

# Clarke dual method for Hamiltonian systems with non standard grow

(In alphabetical order)

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## Abstract

In this paper we consider the problem of finding periodic solutions of  
certain Hamiltonian systems .....blablabla

## 1 Main problem

Let  $H : [0, T] \times \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}$ . We are looking for periodic solutions of the  
Hamiltonian system

$$\begin{cases} \dot{q}(t) &= D_p H(t, q(t), p(t)) \\ \dot{p}(t) &= -D_q H(t, q(t), p(t)) \\ p(0) &= p(T), q(0) = q(T) \end{cases} \quad (1)$$

for  $t \in [0, T]$ . I think that, like in [4], is better to present the Hamiltonian  
problem as the main problem

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**2010 AMS Subject Classification.** Primary: 34C25. Secondary: 34B15.

**Keywords and phrases.** Periodic Solutions, Orlicz Spaces, Euler-Lagrange, Critical Points.

An alternative writing of (1) using the combined variable  $u = (q, p)$  and the canonical symplectic matrix

$$J = \begin{pmatrix} 0 & I_{d \times d} \\ -I_{d \times d} & 0 \end{pmatrix}$$

is the following

$$\dot{u} = J \nabla H(t, u(t)) \quad (2)$$

or equivalently

$$J \dot{u} = -\nabla H(t, u(t)) \quad (3)$$

where  $\nabla H$  is the gradient of  $H$  with respect to the combined variable.

## 2 Preliminaries

We will use some basic concepts of convex analysis that we list below.

Let  $\Gamma_0(\mathbb{R}^d) = \{F : \mathbb{R}^d \rightarrow (-\infty, +\infty] \text{ convex, lower semicontinuous functions with non-empty effective domain.}\}$

The Fenchel conjugate of  $F$  is given by

$$F^*(p) = \sup_{q \in \mathbb{R}^d} \langle p, q \rangle - F(q)$$

The Fenchel conjugate satisfies the following properties:

1.  $F^* \in \Gamma_0(\mathbb{R}^d)$
2. If  $F \leq G$ , then  $G^* \leq F^*$ .
3. If  $G(q) = \alpha F(\beta q) + \sigma$  with  $\alpha, \beta, \sigma > 0$  then  $G^*(p) = \alpha F^*(\frac{p}{\beta \alpha}) - \sigma$

Let  $\Phi : \mathbb{R}^d \rightarrow [0, +\infty)$  be a differentiable, convex function such that  $\Phi(0) = 0$ ,  $\Phi(q) > 0$  if  $q \neq 0$ ,  $\Phi(-q) = \Phi(q)$ , and

$$\lim_{|q| \rightarrow \infty} \frac{\Phi(q)}{|q|} = +\infty, \quad (4)$$

where  $|\cdot|$  denotes the euclidean norm on  $\mathbb{R}^d$ . From now on, we say that  $\Phi$  is an  $G$ -function if  $\Phi$  satisfies the previous properties.

We write  $\Psi$  for the Fenchel conjugate of  $\Phi$ .

We do not assume that  $\Phi$  and  $\Phi'$  satisfy the  $\Delta_2$ -condition.

We denote by  $\partial F(q)$  the subdifferential of  $F$  in the sense of convex analysis (see [1, 2])

The next result is a generalization of [3, Prop. 2.2, p.34]

**Proposition 2.1.** *Let  $F \in \Gamma_0(\mathbb{R}^d)$ . Suppose that there exist an anisotropic function  $\Phi$  and non negative constants  $\beta, \gamma$  such that*

$$-\beta \leq F(q) \leq \Phi(q) + \gamma, \text{ for all } q \in \mathbb{R}^d. \quad (5)$$

Now, if  $p \in \partial F(q)$  then

$$\Psi(p) \leq \Phi(2q) + 2(\beta + \gamma). \quad (6)$$

*Proof.* If  $p \in \partial F(q)$ , from [3, Thm. 2.2, p.33],

$$F^*(p) = \langle p, q \rangle - F(q) \quad (7)$$

Conjugating (5), we have

$$F^*(p) \geq \Psi(p) - \gamma. \quad (8)$$

From Young's inequality, we get

$$\langle p, q \rangle = \frac{1}{2} \langle p, 2q \rangle \leq \frac{1}{2} \Psi(p) + \frac{1}{2} \Phi(2q) \quad (9)$$

By eqs. (5) and (7) to (9), we get

$$\Psi(p) \leq \frac{1}{2} \Psi(p) + \frac{1}{2} \Phi(2q) + \beta + \gamma$$

which implies (6) □

### 3 Optimal bounds for a symplectic bilinear form

We consider the Euclidean space  $\mathbb{R}^{2d}$  equipped with the standard symplectic structure given by bilinear canonical symplectic 2-form

$$\Omega(u, v) := \langle Ju, v \rangle.$$

Let  $\Phi$  a  $G$ -function. We consider the *symplectic  $G$ -function*  $\bar{\Phi}$  defined symplectic manifold  $\mathbb{R}^{2d}$

$$\bar{\Phi}(u) = \bar{\Phi}(q, p) := \Phi(q) + \Psi(p).$$

**I think the  $\bar{\Phi}$  is the appropriate form of the  $G$ -function defined on the symplectic manifold  $\mathbb{R}^{2n}$**

The  $G$ -function  $\bar{\Phi}$  has the following important property

$$\bar{\Phi}(Ju) = \bar{\Psi}(u). \quad (10)$$

and

$$\nabla \bar{\Phi}(Ju) = J\bar{\Psi}(u). \quad (11)$$

Here we are agreeing that  $\nabla \Phi$  is a column vector.

As a consequence of (10), the matrix  $J$  induce a isometry between the spaces  $L^{\bar{\Phi}}([0, T], \mathbb{R}^{2d})$  and  $L^{\bar{\Psi}}([0, T], \mathbb{R}^{2d})$ . Therefore we can extend  $\Omega$  to a bilinear form  $\bar{\Omega}$  on  $L^{\bar{\Phi}}([0, T], \mathbb{R}^{2d})$  of the following way

$$\overline{\Omega}(u, v) := \int_0^T \Omega(u, v) dt, \quad u, v \in L^{\overline{\Phi}}([0, T], \mathbb{R}^{2d})$$

We consider the following functional

$$\Theta(u) := \overline{\Omega}(u, \dot{u}).$$

We are interested in to find bounds of the quadratic functional  $\Theta$  of the following type

$$\theta(u) \geq -C \int_0^T \overline{\Phi}(\dot{u}) dt, \quad (12)$$

for  $u \in W^1 L^{\overline{\Phi}}([0, T], \mathbb{R}^{2d})$ . It is important to get the best constant  $C$  in previous inequality because this constant imposes restrictions to the Hamiltonian  $H$ . We call to the best constant in (12)  $C_{\overline{\Phi}}$ .

If  $\Phi(q) = |q|^2/2$  was proved in [3, Prop. 3.2] that  $C_{\Phi,1} = T/\pi$ . Below we prove that this is the optimal one. In [6, Lem. 3.3] was proved that (12) holds for  $\Phi(q) = |q|^\alpha/\alpha$ ,  $1 < \alpha < \infty$  and  $C_{\Phi} = 2T$ . Since this constant is not equal to  $T/\pi$  when  $\alpha = 2$  it is not optimal.

**Proposition 3.1.** *Let  $\Phi : \mathbb{R}^d \rightarrow [0, +\infty)$  be any  $G$ -function and  $\overline{\Phi}$ . Then (12) holds for and  $C = 2T^{-1}$  for every  $u \in W^1 L^{\overline{\Phi}}([0, T], \mathbb{R}^{2d})$ .*

*Proof.* Let  $u \in W^1 L^{\overline{\Phi}}([0, T], \mathbb{R}^{2d})$ . As is usual we write  $u = \tilde{u} + \bar{u}$  where

$$\bar{u} = \frac{1}{T} \int_0^T u(t) dt.$$

From [5, Lem. 2.4] we have that

$$\int_0^T \overline{\Phi}(\tilde{u}) dt \leq \int_0^T \overline{\Phi}(T\tilde{u}) dt.$$

Then by Young's inequality and using (10)

$$\begin{aligned} \int_0^T \Omega(\dot{u}, u) dt &= T \int_0^T \langle J\dot{u}, T^{-1}\tilde{u} \rangle dt \\ &\geq -T \left\{ \int_0^T \overline{\Psi}(J\dot{u}) dt + \int_0^T \overline{\Phi}(T^{-1}\tilde{u}) dt \right\} \\ &\geq -2T \left\{ \int_0^T \overline{\Phi}(\dot{u}) dt \right\} \end{aligned}$$

□

Clearly the constant  $2/T$  is far from optimal. A possible way to improve  $C$  is to consider other averages  $\bar{u}$ . The mean value that it was used is the standard one considered in the literature. But this value is appropriate for the Hilbert setting  $\Phi(q) = |q|^2/2$ . In this case, the value of  $\bar{u}$  is the nearest (in the  $L^2$ -norm) constant vector to  $u$ . For an arbitrary  $G$  function, it seems more reasonable to consider the nearest constant vector to  $u$  with respect to the  $\bar{\Phi}$ -integral, i.e.

$$\int_0^T \bar{\Phi}(u - \bar{u}) dt \leq \int_0^T \bar{\Phi}(u - u_0) dt, \quad \text{for every } u_0 \in \mathbb{R}^{2n}$$

Equivalently  $\bar{u}$  is characterized by

$$\int_0^T \nabla \bar{\Phi}(u - \bar{u}) dt = 0.$$

There is not an explicit formula as in the Hilbert setting.

**PROBLEM 1.** We can get a better constant taking this  $\bar{u}$ ???

Now, we produce a generalization of [3, Thm. 2.3, pp.].

**Theorem 3.2.** *Suppose that*

1.  $H : [0, T] \times \mathbb{R}^{2d} \rightarrow \mathbb{R}$  is measurable in  $t$ , continuously differentiable with respect to  $u$ .
2. there exist  $\beta, \gamma \in L^1([0, T], \mathbb{R})$ ,  $\Lambda > \lambda > 0$  such that

$$\Phi^*\left(\frac{u}{\Lambda}\right) - \beta(t) \leq H(t, u) \leq \Phi^*\left(\frac{u}{\lambda}\right) + \gamma(t) \quad (13)$$

Then there exists  $\Lambda_0$  such that the dual action

$$\chi(v) = \int_0^T \frac{1}{2} \langle J\dot{v}, v \rangle + H^*(t, \dot{v}) dt$$

is continuously differentiable in  $W^1 L_T^\Phi([0, T], \mathbb{R}^d) \cap \{u | d(u, L^\infty) < \Lambda_0\}$ .

If  $v$  is a critical point of  $\chi$  with  $d(\dot{v}, L^\infty) < \Lambda_0$ , the function defined by  $u(t) = \nabla H^*(t, \dot{v})$  solves

$$\begin{cases} \dot{u} &= J \nabla H(t, u) \\ u(T) &= u(T) \end{cases}$$

*Proof.* Conjugating 2 we obtain

$$\Phi(\lambda u) - \gamma(t) \leq H^*(t, v) \leq \Phi(\Lambda v) + \beta(t) \quad (14)$$

Since  $H^*$  is smooth, we have  $\partial_v H^*(t, v) = \{\nabla_v H^*(t, v)\}$ . Applying Proposition 2.1 with  $F = H^*$ ,  $\Phi(\Lambda v)$  instead of  $\Phi(u)$  and  $u = \nabla H^*(t, v) \in \partial_v H(t, v)$ , inequality (13) becomes

$$\Phi^*\left(\frac{\nabla H^*(t, v)}{\Lambda}\right) \leq \Phi(2\Lambda v) + 2(\beta + \gamma). \quad (15)$$

which will be the main inequality in the proof.

We are planning to obtain the structure condition (??) of [5] which guarantees differentiability.

We consider the Lagrangian

$$\mathcal{L}(t, v, \xi) = \frac{1}{2} \langle J\xi, v \rangle + H^*(t, \xi) \quad (16)$$

and we have to prove that there exist  $\Lambda_0 > \lambda_0 > 0$  such that

$$|\mathcal{L}| + |\nabla_v \mathcal{L}| + \Phi^* \left( \frac{\nabla_\xi \mathcal{L}}{\lambda_0} \right) \leq a(v) \left( b(t) + \Phi \left( \frac{\xi}{\Lambda_0} \right) \right) \quad (17)$$

We start with  $|\mathcal{L}|$ . From (14),

$$|\mathcal{L}| \leq \frac{1}{2} |\langle J\xi, v \rangle| + H^*(t, \xi) \leq \frac{1}{2} |\xi| |v| + \Phi(\Lambda\xi) + \beta(t).$$

Since  $\frac{\Phi(x)}{|x|} \rightarrow \infty$  as  $|x| \rightarrow \infty$ , there exists  $C > 0$  such that  $|x| \leq \Phi(|x|) + C$  for all  $x \in \mathbb{R}^d$ . Then,

$$|\mathcal{L}| \leq \frac{1}{2} \frac{|v|}{\Lambda} (\Phi(\Lambda\xi) + C) + \Phi(\Lambda\xi) + \beta(t) \leq \max \left\{ \Lambda, \frac{|v|}{2\Lambda} \right\} [\Phi(\Lambda\xi) + C + \beta(t)]$$

which is an estimate like the second member of (17).

Now, we treat  $|\nabla_v \mathcal{L}|$  and we get

$$|\nabla_v \mathcal{L}| = \frac{1}{2} |J\xi| \leq |\xi| \leq \frac{1}{2\Lambda} (\Phi(\Lambda\xi) + C). \quad (18)$$

which is also an estimate of the desired type.

Finally, we deal with  $\Phi(\nabla_\xi \mathcal{L} \lambda_0)$ . As  $\Phi^*$  is a convex, pair??? function, we have

$$\Phi^* \left( \frac{\nabla_\xi \mathcal{L}}{\lambda_0} \right) = \Phi^* \left( \frac{-\frac{1}{2} Jv}{\lambda_0} + \frac{\nabla H^*(t, \xi)}{\lambda_0} \right) \leq \frac{1}{2} \Phi^* \left( \frac{Jv}{\lambda_0} \right) + \frac{1}{2} \Phi^* \left( \frac{2\nabla H^*(t, \xi)}{\lambda_0} \right).$$

We choose  $\frac{2}{\lambda_0} = \frac{1}{\Lambda}$  with  $\Lambda$  as in (15) and we finally have

$$\begin{aligned} \Phi^* \left( \frac{\nabla_\xi \mathcal{L}}{\lambda_0} \right) &\leq \Phi^* \left( \frac{Jv}{2\Lambda} \right) + \Phi(2\Lambda\xi) + 2(\beta + \gamma) = \\ &\max \left\{ \Phi^* \left( \frac{Jv}{2\Lambda} \right), 1 \right\} [\Phi(2\Lambda\xi) + 2(\beta + \gamma)] \end{aligned} \quad (19)$$

which is a bound like the second member of (17).

Therefore, from (17), (18), (19) and choosing the worst functions  $a$  and  $b$ , we obtain condition (??).

Next, [5, Thm. 4.5] implies differentiability of  $\chi$  in a set like  $W^1 L_T^\Phi([0, T], \mathbb{R}^d) \cap \{u | d(\dot{u}, L^\infty) < \lambda_0\}$ .

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## References

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If  $v \in W^1 L_T^\Phi([0, T], \mathbb{R}^d)$  is a critical point of  $\chi$  with  $d(\dot{v}, L^\infty) < \lambda_0$  then, from equations (21) of [5] we obtain

$$0 = \int_0^T \frac{1}{2} \langle J\dot{v}, h \rangle - \frac{1}{2} \langle \dot{h}, Jv \rangle + \langle \nabla H^*(t, \dot{v}), \dot{h} \rangle.$$

The rest of the proof follows as in [3].

□

## Acknowledgments

The authors are partially supported by UNRC and UNLPam grants. The second author is partially supported by a UNSL grant.

## References

- [1] F. Clarke. *Optimization and Nonsmooth Analysis*. Classics in Applied Mathematics. SIAM, Philadelphia, 1990.
- [2] F. Clarke. *Functional Analysis, Calculus of Variations and Optimal Control*. Graduate Texts in Mathematics. 2013.
- [3] J. Mawhin and M. Willem. *Critical point theory and Hamiltonian systems*. Springer-Verlag, New York, 1989.
- [4] J. Mawhin and M. Willem. *Critical Point Theory and Hamiltonian Systems*. Applied Mathematical Sciences. Springer, 2010.
- [5] Fernando D Mazzone and Sonia Acinas. Periodic solutions of euler-lagrange equations in an anisotropic orlicz-sobolev space setting. *arXiv preprint arXiv:1708.06657*, 2017.
- [6] Y. Tian and W. Ge. Periodic solutions of non-autonomous second-order systems with a  $p$ -Laplacian. *Nonlinear Anal.*, 66(1):192–203, 2007.