right) \leq \frac{F(a\lambda, ah)}{a^2} \leq \tilde{F}^\alpha((1+\epsilon)\lambda, (1-\epsilon)h)$$

# Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof



# Strategy of the proof

**Goal:**  $\forall \lambda, h > 0, \epsilon \in (0, 1)$  one has for  $a \ll 1$

$$\frac{1}{a^2} F(a\lambda, ah) \leq \frac{1}{a^2} \tilde{F}((1 + \epsilon)a\lambda, (1 - \epsilon)ah)$$

(and viceversa, with  $F \leftrightarrow \tilde{F}$ ).

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# Heuristics

Recall that  $\Delta_n = \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$  and

$$-H_N^0(a\lambda, ah) = -2a\lambda \sum_{n=1}^N (\omega_n + ah) \Delta_n$$

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Recall that  $\Delta_n = \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$  and

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Since  $a \ll 1$ , for  $H_N(a\lambda, ah) \approx 1$  we need  $N \approx t/a^2$  steps, therefore

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We need to show that  $\approx$  can be made  $\asymp$ .

# The proof

## Step 1: Coarse-graining of the renewal process.

Recall that  $\Delta_n = \mathbf{1}_{\{(S_{n-1}, S_n) < 0\}}$  and

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where in  $\Delta_n^{\eta, \delta}$  we replace the microscopic return times  $\tau_n$  by **coarse-grained** return times on blocks of size  $\eta/a^2$ , skipping  $\delta/\eta \gg 1$  blocks between consecutive coarse-grained returns.

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Showing that  $H_N^0 \asymp H_N^1$  is delicate and very technical.

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## Step 2: From discrete charges to the white noise.

$H_N^2$  is obtained from  $H_N^1$  by replacing the charges  $\omega_n$  by i.i.d.  $N(0, 1)$  (discrete white noise).

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$H_N^4 = \tilde{H}_N$  is obtained from  $H_N^3$  by replacing  $\tilde{\Delta}_t^{\eta, \delta}$  by the original (non coarse-grained) continuous-time process  $\tilde{\Delta}_t$ .

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This step is analogous to step 1.

Thank you.