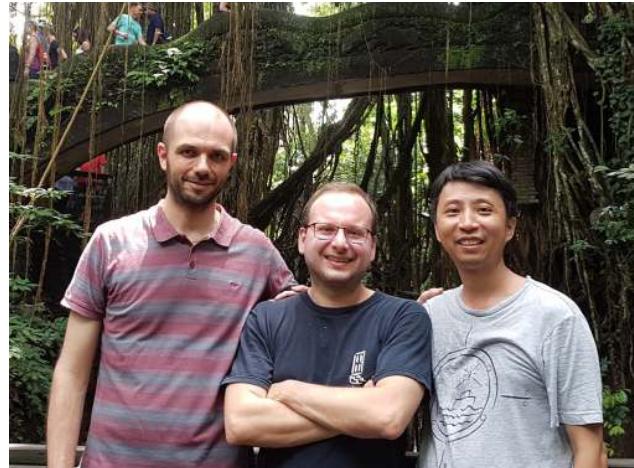


# The critical 2d Stochastic Heat Flow (is not a G.M.C.)

Francesco Caravenna  
Università di Milano-Bicocca

CIRM, Marseille ~ 2 MAY 2022



Based on joint works with  
Nikos Zygouras and Rongfeng Sun

## OVERVIEW

### I. THE CRITICAL 2D STOCHASTIC HEAT FLOW

Motivations, construction, Key features

### II. MAIN RESULTS

Gaussian Multiplicative Chaos, moment bounds

### III. IDEAS FROM THE PROOF

## REFERENCES

- [CSZ 21] F. Caravenna, R. Sun, N. Zygouras  
THE CRITICAL 2D STOCHASTIC HEAT FLOW  
arXiv (2021)
- [CSZ 22] F. Caravenna, R. Sun, N. Zygouras  
... IS NOT A GAUSSIAN MULTIPLICATIVE CHAOS  
In preparation

## 1. WHAT IS THE STOCHASTIC HEAT FLOW ?

It is a "universal" stochastic process of measures on  $\mathbb{R}^2$

$$\mathcal{L}_{s,t}^\vartheta(dy, dx) \quad 0 \leq s < t < \infty \quad x, y \in \mathbb{R}^2 \quad \vartheta \in \mathbb{R}$$

Heuristics: "natural candidate" solution of the (ill-defined)  
"critical" 2d Stochastic Heat Equation:

$$\partial_t u(t,x) = \frac{1}{4} \Delta u(t,x) + \beta \underbrace{\xi(t,x)}_{\rightarrow \text{SPACE-TIME WHITE NOISE}} u(t,x) \quad (\text{SHE})$$

Fix starting time  $s$ , initial condition  $u(s,x) = \psi(x)$ :

$$\rightsquigarrow \int_{\mathbb{R}^2} \psi(y) \mathcal{L}_{s,t}^\vartheta(dy, \cdot) = \text{"natural candidate" for } u(t, \cdot)$$

( $u(t, \cdot)$  expected not to be a function, but a measure on  $\mathbb{R}^2$ )

What do we mean by "natural candidate"? We will see by  
REGULARIZATION of the eq. & RENORMALIZATION of  $\beta \leftrightarrow \vartheta$

Henceforth we fix  $s=0$ ,  $\psi \equiv 1$  and call

$$\int_{y \in \mathbb{R}^2} \mathcal{L}_{s,t}^\vartheta(dy, dx) = \mathcal{L}_t^\vartheta(dx) = \text{STOCHASTIC HEAT FLOW (SHF)}$$

## 2. A LINK WITH DIRECTED POLYMERS

We can **REGULARIZE** the Stochastic Heat Equation by  
**discretizing time and space** (or by mollification  $\rightarrow$  later)

$$\underbrace{\partial_t^N u(t,x)}_{N \cdot \{u(t,x) - u(t-\frac{1}{N},x)\}} = \frac{1}{4} \underbrace{\Delta^N u(t,x)}_{\text{SCALED LATTICE LAPLACIAN}} + \beta N \underbrace{\langle \eta(t,x) u(t,x) \rangle}_{\text{SPACE AVERAGE } \frac{1}{4} \sum_{x' \sim x}}$$

I.I.D. ZERO MEAN, UNIT VARIANCE

$$\text{For } (t,x) = \left( \frac{n}{N}, \frac{z}{\sqrt{N}} \right) \in \frac{N}{N} \times \frac{\mathbb{Z}^2}{\sqrt{N}}$$

$$\text{Special choice of noise: } \eta(t,x) = \frac{e^{\beta \omega(n,z) - \frac{\beta^2}{2}} - 1}{\beta}$$

I.I.D.  $N(0,1)$

$$u(t,x) = Z_N(n,z) = E \left[ e^{\sum_{i=n+1}^N \beta \omega(i,S_i) - \frac{\beta^2}{2}} \mid S_n = z \right]$$

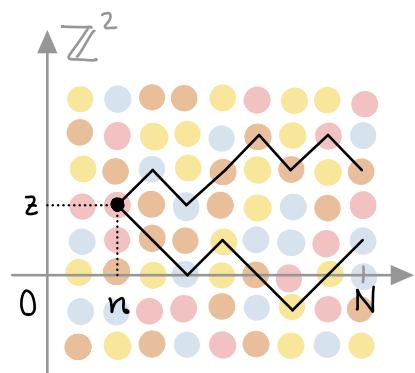
actually  $1-t$

$\uparrow$

$(S_i)$  SIMPLE RANDOM WALK ON  $\mathbb{Z}^2$

Partition function of the

DIRECTED POLYMER  
IN RANDOM ENVIRONMENT



### 3. MAIN RESULTS: EXISTENCE OF THE STOCHASTIC HEAT FLOW

The solution of the **discretized** Stochastic Heat Equation is

$$U(t, x) = \sum_N (\lfloor tN \rfloor, \lfloor x\sqrt{N} \rfloor) \quad (t, x) = \left( \frac{u}{N}, \frac{z}{\sqrt{N}} \right)$$

INTEGER PARTS

A "natural candidate" solution of the original equation (ill-defined) is obtained **removing the discretization**: do we have

$$\int \Psi(x) \sum_N (\lfloor tN \rfloor, \lfloor x\sqrt{N} \rfloor) dx \xrightarrow[N \rightarrow \infty]{d} ?$$

Not yet! As we remove the discretization, we need to **RENORMALIZE** the noise strength  $\beta$  (intermediate disorder regime):

$$\beta \sim \hat{\beta} \frac{\sqrt{\pi}}{\sqrt{\log N}} \quad \text{for fixed } 0 < \hat{\beta} < \infty$$

Indeed, for  $x \neq y \in \mathbb{R}^2$

$$\mathbb{E} [\sum_N (\lfloor tN \rfloor, \lfloor x\sqrt{N} \rfloor), \sum_N (\lfloor tN \rfloor, \lfloor y\sqrt{N} \rfloor)] \xrightarrow[N \rightarrow \infty]{} \begin{cases} 0 & \text{if } \hat{\beta} < 1 \\ K(x, y) \in (0, \infty) & \text{if } \hat{\beta} = 1 \\ \infty & \text{if } \hat{\beta} > 1 \end{cases}$$

$$(\mathbb{E} [\sum_N (\lfloor tN \rfloor, \lfloor x\sqrt{N} \rfloor)] \equiv 1)$$

[Bertini-Cancrini '98]

The 2d directed polymer model undergoes a phase transition from weak to strong disorder. (see [CSZ, AAP 2017])

We fix the critical point  $\hat{\beta} = 1$ , in fact a critical window:

$$\text{⑥} \quad \beta = \left( 1 + \frac{g + o(1)}{\log N} \right) \frac{\sqrt{\pi}}{\sqrt{\log N}} \quad \text{for } -\infty < g < +\infty$$

Our first main result:

Theorem 1. [CSZ 21]

Fix  $g \in \mathbb{R}$ . As  $N \rightarrow \infty$  we have the convergence in f.d.d. of

$$(Z_N(tN, x\sqrt{N}) dx)_{t \in [0, 1]}$$

(as random measures on  $\mathbb{R}^2$ ) to a unique limit

$$(\mathcal{L}_t(dx))_{t \in [0, 1]}$$

which we call the CRITICAL 2D STOCHASTIC HEAT FLOW.

The SHF encodes asymptotic properties of directed polymers, and it gives a meaning to the solution of the Stochastic Heat Equation.

The rest of the talk is devoted to investigating the SHF.

## 4. GAUSSIAN MULTIPLICATIVE CHAOS

Let  $X = (X(x))_{x \in \mathbb{R}^2}$  be a (generalized) Gaussian field on  $\mathbb{R}^2$ .  
 Fix a reference measure  $\mu(dx)$  on  $\mathbb{R}^2$ .

Gaussian Multiplicative Chaos (GMC) is a random measure on  $\mathbb{R}^2$ :

$$\text{formally } \mathcal{M}(dx) = e^{X(x) - \frac{1}{2} \text{Var}[X(x)]} \mu(dx)$$

Often  $X$  is specified by (cf. GFF)

$$\text{Cov}[X(x), X(y)] =: \kappa(x, y) \rightarrow \infty \quad \text{as } |y-x| \rightarrow 0$$

$X(x)$  is not defined pointwise:  $X(\varphi) = \int \varphi(x) X(x) dx \sim N(0, \sigma_\varphi^2)$

$$\sigma_\varphi^2 = \text{Var}[X(\varphi)] = \iint \varphi(x) \varphi(y) \kappa(x, y) dx dy < \infty$$

Approximate  $\kappa(x, y) = \lim_{\varepsilon \downarrow 0} \kappa_\varepsilon(x, y)$  by smooth covariance functions.

Then  $X_\varepsilon \sim N(0, \kappa_\varepsilon)$  is defined pointwise,  $X_\varepsilon(x)$ , and we set

$$\mathcal{M}(dx) := \lim_{\varepsilon \downarrow 0} e^{X_\varepsilon(x) - \frac{1}{2} \text{Var}[X_\varepsilon(x)]} \mu(dx)$$

## 5. MAIN RESULTS: SHF & GMC

SUB-CRITICAL REGIME

For  $\hat{\beta} < 1$ , it is known that  $\log Z_N$  has Gaussian fluctuations:

$$\int \Psi(x) \sqrt{\log N} \left( \log Z_N(tN, x\sqrt{N}) - \mathbb{E}[\log Z_N] \right) dx \xrightarrow[N \rightarrow \infty]{d} \text{GAUSSIAN}$$

EDWARDS-WILKINSON

This can be roughly rephrased as:

$$Z_N(tN, x\sqrt{N}) \approx \exp(\text{GAUSSIAN})$$

Does this picture apply at criticality, i.e. for  $\hat{\beta} = 1$ ?

We can state our second main result:

Theorem 2. [CSZ 22]

The critical 2d SHF is **not** a GMC.

KPZ EQUATION

This suggests that  $\log Z_N(tN, x\sqrt{N})$  is **not** asymptotically Gaussian at criticality  $\hat{\beta} = 1$ .

Also: the SHF  $\mathcal{L}_t^\psi(dx)$  is a new class of random measures on  $\mathbb{R}^2$ .

## 6. MAIN RESULTS: MOMENT BOUNDS

Consider a GMC  $\mathcal{M}(dx) = e^{X(x) - \frac{1}{2} \text{VAR}[X(x)]} \mu(dx)$ .

We can compute the moments

$$\mathbb{E}[\mathcal{M}(dx)] = \mu(dx)$$

$$\mathbb{E}[\mathcal{M}(dx) \mathcal{M}(dy)] = e^{\kappa(x,y)} \mu(dx) \mu(dy)$$

Henceforth we fix  $t > 0$ ,  $\theta \in \mathbb{R}_+$ . For the SHF  $\mathcal{Z}_t^\theta(dx)$  we have

$$\mathbb{E}[\mathcal{Z}_t^\theta(dx)] = dx$$

$$\mathbb{E}[\mathcal{Z}_t^\theta(dx) \mathcal{Z}_t^\theta(dy)] = \underbrace{\kappa(x,y)}_{\sim C \log \frac{1}{|y-x|}} dx dy$$

$$\lim_{N \rightarrow \infty} \mathbb{E}[Z_N(tN, x\sqrt{N}), Z_N(tN, y\sqrt{N})] \sim C \log \frac{1}{|y-x|} \text{ as } |y-x| \rightarrow 0$$

We can match 1<sup>st</sup> & 2<sup>nd</sup> moments of  $\mathcal{M}$  and  $\mathcal{Z}_t^\theta$  by choosing

$$\mu(dx) = dx \quad \kappa(x,y) = \log \kappa(x,y) \sim \log \log \frac{1}{|y-x|}$$

The GMC  $\mathcal{M}$  is now completely determined.

$$\text{Set } \mathcal{L}_t^\varphi(\varphi) := \int \varphi(x) \mathcal{L}_t^\varphi(dx), \quad \mathcal{M}(\varphi) := \int \varphi(x) \mathcal{M}(dx)$$

We can prove that  $\mathcal{L}_t^\varphi \neq \mathcal{M}$  by comparing third moments.

Theorem 3. (3<sup>rd</sup> moment bound)

[CSZ 22]

If  $\varphi$  is the indicator of a ball, or the heat Kernel, then

$$\mathbb{E}[\mathcal{L}_t^\varphi(\varphi)^3] > \mathbb{E}[\mathcal{M}(\varphi)^3].$$

For the heat Kernel  $g_\delta(x) := \frac{e^{-\frac{|x|^2}{2\delta}}}{2\pi\delta}$  we bound higher moments.

Theorem 4. (higher moments bound)

[CSZ 22]

There is  $\gamma = \gamma_{t,\alpha} > 0$  such that  $\forall h \in \mathbb{N}$

$$\mathbb{E}[\mathcal{L}_t^\varphi(g_\delta)^h] \geq (1+\gamma) \mathbb{E}[\mathcal{L}_t^\varphi(g_\delta)^2]^{\binom{h}{2}} \quad \forall \delta \in (0,1),$$

while  $\mathbb{E}[\mathcal{M}(g_\delta)^h] \sim \mathbb{E}[\mathcal{M}(g_\delta)^2]^{\binom{h}{2}}$  as  $\delta \downarrow 0$ .

Moments upper bounds for  $\beta < 1$  by [Cosco, Zeitouni], [Lygkoni, Zygouras]

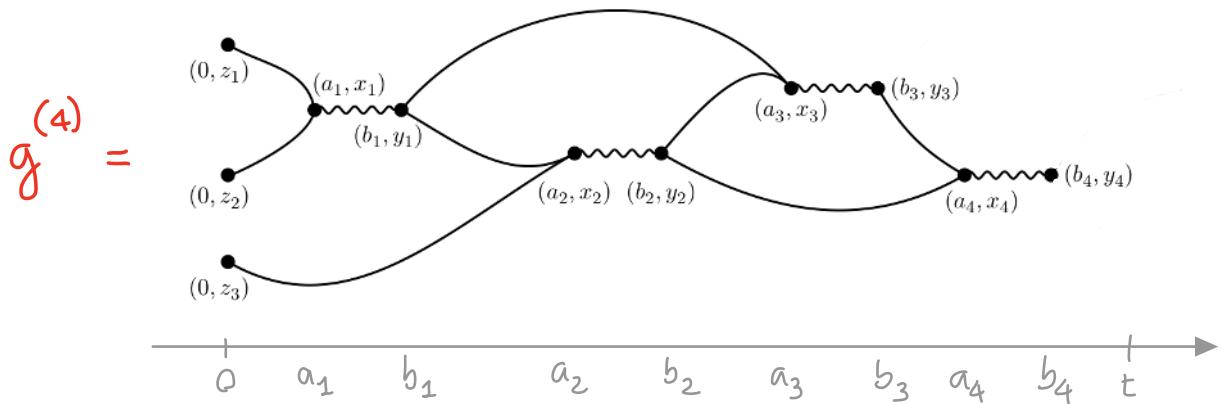
## 7. IDEAS FROM THE PROOF (I): MOMENT FORMULAS

We prove Theorem 3 exploiting an exact formula for the third moment, derived in [CSZ 19]:

$$\mathbb{E} \left[ \mathcal{L}_t^{\vartheta} (dx) \mathcal{L}_t^{\vartheta} (dy) \mathcal{L}_t^{\vartheta} (dz) \right] = K^{(3)}(x, y, z) dx dy dz$$

$$K^{(3)}(z_1, z_2, z_3) = \sum_{m \geq 2} 3 \cdot 2^{m-1} (2\pi)^m \int \dots \int d\vec{a} d\vec{b} d\vec{x} d\vec{y} g_{\vec{z}}^{(m)}(\vec{a}, \vec{b}, \vec{x}, \vec{y})$$

$0 < a_1 < b_1 < \dots < a_m < b_m < t$   
 $x_1, y_1, \dots, x_m, y_m \in \mathbb{R}^2$



where:

$$g_{\frac{a-b}{2}}(x-y)$$

$$g_{\frac{b-a}{4}}(y-x) \cdot G_g(b-a)$$

By an explicit Gaussian integration, we can integrate out all space variables  $\vec{x}, \vec{y}$  to get the equality

$$\int \dots \int g_{\vec{z}}^{(m)}(\vec{a}, \vec{b}, \vec{x}, \vec{y}) d\vec{x} d\vec{y} = \prod_{i=1}^m G_g(b_i - a_i).$$

$$\cdot g_{\bar{a}_1}(z_1 - z_2) \cdot g_{\bar{a}_2}\left(z_3 - \frac{z_1 + z_2}{2}\right) \cdot \prod_{i=3}^m g_{\overline{a_i - b_{i-2}}}^{(0)}$$

For suitable  $\bar{a}_2 < a_2$  and  $\overline{a_i - b_{i-2}} < a_i - b_{i-2}$ .

Exploiting monotonicity of the heat Kernel, we replace the quantities  $\bar{a}_2$  by  $a_2$  and  $\overline{a_i - b_{i-2}}$  by  $a_i - b_{i-2}$ .

This yields the bound  $\mathbb{E}[\mathcal{L}_t^\vartheta(\varphi)^3] > \mathbb{E}[\mathcal{M}(\varphi)^3]$ .  $\square$

Remark. Exact formulas are available for higher moments:

$$\mathbb{E}\left[\prod_{i=1}^h \mathcal{L}_t^\vartheta(dz_i)\right] = K^{(h)}(z_1, \dots, z_h) dz_1 \dots dz_h$$

$$K^{(h)}(z_1, \dots, z_h) = \sum_m \text{EXPLICIT (COMPLICATED) INTEGRALS ON } (\text{TIME} \times \text{SPACE}^h)^{2m}$$

This was proved by [Gu, Quastel, Tsai, PMP'21]

## 8. IDEAS FROM THE PROOF (II): CORRELATION INEQUALITY

Recall that

$$Z_N(\phi, z) = E \left[ e^{\sum_{l=1}^N \beta \omega(l, S_l) - \frac{\beta^2}{2}} \mid S_0 = z \right]$$

Then

$$\star \quad E \left[ \prod_{i=1}^h Z_N(z_i) \right] = E \left[ \prod_{1 \leq i < j \leq h} e^{\beta^2 L_N^{(i,j)}} \mid S_0^{(i)} = z_i \forall i \right]$$

$$L_N^{(i,j)} = \sum_{l=1}^N \mathbb{1}_{\{S_l^{(i)} = S_l^{(j)}\}} \quad \text{OVERLAP OF } S^{(i)} \text{ AND } S^{(j)}$$

Our goal is  $E[\mathcal{L}_t^\phi(g_\delta)^h] \geq (1+\gamma) E[\mathcal{L}_t^\phi(g_\delta)^2]^{\binom{h}{2}}$

Set  $\gamma = 0$  (for simplicity) and replace  $\mathcal{L}_t^\phi$  by  $Z_N$ .

We need to exchange  $E$  and  $\prod_{1 \leq i < j \leq h}$  in  $\star$ :

are  $\exp(\beta^2 L_N^{(i,j)})$  positively correlated? (at least as  $N \rightarrow \infty$ )

We prove that this holds in a continuum setting, replacing random walks  $S^{(i)}$   $\rightsquigarrow$  Brownian motions  $B^{(i)}$

Based on the GAUSSIAN CORRELATION INEQUALITY [Rogers '14].

## CONCLUSIONS

We introduced the CRITICAL 2D STOCHASTIC HEAT FLOW as a scaling limit of directed polymer partition functions

It is a universal process of random measures on  $\mathbb{R}^2$ , different from GMC, which is a natural candidate for the solution of the critical 2d Stochastic Heat Equation.

It has some explicit (but challenging!) features.

Many interesting questions are open:

- SINGULARITY W.R.T. LEBESGUE MEASURE
- FLOW PROPERTY
- CHARACTERIZING PROPERTIES
- TAKING LOG  $\rightsquigarrow$  KPZ

Thanks!