# Periodic solutions of Euler-Lagrange equations in an anisotropic Orlicz-Sobolev space setting

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#### **Abstract**

#### 1 Introduction

In this paper we obtain existence of solutions for systems of equations of the type:

$$\begin{cases} \frac{d}{dt} D_y \mathcal{L}(t, u(t), u'(t)) = D_x \mathcal{L}(t, u(t), u'(t)) & \text{a.e. } t \in (0, T), \\ u(0) - u(T) = u'(0) - u'(T) = 0, \end{cases}$$
 (P)

where the function  $\mathcal{L}:[0,T]\times\mathbb{R}^d\times\mathbb{R}^d\to\mathbb{R},\ d\geqslant 1$  (called the *Lagrange function* or *lagrangian*) satisfying that it is measurable in t for each  $(x,y)\in\mathbb{R}^d\times\mathbb{R}^d$  and continuously differentiable in (x,y) for almost every  $t\in[0,T]$ . The unknown function  $u:[0,T]\to\mathbb{R}^d$  is assumed absolutely continuous.

Our approach involves the direct method of the calculus of variations in the framework of *anisotropic Orlicz-Sobolev spaces*. We suggest the articles [16] for definitions and main results on anisotropic Orlicz spaces, see also [2]. These spaces allow us to unify and extend previous results on existences of solutions for systems like (P).

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Through this article we said that a function  $\Phi : \mathbb{R}^d \to [0, +\infty)$  is of  $N_\infty$  class if  $\Phi$  is convex,  $\Phi(0) = 0$ ,  $\Phi(x) > 0$  if  $x \neq 0$  and  $\Phi(-y) = \Phi(y)$ , and

$$\lim_{|y| \to \infty} \frac{\Phi(y)}{|y|} = +\infty. \tag{1}$$

where  $|\cdot|$  denotes the euclidean norm on  $\mathbb{R}^d$ . From [3, Cor. 2.35] a  $N_{\infty}$  function is continuous.

Associated to  $\Phi$  we have the *complementary function*  $\Psi$  which is defined in  $\xi \in \mathbb{R}^d$  as

$$\Psi(\xi) = \sup_{y \in \mathbb{R}^d} y \cdot \xi - \Phi(y) \tag{2}$$

then, from the continuity of  $\Phi$  and (1), we have that  $\Psi: \mathbb{R}^d \to [0, \infty)$ . Moreover, it is easy to see that  $\Psi$  is a convex function such that  $\Psi(0) = 0$ ,  $\Psi(-\xi) = \Psi(\xi)$  and  $\Psi(\xi) \to \infty$  when  $|\xi| \to \infty$ , see [3, Section 4.2], [8, Chapter 2]. Moreover  $\Psi$  satisfies (1) (see [16, Th. 2.2]). i.e.  $\Psi$  is  $N_{\infty}$  function.

Some examples of  $N_{\infty}$  functions are the following.

Example 1.1.  $\Phi_p(y) \coloneqq |y|^p/p$ , for  $1 . In this case <math>\Psi(\xi) = |\xi|^q/q$ , q = p/(p-1). Example 1.2. If  $\Phi : \mathbb{R} \to [0, +\infty)$  is a  $N_\infty$  function on  $\mathbb{R}$  then  $\overline{\Phi}(y) = \Phi(|y|)$  is a  $N_\infty$  function on  $\mathbb{R}^d$ . In this example, as in the previous one, the function  $\Phi$  is *radial*, i.e. the value of  $\Phi(y)$  depends on the norm of y and not on its direction. These cases are not authentically anisotropic.

*Example 1.3.* An anisotropic function  $\Phi(y)$  depends on the direction of y. For example, if  $1 < p_1, p_2 < \infty$ , we define  $\Phi_{p_1, p_2} : \mathbb{R}^d \times \mathbb{R}^d \to [0, +\infty)$  by

$$\Phi_{p_1,p_2}(y_1,y_2) \coloneqq \frac{|y_1|^{p_1}}{p_1} + \frac{|y_2|^{p_2}}{p_2}.$$

Then  $\Phi_{p_1,p_2}$  is a  $N_{\infty}$  function. In this case the complementary function is the function  $\Phi_{q_1,q_2}$  with  $q_i = p_i/(p_i-1)$ .

More generally, if  $\Phi_k : \mathbb{R}^d \to [0, +\infty)$ , k = 1, ..., n, are  $N_\infty$  functions, then  $\Phi : \mathbb{R}^d \times \cdots \times \mathbb{R}^d \to [0, +\infty)$  defined by  $\Phi(y_1, ..., y_n) = \Phi_1(y_1) + \cdots + \Phi_n(y_n)$  is a  $N_\infty$  function. These functions are truly anistropic, i.e. |x| = |y| does not implies that  $\Phi(x) = \Phi(y)$ .

Example 1.4. If  $\Phi : \mathbb{R} \to [0, +\infty)$  is a  $N_{\infty}$  function and  $O \in GL(d, \mathbb{R})$ , then  $\Phi(y) = \Phi(Oy)$  is a  $N_{\infty}$  function.

Example 1.5. An anisotropic  $N_{\infty}$  function is not necessarily controlled by powers if it does not satisfy the  $\Delta_2$  condition (see xxxxx). For example  $\Phi: \mathbb{R}^d : \to [0, +\infty)$  defined by  $\Phi(y) = \exp(|y|) - 1$  is  $N_{\infty}$  function.

The occurrence of Orlicz Spaces in this paper obeys to we will consider the following structure condition on the lagrangian:

$$|\mathcal{L}| + |\nabla_x \mathcal{L}| + \Psi(\nabla_y \mathcal{L}) \le a(x) \left\{ b(t) + \Phi\left(\frac{y}{\lambda}\right) \right\},$$
 (S)

for a.e.  $t \in [0,T]$ , where  $a \in C(\mathbb{R}^d, [0,+\infty)), b \in L^1([0,T], [0,+\infty))$ .

Our condition (S) includes structure conditions that have previously been considered in the literature. For example, it is easy to see that, when  $\Phi(x)$  is as in Example

1.1, then the condition (S) is equivalent to the structure condition in [8, Th. 1.4]. If  $\Phi$  is a radial  $N_{\infty}$  function such that  $\Psi$  satisfies tha  $\Delta_2$  function (see xxxxxx) then (S) is essentially equivalent to conditions [1, Eq. (2)-(4)]. If  $\Phi$  is as in Example 1.3 and  $\mathcal{L} = \mathcal{L}(t, x_1, x_2, y_1, y_2)$  then Inequality (S) is related to estructure conditions like [22, Lemma 3.1, Eq. (3.1)]. As can be seen, condition (S) is a more compact expression than [22, Lemma 3.1, Eq. (3.1)] and moreover is more weaker, because (S) does not imply a control of  $|D_{y_1}L|$  independent of  $y_2$ . We will return to this point later.

An important example of lagrangian is giving by:

$$\mathcal{L}_{\Phi,F}(t,x,y) \coloneqq \Phi(y) + F(t,x). \tag{3}$$

Here the function F(t,x), which is often referred to potential, be differentiable with respect to x for a.e.  $t \in [0,T]$ . Moreover F satisfies the following conditions:

- (C) F and its gradient  $\nabla_x F$ , with respect to  $x \in \mathbb{R}^d$ , are Carathéodory functions, i.e. they are measurable functions with respect to  $t \in [0,T]$ , for every  $x \in \mathbb{R}^d$ , and they are continuous functions with respect to  $x \in \mathbb{R}^d$  for a.e.  $t \in [0,T]$ .
- (A) For a.e.  $t \in [0,T]$ , it holds that

$$|F(t,x)| + |\nabla_x F(t,x)| \le a(x)b(t). \tag{4}$$

where 
$$a \in C(\mathbb{R}^d, [0, +\infty))$$
 and  $0 \le b \in L^1([0, T], \mathbb{R})$ .

The lagrangian  $\mathcal{L}_{\Phi,F}$  satisfies condition (S). In order to prove this, the only non trivial fact that we should to establish is is that  $\Psi(\nabla_y \mathcal{L}) \leq a(x) \{b(t) + \Phi(y/\lambda)\}$ . But, from inequality xxxx below,  $\Psi(\nabla_y \mathcal{L}) = \Psi(\nabla \Phi(y)) \leq \Phi(2y)$ .

The laplacian  $\mathcal{L}_{\Phi,F}$  leads to the system

$$\begin{cases} \frac{d}{dt} \nabla \Phi(u'(t)) = \nabla_x F(t, u(t)) & \text{a.e. } t \in (0, T), \\ u(0) - u(T) = u'(0) - u'(T) = 0, \end{cases}$$
  $(\mathbf{P}_{\Phi})$ 

Problem  $(P_{\Phi})$  contains, as a particular case, many problems that are usually considered in the literature. For example, the classic book [8] deals mainly with problem (P), for the lagrangian  $\mathcal{L}_{\Phi,F}$ , with  $\Phi(x)=|x|^2/2$ , through various methods: direct, dual action, minimax, etc. The results in [8] were extended and improved in several articles, see [20, 18, 23, 19, 26] to cite some examples. The case  $\Phi(y)=|y|^p/p$ , for arbitrary  $1 were considered in [22, 21], among other papers, and in this case <math>(P_{\Phi})$  is reduced to the p-laplacian system

$$\begin{cases} \frac{d}{dt} \left( u'(t) |u'|^{p-2} \right) = \nabla F(t, u(t)) & \text{a.e. } t \in (0, T) \\ u(0) - u(T) = u'(0) - u'(T) = 0. \end{cases}$$
  $(\mathbf{P_p})$ 

If  $\Phi$  is as in Example 1.3 and  $F:[0,T]\times\mathbb{R}^d\times\mathbb{R}^d\to\mathbb{R}$  is a Carathéodory function, then the equations  $(P_{\Phi})$  become

$$\begin{cases} \frac{d}{dt} \left( |u_1'|^{p_1 - 2} u_1' \right) = F_{x_1}(t, u) & \text{a.e. } t \in (0, T) \\ \frac{d}{dt} \left( |u_2'|^{p_2 - 2} u_2' \right) = F_{x_2}(t, u) & \text{a.e. } t \in (0, T) \\ u(0) - u(T) = u'(0) - u'(T) = 0, \end{cases}$$
  $(\boldsymbol{P_{p_1, p_2}})$ 

where  $x = (x_1, x_2) \in \mathbb{R}^d \times \mathbb{R}^d$  and  $u(t) = (u_1(t), u_2(t)) \in \mathbb{R}^d \times \mathbb{R}^d$ . In the literature, these equations are known as  $(p_1, p_2)$ -Laplacian system, see [25, 12, 24, 9, 10, 11, 6].

In conclusion, the problem (P) with conditions (S) contains several problems that have been considered by many authors in the past.

## 2 Anisotropic Orlicz and Orlicz-Sobolev spaces

In this section, we give a short introduction to Orlicz and Orlicz-Sobolev spaces of vector valued functions associated to anisotropic  $N_{\infty}$  functions  $\Phi: \mathbb{R}^n \to [0, +\infty)$ . References for these topics are [4, 16, 17].

We say that  $\Phi: \mathbb{R}^d \to [0, +\infty)$  satisfies the  $\Delta_2^{\infty}$ -condition, denoted by  $\Phi \in \Delta_2^{\infty}$ , if there exist constants K > 0 and  $M \geqslant 0$  such that

$$\Phi(2x) \leqslant K\Phi(x),\tag{5}$$

for every  $|x| \ge M$ .

We denote by  $\mathcal{M} := \mathcal{M}([0,T],\mathbb{R}^d)$ , with  $d \ge 1$ , the set of all measurable functions (i.e. functions which are limits of simple functions) defined on [0,T] with values on  $\mathbb{R}^d$  and we write  $u = (u_1, \dots, u_d)$  for  $u \in \mathcal{M}$ .

Given an  $N_{\infty}$  function  $\Phi$  we define the modular function  $\rho_{\Phi}: \mathcal{M} \to \mathbb{R}^+ \cup \{+\infty\}$  by

$$\rho_{\Phi}(u) \coloneqq \int_0^T \Phi(u) \ dt.$$

Now, we introduce the *Orlicz class*  $C^{\Phi} = C^{\Phi}([0,T],\mathbb{R}^d)$  by setting

$$C^{\Phi} := \{ u \in \mathcal{M} | \rho_{\Phi}(u) < \infty \}. \tag{6}$$

The Orlicz space  $L^{\Phi}$  =  $L^{\Phi}$  ([0, T],  $\mathbb{R}^d$ ) is the linear hull of  $C^{\Phi}$ ; equivalently,

$$L^{\Phi} := \{ u \in \mathcal{M} | \exists \lambda > 0 : \rho_{\Phi}(\lambda u) < \infty \}. \tag{7}$$

The Orlicz space  $L^{\Phi}$  equipped with the Luxemburg norm

$$\|u\|_{L^\Phi}\coloneqq\inf\left\{\lambda\left|\rho_\Phi\left(\frac{v}{\lambda}\right)dt\leqslant1\right\},$$

is a Banach space.

The subspace  $E^{\Phi} = E^{\Phi}\left([0,T],\mathbb{R}^d\right)$  is defined as the closure in  $L^{\Phi}$  of the subspace  $L^{\infty}\left([0,T],\mathbb{R}^d\right)$  of all  $\mathbb{R}^d$ -valued essentially bounded functions. It is shown that (see [16, Thm. 5.1])  $u \in E^{\Phi}$  if and only if  $\rho_{\Phi}(\lambda u) < \infty$  for any  $\lambda > 0$ . The equality  $L^{\Phi} = E^{\Phi}$  is true if and only if  $\Phi \in \Delta_2^{\infty}$  (see [16, Thm. 5.2]).

A generalized version of *Hölder's inequality* holds in Orlicz spaces (see [16, Thm. 7.2]). Namely, if  $u \in L^{\Phi}$  and  $v \in L^{\Psi}$  then  $u \cdot v \in L^{1}$  and

$$\int_{0}^{T} v \cdot u \, dt \le 2 \|u\|_{L^{\Phi}} \|v\|_{L^{\Psi}}. \tag{8}$$

By  $u\cdot v$  we denote the usual dot product in  $\mathbb{R}^d$  between u and v. We consider the subset  $\Pi(E^\Phi,r)$  of  $L^\Phi$  given by

$$\Pi(E^{\Phi}, r) := \{ u \in L^{\Phi} | d(u, E^{\Phi}) < r \}.$$

This set is related to the Orlicz class  $C^{\Phi}$  by means of inclusions, namely,

$$\Pi(E^{\Phi}, r) \subset rC^{\Phi} \subset \overline{\Pi(E^{\Phi}, r)} \tag{9}$$

for any positive r. This relation is a trivial generalization of [16, Thm. 5.6]. If  $\Phi \in \Delta_2^{\infty}$ , then the sets  $L^{\Phi}$ ,  $E^{\Phi}$ ,  $\Pi(E^{\Phi}, r)$  and  $C^{\Phi}$  are equal.

As usual, if  $(X, \|\cdot\|_X)$  is a normed space and  $(Y, \|\cdot\|_Y)$  is a linear subspace of X, we write  $Y \to X$  and we say that Y is *embedded* in X when there exists  $\hat{C} > 0$  such that  $||y||_X \le C||y||_Y$  for any  $y \in Y$ . With this notation, Hölder's inequality states that  $L^{\Phi} \hookrightarrow [L^{\Psi}]^*$ , where a function  $v \in L^{\Phi}$  is associated to  $\xi_v \in [L^{\Psi}]^*$  being

$$\xi_v(u) = \langle \xi_v, u \rangle = \int_0^T v \cdot u \, dt, \tag{10}$$

We highlight the following result that is a consequence of Theorems 7.1 and 7.3 in [16].

**Proposition 2.1.** If  $\Psi$  satisfies the  $\Delta_2^{\infty}$ -condition then  $L^{\Phi}([0,T],\mathbb{R}^d) = [L^{\Psi}([0,T],\mathbb{R}^d)]^*$ .

We define the Sobolev-Orlicz space  $W^1L^{\Phi}$  by

$$W^{1}L^{\Phi}\left([0,T],\mathbb{R}^{d}\right)\coloneqq\left\{u|u\in AC\left([0,T],\mathbb{R}^{d}\right)\text{ and }u'\in L^{\Phi}\left([0,T],\mathbb{R}^{d}\right)\right\},$$

where  $AC([0,T],\mathbb{R}^d)$  denotes the space of all  $\mathbb{R}^d$  valued absolutely continuous functions defined on [0,T]. The space  $W^1L^{\Phi}([0,T],\mathbb{R}^d)$  is a Banach space when equipped with the norm

$$||u||_{W^1L^{\Phi}} = ||u||_{L^{\Phi}} + ||u'||_{L^{\Phi}}. \tag{11}$$

We introduce the following subspaces of  $W^1L^\Phi$ 

$$W^{1}E^{\Phi} = \{u \in W^{1}L^{\Phi} | u' \in E^{\Phi}\},\$$

$$W^{1}E^{\Phi}_{T} = \{u \in W^{1}E^{\Phi} | u(0) = u(T)\}.$$
(12)

Given a Young's function  $\Phi$ , we define function  $A_{\Phi}: \mathbb{R}^+ \to \mathbb{R}^+$  by

$$A_{\Phi}(s) = \min\left\{\Phi(x) \,\middle|\, |x| = s\right\},\tag{13}$$

Let us establish some elementary properties of  $A_{\Phi}$  that we will use in this article.

**Proposition 2.2.** The function  $A_{\Phi}$  has the following properties:

- 1.  $A_{\Phi}$  is continuous,
- 2.  $A_{\Phi}(s)/s$  is increasing,

- 3.  $A_{\Phi}(|x|)$  is the greatest radial minorant of  $\Phi(x)$ ,
- 4.  $\Phi$  is  $N_{\infty}$  if and only if  $\lim_{s\to+\infty} A_{\Phi}(s)/s = +\infty$ .

*Proof.* It is well known that finite and convex functions defined on finite dimensional vector spaces are locally Lipschitz functions (see [3]). This fact implies item 1 immediately.

In order to prove item 2, suppose 0 < r < s and  $x \in \mathbb{R}^d$  with  $A_{\Phi}(s) = \Phi(x)$ . Then, from the definition of  $A_{\Phi}$  and the convexity of  $\Phi$ ,

$$\frac{A_{\Phi}(r)}{r} \leqslant \frac{\Phi\left(\frac{r}{s}x\right)}{r} \leqslant \frac{\Phi\left(x\right)}{s} = \frac{A_{\Phi}(s)}{s}.$$

Property in items 3 and 4 are obtained easily.

Example 2.1. We compute  $A_{\Phi}$  for the function  $\Phi = \Phi_{p_1,p_2}$  given in Example (1.3). We apply the method of Lagrange multipliers (see [7, Ch. 11]) to solve the next minimization problem subject to constraints

$$\begin{cases} \text{ minimize } \Phi_{p_1, p_2}(y_1, y_2) \\ \text{ subject to } |y_1|^2 + |y_2|^2 = r^2 \end{cases}.$$

The first order conditions are

$$\begin{cases} |y_1|^{p_1-2}y_1 + \lambda y_1 &= 0\\ |y_2|^{p_2-2}y_2 + \lambda y_2 &= 0\\ |y_1|^2 + |y_2|^2 &= r^2 \end{cases}$$
(14)

These equations are solved, among others, by the following two sets of citical points: a) |x| = r, y = 0 and  $\lambda = -r^{p_1-2}$  and b) x = 0, |y| = r and  $\lambda = -r^{p_2-2}$ . These sets are infinite when d > 1. Associated with these critical points we have the following critical values: a)  $r^{p_1}/p_1$  and b)  $r^{p_2}/p_2$ .

We deal with  $p_1 \le 2$  and  $p_2 \le 2$  being one of them (suppose  $p_2$ ) different from 2. The remaining cases can be treated with similar techniques.

If  $(y_1,y_2)$  solve (14) with  $y_1 \neq 0$  and  $y_2 \neq 0$  then  $|y_2| = |y_1|^{\frac{p_1-2}{p_2-2}}$  and  $\lambda = -|y_1|^{p_1-2}$ . We use second order conditions for constrained problems. We have to consider the tangent plane at the point  $(y_1,y_2) \in \mathbb{R}^{2n}$ , i.e.  $M = \{(\xi,\eta) \in \mathbb{R}^{2n} : \xi y_1^t + \eta y_2^T = 0\}$ . Let L be the Lagrangian associated to the constrained problem:  $L(y_1,y_2,\lambda) = \Phi(y_1,y_2) + \lambda H(y_1,y_2)$  being H = 0 the constraint. We must analyze the positivity of the quadratic form associated to the matrix of second partial derivatives  $\mathcal{H} = D^2 \Phi + \lambda D^2 H$  on the subspace M. By elementary computations we have for  $(\xi,\eta) \in M$ 

$$(\xi, \eta)^t \mathcal{H}(\xi, \eta) = |\lambda| (\xi^t x)^2 [|y_1|^{-2} (p_1 - 2) + (p_2 - 2)|y_2|^{-2}],$$

on the subspace M. We note that  $(-y_2, y_1) \in M$  and  $(-y_2, y_1)^t \mathcal{H}(-y_2, y_1) < 0$ . Then, by second order necessary conditions [7, p.333], at  $(y_1, y_2)$  there cannot be a minimum. Therefore, the only minima occur at  $y_1 = 0$  or  $y_2 = 0$ , then

$$A_{\Phi}(x,y) = \min\{r^{p_1}/p_1, r^{p_2}/p_2\}.$$

More generally, it holds that

$$K_1 \min\{r^{p_1}, r^{p_2}\} \le A_{\Phi} \le K_2 \min\{r^{p_1}, r^{p_2}\}$$

with  $K_1, K_2 > 0$ , for every  $1 < p_1, p_2 < \infty$ .

As is customary, we will use the decomposition  $u = \overline{u} + \widetilde{u}$  for a function  $u \in L^1([0,T])$  where  $\overline{u} = \frac{1}{T} \int_0^T u(t) \ dt$  and  $\widetilde{u} = u - \overline{u}$ .

The following lemma is an elementary generalization to anisotropic Sobolev-Orlicz spaces of known results of Sobolev spaces.

**Lemma 2.3.** Let  $\Phi : \mathbb{R}^d \to [0, +\infty)$  be a Young's function and let  $u \in W^1L^{\Phi}([0, T], \mathbb{R}^d)$ . Let  $A_{\Phi} : \mathbb{R}^+ \to \mathbb{R}^+$  be the function defined by (13). Then

1. For every  $s, t \in [0, T]$ ,  $s \neq t$ ,

$$|u(t) - u(s)| \le ||u'||_{L^{\Phi}} |s - t| A_{\Phi}^{-1} \left( \frac{1}{|s - t|} \right)$$
 (Morrey's inequality)

$$||u||_{L^{\infty}} \leqslant A_{\Phi}^{-1}\left(\frac{1}{T}\right) \max\{1, T\} ||u||_{W^1L^{\Phi}}$$
 (Sobolev's inequality)

2. We have  $\widetilde{u} \in L^{\infty}([0,T],\mathbb{R}^d)$  and

$$\|\widetilde{u}\|_{L^{\infty}} \leqslant TA_{\Phi}^{-1}\left(\frac{1}{T}\right)\|u'\|_{L^{\Phi}}$$
 (Sobolev-Wirtinger's inequality)

3. If  $\Phi$  is  $N_{\infty}$  then the space  $W^1L^{\Phi}([0,T],\mathbb{R}^d)$  is compactly embedded in the space of continuous functions  $C([0,T],\mathbb{R}^d)$ .

*Proof.* By the absolutely continuity of u, Jensen's inequality and the definition of the Luxemburg norm, we have

$$\Phi\left(\frac{u(t) - u(s)}{\|u'\|_{L^{\Phi}}|s - t|}\right) \leqslant \Phi\left(\frac{1}{|s - t|} \int_{s}^{t} \frac{u'(r)}{\|u'\|_{L^{\Phi}}} dr\right) 
\leqslant \frac{1}{|s - t|} \int_{s}^{t} \Phi\left(\frac{u'(r)}{\|u'\|_{L^{\Phi}}}\right) dr \leqslant \frac{1}{|s - t|}.$$

By Proposition 2.2(3) we have  $A_{\Phi}^{-1}\Phi(x) \ge |x|$ , therefore we get

$$\frac{|u(t) - u(s)|}{\|u'\|_{L^{\Phi}}|s - t|} \leqslant A_{\Phi}^{-1} \left(\frac{1}{|s - t|}\right),$$

then 1 holds.

Now, we use Morrey's inequality and Proposition 2.2 (2) and we have

$$|u(t) - \overline{u}| = \left| \frac{1}{T} \int_0^T u(t) - u(s) \, ds \right|$$

$$\leq \frac{1}{T} \int_0^T |u(t) - u(s)| \, ds$$

$$\leq \|u'\|_{L^{\Phi}} T A_{\Phi}^{-1} \left( \frac{1}{T} \right)$$

In order to prove the Sobolev's inequality, we note that, using Jensen's inequality and the definition of  $\|u\|_{L^{\Phi}}$ , we obtain

$$\Phi\left(\frac{\overline{u}}{\|u\|_{L^{\Phi}}}\right) \leqslant \frac{1}{T} \int_{0}^{T} \Phi\left(\frac{u(s)}{\|u\|_{L^{\Phi}}}\right) ds \leqslant \frac{1}{T}$$

Then by By Proposition 2.2(3)

$$|\overline{u}| \leqslant A_{\Phi}^{-1} \left(\frac{1}{T}\right) \|u\|_{L^{\Phi}}.$$

Therefore, from this and (Sobolev-Wirtinger's inequality) we get

$$\begin{split} \|u\|_{L^{\infty}} &\leqslant |\overline{u}| + \|\tilde{u}\|_{L^{\infty}} \\ &\leqslant A_{\Phi}^{-1} \left(\frac{1}{T}\right) \|u\|_{L^{\Phi}} + T A_{\Phi}^{-1} \left(\frac{1}{T}\right) \|u'\|_{L^{\Phi}} \\ &\leqslant A_{\Phi}^{-1} \left(\frac{1}{T}\right) \max\{1, T\} \|u\|_{W^{1}L^{\Phi}} \end{split}$$

In order to prove item 3, we take a bounded sequence  $u_n$  in  $W^1L^{\Phi}\left([0,T],\mathbb{R}^d\right)$ . Since  $\Phi$  is  $N_{\infty}$ , from Proposition 2.2(4) we obtain  $sA_{\Phi}^{-1}(1/s) \to 0$  when  $s \to 0$ . Therefore (Morrey's inequality) implies that  $u_n$  are equicontinuous. Furthermore (??) implies that  $u_n$  is bounded in  $C\left([0,T],\mathbb{R}^d\right)$ . Therefore by the Arzela-Ascoli Theorem we obtain a subsequence  $n_k$  and  $u \in C\left([0,T],\mathbb{R}^d\right)$  with  $u_{n_k} \to u$  in  $C\left([0,T],\mathbb{R}^d\right)$ .

**Lemma 2.4.** Let  $\{u_n\}_{n\in\mathbb{N}}$  be a sequence of functions in  $\Pi(E^{\Phi},1)$  converging to  $u\in\Pi(E^{\Phi},1)$  in the  $L^{\Phi}$ -norm. Then, there exist a subsequence  $u_{n_k}$  and a real valued function  $h\in L^1([0,T],\mathbb{R})$  such that  $u_{n_k}\to u$ —a.e. and  $\Phi(u_{n_k})\leqslant h$ —a.e.

*Proof.* Since  $d(u, E^{\Phi}) < 1$  and  $u_n$  converges to u, there exists  $u_0 \in E^{\Phi}$ , a subsequence of  $u_n$  (again denoted  $u_n$ ) and 0 < r < 1 such that  $d(u_n, u_0) < r$ . Let  $\lambda_0 \in (r, 1)$ . By extracting more subsequences, if necessary, we can assume that  $u_n \to u$  a.e. and

$$\lambda_n := \|u_{n+1} - u_n\|_{L^{\Phi}} < \frac{1 - \lambda_0}{2^n}, \quad \text{for } n \geqslant 1.$$

We can assume  $\lambda_n > 0$  for every  $n = 0, \ldots$ 

Let  $\lambda := 1 - \sum_{n=0}^{\infty} \lambda_n$  and define  $h : [0, T] \to \mathbb{R}$  by

$$h(x) = \lambda \Phi\left(\frac{u_0}{\lambda}\right) + \sum_{n=0}^{\infty} \lambda_n \Phi\left(\frac{u_{n+1} - u_n}{\lambda_n}\right). \tag{15}$$

Note that  $\sum_{n=0}^{\infty} \lambda_n + \lambda = 1$ , therefore for any  $n = 1, \dots$ 

$$\Phi(u_n) = \Phi\left(\lambda \frac{u_0}{\lambda} + \sum_{j=0}^{n-1} \lambda_j \frac{u_{j+1} - u_j}{\lambda_j}\right)$$

$$\leq \lambda \Phi\left(\frac{u_0}{\lambda}\right) + \sum_{j=0}^{n-1} \lambda_j \Phi\left(\frac{u_{j+1} - u_j}{\lambda_j}\right) \leq h$$

Since  $u_0 \in E^{\Phi} \subset C^{\Phi}$  and  $E^{\Phi}$  is a subspace we have that  $\Phi(u_0/\lambda) \in L^1([0,T],\mathbb{R})$ . On the other hand  $||u_{n+1} - u_n||_{L^{\Phi}} \leq \lambda_n$ , therefore

$$\int_0^T \Phi\left(\frac{u_{j+1} - u_j}{\lambda_j}\right) dt \le 1.$$

Then  $h \in L^1([0,T],\mathbb{R})$ .

# 3 Differentiability Gateâux of action integrals in anisotropic Orlicz spaces

In this section we give a brief introduction to superposition operators between anistropic Orlicz Spaces. We apply these results to obtain Gateâux differentiability of action integrals associated to lagrangian functions defined on Sobolev-Orlicz spaces.

Henceforth we assume that  $f:[0,T]\times\mathbb{R}^d\to\mathbb{R}^d$  is a Carathéodory function, i.e.

(C) f is measurable with respect to  $t \in [0,T]$  for every  $x \in \mathbb{R}^d$ , and f is a continuous function with respect to  $x \in \mathbb{R}^d$  for a.e.  $t \in [0,T]$ .

**Definition 3.1.** For  $f:[0,T]\times\mathbb{R}^d\to\mathbb{R}^d$  we denote by  $\boldsymbol{f}$  the Nemytskii (o superposition) operator defined for functions  $u:[0,T]\to\mathbb{R}^d$  by

$$fu(t) = f(t, u(t))$$

In the following Theorem we enumerate some known properties for superposition operators defined on anisotropic Orlicz spaces of vector functions. For the proofs see [5] for scalar functions and [15, 14, 13] for the generalization to  $\mathbb{R}^d$ -valued (moreover Banach spaces valued) functions in a anisotropic Orlicz Spaces (moreover modular anisotropic spaces).

**Theorem 3.2.** We assume that f satisfies condition ((C)) and that  $\Phi_1, \Phi_2 : \mathbb{R}^d \to [0, +\infty)$  are anisotropic Young functions. Then

- Measurability. The operator f maps measurable function into measurable functions
- 2. Extensibility. If the operator  ${\bf f}$  acts from the ball  $B_{L^{\Phi_1}}(r)\coloneqq\{u\in L^{\Phi_1}|\|u\|_{L^{\Phi_1}}< r\}$  into the space  $L^{\Phi_2}$  or the space  $E^{\Phi_2}$  then  ${\bf f}$  can be extended from  $\Pi(E^{\Phi_1},r)$  into space  $L^{\Phi_2}$  or  $E^{\Phi_2}$ , respectively.
- 3. Continuity. If the operator f acts from  $\Pi(E^{\Phi_1}, r)$  into space  $E^{\Phi_2}$ , then f is continuous.

Given a continuous function  $a \in C(\mathbb{R}^n, \mathbb{R}^+)$ , we define the composition operator  $a : \mathcal{M}_d \to \mathcal{M}_d$  by a(u)(x) = a(u(x)).

We will often use the following result whose proof can be performed as that of Corollary 2.3 in [1].

**Lemma 3.3.** If  $a \in C(\mathbb{R}^d, \mathbb{R}^+)$  then  $\mathbf{a} : W^1 L^{\Phi} \to L^{\infty}([0,T])$  is bounded. More concretely, there exists a non decreasing function  $A : \mathbb{R}^+ \to \mathbb{R}^+$  such that  $\|\mathbf{a}(u)\|_{L^{\infty}([0,T])} \le A(\|u\|_{W^1 L^{\Phi}})$ .

*Proof.* Let  $A \in C(\mathbb{R}^+, \mathbb{R}^+)$  be a non decreasing, continuous function defined by  $\alpha(s) := \sup_{\|x\| \le s, x \in \mathbb{R}^d} |a(x)|$ . If  $u \in W^1L_d^{\Phi}$  then, by Sobolev's inequality, for a.e.  $t \in [0, T]$ 

$$a(u(t)) \le \alpha(\|u\|_{L^{\infty}}) \le \alpha\left(A_{\Phi}^{-1}\left(\frac{1}{T}\right)\max\{1,T\}\|u\|_{W^{1}L^{\Phi}}\right) =: A(\|u\|_{W^{1}L^{\Phi}}).$$

HABRÍA QUE VER DÓNDE SE UBICA LA CONDICIÓN DE ESTRUCTURA...QUIZÁS EN LA INTRODUCCIÓN?....

Next, we deal with the differentiability of the action integral

$$I(u) = \int_0^T \mathcal{L}(t, u(t), \dot{u}(t)) dt.$$
 (16)

**Theorem 3.4.** Let  $\mathcal{L}$  be a differentiable Carathéodory function satisfying (S). Then the following statements hold:

- 1. The action integral given by (16) is finitely defined on  $\mathcal{E}^{\Phi} := W^1 L^{\Phi} \cap \{u | \dot{u} \in \Pi(E^{\Phi}, 1)\}.$
- 2. The function I is Gâteaux differentiable on  $\mathcal{E}^{\Phi}$  and its derivative I' is demicontinuous from  $\mathcal{E}^{\Phi}$  into  $\left[W^1L^{\Phi}\right]^*$ . Moreover, I' is given by the following expression

$$\langle I'(u), v \rangle = \int_0^T \left\{ D_x \mathcal{L}(t, u, \dot{u}) \cdot v + D_y \mathcal{L}(t, u, \dot{u}) \cdot \dot{v} \right\} dt. \tag{17}$$

3. If  $\Psi \in \Delta_2$  then I' is continuous from  $\mathcal{E}^{\Phi}$  into  $\left[W^1L^{\Phi}\right]^*$  when both spaces are equipped with the strong topology.

*Proof.* Let  $u \in \mathcal{E}^{\Phi}$ . As

$$\dot{u} \in \Pi(E^{\Phi}, 1) \subset C_1^{\Phi} \tag{18}$$

and (9), then  $\Phi(\dot{u}(t)) \in L^1$ . Now,

$$|\mathcal{L}(\cdot, u, \dot{u})| + |D_x \mathcal{L}(\cdot, u, \dot{u})| + \Psi(D_y \mathcal{L}(\cdot, u, \dot{u})) \leq A(\|u\|_{W^1 L^{\Phi}})(b + \Phi(\dot{u})) \in L^1,$$
(19)

by (S) and Lemma 3.3. Thus item (1) is proved.

We split up the proof of item 2 into four steps.

Step 1. The non linear operator  $u \mapsto D_x \mathcal{L}(t, u, \dot{u})$  is continuous from  $\mathcal{E}^{\Phi}$  into  $L^1([0, T])$  with the strong topology on both sets.

Let  $\{u_n\}_{n\in\mathbb{N}}$  be a sequence of functions in  $\mathcal{E}^{\Phi}$  and let  $u\in\mathcal{E}^{\Phi}$  such that  $u_n\to u$  in  $W^1L^{\Phi}$ . By (Sobolev's inequality), we have

$$|u_n(t) - u(t)| \le TA_{\Phi}^{-1}\left(\frac{1}{T}\right) ||u_n - u||_{L^{\Phi}}$$

then  $u_n \to u$  uniformly. As  $\dot{u}_n \to \dot{u} \in \mathcal{E}^{\Phi}$ , by Lemma 2.4, there exist a subsequence of  $\dot{u}_{n_k}$  (again denoted  $\dot{u}_{n_k}$ ) and a function  $h \in L^1([0,T],\mathbb{R})$  such that  $\dot{u}_{n_k} \to \dot{u}$  a.e. and  $\Phi(\dot{u}_{n_k}) \leq h$  a.e.

Since  $u_{n_k}$ , k = 1, 2, ..., is a strong convergent sequence in  $W^1L^{\Phi}$ , it is a bounded sequence in  $W^1L^{\Phi}$ . According to item (3) of Lemma 2.3, there exists M > 0 such that  $\|a(u_{n_k})\|_{L^{\infty}} \leq M$ , k = 1, 2, ... From the previous facts and (19), we get

$$|D_x \mathcal{L}(\cdot, u_{n_k}, \dot{u}_{n_k})| \leq a(|u_{n_k}|)(b + \Phi(\dot{u}_{n_k})) \leq M(b + h) \in L^1.$$

On the other hand, by the continuous differentiability of  $\mathcal{L}$ , we have

$$D_x \mathcal{L}(t, u_{n_k}(t), \dot{u}_{n_k}(t)) \to D_x \mathcal{L}(t, u(t), \dot{u}(t))$$
 for a.e.  $t \in [0, T]$ .

Applying the Dominated Convergence Theorem we conclude the proof of step 1. Step 2. The non linear operator  $u \mapsto D_y \mathcal{L}(t,u,\dot{u})$  is continuous from  $\mathcal{E}^{\Phi}$  with the strong topology into  $\left[L^{\Phi}\right]^*$  with the weak\* topology.

Let  $u \in \mathcal{E}^{\Phi}$ . From (19) it follows that

$$D_{\nu}\mathcal{L}(\cdot, u, \dot{u}) \in C^{\Psi}.$$
 (20)

Así? o conviene poner la cota de  $\Psi(D_u)$  explícitamente???

Note that (19), (20) and the imbeddings  $W^1L^{\Phi} \to L^{\infty}$  and  $L^{\Psi} \to \left[L^{\Phi}\right]^*$  imply that the second member of (17) defines an element of  $\left[W^1L^{\Phi}\right]^*$ .

Let  $u_n, u \in \mathcal{E}^{\Phi}$  such that  $u_n \to u$  in the norm of  $W^1L^{\Phi}$ . We must prove that  $D_y\mathcal{L}(\cdot, u_n, \dot{u}_n) \stackrel{w^*}{\rightharpoonup} D_y\mathcal{L}(\cdot, u, \dot{u})$ . On the contrary, there exist  $v \in L^{\Phi}$ ,  $\epsilon > 0$  and a subsequence of  $\{u_n\}$  (denoted  $\{u_n\}$  for simplicity) such that

$$|\langle D_{\nu} \mathcal{L}(\cdot, u_n, \dot{u}_n), v \rangle - \langle D_{\nu} \mathcal{L}(\cdot, u, \dot{u}), v \rangle| \ge \epsilon. \tag{21}$$

We have  $u_n \to u$  in  $L^{\Phi}$  and  $\dot{u}_n \to \dot{u}$  in  $L^{\Phi}$ . By Lemma 2.4, there exist a subsequence of  $\{u_n\}$  (again denoted  $\{u_n\}$  for simplicity) and a function  $h \in L^1([0,T],\mathbb{R})$  such that

 $u_n \to u$  uniformly,  $\dot{u}_n \to \dot{u}$  a.e. and  $\Phi(\dot{u}_n) \leqslant h$  a.e. As in the previous step, since  $u_n$  is a convergent sequence, Lemma 3.3 implies that  $a(|u_n(t)|)$  is uniformly bounded by a certain constant M > 0. Therefore, from inequality (19) with  $u_n$  instead of u, we have

$$\Psi(D_{\nu}\mathcal{L}(\cdot, u_n, \dot{u}_n)) \leq M(b+h) \in L^1.$$
(22)

As  $v \in L^{\Phi}$  there exists  $\lambda > 0$  such that  $\Phi(\frac{v}{\lambda}) \in L^1$ . Now, by Young inequality and (22), we have

$$\lambda D_{y} \mathcal{L}(\cdot, u_{n_{k}}, \dot{u}_{n_{k}}) \cdot \frac{v(t)}{\lambda}$$

$$\leq \lambda \left[ \Psi(D_{y} \mathcal{L}(\cdot, u_{n_{k}}, \dot{u}_{n_{k}})) + \Phi\left(\frac{v}{\lambda}\right) \right]$$

$$\leq \lambda M(b+h) + \lambda \Phi\left(\frac{v}{\lambda}\right) \in L^{1}$$
(23)

Finally, from the Lebesgue Dominated Convergence Theorem, we deduce

$$\int_0^T D_y \mathcal{L}(t, u_{n_k}, \dot{u}_{n_k}) \cdot v \, dt \to \int_0^T D_y \mathcal{L}(t, u, \dot{u}) \cdot v \, dt \tag{24}$$

which contradicts the inequality (21). This completes the proof of step 2.

Step 3. We will prove (17). For  $u \in \mathcal{E}^{\Phi}$  and  $0 \neq v \in W^1L^{\Phi}$ , we define the function

$$H(s,t) \coloneqq \mathcal{L}(t,u(t) + sv(t), \dot{u}(t) + s\dot{v}(t)).$$

For  $|s| \leq s_0 := \min\{\left(1 - d(\dot{u}, E^{\Phi})\right) / \|v\|_{W^1L^{\Phi}}, 1 - d(\dot{u}, E^{\Phi})\}$ , using triangle inequality we get  $d\left(\dot{u} + s\dot{v}, E^{\Phi}\right) < 1$  and thus  $\dot{u} + s\dot{v} \in \Pi(E^{\Phi}, 1)$ . These facts imply, in virtue of Theorem 3.4 item 1, that I(u + sv) is well defined and finite for  $|s| \leq s_0$ .

We also have  $\|u+sv\|_{W^1L^\Phi} \le \|u\|_{W^1L^\Phi} + s_0\|v\|_{W^1L^\Phi}$ ; then, by Lemma 3.3, there exists M>0 such that  $\|a(u+sv)\|_{L^\infty} \le M$ .

Let  $\lambda > 0$  such that  $\Phi(\frac{\dot{v}}{\lambda}) \in L^1$ . On the other hand, if  $\dot{v} \in L^{\Phi}$  and  $|s| \leq s_0 \lambda^{-1}$ , from the convexity and the parity of  $\Phi$ , we get

$$\Phi(\dot{u} + s\dot{v}) = \Phi\left((1 - s_0)\frac{\dot{u}}{1 - s_0} + s_0\frac{s}{s_0}\dot{v}\right) \leqslant (1 - s_0)\Phi\left(\frac{\dot{u}}{1 - s_0}\right) + s_0\Phi\left(\frac{s}{s_0}\dot{v}\right)$$

$$\leqslant (1 - s_0)\Phi\left(\frac{\dot{u}}{1 - s_0}\right) + s_0\Phi\left(\frac{\dot{v}}{\lambda}\right) \in L^1$$

As  $\dot{u} \in \Pi(E^{\Phi}, 1)$  then

$$d\left(\frac{\dot{u}}{1-s_0}, E^{\Phi}\right) = \frac{1}{1-s_0}d(\dot{u}, E^{\Phi}) < 1$$

and therefore  $\frac{\dot{u}}{1-s_0} \in C^{\Phi}$ .

Now, applying (19), (23), the fact that  $v \in L^{\infty}$  and  $\dot{v} \in L^{\Phi}$ , we get

$$|D_{s}H(s,t)| = \left| D_{x}\mathcal{L}(t,u+sv,\dot{u}+s\dot{v}) \cdot v + \lambda D_{y}\mathcal{L}(t,u+sv,\dot{u}+s\dot{v}) \cdot \frac{\dot{v}}{\lambda} \right|$$

$$\leq M \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] |v|$$

$$+ \lambda \left[ \Psi(D_{y}\mathcal{L}(t,u+sv,\dot{u}+s\dot{v})) + \Phi\left(\frac{\dot{v}}{\lambda}\right) \right]$$

$$\leq M \left\{ \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] |v| \right\} + \lambda M \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] + \lambda \Phi\left(\frac{\dot{v}}{\lambda}\right)$$

$$= M \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] (|v| + \lambda) + \lambda \Phi\left(\frac{\dot{v}}{\lambda}\right) \in L^{1}.$$
(25)

Consequently, I has a directional derivative and

$$\langle I'(u), v \rangle = \frac{d}{ds} I(u + sv) \Big|_{s=0} = \int_0^T \left\{ D_x \mathcal{L}(t, u, \dot{u}) \cdot v + D_y \mathcal{L}(t, u, \dot{u}) \cdot \dot{v} \right\} dt.$$

Moreover, from the previous formula, (19), (20), and Lemma 2.3, we obtain

$$|\langle I'(u), v \rangle| \leq ||D_x \mathcal{L}||_{L^1} ||v||_{L^{\infty}} + ||D_y \mathcal{L}||_{L^{\Psi}} ||\dot{v}||_{L^{\Phi}} \leq C ||v||_{W^1 L^{\Phi}}$$

with a appropriate constant C.

This completes the proof of the Gâteaux differentiability of I.

Step 4. The operator  $I': \mathcal{E}^{\Phi} \to \left[W^1L_d^{\Phi}\right]^*$  is demicontinuous. This is a consequence of the continuity of the mappings  $u \mapsto D_x \mathcal{L}(t,u,\dot{u})$  and  $u \mapsto D_y \mathcal{L}(t,u,\dot{u})$ . Indeed, if  $u_n, u \in \mathcal{E}^{\Phi}$  with  $u_n \to u$  in the norm of  $W^1L^{\Phi}$  and  $v \in W^1L^{\Phi}$ , then

$$\langle I'(u_n), v \rangle = \int_0^T \{ D_x \mathcal{L}(t, u_n, \dot{u}_n) \cdot v + D_y \mathcal{L}(t, u_n, \dot{u}_n) \cdot \dot{v} \} dt$$

$$\to \int_0^T \{ D_x \mathcal{L}(t, u, \dot{u}) \cdot v + D_y \mathcal{L}(t, u, \dot{u}) \cdot \dot{v} \} dt$$

$$= \langle I'(u), v \rangle.$$

In order to prove item 3, it is necessary to see that the maps  $u \mapsto D_x \mathcal{L}(t, u, \dot{u})$  and  $u \mapsto D_u \mathcal{L}(t, u, \dot{u})$  are norm continuous from  $\mathcal{E}^{\Phi}$  into  $L^1$  and  $L^{\Psi}$ , respectively.

The continuity of the first map has already been proved in step 1.

Si eliminamos la demicontinuidad del segundo item, hay que copiar la continuidad de  $D_x$  aquí!!!

Let  $u_n, u \in \mathcal{E}^{\Phi}$  with  $||u_n - u||_{W^1L^{\Phi}} \to 0$ .

Applying Lemma 2.4 to  $\dot{u}_n$ , there exists a subsequence (denoted  $\dot{u}_n$  for simplicity) such that  $\dot{u}_n \in L^{\Phi}$  and a function  $h \in L^1$  such that  $\Psi(\dot{u}_n) \leq h$  and  $\dot{u}_n \to \dot{u}$  a.e.

Then, by (23) we have  $\Psi(v_n) \leq m(t) \in L^1$  being  $v_n := D_y \mathcal{L}(\cdot, u_n, \dot{u}_n)$  and m(t) :=M(b+h). In addition, from the continuous differentiability of  $\mathcal{L}$ , we have that  $v_n \to v$ a.e. where  $D_{\nu}\mathcal{L}(\cdot, u, \dot{u})$ .

As  $\Psi \in \Delta_2$ , there exists  $c : \mathbb{R}^+ \to ???$  such that  $\Psi(\lambda x) \leq c(|\lambda|)\Psi(x)$ . Then,

$$\begin{split} &\Psi(\frac{v_n-v}{\lambda})\leqslant c(|\lambda|^{-1})\Psi(v_n-v) \text{ for every } \lambda\in\mathbb{R}.\\ &\quad\text{Therefore, } \Psi(\frac{v_n-v}{\lambda})\to 0 \text{ a.e. as } n\to\infty \text{ and } \Psi(\frac{v_n-v}{\lambda})\leqslant c(|\lambda|^{-1})K\Psi(v_n)+\\ &\Psi(v))\leqslant c(|\lambda|^{-1})K\big[m(t)+\Psi(v)\big])\in L^1. \end{split}$$

Now, by Dominated Convergence Theorem, we get  $\int \Psi(\frac{v_n - v}{\lambda}) dt \to 0$  for every  $\lambda > 0$ . Thus,  $v_n \to v$  in  $L^{\Psi}$ .

The continuity of I' follows from the continuity of  $D_x \mathcal{L}$  and  $D_y \mathcal{L}$  using the formula (17).

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