# Periodic solutions of Euler-Lagrange equations with "sublinear nonlinearity" in an Orlicz-Sobolev space setting

Sonia Acinas \*
Instituto de Matemática Aplicada San Luis (CONICET-UNSL)
(5700) San Luis, Argentina
Universidad Nacional de La Pampa
(6300) Santa Rosa, La Pampa, Argentina
sonia.acinas@gmail.com

Fernando D. Mazzone †

Dpto. de Matemática, Facultad de Ciencias Exactas, Físico-Químicas y Naturales Universidad Nacional de Río Cuarto (5800) Río Cuarto, Córdoba, Argentina,

fmazzone@exa.unrc.edu.ar

#### Abstract

In this paper we....

ALGO ASÍ HABÍA QUE ESCRIBIR ACÁ O EN LA INTRO....

The results of this paper improve on the classic ones obtained in [1] and [2] with  $|\nabla F(t, \boldsymbol{x})| \leq f(t)\boldsymbol{x}^{\alpha} + g(t)$  for  $\alpha < 1$  and  $\alpha < p$  respectively. Here, the bounds of the sublinearity are bigger than those considered in the aforementioned works.

#### 1 Introduction

This paper is concerned with the existence of periodic solutions of the problem

$$\begin{cases} \frac{d}{dt}D_{y}\mathcal{L}(t,\boldsymbol{u}(t),\dot{\boldsymbol{u}}(t)) = D_{\boldsymbol{x}}\mathcal{L}(t,\boldsymbol{u}(t),\dot{\boldsymbol{u}}(t)) & \text{a.e. } t \in (0,T) \\ \boldsymbol{u}(0) - \boldsymbol{u}(T) = \dot{\boldsymbol{u}}(0) - \dot{\boldsymbol{u}}(T) = 0 \end{cases}$$
(1)

2010 AMS Subject Classification. Primary: . Secondary: .

Keywords and phrases. .

<sup>\*</sup>SECyT-UNRC, UNSL and CONICET

<sup>†</sup>SECyT-UNRC and CONICET

where T>0,  $\boldsymbol{u}:[0,T]\to\mathbb{R}^d$  is absolutely continuous and the Lagrangian  $\mathcal{L}:[0,T]\times\mathbb{R}^d\times\mathbb{R}^d\to\mathbb{R}$  is a Carathéodory function satisfying the conditions

$$|\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y})| \leq a(|\boldsymbol{x}|) \left(b(t) + \Phi\left(\frac{|\boldsymbol{y}|}{\lambda} + f(t)\right)\right),$$
 (2)

$$|D_{\boldsymbol{x}}\mathcal{L}(t,\boldsymbol{x},\boldsymbol{y})| \leq a(|\boldsymbol{x}|)\left(b(t) + \Phi\left(\frac{|\boldsymbol{y}|}{\lambda} + f(t)\right)\right),$$
 (3)

$$|D_{\boldsymbol{y}}\mathcal{L}(t,\boldsymbol{x},\boldsymbol{y})| \leq a(|\boldsymbol{x}|)\left(c(t) + \varphi\left(\frac{|\boldsymbol{y}|}{\lambda} + f(t)\right)\right).$$
 (4)

In these inequalities we assume that  $a\in C(\mathbb{R}^+,\mathbb{R}^+)$ ,  $\lambda>0$ ,  $\Phi$  is an N-function (see section Preliminaries for definitions),  $\varphi$  is the right continuous derivative of  $\Phi$ . The non negative functions b,c and f satisfy that  $b\in L^1_1([0,T])$ ,  $c\in L^\Psi_1([0,T])$  and  $f\in E^\Phi_1([0,T])$ , where the Banach spaces  $L^1_1([0,T])$ ,  $L^\Psi_1([0,T])$  and  $E^\Phi_1([0,T])$  will be defined later.

It is well known that problem (1) comes from a variational one, that is, a solution of (1) is a critical point of the *action integral* 

$$I(\boldsymbol{u}) = \int_0^T \mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) dt.$$
 (5)

#### 2 Preliminaries

For reader convenience, we give a short introduction to Orlicz and Orlicz-Sobolev spaces of vector valued functions and a list of results that we will use throughout the article. Classic references for Orlicz spaces of real valued functions are [3, 4, 5]. For Orlicz spaces of vector valued functions, see [6] and the references therein.

Hereafter we denote by  $\mathbb{R}^+$  the set of all non negative real numbers. A function  $\Phi: \mathbb{R}^+ \to \mathbb{R}^+$  is called an *N-function* if  $\Phi$  is given by

$$\Phi(t) = \int_0^t \varphi(\tau) \ d\tau, \quad \text{for } t \ge 0,$$

where  $\varphi: \mathbb{R}^+ \to \mathbb{R}^+$  is a right continuous non decreasing function satisfying  $\varphi(0) = 0$ ,  $\varphi(t) > 0$  for t > 0 and  $\lim_{t \to \infty} \varphi(t) = +\infty$ .

Given a function  $\varphi$  as above, we consider the so-called right inverse function  $\psi$  of  $\varphi$  which is defined by  $\psi(s)=\sup_{\varphi(t)\leqslant s}t$ . The function  $\psi$  satisfies the same properties as the function  $\varphi$ , therefore we have an N-function  $\Psi$  such that  $\Psi'=\psi$ . The function  $\Psi$  is called the *complementary function* of  $\Phi$ .

We say that  $\Phi$  satisfies the  $\Delta_2$ -condition, denoted by  $\Phi \in \Delta_2$ , if there exist constants K > 0 and  $t_0 \ge 0$  such that

$$\Phi(2t) \leqslant K\Phi(t) \tag{6}$$

for every  $t \ge t_0$ . If  $t_0 = 0$ , we say that  $\Phi$  satisfies the  $\Delta_2$ -condition globally ( $\Phi \in \Delta_2$  globally).

Let d be a positive integer. We denote by  $\mathcal{M}_d := \mathcal{M}_d([0,T])$  the set of all measurable functions defined on [0,T] with values on  $\mathbb{R}^d$  and we write  $\boldsymbol{u}=(u_1,\ldots,u_d)$  for  $\boldsymbol{u}\in\mathcal{M}_d$ . In this paper we adopt the convention that bold symbols denote points in  $\mathbb{R}^d$ .

Given an N-function  $\Phi$  we define the modular function  $\rho_{\Phi}: \mathcal{M}_d \to \mathbb{R}^+ \cup \{+\infty\}$ by

$$\rho_{\Phi}(\boldsymbol{u}) := \int_0^T \Phi(|\boldsymbol{u}|) \ dt.$$

Here  $|\cdot|$  is the euclidean norm of  $\mathbb{R}^d$ . The Orlicz class  $C_d^\Phi=C_d^\Phi([0,T])$  is given by

$$C_d^{\Phi} := \{ \boldsymbol{u} \in \mathcal{M}_d | \rho_{\Phi}(\boldsymbol{u}) < \infty \}. \tag{7}$$

The Orlicz space  $L_d^\Phi=L_d^\Phi([0,T])$  is the linear hull of  $C_d^\Phi;$  equivalently,

$$L_d^{\Phi} := \{ \boldsymbol{u} \in \mathcal{M}_d | \exists \lambda > 0 : \rho_{\Phi}(\lambda \boldsymbol{u}) < \infty \}.$$
 (8)

The Orlicz space  $L_d^{\Phi}$  equipped with the  $\mathit{Orlicz}$   $\mathit{norm}$ 

$$\|oldsymbol{u}\|_{L^\Phi} := \sup \left\{ \int_0^T oldsymbol{u} \cdot oldsymbol{v} \; dt ig| 
ho_\Psi(oldsymbol{v}) \leqslant 1 
ight\},$$

is a Banach space. By  $u \cdot v$  we denote the usual dot product in  $\mathbb{R}^d$  between u and v. The following alternative expression for the norm, known as Amemiya norm, will be useful (see [4, Thm. 10.5] and [7]). For every  $u \in L^{\Phi}$ ,

$$\|\mathbf{u}\|_{L^{\Phi}} = \inf_{k>0} \frac{1}{k} \{1 + \rho_{\Phi}(k\mathbf{u})\}.$$
 (9)

The subspace  $E_d^\Phi=E_d^\Phi([0,T])$  is defined as the closure in  $L_d^\Phi$  of the subspace  $L_d^\infty$  of all  $\mathbb{R}^d$ -valued essentially bounded functions. It is shown that  $E_d^\Phi$  is the only one maximal subspace contained in the Orlicz class  $C_d^\Phi$ , i.e.  $u\in E_d^\Phi$  if and only if  $\rho_{\Phi}(\lambda \boldsymbol{u}) < \infty$  for any  $\lambda > 0$ .

A generalized version of Hölder's inequality holds in Orlicz spaces (see [4, Th. 9.3]). Namely, if  $u \in L_d^{\Phi}$  and  $v \in L_d^{\Psi}$  then  $u \cdot v \in L_1^1$  and

$$\int_0^T \boldsymbol{v} \cdot \boldsymbol{u} \, dt \leqslant \|\boldsymbol{u}\|_{L^{\Phi}} \|\boldsymbol{v}\|_{L^{\Psi}}. \tag{10}$$

If X and Y are Banach spaces such that  $Y \subset X^*$ , we denote by  $\langle \cdot, \cdot \rangle : Y \times X \to \mathbb{R}$ the bilinear pairing map given by  $\langle x^*, x \rangle = x^*(x)$ . Hölder's inequality shows that  $L_d^\Psi \subset \left[L_d^\Phi\right]^*$ , where the pairing  $\langle {m v}, {m u} \rangle$  is defined by

$$\langle \boldsymbol{v}, \boldsymbol{u} \rangle = \int_0^T \boldsymbol{v} \cdot \boldsymbol{u} \, dt \tag{11}$$

with  $u \in L_d^{\Phi}$  and  $v \in L_d^{\Psi}$ . Unless  $\Phi \in \Delta_2$ , the relation  $L_d^{\Psi} = \left[L_d^{\Phi}\right]^*$  will not hold. In general, it is true that  $\left[E_d^\Phi\right]^*=L_d^\Psi.$  Like in [4], we will consider the subset  $\Pi(E_d^\Phi,r)$  of  $L_d^\Phi$  given by

$$\Pi(E_d^\Phi,r):=\{\boldsymbol{u}\in L_d^\Phi|d(\boldsymbol{u},E_d^\Phi)< r\}.$$

This set is related to the Orlicz class  $C_d^{\Phi}$  by means of inclusions, namely,

$$\Pi(E_d^{\Phi}, r) \subset rC_d^{\Phi} \subset \overline{\Pi(E_d^{\Phi}, r)}$$
(12)

for any positive r. If  $\Phi \in \Delta_2$ , then the sets  $L_d^\Phi$ ,  $E_d^\Phi$ ,  $\Pi(E_d^\Phi,r)$  and  $C_d^\Phi$  are equal. We define the *Sobolev-Orlicz space*  $W^1L_d^\Phi$  (see [3]) by

$$W^1L_d^{\Phi}:=\{oldsymbol{u}|oldsymbol{u} ext{ is absolutely continuous and } oldsymbol{\dot{u}}\in L_d^{\Phi}\}.$$

 $W^1L_d^{\Phi}$  is a Banach space when equipped with the norm

$$\|\mathbf{u}\|_{W^{1}L^{\Phi}} = \|\mathbf{u}\|_{L^{\Phi}} + \|\dot{\mathbf{u}}\|_{L^{\Phi}}. \tag{13}$$

For a function  $\boldsymbol{u} \in L^1_d([0,T])$ , we write  $\boldsymbol{u} = \overline{\boldsymbol{u}} + \widetilde{\boldsymbol{u}}$  where  $\overline{\boldsymbol{u}} = \frac{1}{T} \int_0^T \boldsymbol{u}(t) \ dt$  and  $\widetilde{\boldsymbol{u}} = \boldsymbol{u} - \overline{\boldsymbol{u}}$ .

As usual, if  $(X,\|\cdot\|_X)$  is a Banach space and  $(Y,\|\cdot\|_Y)$  is a subspace of X, we write  $Y\hookrightarrow X$  and we say that Y is *embedded* in X when the restricted identity map  $i_Y:Y\to X$  is bounded. That is, there exists C>0 such that for any  $y\in Y$  we have  $\|y\|_X\leqslant C\|y\|_Y$ . With this notation, Hölder's inequality states that  $L_d^\Psi\hookrightarrow \left[L_d^\Phi\right]^*$ ; and, it is easy to see that for every N-function  $\Phi$  we have that  $L_d^\infty\hookrightarrow L_d^\Phi\hookrightarrow L_d^1$ . Recall that a function  $w:\mathbb{R}^+\to\mathbb{R}^+$  is called a *modulus of continuity* if w is a

Recall that a function  $w:\mathbb{R}^+\to\mathbb{R}^+$  is called a *modulus of continuity* if w is a continuous increasing function which satisfies w(0)=0. For example, it can be easily shown that  $w(s)=s\Phi^{-1}(1/s)$  is a modulus of continuity for every N-function  $\Phi$ . We say that  $u:[0,T]\to\mathbb{R}^d$  has modulus of continuity w when there exists a constant C>0 such that

$$|\boldsymbol{u}(t) - \boldsymbol{u}(s)| \leqslant Cw(|t - s|). \tag{14}$$

We denote by  $C^w([0,T],\mathbb{R}^d)$  the space of w-Hölder continuous functions. This is the space of all functions satisfying (14) for some C>0 and it is a Banach space with norm

$$\|m{u}\|_{C^w([0,T],\mathbb{R}^d)} := \|m{u}\|_{L^\infty} + \sup_{t 
eq s} rac{|m{u}(t) - m{u}(s)|}{w(|t-s|)}.$$

An important aspect of the theory of Sobolev spaces is related to embedding theorems. There is an extensive literature on this question in the Orlicz-Sobolev space setting, see for example [8, 9, 10, 11, 12]. The next simple lemma, whose proof can be found in [13], will be used systematically.

**Lemma 2.1.** Let  $w(s) := s\Phi^{-1}(1/s)$ . Then, the following statements hold:

1.  $W^1L^\Phi\hookrightarrow C^w([0,T],\mathbb{R}^d)$  and for every  ${\boldsymbol u}\in W^1L^\Phi$ 

$$|\boldsymbol{u}(t) - \boldsymbol{u}(s)| \leq \|\dot{\boldsymbol{u}}\|_{L^{\Phi}} w(|t - s|), \tag{15}$$

$$\|u\|_{L^{\infty}} \leqslant \Phi^{-1}\left(\frac{1}{T}\right) \max\{1, T\} \|u\|_{W^{1}L^{\Phi}}$$
 (16)

2. For every  $u \in W^1L^{\Phi}$  we have  $\widetilde{u} \in L^{\infty}_d$  and

$$\|\widetilde{\boldsymbol{u}}\|_{L^{\infty}} \leqslant T\Phi^{-1}\left(\frac{1}{T}\right)\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}$$
 (Sobolev's inequality). (17)

The following result is analogous to some lemmata in  $W^1L_d^p$ , see [14].

**Lemma 2.2.** If  $\|u\|_{W^1L^{\Phi}} \to \infty$ , then  $(|\overline{u}| + \|\dot{u}\|_{L^{\Phi}}) \to \infty$ .

*Proof.* By the decomposition  $u=\overline{u}+\tilde{u}$  and some elementary operations, we get

$$\|\boldsymbol{u}\|_{L^{\Phi}} = \|\overline{\boldsymbol{u}} + \tilde{\boldsymbol{u}}\|_{L^{\Phi}} \leqslant \|\overline{\boldsymbol{u}}\|_{L^{\Phi}} + \|\tilde{\boldsymbol{u}}\|_{L^{\Phi}} = |\overline{\boldsymbol{u}}| \|1\|_{L^{\Phi}} + \|\tilde{\boldsymbol{u}}\|_{L^{\Phi}}. \tag{18}$$

It is known that  $L_d^\infty\hookrightarrow L_d^\Phi$ , i.e., there exists  $C_1=C_1(T)>0$  such that for any  $\tilde{\boldsymbol{u}}\in L_d^\infty$  we have

$$\|\tilde{\boldsymbol{u}}\|_{L^{\Phi}} \leqslant C_1 \|\tilde{\boldsymbol{u}}\|_{L^{\infty}};$$

and, applying Sobolev's inequality, we obtain the Wirtinger's inequality, that is there exists  $C_2=C_2(T)>0$  suvh that

$$\|\tilde{\boldsymbol{u}}\|_{L^{\Phi}} \leqslant C_2 \|\dot{\boldsymbol{u}}\|_{L^{\Phi}}. \tag{19}$$

Therefore, from (18), (19) and (13), we get

$$\|\boldsymbol{u}\|_{W^1L^{\Phi}} \leqslant C_3(|\overline{\boldsymbol{u}}| + \|\dot{\boldsymbol{u}}\|_{L^{\Phi}})$$

where  $C_3=C_3(T)$ . Finally, as  $\|\boldsymbol{u}\|_{W^1L^\Phi}\to\infty$  we conclude that  $(|\overline{\boldsymbol{u}}|+\|\dot{\boldsymbol{u}}\|_{L^\Phi})\to\infty$ .

We present a definition that will be useful later.

**Definition 2.3.** A function  $\mathcal{L}: [0,T] \times \mathbb{R}^d \times \mathbb{R}^d \to \mathbb{R}$  is a Carathéodory function if for fixed  $(\boldsymbol{x},\boldsymbol{y})$  the map  $t \mapsto \mathcal{L}(t,\boldsymbol{x},\boldsymbol{y})$  is measurable and for fixed t the map  $(\boldsymbol{x},\boldsymbol{y}) \mapsto \mathcal{L}(t,\boldsymbol{x},\boldsymbol{y})$  is continuous for almost everywhere  $t \in [0,T]$ . We say that  $\mathcal{L}(t,\boldsymbol{x},\boldsymbol{y})$  is differentiable Carathéodory if in addition  $\mathcal{L}(t,\boldsymbol{x},\boldsymbol{y})$  is continuously differentiable with respect to  $\boldsymbol{x}$  and  $\boldsymbol{y}$  for almost everywhere  $t \in [0,T]$ .

In [13] we proved the next results.

**Theorem 2.4.** Let  $\mathcal{L}$  be a differentiable Carathéodory function satisfying (2), (3) and (4). Then the following statements hold:

- 1. The action integral given by (5) is finitely defined on  $\mathcal{E}_d^{\Phi}(\lambda) := W^1 L_d^{\Phi} \cap \{ \boldsymbol{u} | \boldsymbol{\dot{u}} \in \Pi(E_d^{\Phi}, \lambda) \}.$
- 2. The function I is Gâteaux differentiable on  $\mathcal{E}_d^{\Phi}(\lambda)$  and its derivative I' is demicontinuous from  $\mathcal{E}_d^{\Phi}(\lambda)$  into  $\left[W^1L_d^{\Phi}\right]^*$ . Moreover, I' is given by the following expression

$$\langle I'(\boldsymbol{u}), \boldsymbol{v} \rangle = \int_0^T \left\{ D_{\boldsymbol{x}} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \boldsymbol{v} + D_{\boldsymbol{y}} \mathcal{L}(t, \boldsymbol{u}, \dot{\boldsymbol{u}}) \cdot \dot{\boldsymbol{v}} \right\} dt.$$
 (20)

3. If  $\Psi \in \Delta_2$  then I' is continuous from  $\mathcal{E}_d^{\Phi}(\lambda)$  into  $\left[W^1L_d^{\Phi}\right]^*$  when both spaces are equipped with the strong topology.

In [13] we derive the Euler-Lagrange equations associated to critical points of action integrals on the subspace of T-periodic functions. We denote by  $W^1L_T^{\Phi}$  the subspace of  $W^1L_d^{\Phi}$  containing all T-periodic functions. As usual, when Y is a subspace of the Banach space X, we denote by  $Y^{\perp}$  the *annihilator subspace* of  $X^*$ , i.e. the subspace that consists of all bounded linear functions which are identically zero on Y.

We recall that a function  $f: \mathbb{R}^d \to \mathbb{R}$  is called *strictly convex* if  $f\left(\frac{x+y}{2}\right) < \frac{1}{2}\left(f\left(x\right) + f\left(y\right)\right)$  for  $x \neq y$ . It is well known that if f is a strictly convex and differentiable function, then  $D_{\boldsymbol{x}}f: \mathbb{R}^d \to \mathbb{R}^d$  is a one-to-one map (see, e.g. [15, Thm. 12.17]).

**Theorem 2.5.** Let  $u \in \mathcal{E}_d^{\Phi}(\lambda)$  be a T-periodic function. The following statements are equivalent:

- 1.  $I'(\boldsymbol{u}) \in (W^1 L_T^{\Phi})^{\perp}$
- 2.  $D_{\boldsymbol{y}}\mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t))$  is an absolutely continuous function and  $\boldsymbol{u}$  solves the following boundary value problem

$$\begin{cases}
\frac{d}{dt}D_{y}\mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) = D_{\boldsymbol{x}}\mathcal{L}(t, \boldsymbol{u}(t), \dot{\boldsymbol{u}}(t)) & a.e. \ t \in (0, T) \\
\boldsymbol{u}(0) - \boldsymbol{u}(T) = D_{\boldsymbol{y}}\mathcal{L}(0, \boldsymbol{u}(0), \dot{\boldsymbol{u}}(0)) - D_{\boldsymbol{y}}\mathcal{L}(T, \boldsymbol{u}(T), \dot{\boldsymbol{u}}(T)) = 0.
\end{cases} (21)$$

Moreover if  $D_{\boldsymbol{y}}\mathcal{L}(t,x,y)$  is T-periodic with respect to the variable t and strictly convex with respect to  $\boldsymbol{y}$ , then  $D_{\boldsymbol{y}}\mathcal{L}(0,\boldsymbol{u}(0),\dot{\boldsymbol{u}}(0)) - D_{\boldsymbol{y}}\mathcal{L}(T,\boldsymbol{u}(T),\dot{\boldsymbol{u}}(T)) = 0$  is equivalent to  $\dot{\boldsymbol{u}}(0) = \dot{\boldsymbol{u}}(T)$ .

### Habría que ver si el lugar de los índices es el adecuado. Copié lo que teníamos en el primer trabajo.

Next, we enumerate some definitions and results from the theory of convex functions. We suggest [16, 17, 4, 18, 5] for definitions, proofs and additional details.

We denote by  $\alpha_{\varphi}$  and  $\beta_{\varphi}$  the so called *Matuszewska-Orlicz indices* of the function  $\varphi$ , which are defined next. Given an increasing, unbounded, continuous function  $\varphi:[0,+\infty)\to[0,+\infty)$  such that  $\varphi(0)=0$  we define

$$\alpha_{\varphi} := \lim_{t \to 0^{+}} \frac{\log \left( \sup_{u > 0} \frac{\varphi(tu)}{\varphi(u)} \right)}{\log(t)}, \quad \beta_{\varphi} := \lim_{t \to +\infty} \frac{\log \left( \sup_{u > 0} \frac{\varphi(tu)}{\varphi(u)} \right)}{\log(t)}. \tag{22}$$

We have that  $0 \le \alpha_{\varphi} \le \beta_{\varphi} \le +\infty$ . The relation  $\beta_{\varphi} < \infty$  holds true if and only if  $\varphi$  is a  $\Delta_2$ -function. If  $\varphi$  is a homeomorphism we have that

$$\alpha_{\varphi^{-1}} = \frac{1}{\beta_{\varphi}}.\tag{23}$$

Moreover  $\varphi \in \mathcal{F}$  implies  $\alpha_{\varphi} \geq 1$ . As a consequence,  $\varphi^{-1}$  is a  $\Delta_2$ -function.

It is well known that if  $\varphi$  is an increasing  $\Delta_2$ -function,  $\varphi$  is controlled by above and below by power functions. More concretely, for every  $\epsilon>0$  there exists a constant  $K=K(\varphi,\epsilon)$  such that, for every  $t,u\geq 0$ ,

$$K^{-1} \min \big\{ t^{\beta_{\varphi} + \epsilon}, t^{\alpha_{\varphi} - \epsilon} \big\} \varphi(u) \leqslant \varphi(tu) \leqslant K \max \big\{ t^{\beta_{\varphi} + \epsilon}, t^{\alpha_{\varphi} - \epsilon} \big\} \varphi(u). \tag{24}$$

**Lemma 2.6.** Let  $\Phi, \Psi$  complementary functions. The next statements are equivalent:

- 1.  $\Psi \in \Delta_2$  globally.
- 2. There exists an N-function  $\Phi_1 \in \Delta_2$  such that

$$\Phi(rs) \ge \Phi_1(r)\Phi(s) \text{ for every } r \ge 1, \quad s \ge 0.$$
(25)

*Proof.* 1) $\Rightarrow$ 2) As  $\Psi \in \Delta_2$  globally, there exist k > 0 and  $\nu > 1$  such that

$$\Phi(rs) \ge kr^{\nu}\Phi(s) \quad r \ge 1, \ s > 0,$$

which is (25) with  $\Phi_1(r) = kr^{\nu}$  that is an N-function that satisfies the  $\Delta_2$ -condition. 2) $\Rightarrow$ 1) Next, we follow [5, p. 32, Prop. 13] and [5, p. 29, Prop. 9]. Assume that

$$\Phi_1(r)\Phi(s) \leqslant \Phi(rs) \ r > 1, \ s \ge 0.$$

Let  $u=\Phi_1(r)\geq \Phi_1(1)$  and  $v=\Phi(s)\geq 0$ . By the well known inequality [5, p. 13, Prop. 1] and (25) we have for  $u\geq \Phi_1(1)$  and  $v\geq 0$ 

$$\frac{uv}{\Psi^{-1}(uv)} \leqslant \Phi^{-1}(uv) \leqslant \Phi_1^{-1}(u)\Phi^{-1}(v) \leqslant \frac{4uv}{\Psi_1^{-1}(u)\Psi^{-1}(v)},$$

then

$$\Psi_1^{-1}(u)\Psi^{-1}(v) \leqslant 4\Psi^{-1}(uv).$$

If we take  $x = \Psi_1^{-1}(u) \ge \Psi_1^{-1}(\Phi_1(1))$  and  $y = \Psi^{-1}(v) \ge 0$ , then

$$\Psi\left(\frac{xy}{4}\right) \leqslant \Psi_1(x)\Psi(y).$$

Now, taking  $x \ge \max\{8, \Psi_1^{-1}(\Phi_1(1))\}$  we get that  $\Psi \in \Delta_2$  globally.

The following Lemma generalizes [13, Lemma 5.2].

**Lemma 2.7.** Let  $\Phi$  be N-functions and suppose that  $\Psi \in \Delta_2$  globally. Then

$$\lim_{\|u\|_{L^{\Phi}} \to \infty} \frac{\int_{0}^{T} \Phi(|u|) dt}{\Phi_{0}(\|u\|_{L^{\Phi}})} = \infty, \tag{26}$$

for every  $\Phi_0$  with  $\Phi_0 = o(\Phi_1)$  at  $\infty$  where  $\Phi_1$  is any N-function satisfying (25). Reciprocally if (26) holds for some N-function  $\Phi_0$  then  $\Psi \in \Delta_2$  (at  $\infty$ ).

*Proof.* By the assumptions on  $\Phi$  and  $\Phi_1$  and identity (9), we have

$$\frac{\int_0^T \Phi(|u|) \, dt}{\Phi_0(\|u\|_{L^\Phi})} \ge \Phi_1(r) \frac{\int_0^T \Phi(r^{-1}|u|) \, dt}{\Phi_0(\|u\|_{L^\Phi})} \ge \frac{\Phi_1(r)}{\Phi_0(\|u\|_{L^\Phi})} \{r^{-1} \|u\|_{L^\Phi} - 1\}.$$

Now, we choose  $r=\frac{\|u\|_{L^\Phi}}{2}$ , as  $\|u\|_{L^\Phi}\to\infty$  we can assume r>1, we use the fact that  $\Phi_1\in\Delta_2$  and  $\Phi_0=o(\Phi_1)$  at  $\infty$ , and therefore we get

$$\lim_{\|u\|_{L^{\Phi}} \to \infty} \frac{\int_{0}^{T} \Phi(|u|) dt}{\Phi_{0}(\|u\|_{L^{\Phi}})} \ge \lim_{\|u\|_{L^{\Phi}} \to \infty} \frac{\Phi_{1}\left(\frac{\|u\|_{L^{\Phi}}}{2}\right)}{\Phi_{0}(\|u\|_{L^{\Phi}})} \ge C \lim_{\|u\|_{L^{\Phi}} \to \infty} \frac{\Phi_{1}(\|u\|_{L^{\Phi}})}{\Phi_{0}(\|u\|_{L^{\Phi}})} = \infty.$$

The last affirmation of the Lemma follows from the fact that if  $\Phi_0$  is N-function then  $\Phi_0(u) \geq ku$  for k small enough. Therefore (26) holds for  $\Phi_0(u) = |u|$ , then [13, Lemma 5.2] implies  $\Psi \in \Delta_2$  at  $\infty$ .

Remark 1. We point out that this Lemma can be applied to more cases than [13, Lemma 5.2]. For example, if  $\Phi(u) = u^2$ ,  $\Phi_1$  and  $\Phi_0$  are a N-functions with principal parts equal to  $u^2/\log u$  and  $u^2/(\log u)^2$  respectively (see [4, p. 16] and [4, Section 7] for the definition and properties of principal part). Then (26) holds for  $\Phi_0$ , however  $\Phi_0(u)$  is not dominated for any power function  $|u|^\alpha$  for every  $\alpha < 2$ .

We define the following functionals  $J_{C,\mu}:L^{\Phi}\to (-\infty,+\infty]$  and  $H_{C,\sigma}:\mathbb{R}^n\to\mathbb{R}$ , with  $C,\nu,\sigma>0$ , by

$$J_{C,\nu}(\boldsymbol{u}) := \rho_{\Phi}(\boldsymbol{u}) - C \|\boldsymbol{u}\|_{L^{\Phi}}^{\nu}, \tag{27}$$

and

$$H_{C,\sigma}(\boldsymbol{x}) = \int_0^T F(t, \boldsymbol{x}) dt - C|\boldsymbol{x}|^{\sigma}, \tag{28}$$

Debería ir algo así, en lugar de la anterior!!!

$$H_{C,\Phi_0}(\boldsymbol{x}) = \int_0^T F(t, \boldsymbol{x}) dt - C\Phi_0(|\boldsymbol{x}|)$$
(29)

respectively.

## 3 Lagrangians satisfying sublinear nonlinearity type conditions

Like in [13] we consider Lagrangians  $\mathcal{L}$  which are lower bounded as follows

$$\mathcal{L}(t, \boldsymbol{x}, \boldsymbol{y}) \ge \alpha_0 \Phi\left(\frac{|\boldsymbol{y}|}{\Lambda}\right) + F(t, \boldsymbol{x}).$$
 (30)

If  $\mathcal{L}$  is given by right hand side in (3) and  $\Phi(u) = |u|^2$  then the ode in (1) is quasilinear, which is is linear in  $\dot{\boldsymbol{u}}$  and non linear in  $\boldsymbol{u}$ , being  $F(t, \boldsymbol{u}(t))$  the nonlinearity. Following the literature, we refer to F as the non linearity yet we assume in just the inequality. In [1] and [2] the authors considered for p-laplacian non linearities satisfying the inequality

$$|\nabla F(t, \boldsymbol{x})| \leq b_1(t)|\boldsymbol{x}|^{\alpha} + b_2(t),$$

where  $b_1, b_2 \in L_1^1$  and  $\alpha$  is any power less than p.

Based on [19] we say that F satisfies the condition (A) if F(t, x) is a Carathéodory function and F is continuously differentiable with respect to x. Moreover, the next inequality holds

8

$$|F(t, \boldsymbol{x})| + |D_{\boldsymbol{x}}F(t, \boldsymbol{x})| \le a(|\boldsymbol{x}|)b_0(t), \quad \text{for a.e. } t \in [0, T], \forall \boldsymbol{x} \in \mathbb{R}^d.$$
 (31)

Now, we have another result about coercivity of I assuming some conditions on the nonlinearity  $\nabla F$ .

**Theorem 3.1.** Let  $\mathcal{L}$  be a lagrangian function satisfying (2), (3), (4), (3) and F satisfies condition (A). We assume the following conditions:

- 1.  $\Psi \in \Delta_2$ .
- 2. There exist non negative functions  $b_1, b_2 \in L_1^1$  such that for any  $x \in \mathbb{R}^d$  and a.e.  $t \in [0,T]$

$$|\nabla F(t, \boldsymbol{x})| \leqslant b_1(t)\varphi_0(|\boldsymbol{x}|) + b_2(t) \tag{32}$$

being  $\varphi_0 = \Phi_0'$  and  $\Phi_0 \in \Delta_2$  globally that satisfies  $\Phi_0 = o(\Phi_1)$  at  $\infty$  and where  $\Phi(rs) \ge \Phi_1(r)\Phi(s)$ .

3.

$$\Phi_0(|\boldsymbol{x}|) = o\left(\int_0^T F(t, \boldsymbol{x}) \ dt\right) \quad as \quad |\boldsymbol{x}| \to \infty. \tag{33}$$

Then the action integral I is coercive.

*Proof.* By the decomposition  $u = \overline{u} + \tilde{u}$ , Mean Value Theorem, Cauchy-Schwarz's inequality and (43), we have

$$\left| \int_{0}^{T} F(t, \boldsymbol{u}) - F(t, \overline{\boldsymbol{u}}) dt \right| = \left| \int_{0}^{T} \int_{0}^{1} \nabla F(t, \overline{\boldsymbol{u}} + s \tilde{\boldsymbol{u}}(t)) \cdot \tilde{\boldsymbol{u}}(t) ds dt \right|$$

$$\leq \int_{0}^{T} \int_{0}^{1} b_{1}(t) f(|\overline{\boldsymbol{u}} + s \tilde{\boldsymbol{u}}(t)|) |\tilde{\boldsymbol{u}}(t)| ds dt + \int_{0}^{T} \int_{0}^{1} b_{2}(t) |\tilde{\boldsymbol{u}}(t)| ds dt$$

$$= I_{1} + I_{2}.$$
(34)

On the one hand, by Hölder's inequality and Sobolev's inequality, we estimate  $I_2$  as follows

$$I_2 \leqslant ||b_2||_{L^1} ||\tilde{\boldsymbol{u}}||_{L^{\infty}} \leqslant C_1 ||\dot{\boldsymbol{u}}||_{L^{\Phi}}.$$
 (35)

where  $C_1 = C_1(\|b_2\|_{L^1}, T)$ .

On the other hand, since  $\Phi_0 \in \Delta_2$  globally, then  $\varphi \in \Delta_2$  globally and consequently  $\varphi_0$  is a subadditive function, i.e. there exists k>0 such that  $\varphi_0(a+b) \leqslant k(\varphi_0(a)+\varphi_0(b))$  for every  $a,b\geq 0$ . In this way, we have

$$\varphi_0(|\overline{\boldsymbol{u}} + s\tilde{\boldsymbol{u}}(t)|) \leqslant C(\varphi_0)[\varphi_0(|\overline{\boldsymbol{u}}|) + \varphi_0(||\tilde{\boldsymbol{u}}||_{L^{\infty}})]. \tag{36}$$

where  $s \in [0, 1]$  and  $C(\varphi_0)$  springs from the subadditivity of  $\varphi_0$ .

Now, inequality (47), Hölder's inequality, Sobolev's inequality, and the monotonicity and the subadditivity of  $\varphi_0$  imply that

$$I_{1} \leqslant C(\varphi_{0}) \left\{ \varphi_{0}(|\overline{\boldsymbol{u}}|) \|b_{1}\|_{L^{1}} \|\tilde{\boldsymbol{u}}\|_{L^{\infty}} + \|b_{1}\|_{L^{1}} \varphi_{0}(\|\tilde{\boldsymbol{u}}\|_{L^{\infty}}) \|\tilde{\boldsymbol{u}}\|_{L^{\infty}} \right\}$$

$$\leqslant C_{2} \left\{ \varphi_{0}(|\overline{\boldsymbol{u}}|) \|\dot{\boldsymbol{u}}\|_{L^{\Phi}} + \varphi_{0}(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}) \|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \right\}$$

$$(37)$$

where  $C_2 = C_2(\varphi_0, T, ||b_1||_{L^1}).$ 

Next, by Young's inequality with complementary functions  $\Phi_0$  and  $\Psi_0$  and the fact that  $\Phi_0 \in \Delta_2$  globally, we get

$$\varphi_{0}(|\overline{\boldsymbol{u}}|)\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \leqslant \Psi_{0}(\varphi_{0}(|\overline{\boldsymbol{u}}|) + \Phi_{0}(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}) 
= |\overline{\boldsymbol{u}}|\varphi_{0}(|\overline{\boldsymbol{u}}|) + \Phi_{0}(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}) 
\leqslant C(\Phi_{0})\Phi_{0}(|\overline{\boldsymbol{u}}|) + \Phi_{0}(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}})$$
(38)

and

$$\varphi_0(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}})\|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \leqslant C(\Phi_0)\Phi_0(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}),\tag{39}$$

with  $C(\Phi_0)$  the constant that comes from the  $\Delta_2$  condition on  $\Phi_0$ .

From (48), (49), (39) and (46), we have

$$I_1 + I_2 \leqslant C_3 \left\{ \Phi_0(|\overline{\boldsymbol{u}}|) + \Phi_0(||\dot{\boldsymbol{u}}||_{L^{\Phi}}) + ||\dot{\boldsymbol{u}}||_{L^{\Phi}} \right\}$$
 (40)

with  $C_3 = C_3(\Phi_0, T, ||b_1||_{L^1}, ||b_2||_{L^1}).$ 

We recall that as  $\Phi_0$  is an N-function, there exists C>0 such that  $\Phi_0(x)\geq Cx$  for every  $x\geq 1$ ; and for  $0\leqslant x\leqslant 1$ , we also have  $x\leqslant 1+C\Phi_0(x)$ . Thus,

$$x \leqslant C\Phi_0(x) + 1$$
 for every  $x \ge 0$ . (41)

#### $C = \Phi_0(1)$ ????, en las notas pero no lo veo claro

In the subsequent estimates, we use the decomposition  $u = \overline{u} + \tilde{u}$ , (3), (45), (50), (41), the fact that  $\Phi_0 \in \Delta_2$  and we get

$$I(\boldsymbol{u}) \geq \alpha_{0}\rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) + \int_{0}^{T} F(t,\boldsymbol{u}) dt$$

$$= \alpha_{0}\rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) + \int_{0}^{T} \left[F(t,\boldsymbol{u}) - F(t,\overline{\boldsymbol{u}})\right] dt + \int_{0}^{T} F(t,\overline{\boldsymbol{u}}) dt$$

$$\geq \alpha_{0}\rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) - C_{3}\Phi_{0}(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}) + \int_{0}^{T} F(t,\overline{\boldsymbol{u}}) dt - C_{3}\Phi_{0}(|\overline{\boldsymbol{u}}|) - C_{3}\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}$$

$$\geq \alpha_{0}\rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) - C_{4}\Phi_{0}(\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}) + H_{C_{3},\Phi_{0}}(\overline{\boldsymbol{u}}) - C_{3}$$

$$\geq \alpha_{0}\rho_{\Phi}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) - C_{5}\Phi_{0}\left(\frac{\|\dot{\boldsymbol{u}}\|_{L^{\Phi}}}{\Lambda}\right) + H_{C_{3},\Phi_{0}}(\overline{\boldsymbol{u}}) - C_{3}$$

$$= \alpha_{0}J_{C_{6},\Phi_{0}}\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) + H_{C_{3},\Phi_{0}}(\overline{\boldsymbol{u}}) - C_{3}$$

$$(42)$$

where  $C_4=(C+1)C_3, C_5=C_5(\Phi_0,\Lambda,C_4)$  and  $C_6=\frac{C_5}{\alpha_0}$ .

De acá en adelante, es copia del final de la otra demo. Debe adaptarse a la fórmula anterior, si es que ella logra sobrevivir....

Let  $u_n$  be a sequence in  $\mathcal{E}_d^{\Phi}(\lambda)$  with  $\|u_n\|_{W^1L^{\Phi}} \to \infty$  and we have to prove that  $I(u_n) \to \infty$ . On the contrary, suppose that for a subsequence, still denoted by

 $u_n,\ I(u_n)$  is upper bounded, i.e., there exists M>0 such that  $|I(u_n)|\leqslant M$ . As  $\|u_n\|_{W^1L^\Phi}\to\infty$ , from Lemma 2.2, we have  $|\overline{u}_n|+\|\dot{u}_n\|_{L^\Phi}\to\infty$ . Then, there exists a subsequence of  $\{u_n\}$ , still denoted by  $u_n$ , which is not bounded. Then,  $|\overline{u}_n|\to\infty$  or  $\|\dot{u}_n\|_{L^\Phi}\to\infty$ . Now, Lemma  $\ref{lem:harmonic_norm}$  implies that the functional  $J_{C_4,\mu'}(\frac{\dot{u}}{\Lambda})$  is coercive, and, by (44), the functional  $H(\overline{u})$  is also coercive, then  $J_{C_4,\mu'}(\frac{\dot{u}_n}{\Lambda})\to\infty$  or  $H(\overline{u}_n)\to\infty$ . From (31), we have that on a bounded set the functional  $H(\overline{u}_n)$  is lower bounded; and,  $J_{C_4,\mu'}(\frac{\dot{u}_n}{\Lambda})$  is also lower bounded because the modular  $\rho_\Phi\left(\frac{\dot{u}}{\Lambda}\right)$  is always bigger than zero. Therefore,  $I(u_n)\to\infty$  as  $\|u_n\|_{W^1L^\Phi}\to\infty$  which contradicts the initial assumption on the behavior of  $I(u_n)$ .

#### 3.1 versión potencias por si acaso...

**Theorem 3.2.** Let  $\mathcal{L}$  be a lagrangian function satisfying (2), (3), (4), (3) and F satisfies condition (A). We assume the following conditions:

- 1.  $\Psi \in \Delta_2$ .
- 2. There exist non negative functions  $b_1, b_2 \in L_1^1$  and a constant  $1 < \mu < \alpha_{\Phi}$  such that for any  $\mathbf{x} \in \mathbb{R}^d$  and a.e.  $t \in [0, T]$

$$|\nabla F(t, \mathbf{x})| \le b_1(t)|\mathbf{x}|^{\mu - 1} + b_2(t).$$
 (43)

3. There exists a real positive number  $\sigma$  such that  $\sigma > (\mu - 1)\beta_{\Psi}$  and

$$|\boldsymbol{x}|^{\sigma} = o\left(\int_0^T F(t, \boldsymbol{x}) dt\right) \quad as \quad |\boldsymbol{x}| \to \infty.$$
 (44)

*Then the action integral I is coercive.* 

*Proof.* By the decomposition  $u = \overline{u} + \tilde{u}$ , Mean Value Theorem, Cauchy-Schwarz's inequality and (43), we have

$$\left| \int_{0}^{T} F(t, \boldsymbol{u}) - F(t, \overline{\boldsymbol{u}}) dt \right| = \left| \int_{0}^{T} \int_{0}^{1} \nabla F(t, \overline{\boldsymbol{u}} + s \tilde{\boldsymbol{u}}(t)) \cdot \tilde{\boldsymbol{u}}(t) ds dt \right|$$

$$\leq \int_{0}^{T} \int_{0}^{1} b_{1}(t) |\overline{\boldsymbol{u}} + s \tilde{\boldsymbol{u}}(t)|^{\mu - 1} |\tilde{\boldsymbol{u}}(t)| ds dt + \int_{0}^{T} \int_{0}^{1} b_{2}(t) |\tilde{\boldsymbol{u}}(t)| ds dt$$

$$= I_{1} + I_{2}$$

$$(45)$$

On the one hand, by Hölder's inequality and Sobolev's inequality, we estimate  $I_2$  as follows

$$I_2 \leqslant \|b_2\|_{L^1} \|\tilde{\boldsymbol{u}}\|_{L^{\infty}} \leqslant C_1 \|\dot{\boldsymbol{u}}\|_{L^{\Phi}}.$$
 (46)

where  $C_1 = C_1(\|b_2\|_{L^1}, T)$ . On the other hand, as  $s \in [0, 1]$ , we have

$$|\overline{\boldsymbol{u}} + s\tilde{\boldsymbol{u}}(t)|^{\mu - 1} \leqslant C(\mu)(|\overline{\boldsymbol{u}}|^{\mu - 1} + ||\tilde{\boldsymbol{u}}||_{L^{\infty}}^{\mu - 1}). \tag{47}$$

where  $C(\mu) = 2^{\mu-2}$ , for  $\mu \ge 2$  and  $C(\mu) = 1$ , for  $1 < \mu < 2$ . Now, inequality (47), Hölder's inequality and Sobolev's inequality imply that

$$I_{1} \leqslant C(\mu) \left( |\overline{\boldsymbol{u}}|^{\mu-1} \int_{0}^{T} b_{1}(t) |\tilde{\boldsymbol{u}}(t)| dt + ||\tilde{\boldsymbol{u}}||_{L^{\infty}}^{\mu-1} \int_{0}^{T} b_{1}(t) |\tilde{\boldsymbol{u}}(t)| dt \right)$$

$$\leqslant C(\mu) \left\{ |\overline{\boldsymbol{u}}|^{\mu-1} ||b_{1}||_{L^{1}} ||\tilde{\boldsymbol{u}}||_{L^{\infty}} + ||b_{1}||_{L^{1}} ||\tilde{\boldsymbol{u}}||_{L^{\infty}}^{\mu} \right\}$$

$$\leqslant C_{2} \left\{ |\overline{\boldsymbol{u}}|^{\mu-1} ||\dot{\boldsymbol{u}}||_{L^{\Phi}} + ||\dot{\boldsymbol{u}}||_{L^{\Phi}}^{\mu} \right\}, \tag{48}$$

where  $C_2=C_2(\mu,T,\|b_1\|_{L^1})$ . Let  $\mu'$  be a positive constant such that  $1<\mu\leqslant\mu'<\alpha_\Phi$ . Next, using Young's inequality with conjugate exponents  $\mu'$  and  $\frac{\mu'}{\mu'-1}$  we get

$$|\overline{\boldsymbol{u}}|^{\mu-1} \|\dot{\boldsymbol{u}}\|_{L^{\Phi}} \leqslant \frac{(\mu'-1)}{\mu'} |\overline{\boldsymbol{u}}|^{\sigma} + \frac{1}{\mu'} \|\dot{\boldsymbol{u}}\|_{L^{\Phi}}^{\mu'}$$
 (49)

where  $\sigma = \frac{(\mu-1)\mu'}{\mu'-1}$ . We point out that  $\sigma$  is an arbitrary positive constant bigger than  $(\mu-1)b_{\Psi}$ .

From (48), (49), (46) and the inequality  $x^{r_1} \leqslant x^{r_2} + 1$ , for any  $x \ge 0$  and  $r_1 \leqslant r_2$ , we have

$$I_{1} + I_{2} \leqslant C_{3} \left\{ |\overline{\boldsymbol{u}}|^{\sigma} + ||\dot{\boldsymbol{u}}||_{L^{\Phi}}^{\mu'} + ||\dot{\boldsymbol{u}}||_{L^{\Phi}}^{\mu} + ||\dot{\boldsymbol{u}}||_{L^{\Phi}} \right\}$$

$$\leqslant C_{3} \left\{ |\overline{\boldsymbol{u}}|^{\sigma} + ||\dot{\boldsymbol{u}}||_{L^{\Phi}}^{\mu'} + 1 \right\}$$
(50)

with  $C_3 = C_3(\mu, T, ||b_1||_{L^1}, \mu')$ . In the subsequent estimates, we use the decomposition  $u = \overline{u} + \tilde{u}$ , (3), (45), (50) and we get

$$I(\boldsymbol{u}) \geq \alpha_{0} \rho_{\Phi} \left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) + \int_{0}^{T} F(t, \boldsymbol{u}) dt$$

$$= \alpha_{0} \rho_{\Phi} \left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) + \int_{0}^{T} \left[F(t, \boldsymbol{u}) - F(t, \overline{\boldsymbol{u}})\right] dt + \int_{0}^{T} F(t, \overline{\boldsymbol{u}}) dt$$

$$\geq \alpha_{0} \rho_{\Phi} \left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) - C_{3} \|\dot{\boldsymbol{u}}\|_{L^{\Phi}}^{\mu'} + \int_{0}^{T} F(t, \overline{\boldsymbol{u}}) dt - C_{3} |\overline{\boldsymbol{u}}|^{\sigma} - C_{3}$$

$$= \alpha_{0} J_{C_{4}, \mu'} \left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right) + H_{C_{3}, \sigma}(\overline{\boldsymbol{u}}) - C_{3},$$
(51)

where  $C_4 = \Lambda^{\mu'} C_3 / \alpha_0$ .

Let  $u_n$  be a sequence in  $\mathcal{E}_d^\Phi(\lambda)$  with  $\|u_n\|_{W^1L^\Phi}\to\infty$  and we have to prove that  $I(u_n)\to\infty$ . On the contrary, suppose that for a subsequence, still denoted by  $u_n$ ,  $I(u_n)$  is upper bounded, i.e., there exists M>0 such that  $|I(u_n)|\leqslant M$ . As  $\|u_n\|_{W^1L^\Phi}\to\infty$ , from Lemma 2.2, we have  $|\overline{u}_n|+\|\dot{u}_n\|_{L^\Phi}\to\infty$ . Then, there exists a subsequence of  $\{u_n\}$ , still denoted by  $u_n$ , which is not bounded. Then,  $|\overline{u}_n|\to\infty$  or

 $\|\dot{\boldsymbol{u}}_n\|_{L^\Phi} \to \infty$ . Now, Lemma  $\ref{Lemma 27}$  implies that the functional  $J_{C_4,\mu'}(\frac{\dot{\boldsymbol{u}}}{\Lambda})$  is coercive, and, by (44), the functional  $H(\overline{\boldsymbol{u}})$  is also coercive, then  $J_{C_4,\mu'}(\frac{\dot{\boldsymbol{u}}_n}{\Lambda}) \to \infty$  or  $H(\overline{\boldsymbol{u}}_n) \to \infty$ . From (31), we have that on a bounded set the functional  $H(\overline{\boldsymbol{u}}_n)$  is lower bounded; and,  $J_{C_4,\mu'}(\frac{\dot{\boldsymbol{u}}_n}{\Lambda})$  is also lower bounded because the modular  $\rho_\Phi\left(\frac{\dot{\boldsymbol{u}}}{\Lambda}\right)$  is always bigger than zero. Therefore,  $I(\boldsymbol{u}_n) \to \infty$  as  $\|\boldsymbol{u}_n\|_{W^1L^\Phi} \to \infty$  which contradicts the initial assumption on the behavior of  $I(\boldsymbol{u}_n)$ .

Leer y ver si es coherente lo anterior, si conviene trabajar siempre con  $u_n$  o habría que usar la notación de subsucesiones explícita!!!

Falta leer y corregir la sección que sigue del caso límite!!!

#### 4 Limit case $\mu = \alpha_{\Phi}$

Coercivity was obtained even in the limit case  $\mu=1$  and  $\mu=p$  (see [20, 2]) assuming additional conditions on one of the  $L^1$  functions that bound the nonlinearity. These results lean on the fact that in  $L^p$  spaces the norm and the modular coincides, that is,  $\|\cdot\|_p^p = O(\int_0^T |\cdot|^p dt)$ . In Orlicz spaces,  $\|\cdot\|_{L^\Phi}^\mu$  can be upper controlled by a modular provided that  $\mu<\alpha_\Phi$  for any N-function  $\Phi$ . Nevertheless, the limit case does not hold for any  $\Phi$ , i.e. in general  $\|\cdot\|_{L^\Phi}^{\alpha_\Phi} = O(\int_0^T \Phi(|u|) dt)$  is false as can be seen as follows.

Let  $\Phi$  be an N-function that satisfies the  $\Delta_2$ -condition. We claim that the inequality  $\Phi(tu) \geq t^{\alpha_{\Phi}}\Phi(u)$  for any u>0 and for any  $t\geq 1$  is false.

With the aim of proving the above assertion, we define the function

$$\Phi(u) = \begin{cases} \frac{p-1}{p} u^p & u \leqslant e \\ \frac{u^p}{\log u} - \frac{e^p}{p} & u > e \end{cases}$$

with p > 1.

Me parece que habría que decir algo que aclare los que se hace a continuación. Por ejemplo:

Next, we will study the behaviour of  $\Phi$ / establish some properties of  $\Phi$ /...

**Theorem 4.1.** If  $p \ge \frac{1+\sqrt{2}}{2}$ , then  $\Phi$  is an N-function.

Proof. We have

$$\Phi'(u) = \begin{cases} (p-1)u^{p-1} & := & \varphi_1(u) & \text{if } u < e \\ \frac{u^{p-1}}{\log u} (p - \frac{1}{\log u}) & := & \varphi_2(u) & \text{if } u > e \end{cases}$$

and  $\Phi$  is differentiable at e because  $\varphi_1(e) = \varphi_2(e) = (p-1)e^{p-1}$ .

We will see that  $\Phi'$  is increasing when  $p \geq \frac{1+\sqrt{2}}{2}$ . Then, it is enough to see that  $\varphi_1$  is increasing on [0,e] and  $\varphi_2$  is increasing on  $[e,\infty)$  for  $p \geq \frac{1+\sqrt{2}}{2}$ .

Tendríamos que ver que  $\varphi_1, \varphi_2$  son crecientes y que  $\varphi_2 \to \infty$  cuando  $u \to \infty$ ????? o basta con ver que  $\varphi_2$  es creciente????

 $\varphi_1$  is an increasing function for p > 1 and  $\varphi_1(u) \to 0$  as  $u \to 0$ .

On the other hand,  $\varphi_2(u) \to \infty$  as  $u \to \infty$  provided that p > 1. And,  $\varphi_2'(u) > 0$ on  $[e, \infty)$  if and only if

$$\left(p^2 - p - \frac{2p}{\log u} + \frac{1}{\log u} + \frac{2}{\log^2 u}\right) > 0.$$
 (52)

Now, as (52) holds if and only if  $p \notin (\frac{1-\sqrt{2}}{2}, \frac{1+\sqrt{2}}{2})$ , then  $\varphi_2$  is an icreasing function when  $p \geq \frac{1+\sqrt{2}}{2}$ .

Conviene agregar más cuentas en la prueba anterior, es decir, que la fórmula (52) es una cuadrática y bla...???

**Theorem 4.2.** There exists a constant C > 0 such that

$$\Phi(tu) \leqslant ct^p \Phi(u) \quad t > 1, u > 0. \tag{53}$$

For every  $\varepsilon > 0$  there exists a constant  $C = C(\varepsilon, p)$  such that

$$\Phi(tu) \ge Ct^{p-\varepsilon}\Phi(u) \quad t \ge 1, u > 0. \tag{54}$$

*Proof.* In order to prove (53), we analyze three cases.

If  $u \le tu \le e$ , then  $\Phi(tu) = t^p \Phi(u)$  and (53) holds with C = 1.

If  $u\leqslant e\leqslant tu$ , as  $\frac{e^p}{p}>0$  and  $\log(tu)\geq 1$ , we have  $\Phi(tu)\leqslant t^pu^p=\frac{p}{p-1}t^p\Phi(u)$ . Thus, (53) holds with  $C=\frac{p}{p-1}$ . If  $e\leqslant u\leqslant tu$ , then

$$\Phi(tu) \leqslant \frac{t^p u^p}{\log(tu)} \leqslant \frac{t^p u^p}{\log(u)} = \frac{pt^p v}{\log v}$$
 (55)

If  $\alpha>0$ , the function  $f(x)=\frac{x}{x-\alpha}$  is decreasing on  $(\alpha,\infty)$ . And, the function  $g(v)=\frac{pv}{\log v}$  is decreasing on  $[e^p,\infty)$ . Therefore,  $f\circ g$  is decreasing on  $[e^p,\infty)$  and we have

$$(f \circ g)(v) = \frac{\frac{pv}{\log v}}{\frac{pv}{\log v} - \frac{e^p}{n}} \leqslant e^p - \frac{e^p}{p} = \frac{p}{p-1}$$

for every  $v \geq e^p$ .

In this way, from (55), we have

$$\Phi(tu) \leqslant \frac{pt^p}{p-1} \left( \frac{pv}{\log v} - \frac{e^p}{p} \right) = \frac{pt^p}{p-1} \left( \frac{u^p}{\log u} - \frac{e^p}{p} \right)$$

and (53) holds with  $C = \frac{p}{p-1}$ . Now, we will prove (54).

If  $u \leqslant tu \leqslant e$ , (54) is immediate because  $t^p \ge t^{p-\varepsilon}$  for every  $t \ge 1$ , p > 1 and  $\varepsilon$ sufficiently small????

If  $u\leqslant e\leqslant tu$ , as  $f(t)=\frac{t}{\log t}$  is increasing on  $[e,\infty)$  then  $f(t)\geq e$  for every  $t\geq e$ . Habría que mirar en  $[e^p,\infty)$  para que  $f(t)\geq \frac{e^p}{p}$ ???? O, como f es creciente y  $e^p \leqslant (tu)^p$  entonces  $f((tu)^p) \ge f(e^p)$ ??? Now,

$$\Phi(tu) = \frac{p(tu)^p}{\log(tu)^p} - \frac{e^p}{p} = \frac{(p-1)(tu)^p}{\log(tu)^p} + \frac{(tu)^p}{\log(tu)^p} - \frac{e^p}{p} \ge \frac{p-1}{p} \frac{(tu)^p}{\log(tu)} \ge \frac{p-1}{p} \frac{t^{\varepsilon}}{\log t + 1} t^{p-\varepsilon} u^p.$$

Since  $f(t)=rac{t^{arepsilon}}{\log t+1}$  attains its minimum value  $arepsilon e^{1-arepsilon}$  at  $e^{rac{1-arepsilon}{arepsilon}}$ , then

$$\Phi(tu) \ge \frac{p-1}{p} \varepsilon e^{1-\varepsilon} t^{p-\varepsilon} u^p.$$

If  $e \leqslant u \leqslant tu$ , then

$$\Phi(tu) = \frac{p-1}{p} \frac{(tu)^p}{\log(tu)} + \frac{1}{p} \frac{(tu)^p}{\log(tu)} - \frac{e^p}{p} \geq \frac{p-1}{p} \frac{(tu)^p}{\log(tu)} = \frac{p-1}{p} \frac{t^\varepsilon \log(u)^\varepsilon}{\log(t^\varepsilon u^\varepsilon)} \frac{t^{p-\varepsilon} u^p}{\log u}$$

Let  $f(s)=\frac{sA}{\log s+1}$  with  $s\geq 1$  and  $A\geq \varepsilon$ . Then, the function f attains a minimum at  $s=e^{1-A}$ ; but, s has to be bigger than 1, then it is necessary that  $\varepsilon\leqslant A\leqslant 1$ . And, the minimum value is  $f(e^{1-A})=Ae^{1-A}\geq \varepsilon$ . If  $A\geq 1$ , f attains the minimum at s=1??? and f(1)=1. Then,  $f\geq \varepsilon$  and therefore

$$\Phi(tu) \geq \frac{p-1}{p} \varepsilon \frac{t^{p-\varepsilon} u^p}{\log u} \geq \frac{p-1}{p} \varepsilon t^{p-\varepsilon} \Phi(u).$$

#### La prueba anterior está rara.

Remark 2. The inequality

$$\Phi(tu) \ge Ct^p\Phi(u)$$

is false for every C because for every  $u \ge e$  we have

$$\lim_{t \to \infty} \frac{\Phi(tu)}{t^p \Phi(u)} = \lim_{t \to \infty} \frac{\frac{u^p}{\log(tu)} - \frac{e^p}{pt^p}}{\frac{u^p}{\log u} - \frac{e^p}{p}} = 0$$

El límite intermedio de la fórmula anterior se podría quitar.

Theorem 4.3.  $\alpha_{\Phi} = \beta_{\Phi} = p$ 

Proof. From (22) and (53), we get

$$\beta_{\Phi} = \lim_{t \to \infty} \frac{\log \left[ \sup_{u > 0} \frac{\Phi(tu)}{\Phi(u)} \right]}{\log t} \leqslant \lim_{t \to \infty} \frac{\log C + p \log t}{\log t} = p.$$

On the other hand, employing (22) and performing some elementary calculations, we obtain

$$\alpha_{\Phi} = \lim_{t \to 0^+} \frac{\log \left[ \sup_{u > 0} \frac{\Phi(tu)}{\Phi(u)} \right]}{\log t} = \lim_{s \to \infty} \frac{\log \left[ \sup_{v > 0} \frac{\Phi(v)}{\Phi(sv)} \right]^{-1}}{\log s} = \lim_{s \to \infty} \frac{\log \left[ \inf_{v > 0} \frac{\Phi(sv)}{\Phi(v)} \right]}{\log s}$$

where v := tu and  $s := \frac{1}{t}$ . Then, using (54), for every  $\varepsilon > 0$  we have

$$\alpha_{\Phi} = \lim_{s \to \infty} \frac{\log \left[ \inf_{v > 0} \frac{\Phi(sv)}{\Phi(v)} \right]}{\log s} \ge \lim_{s \to \infty} \frac{\log C + (p - \varepsilon) \log s}{\log s} \ge p - \varepsilon,$$

therefore  $\alpha_{\Phi} \geq p$ .

Finally, as 
$$\alpha_{\Phi} \leqslant \beta_{\Phi} \leqslant p$$
, we get  $\alpha_{\Phi} = \beta_{\Phi} = p$ .

Now, we are able to see that

$$\rho_{\Phi}(u) = \int_{0}^{T} \Phi(|u|) \, dx \ge C \|u\|_{L^{\Phi}}^{\alpha_{\Phi}} = C \|u\|_{L^{\Phi}}^{p}$$

is false.

In fact, if we take  $u\equiv t>0$ , then  $\|u\|_{L^\Phi}^p=C_1t^p$  where  $C_1=\|1\|_{L^\Phi}$  and  $\int_0^T\Phi(|u|)\,dx=C_2\Phi(t)$  with  $C_2=T$ . Then, if  $\rho_\Phi(u)\geq C\|u\|_{L^\Phi}^p$  were true, then  $\Phi(t)\geq Ct^p$  would also be true; however, this last inequality is false.

#### Acknowledgments

The authors are partially supported by a UNRC grant number 18/C417. The first author is partially supported by a UNSL grant number 22/F223.

#### References

- [1] C.-L. Tang, Periodic solutions for nonautonomous second order systems with sublinear nonlinearity, Proc. Amer. Math. Soc. 126 (11) (1998) 3263–3270.
- [2] X. Tang, X. Zhang, Periodic solutions for second-order Hamiltonian systems with a *p*-Laplacian, Ann. Univ. Mariae Curie-Skłodowska Sect. A 64 (1) (2010) 93–113.
- [3] R. Adams, J. Fournier, Sobolev spaces, Elsevier/Academic Press, Amsterdam, 2003.
- [4] M. A. Krasnosel'skiĭ, J. B. Rutickiĭ, Convex functions and Orlicz spaces, P. Noordhoff Ltd., Groningen, 1961.
- [5] M. M. Rao, Z. D. Ren, Theory of Orlicz spaces, Vol. 146, Marcel Dekker, Inc., New York, 1991.

- [6] G. Schappacher, A notion of Orlicz spaces for vector valued functions, Appl. Math. 50 (4) (2005) 355–386.
- [7] H. Hudzik, L. Maligranda, Amemiya norm equals Orlicz norm in general, Indag. Math. (N.S.) 11 (4) (2000) 573–585.
- [8] A. Cianchi, A fully anisotropic Sobolev inequality, Pacific J. Math. 196 (2) (2000) 283–295.
- [9] A. Cianchi, Some results in the theory of Orlicz spaces and applications to variational problems, in: Nonlinear analysis, function spaces and applications, Vol. 6 (Prague, 1998), Acad. Sci. Czech Repub., Prague, 1999, pp. 50–92.
- [10] N. Clavero, Optimal Sobolev embeddings and Function Spaces, http://www.maia.ub.edu/~soria/sobolev1.pdf, last accessed: 2014-12-22. (2011).
- [11] D. Edmunds, R. Kerman, L. Pick, Optimal Sobolev imbeddings involving rearrangement-invariant quasinorms, J. Funct. Anal. 170 (2) (2000) 307–355.
- [12] R. Kerman, L. Pick, Optimal Sobolev imbeddings, Forum Math. 18 (4) (2006) 535–570.
- [13] S. Acinas, L. Buri, G. Giubergia, F. Mazzone, E. Schwindt, Some existence results on periodic solutions of Euler-Lagrange equations in an Orlicz-Sobolev space setting, Nonlinear Analysis, TMA.(to appear).
- [14] B. Xu, C.-L. Tang, Some existence results on periodic solutions of ordinary *p*-Laplacian systems, J. Math. Anal. Appl. 333 (2) (2007) 1228–1236.
- [15] R. T. Rockafellar, R. Wets, Variational analysis, Springer-Verlag, Berlin, 1998.
- [16] A. Fiorenza, M. Krbec, Indices of Orlicz spaces and some applications, Comment. Math. Univ. Carolin. 38 (3) (1997) 433–451.
- [17] J. Gustavsson, J. Peetre, Interpolation of Orlicz spaces, Studia Mathematica 60 (1) (1977) 33–59.
  URL http://eudml.org/doc/218150
- [18] L. Maligranda, Orlicz spaces and interpolation, Vol. 5 of Seminários de Matemática [Seminars in Mathematics], Universidade Estadual de Campinas, Departamento de Matemática, Campinas, 1989.
- [19] J. Mawhin, M. Willem, Critical point theory and Hamiltonian systems, Springer-Verlag, New York, 1989.
- [20] F. Zhao, X. Wu, Periodic solutions for a class of non-autonomous second order systems, J. Math. Anal. Appl. 296 (2) (2004) 422–434.