Periodic solutions for a Sitnikov restricted $n+1$-body problem with primaries in rigid motion

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Abstract

1 Introduction

In this paper we discuss existence of periodic solutions for the following restricted nonplanar Newtonian $n+1$-body problem P (see figure 1):

P_1 We have n primary bodies of masses $m_1, \dots, m_n$ and an additional masless body.

P_2 The primary bodies are in a central configuration rigid motion (see [11, Section 2.9]). This motion is carried out in a plane Π.

P_3 The massless particle is moving on the perpendicular line to Π passing through the center of masses.

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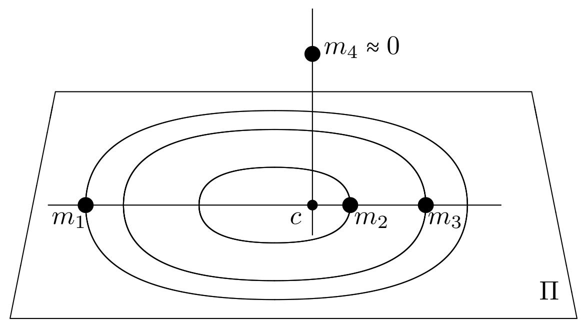

Figure 1: Four-body problem with three primaries

Problems like the one presented have been extensively discussed in the literature. In [16] K. Sitnikov considered the problem of two body in a Keplerian motion and a massless particle moving in the perpendicular line to the orbital plane passing through the center of masses. Sitnikov obtained deep results about existence of solutions, some of them periodic (see [13, III(5)]. Since then many other authors have studied Sitnikov problem, for instance Liu, Zhou, and Sun [9], Hagel and Trenkler [6], Dvorak [5], Dankowicz and Holmes [4], Llibre, Meyer and Soler [10], Chesley [3], Jiménez-Lara a and Escalona-Buendía [7], Llibre and Ortega [12], Pérez, Jiménez and Lacomba [15].

Problems like the Sitnikov problem for four bodies were addressed more recently. In [17] Soulis, Papadakis and Bountis studied existence, linear stability and bifurcations for a problem similar to P, where in place to have a Eulerian colinear configuration they had a Lagrangian equilateral triangle configuration for the primaries bodies, which are supposed to have the same mass $m_1 = m_2 = m_3$. Later, In [1] Baltagiannis nad Papadakis considered more general masses and in [14] Pandey and Ahmad extend the analysis started in [17] to the case when the primaries are oblate (not mass points). In [19], Zhao and Zhang proved existence of periodic solutions for a problem similar to the one dealt with in [17]. They used a variational approach. In the present paper we extend the analysis in [19] to the case of a colinear central configuration for the primaries. In [8] Li, Zhang and Zhao studied a special type of restricted circular $N+1$-body problem with equal masses for the primaries.

Given that in our problem the primaries are no longer equidistant and their relative position is determined by a polynomial equation of fifth degree, the calculations involved here are tedious to reproduce completely and difficult for that the reader to check them by hand. For this reason we have prepared a jupyter-notebook (see [2]) with some of these calculations. With a little knowledge of Python-Sympy (see [18]) the reader can check and reproduce them easily.

2 Preliminaries and Main Results

We start considering n bodies, $n > 2$, of masses $m_1, \dots, m_n$ moving in a Euclidean 3-dimensional space according to Newton's laws of motion. We assume that $q_1(t), \dots, q_n(t)$ are the coordinates (column vectors) of the bodies in some inertial Cartesian coordinate system. We denote by $r_{ij} = |q_i - q_j|$ the Euclidean distance between q_i and q_j. We can suppose without any loos generality, we can assume the center of mass $c := \sum_j m_j q_j / M$ ($M := \sum_j m_j$) is fixed at the origin ($c = 0$).

We assume that these bodies are in a *rigid motion*. We recall that a *rigid motion*, is a solution of motions equations with r_{ij} constant. It is known (see -buscar referencias) that in a rigid motion is performed in a plane Π. We assume that Π is the plane determined by the first two coordinates axes. Then a rigid motion has the form

$$q_j(t) = Q(\nu t) q_j^0,$$

where

$$Q(\nu t) = \begin{pmatrix} \cos(\nu t) & -\sin(\nu t) & 0 \\ \sin(\nu t) & \cos(\nu t) & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

and $q_j^0 \in \Pi$, $j = 1, \dots, n$ are vectors in a planar *central configuration* (CC) in $\mathbb{R}^3$, i.e. there exists $\lambda \in \mathbb{R}$ such that

$$\nabla_j U(q_1^0, \dots, q_n^0) + \lambda m_j q_j^0 = 0, \quad j = 1, \dots, n.$$

where the *potential function* U is defined by:

$$U(x) = \sum_{i<j} \frac{m_i m_j}{r_{ij}}, \tag{1}$$

and ∇_j denotes the 3-dimensional partial gradient with respect to q_j. According to [11, Eq. (2.16)] we have $\nu^2 = \lambda$. The primaries bodies perform a periodic motion with period $T := 2\pi/\nu$.

We suppose that we have a massless particle with coordinates $q(t) = (x(t), y(t), z(t)) \in \mathbb{R}^3$. This particle does not disturb the rigid motion of primaries. We want to find conditions under which this particle perform a T-periodic motion on the third axis of coordinates.

The particle q satisfies the Newtonian equations of motion

$$\ddot{q} = \sum_{i=1}^{n} \frac{m_i (q_i - q)}{|q_i - q|^3}, \tag{2}$$

Lemma 2.1. *For $c > 0$ we define the function $y_c(t) := (c+t)^{-3/2}$. If $0 < t_1 < t_2 < \dots < t_k$ then the functions $y_j(t) := y_{t_j}(t)$ are linearly independent on each open interval $I \subset \mathbb{R}^+$.*

Proof. It is sufficient to prove that Wronskian

$$W := W(y_1,\dots,y_k)(t) = \det\begin{pmatrix} y_1 & \cdots & y_k \\ \frac{dy_1}{dt} & \cdots & \frac{dy_k}{dt} \\ \vdots & \ddots & \vdots \\ \frac{d^{k-1}y_1}{dt^{k-1}} & \cdots & \frac{d^{k-1}y_k}{dt^{k-1}} \end{pmatrix}$$

is not null on I.

Using induction is easy to show that

$$\frac{d^i y_c}{dt^i} = \beta_i y_c^{\frac{2i+3}{3}}, \quad \text{for some } \beta_i \neq 0, \text{ and for all } i = 1,\dots. \tag{3}$$

Fix any $t \in I$. Then, according to (3) and writing $\lambda_j := (t+t_j)^{-1}$, we have

$$\begin{aligned} W(t) &= \det\begin{pmatrix} \lambda_1^{3/2} & \lambda_2^{3/2} & \cdots & \lambda_k^{3/2} \\ \beta_1\lambda_1^{5/2} & \beta_1\lambda_2^{5/2} & \cdots & \beta_1\lambda_k^{5/2} \\ \vdots & \vdots & \ddots & \vdots \\ \beta_{k-1}\lambda_1^{k+1/2} & \beta_{k-1}\lambda_2^{k+1/2} & \cdots & \beta_{k-1}\lambda_k^{k+1/2} \end{pmatrix} \\ &= \beta_1\beta_2\cdots\beta_{k-1}\lambda_1^{3/2}\lambda_2^{3/2}\cdots\lambda_k^{3/2}\det\begin{pmatrix} 1 & 1 & \cdots & 1 \\ \lambda_1 & \lambda_2 & \cdots & \lambda_k \\ \vdots & \vdots & \ddots & \vdots \\ \lambda_1^{k-1} & \lambda_2^{k-1} & \cdots & \lambda_k^{k-1} \end{pmatrix} \\ &= \beta_1\beta_2\cdots\beta_{k-1}\lambda_1^{3/2}\lambda_2^{3/2}\cdots\lambda_k^{3/2}\prod_{1\le i<j\le n}(\lambda_j-\lambda_i), \end{aligned}$$

where the last equality follows of the well known Vandermonde determinant identity. Therefore $W \neq 0$ if and only if $\lambda_i \neq \lambda_j$, $i \neq j$, which in turn is equivalent to $t_i \neq t_j$, $i \neq j$. □

Theorem 2.2. *There exists a non-stationary solution of* (2) *with* $x(t) = y(t) = 0$ *if and only if* $q_1^0,\dots,q_n^0$ *satisfy that for any* $r > 0$ *such that the set*

$$F_r = \{i : |q_i^0| = r\}$$

is non empty, then

$$\sum_{i\in F_r} m_i q_i^0 = 0. \tag{4}$$

Proof. We use a rotating coordinate system where the primaries are fixed. Concretely we put

$$\xi = Q(-\nu t)q.$$

In this system the motion equations are

$$\ddot{\xi} + 2\nu B\dot{\xi} + \nu^2 C\xi = \sum_{i=1}^n \frac{m_i(q_i^0-\xi)}{|q_i^0-\xi|^3}, \tag{5}$$

where

$$B := \begin{pmatrix} J & 0_{2\times 1} \\ 0_{1\times 2} & 0 \end{pmatrix}, \quad J := \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \quad \text{and} \quad C = \begin{pmatrix} -I_2 & 0_{2\times 1} \\ 0_{1\times 2} & 0 \end{pmatrix},$$

where $0_{n\times m}$ and I_n denote the null $n \times m$ matrix and the identity $n \times n$ matrix respectively. Assuming that the masless particle is moving on the z-axis then $\xi = q = (0, 0, z)$ and the Coriolis and centrifugal forces, $2\nu B\dot{\xi}$ and $\nu^2 C\xi$ respectively, are null. Therefore, taking account in the first two equation in (5) and identifying the vectors q_i^0, $i = 1, \dots, n$ with vectors in $\mathbb{R}^2$, we have

$$\sum_{i=1}^{n} \frac{m_i q_i^0}{|q_i^0 - \xi|^3} = 0.$$

Let $D = \{|q_i^0| : i = 1, \dots, n\}$. Suppose that $D = \{r_1, \dots, r_k\}$, with $r_i \neq r_j$ for $i \neq j$, and $\{1, \dots, n\} = F_1 \cup \dots \cup F_k$, where if $i \in F_j$ then $|q_i^0| = r_j$. Then

$$\sum_{j=1}^{k} \left\{ \frac{1}{(r_j^2 + z^2)^{3/2}} \sum_{i \in F_j} m_i q_i^0 \right\} = 0.$$

Since we are considering a non-stationary solution, we have that $z(t)$ is not constant. Therefore there exists an interval $I \subset \mathbb{R}^+$ where

$$\sum_{j=1}^{k} \left\{ \frac{1}{(r_j^2 + s)^{3/2}} \sum_{i \in F_j} m_i q_i^0 \right\} = 0, \quad s \in I.$$

Then, according to Lemma 2.1, we obtain (4). □

Acknowledgments

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