

The weak coupling limit of disordered copolymer models

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Joint work with Giambattista Giacomin (Université Paris Diderot)

Erwin Bolthausen's birthday conference

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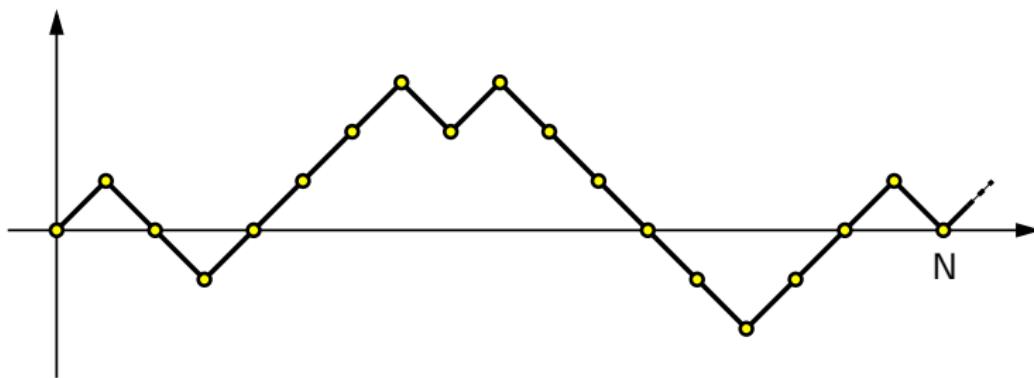
Outline

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2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

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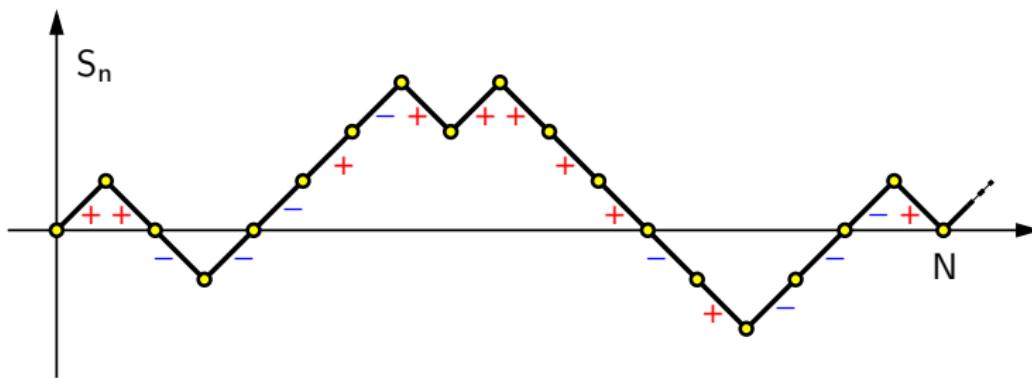
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A random walk with a random potential



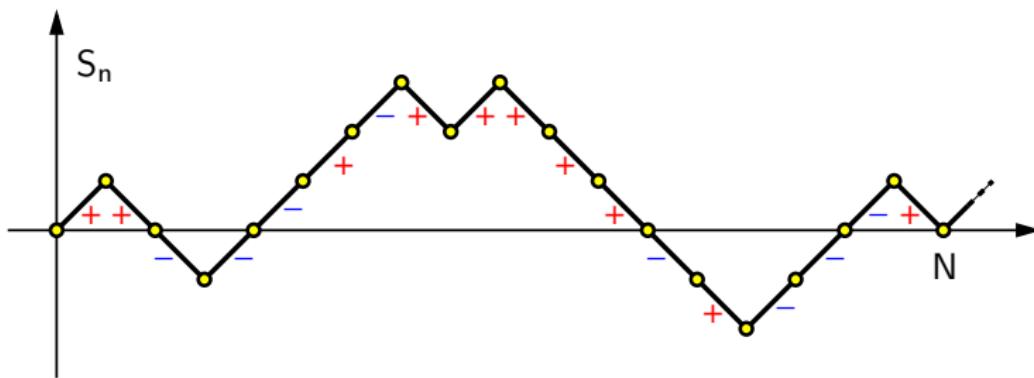
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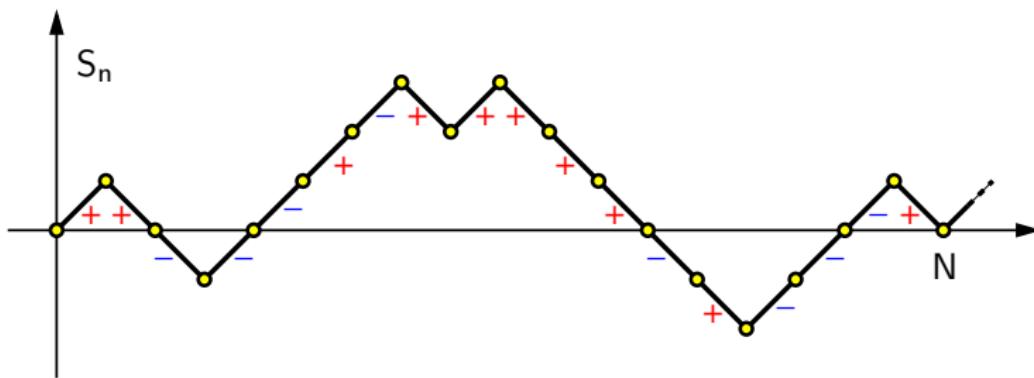
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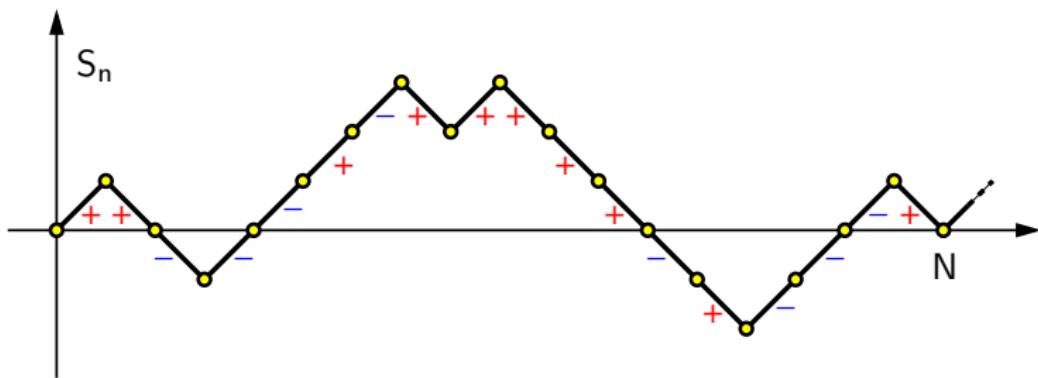


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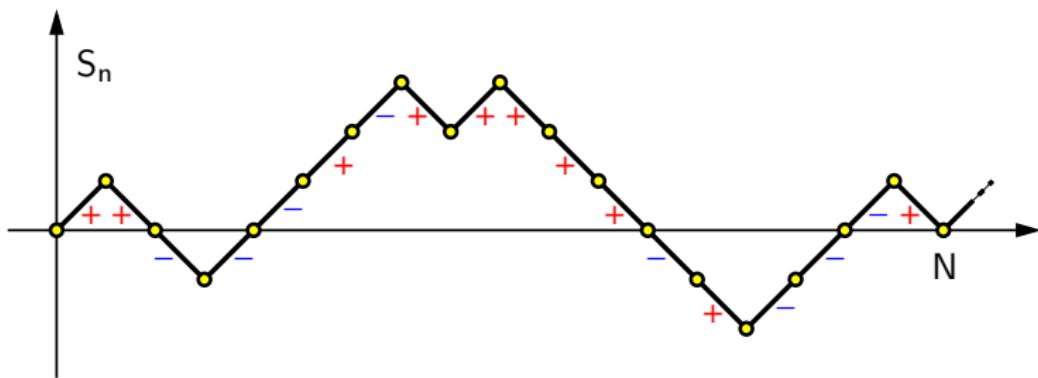


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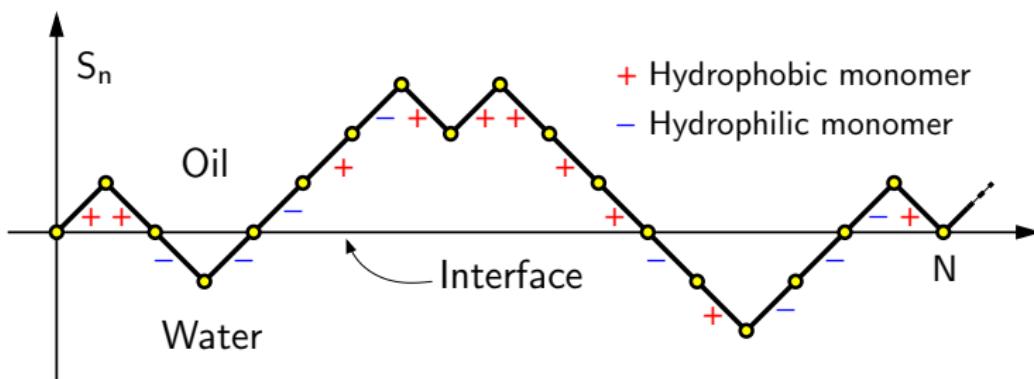
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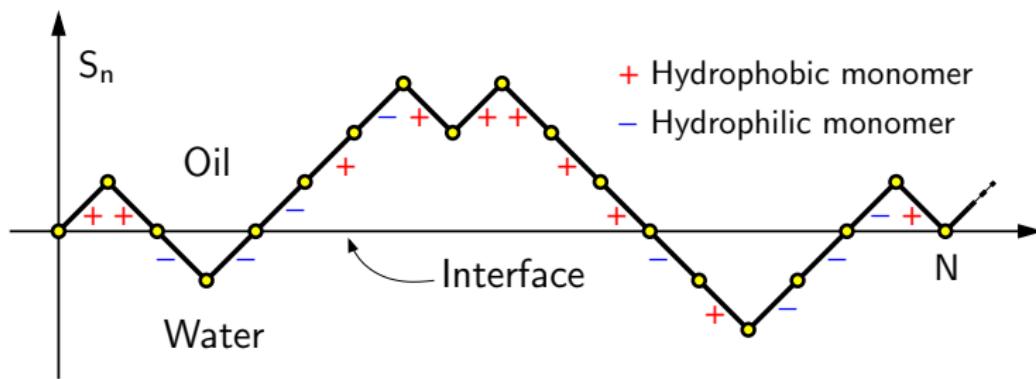
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Localization or Delocalization?

A polymer model interpretation

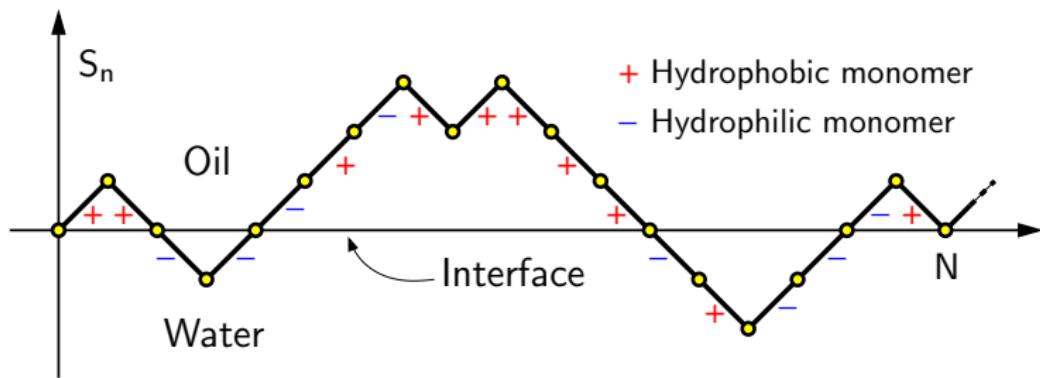


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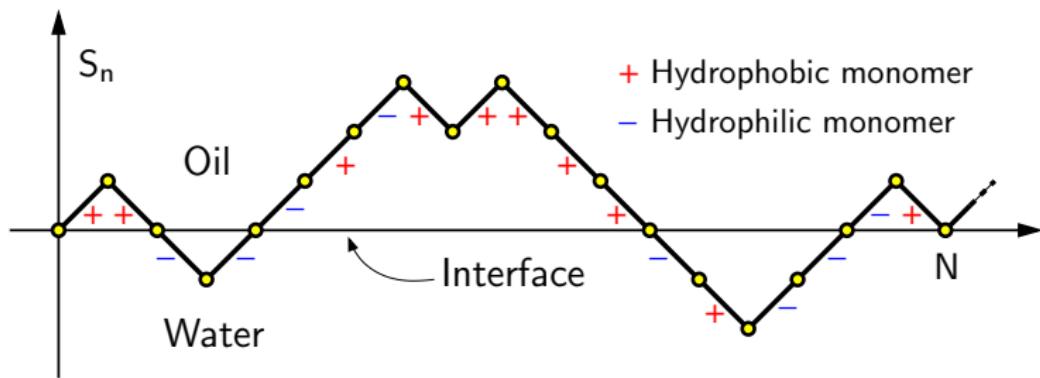
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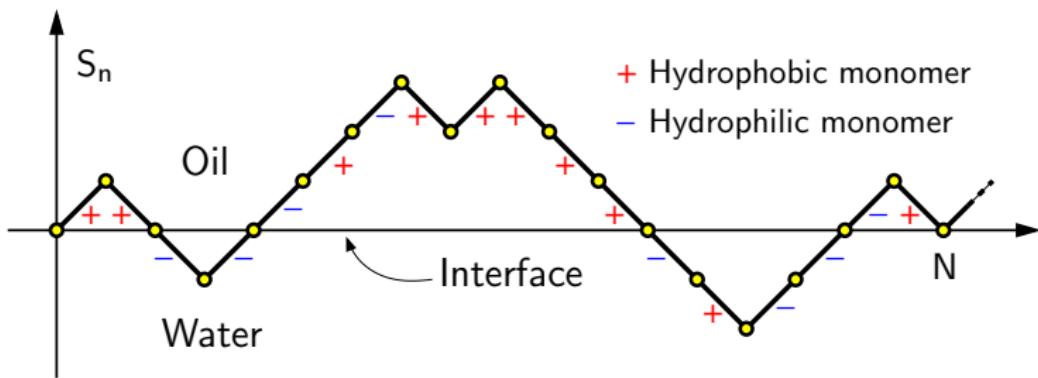
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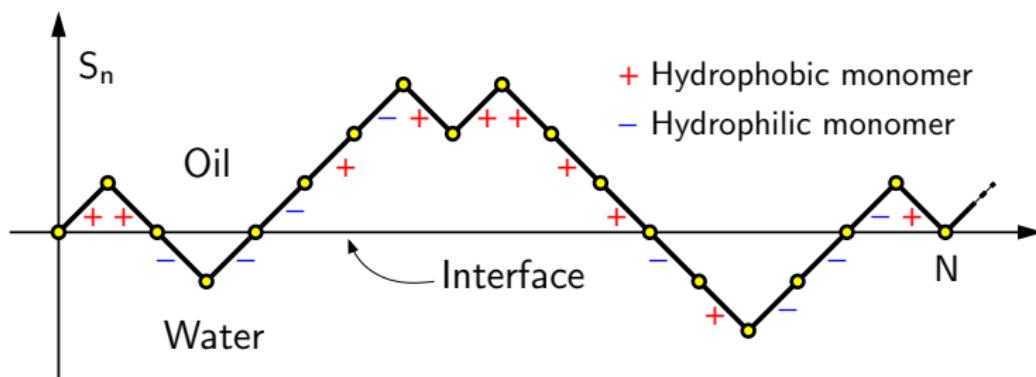
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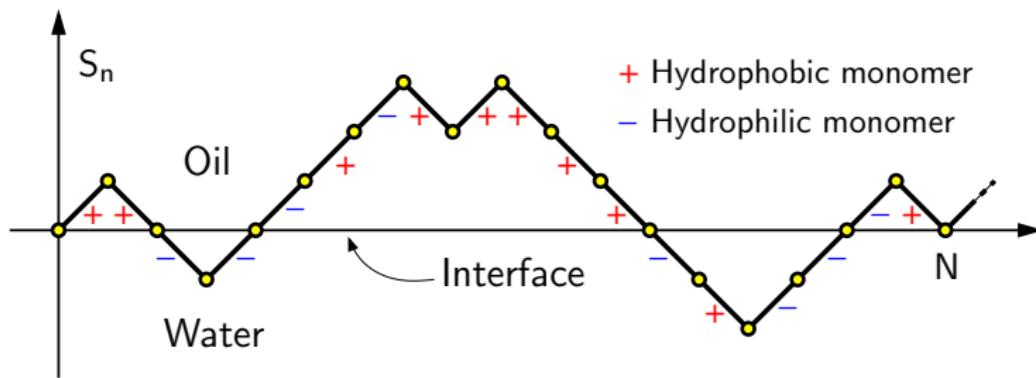
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- ▶ Physics literature: C. Monthus, S. Stepanow, J.-U. Sommer, I. Ya. Erkhimovich, A. Trovato, A. Maritan, S. G. Whittington, C. E. Soteros, A. Rechnitzer, G. K. Iliev, ...

The basic copolymer model

Definition of the model: $\frac{d\mathbf{P}_{N,\omega}}{d\mathbf{P}}(S) := \frac{1}{Z_{N,\omega}} \exp(-H_{N,\omega}(S))$

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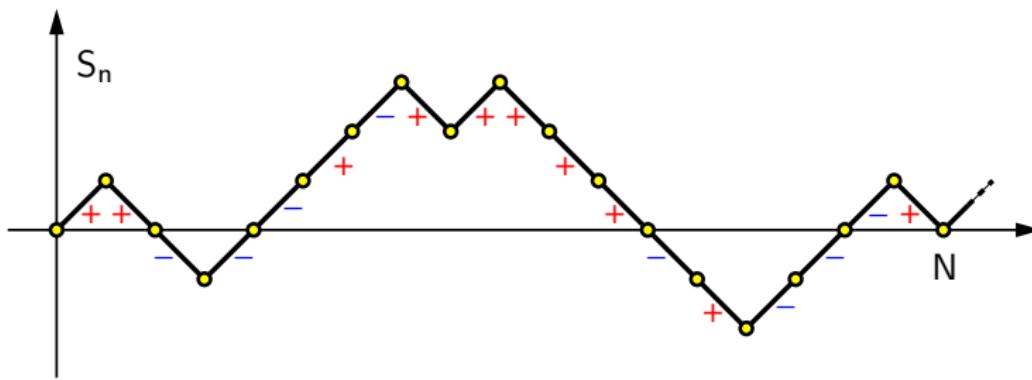
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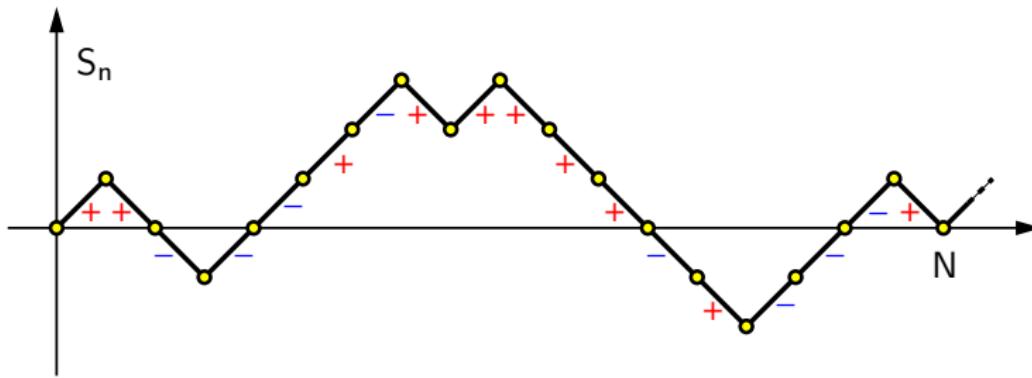
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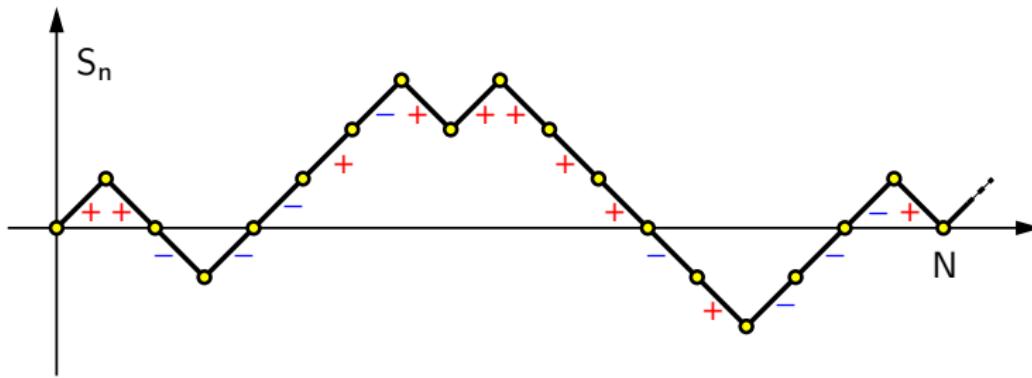
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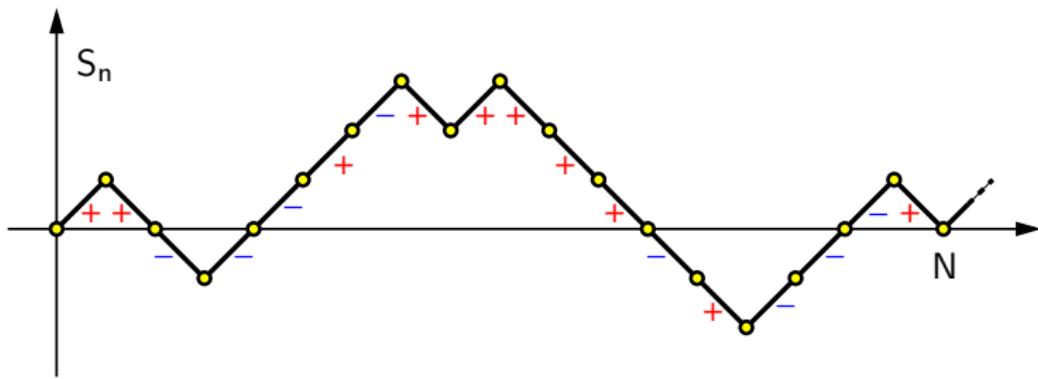
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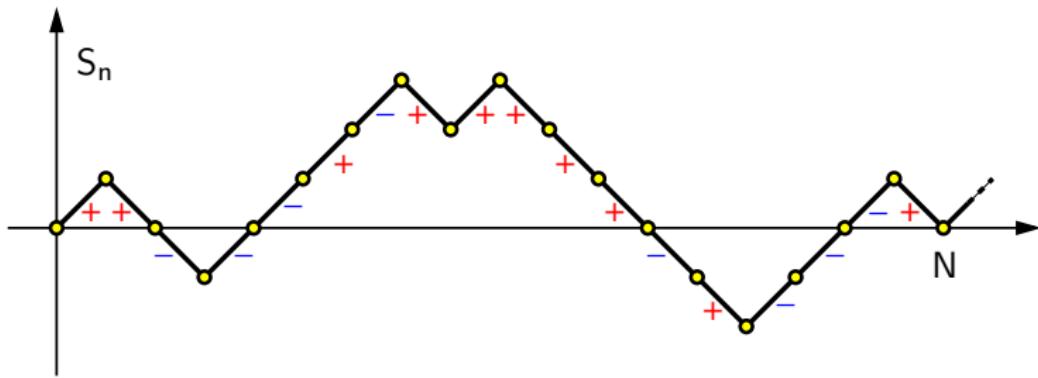


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This definition does correspond to sharply different path behaviors!

The phase diagram: discrete model

Theorem

The regions \mathcal{L} and \mathcal{D} are separated by a strictly increasing, continuous critical line $h_c(\cdot)$:

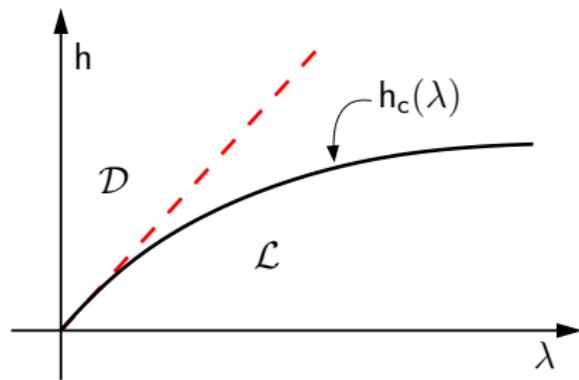
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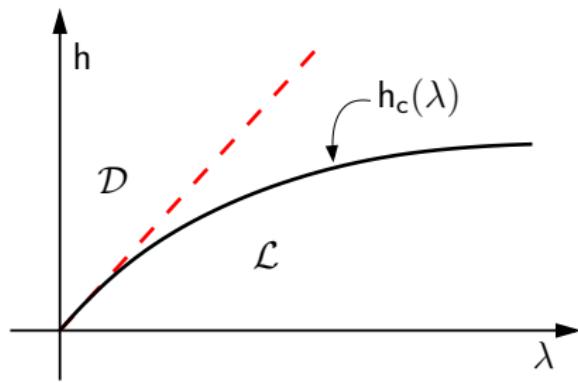


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Theorem

$$h_c(0) = 0 \text{ and as } \lambda \downarrow 0$$

$$h_c(\lambda) = m\lambda + o(\lambda),$$

with

$$\frac{2}{3} \leq m < 1.$$

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The continuum free energy $\tilde{F}(\lambda, h)$ is defined analogously:

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By Brownian scaling $\tilde{F}(a\lambda, ah) = a^2 \tilde{F}(\lambda, h)$ for all $a, \lambda, h \geq 0$.

The phase diagram: continuum model

The continuum free energy $\tilde{F}(\lambda, h)$ is defined analogously:

$$\tilde{F}(\lambda, h) := \lim_{t \rightarrow \infty} \frac{1}{t} \tilde{\mathbb{E}} \left(\log \tilde{Z}_{t,\beta} \right)$$

This time existence is highly non-trivial!

Again $\tilde{F}(\lambda, h) \geq 0$. We then define $\tilde{\mathcal{L}}$ ocalization and $\tilde{\mathcal{D}}$ elocalization:

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Therefore $\tilde{h}_c(\cdot)$ is a straight line: $\tilde{h}_c(\lambda) = \tilde{m}\lambda$.

The weak coupling limit

Theorem ([BdH 97])

For all $\lambda, h \geq 0$

$$\lim_{a \downarrow 0} \frac{1}{a^2} F(a\lambda, ah) = \tilde{F}(\lambda, h).$$

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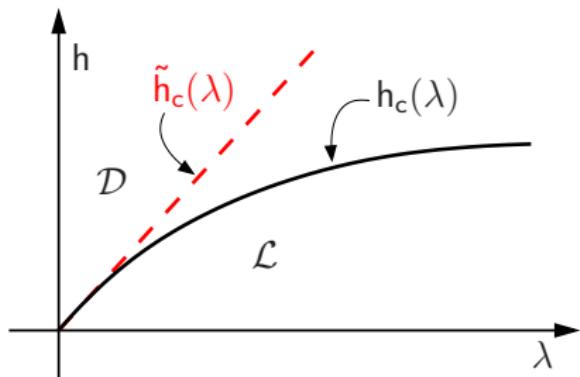
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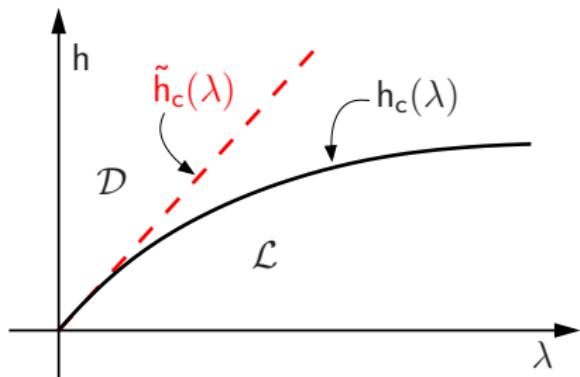
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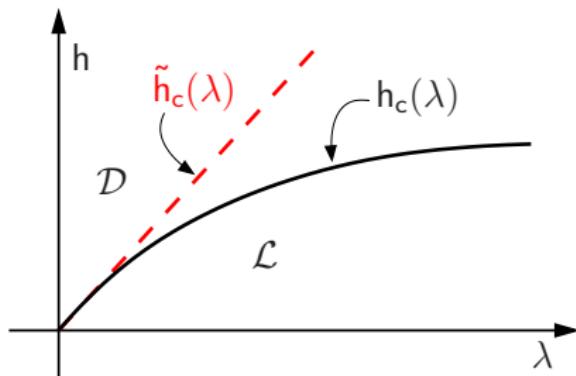
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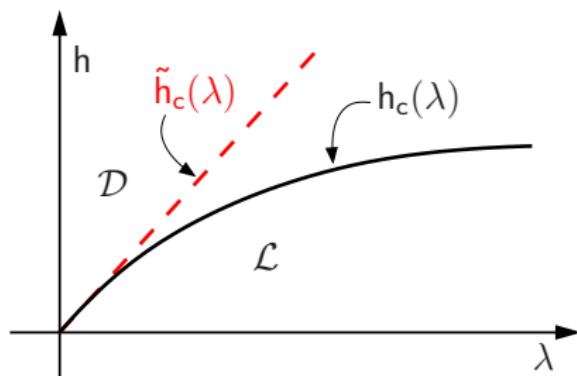
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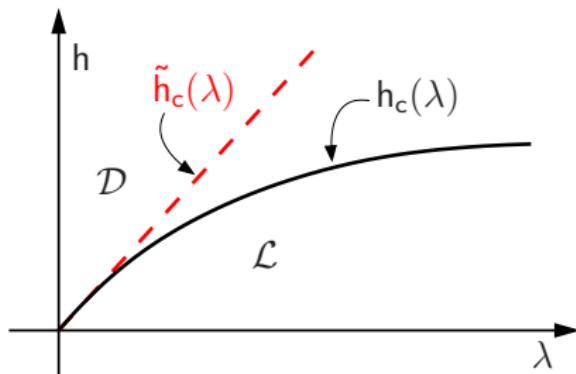
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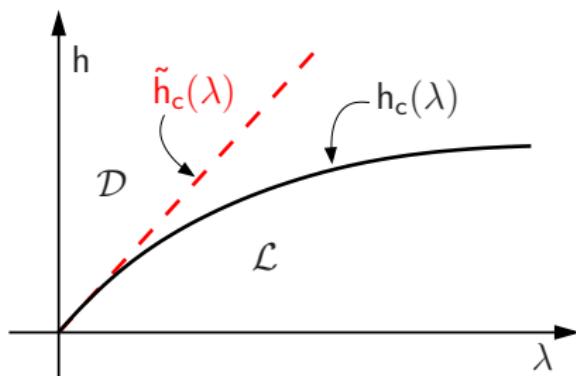
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... unfortunately for just **one** discrete model. Generalization?

Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

Beyond the simple random walk

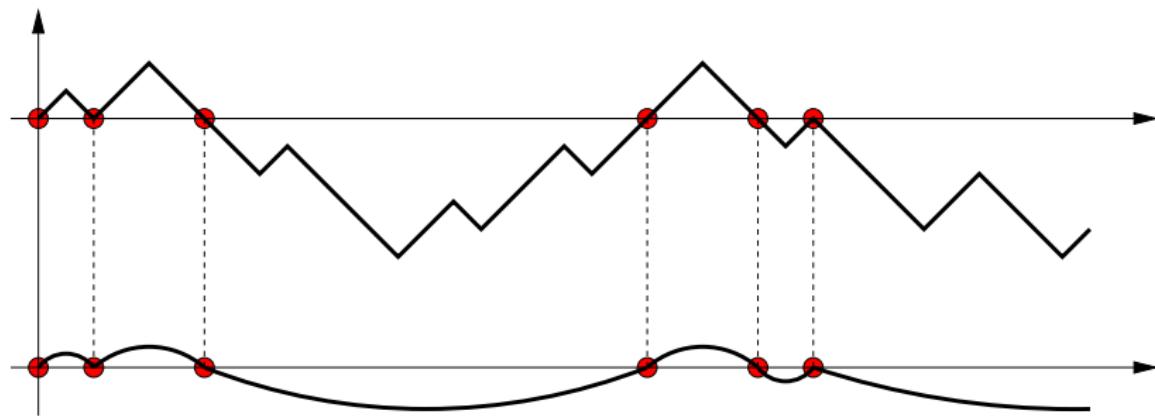
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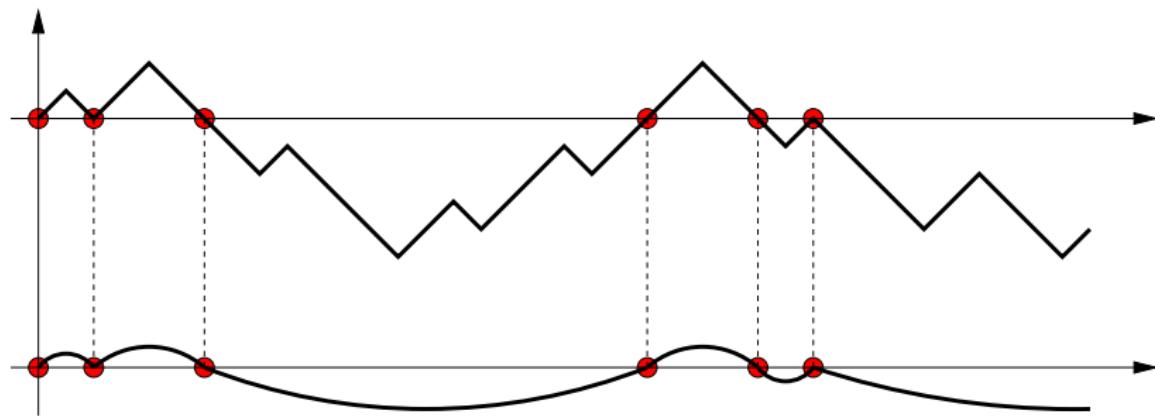
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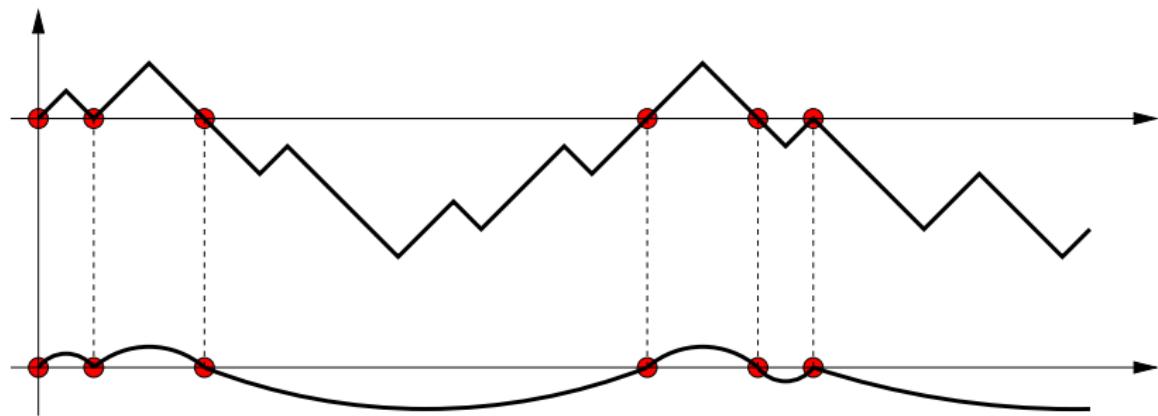


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- ▶ Excursions signs: fair coin tossing (independent of $\{\tau_k\}_{k \geq 0}$)

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Discrete Bessel-like process ($c_\alpha = 1/2 - \alpha$)

$\mathbf{P}(S_{n+1} = x \pm 1 | S_n = x) = \frac{1}{2} \left(1 \pm \frac{c_\alpha}{x} + o\left(\frac{1}{x}\right) \right)$ yields (\star) asymp.

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Note that $F(\cdot, \cdot)$ and $h_c(\cdot)$ depend on the choice of \mathbf{P} and \mathbb{P}

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From $\tilde{\tau}^\alpha$ we obtain $(\tilde{\Delta}^\alpha = \{\tilde{\Delta}_t^\alpha\}_{t \geq 0}, \tilde{\mathbf{P}})$ (For $\alpha = \frac{1}{2}$ we recover BM)

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For all $\lambda, h > 0$ and $\epsilon \in (0, 1)$ there exists $a_0 > 0$ s.t. for all $a < a_0$

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Outline

1. The basic copolymer model
2. The free energy
3. Generalized copolymer models
4. Strategy of the proof

Strategy of the proof

Let $f^i(\lambda, h)$ be the free energy of an Hamiltonian $H_N^i = H_N^i(\lambda, h)$

$$f^i(\lambda, h) := \lim_{N \rightarrow \infty} \frac{1}{N} \mathbb{E} (\log \mathbf{E} [\exp (-H_N^i)])$$

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$H_N^i \simeq H_N^j$ if for all $C > 0$ and $a \ll 1$ [also with $i \leftrightarrow j$]

$$\limsup_{N \rightarrow \infty} \frac{1}{N} \log \mathbb{E} \mathbf{E} \left[\exp \left(-C(H_N^i(a\lambda, ah) - H_N^j(a\lambda, (1-\epsilon)ah)) \right) \right] \leq 0$$

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1. Coarse-graining of the renewal process

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Happy birthday, Erwin!