

# Periodic solutions of Euler-Lagrange equations with “sublinear nonlinearity” in an Orlicz-Sobolev space setting

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## Abstract

In this paper we obtain existence of periodic solutions, in the Orlicz-Sobolev space  $W^1 L^\Phi([0, T])$ , of hamiltonian systems with a potential function  $F$  satisfying the inequality  $|\nabla F(t, x)| \leq b_1(t)\Phi'_0(|x|) + b_2(t)$ , with  $b_1, b_2 \in L^1$  and for certain  $N$ -functions  $\Phi_0$ .

## 1 Introduction

The purpose of this paper is to study the existence of periodic solution for the following non-autonomous second-order systems:

$$\begin{cases} \frac{d}{dt} \left( u'(t) \frac{\Phi'(|u'|)}{|u'|} \right) = \nabla F(t, u(t)) & \text{a.e. } t \in (0, T) \\ u(0) - u(T) = u'(0) - u'(T) = 0 \end{cases} \quad (1)$$

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where  $T > 0$ ,  $u : [0, T] \rightarrow \mathbb{R}^d$  is absolutely continuous and  $\Phi$  is a differentiable  $N$ -function (see section Preliminaries for definitions). Furthermore, the *potential*  $F : [0, T] \times \mathbb{R}^d \rightarrow \mathbb{R}$  satisfies the following conditions:

(C)  $F$  and its gradient  $\nabla F$  are Carathéodory functions, i.e. they are measurable functions with respect to  $t \in [0, T]$ , for every  $x \in \mathbb{R}^d$ , and continuous functions with respect to  $x \in \mathbb{R}^d$  for a.e.  $t \in [0, T]$ .

(A) For a.e.  $t \in [0, T]$ , it holds that

$$|F(t, x)| + |\nabla F(t, x)| \leq a(|x|)b(t). \quad (2)$$

In this inequality we assume that the function  $a : [0, +\infty) \rightarrow [0, +\infty)$  is continuous and nondecreasing and  $0 \leq b \in L^1([0, T], \mathbb{R})$ .

We will call the differential operator

$$L_\Phi[u] = \frac{d}{dt} \left( u'(t) \frac{\Phi'(|u'|)}{|u'|} \right)$$

the  $\Phi$ -laplacian operator. If  $\Phi(x) = |x|^p/p$ ,  $1 < p < \infty$ ,  $L_\Phi$  is the well known  $p$ -laplacian operator. In this case, we have the *Dirichlet problem* for the  $p$ -laplacian

$$\begin{cases} \frac{d}{dt} (u'(t)|u'|^{p-2}) = \nabla F(t, u(t)) & \text{a.e. } t \in (0, T) \\ u(0) - u(T) = u'(0) - u'(T) = 0 \end{cases} \quad (3)$$

The problem (1) comes from a variational one, that is, the equation in (1) is the Euler-Lagrange equation associated to the *action integral*

$$I(u) = \int_0^T \Phi(|u'(t)|) + F(t, u(t)) dt. \quad (4)$$

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The main result of this article is Theorem 3.7 which establishes conditions to guarantee existence of solutions of the problem (1) by minimization of functional (41). We point out that the hypothesis of Theorem 3.7 are generalizations of those given in [1, 2, 3, 4] about the sublinearity.

## 2 Preliminaries

For reader convenience, we give a short introduction to Orlicz and Orlicz-Sobolev spaces of vector valued functions and a list of results that we will use throughout the article. Classic references for Orlicz spaces of real valued functions are [5, 6, 7]. For Orlicz spaces of vector valued functions, see [8] and the references therein.

Hereafter we denote by  $\mathbb{R}^+$  the set of all non negative real numbers. A function  $\Phi : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  is called an  $N$ -function if  $\Phi$  is convex and satisfies that

$$\lim_{t \rightarrow +\infty} \frac{\Phi(t)}{t} = +\infty \quad \text{and} \quad \lim_{t \rightarrow 0} \frac{\Phi(t)}{t} = 0$$

In addition, in this paper we assume that  $\Phi$  is differentiable and we call  $\varphi$  the derivative of  $\Phi$ . On these assumptions,  $\varphi : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  is a homeomorphism whose inverse is  $\psi$ . We denote by  $\Psi$  the primitive of  $\psi$  that satisfies  $\Psi(0) = 0$ . Then,  $\Psi$  is an  $N$ -function which is called the *complementary function* of  $\Phi$ .

There exist several order relations between  $N$ -functions (see [7, Sec. 2.2]). Following [7, Def. 1, p. 15] we say that the  $N$ -function  $\Phi_2$  is *essentially stronger* than the  $N$ -function  $\Phi_1$  ( $\Phi_1 \ll \Phi_2$ ) if and only if there exists  $x_0 \geq 0$  such that  $\Phi_1(x) \leq \Phi_2(ax)$ , for every  $a > 0$  and  $x \geq x_0$ .

We also say that a function  $\eta : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  satisfies the  $\Delta_2$ -condition, denoted by  $\eta \in \Delta_2$ , if there exist constants  $K > 0$  and  $t_0 \geq 0$  such that

$$\eta(2t) \leq K\eta(t), \quad (5)$$

for every  $t \geq t_0$ . If  $t_0 = 0$ , a function  $\eta : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  is said to satisfy the  $\Delta_2$ -condition *globally* ( $\eta \in \Delta_2$  globally).

Let  $d$  be a positive integer. We denote by  $\mathcal{M} := \mathcal{M}([0, T], \mathbb{R}^d)$  the set of all measurable functions defined on  $[0, T]$  with values on  $\mathbb{R}^d$  and we write  $u = (u_1, \dots, u_d)$  for  $u \in \mathcal{M}$ .

Given an  $N$ -function  $\Phi$  we define the *modular function*  $\rho_\Phi : \mathcal{M} \rightarrow \mathbb{R}^+ \cup \{+\infty\}$  by

$$\rho_\Phi(u) := \int_0^T \Phi(|u|) dt.$$

Here  $|\cdot|$  is the euclidean norm of  $\mathbb{R}^d$ . The *Orlicz class*  $C^\Phi = C^\Phi([0, T], \mathbb{R}^d)$  is defined by

$$C_d^\Phi := \{u \in \mathcal{M} | \rho_\Phi(u) < \infty\}. \quad (6)$$

The *Orlicz space*  $L^\Phi = L^\Phi([0, T], \mathbb{R}^d)$  is the linear hull of  $C^\Phi$ ; equivalently,

$$L^\Phi := \{u \in \mathcal{M} | \exists \lambda > 0 : \rho_\Phi(\lambda u) < \infty\}. \quad (7)$$

The Orlicz space  $L^\Phi$  equipped with the *Orlicz norm*

$$\|u\|_{L^\Phi} := \sup \left\{ \int_0^T u \cdot v dt \mid \rho_\Phi(v) \leq 1 \right\},$$

is a Banach space. By  $u \cdot v$  we denote the usual dot product in  $\mathbb{R}^d$  between  $u$  and  $v$ . The following alternative expression for the norm, known as *Amemiya norm*, will be useful (see [6, Thm. 10.5] and [9]). For every  $u \in L^\Phi$ ,

$$\|u\|_{L^\Phi} = \inf_{k>0} \frac{1}{k} \{1 + \rho_\Phi(ku)\}. \quad (8)$$

In particular

$$\|u\|_{L^\Phi} \leq \frac{1}{k} \{1 + \rho_\Phi(ku)\}, \quad \text{for every } k > 0. \quad (9)$$

The subspace  $E^\Phi = E^\Phi([0, T], \mathbb{R}^d)$  is defined as the closure in  $L^\Phi$  of the subspace  $L_d^\infty([0, T], \mathbb{R}^d)$  of all  $\mathbb{R}^d$ -valued essentially bounded functions. It is shown that  $E^\Phi$  is

the only one maximal subspace contained in the Orlicz class  $C^\Phi$ , i.e.  $u \in E^\Phi$  if and only if  $\rho_\Phi(\lambda u) < \infty$  for any  $\lambda > 0$ .

A generalized version of *Hölder's inequality* holds in Orlicz spaces (see [6, Thm. 9.3]). Namely, if  $u \in L^\Phi$  and  $v \in L^\Psi$  then  $u \cdot v \in L^1$  and

$$\int_0^T v \cdot u \, dt \leq \|u\|_{L^\Phi} \|v\|_{L^\Psi}. \quad (10)$$

If  $X$  and  $Y$  are Banach spaces such that  $Y \subset X^*$ , we denote by  $\langle \cdot, \cdot \rangle : Y \times X \rightarrow \mathbb{R}$  the bilinear pairing map given by  $\langle x^*, x \rangle = x^*(x)$ . Hölder's inequality shows that  $L^\Psi \subset [L^\Phi]^*$ , where the pairing  $\langle v, u \rangle$  is defined by

$$\langle v, u \rangle = \int_0^T v \cdot u \, dt, \quad (11)$$

with  $u \in L^\Phi$  and  $v \in L^\Psi$ . Unless  $\Phi \in \Delta_2$ , the relation  $L^\Psi = [L^\Phi]^*$  will not hold. In general, it is true that  $[E^\Phi]^* = L^\Psi$ .

We define the *Sobolev-Orlicz space*  $W^1 L^\Phi$  (see [5]) by

$$W^1 L^\Phi := \{u \mid u \text{ is absolutely continuous on } [0, T] \text{ and } u' \in L^\Phi\}.$$

$W^1 L^\Phi$  is a Banach space when equipped with the norm

$$\|u\|_{W^1 L^\Phi} = \|u\|_{L^\Phi} + \|u'\|_{L^\Phi}. \quad (12)$$

Moreover, we introduce the following subspaces of  $W^1 L^\Phi$

$$\begin{aligned} W^1 E^\Phi &= \{u \in W^1 L^\Phi \mid u' \in E^\Phi\}, \\ W^1 E_T^\Phi &= \{u \in W^1 E^\Phi \mid u(0) = u(T)\}. \end{aligned} \quad (13)$$

For a function  $u \in L_d^1([0, T])$ , we write  $u = \bar{u} + \tilde{u}$  where  $\bar{u} = \frac{1}{T} \int_0^T u(t) \, dt$  and  $\tilde{u} = u - \bar{u}$ .

As usual, if  $(X, \|\cdot\|_X)$  is a Banach space and  $(Y, \|\cdot\|_Y)$  is a subspace of  $X$ , we write  $Y \hookrightarrow X$  and we say that  $Y$  is *embedded* in  $X$  when the restricted identity map  $i_Y : Y \rightarrow X$  is bounded. That is, there exists  $C > 0$  such that for any  $y \in Y$  we have  $\|y\|_X \leq C \|y\|_Y$ . With this notation, Hölder's inequality states that  $L^\Psi \hookrightarrow [L^\Phi]^*$ ; and, it is easy to see that for every  $N$ -function  $\Phi$  we have that  $L_d^\infty \hookrightarrow L^\Phi \hookrightarrow L_d^1$ .

Recall that a function  $w : \mathbb{R}^+ \rightarrow \mathbb{R}^+$  is called a *modulus of continuity* if  $w$  is a continuous increasing function which satisfies  $w(0) = 0$ . For example, it can be easily shown that  $w(s) = s\Phi^{-1}(1/s)$  is a modulus of continuity for every  $N$ -function  $\Phi$ . We say that  $u : [0, T] \rightarrow \mathbb{R}^d$  has modulus of continuity  $w$  when there exists a constant  $C > 0$  such that

$$|u(t) - u(s)| \leq Cw(|t - s|). \quad (14)$$

We denote by  $C^w([0, T], \mathbb{R}^d)$  the space of  $w$ -Hölder continuous functions. This is the space of all functions satisfying (14) for some  $C > 0$  and it is a Banach space with norm

$$\|u\|_{C^w([0, T], \mathbb{R}^d)} := \|u\|_{L^\infty} + \sup_{t \neq s} \frac{|u(t) - u(s)|}{w(|t - s|)}.$$

An important aspect of the theory of Sobolev spaces is related to embedding theorems. There is an extensive literature on this question in the Orlicz-Sobolev space setting, see for example [10, 11, 12, 13, 14]. The next simple lemma, whose proof can be found in [15], will be used systematically.

**Lemma 2.1.** *Let  $w(s) := s\Phi^{-1}(1/s)$ . Then, the following statements hold:*

1.  $W^1L^\Phi \hookrightarrow C^w([0, T], \mathbb{R}^d)$  and for every  $u \in W^1L^\Phi$

$$|u(t) - u(s)| \leq \|u'\|_{L^\Phi} w(|t - s|) \quad (\text{Morrey's inequality}), \quad (15)$$

$$\|u\|_{L^\infty} \leq \Phi^{-1}\left(\frac{1}{T}\right) \max\{1, T\} \|u\|_{W^1L^\Phi} \quad (\text{Sobolev's inequality}). \quad (16)$$

2. For every  $u \in W^1L^\Phi$  we have  $\tilde{u} \in L_d^\infty$  and

$$\|\tilde{u}\|_{L^\infty} \leq T\Phi^{-1}\left(\frac{1}{T}\right) \|u'\|_{L^\Phi} \quad (\text{Sobolev-Wirtinger's inequality}). \quad (17)$$

### 3 Lagrangians satisfying sublinear nonlinearity type conditions

**Lemma 3.1.**  *$E^\Phi$  is weak\* closed in  $L^\Phi$ .*

*Proof.* From [7, Thm. 7, p. 110] we have that  $L^\Phi = [E^\Psi]^*$ . Then,  $L^\Phi$  is a dual and therefore we are allowed to speak about the weak\* topology of  $L^\Phi$ . Besides,  $E^\Phi$  is separable (see [7, Thm. 1, p. 87]). Let  $S = E^\Phi \cap \{u \in L^\Phi \mid \|u\|_{L^\Phi} \leq 1\}$ , then  $S$  is closed in the norm  $\|\cdot\|_{L^\Phi}$ . Now, according to [7, Cor. 5, p. 148]  $S$  is weak\* sequentially compact. Thus,  $S$  is weak\* sequentially closed because if  $u_n \in S$  and  $u_n \xrightarrow{*} u \in L^\Phi$  then the weak\* sequential compactness implies the existence of  $v \in S$  and a subsequence  $u_{n_k}$  such that  $u_{n_k} \xrightarrow{*} v$ . Finally, by the uniqueness of the limit, we get  $u = v \in S$ . As  $E^\Psi$  is separable and  $L^\Phi = [E^\Psi]^*$ , the ball of  $L^\Phi$   $\{u \in L^\Phi \mid \|u\|_{L^\Phi} \leq 1\}$  is weak\* metrizable (see [16, Thm. 5.1, p. 138]). Thus,  $S$  is closed with respect to the weak\* topology. Now, by Krein-Smulian theorem, [16, Cor. 12.6, p. 165] implies that  $E^\Phi$  is weak\* closed.  $\square$

The following result is analogous to some lemmata in  $W^{1,p}$ , see [17].

**Lemma 3.2.** *If  $\|u\|_{W^1L^\Phi} \rightarrow \infty$ , then  $(|\bar{u}| + \|u'\|_{L^\Phi}) \rightarrow \infty$ .*

*Proof.* By the decomposition  $u = \bar{u} + \tilde{u}$  and some elementary operations, we get

$$\|u\|_{L^\Phi} = \|\bar{u} + \tilde{u}\|_{L^\Phi} \leq \|\bar{u}\|_{L^\Phi} + \|\tilde{u}\|_{L^\Phi} = |\bar{u}|1\|_{L^\Phi} + \|\tilde{u}\|_{L^\Phi}. \quad (18)$$

It is known that  $L_d^\infty \hookrightarrow L^\Phi$ , i.e. there exists  $C_1 = C_1(T) > 0$  such that for any  $\tilde{u} \in L_d^\infty$  we have

$$\|\tilde{u}\|_{L^\Phi} \leq C_1 \|\tilde{u}\|_{L^\infty};$$

and, applying Sobolev's inequality, we obtain Wirtinger's inequality, that is there exists  $C_2 = C_2(T) > 0$  such that

$$\|\tilde{u}\|_{L^\Phi} \leq C_2 \|u'\|_{L^\Phi}. \quad (19)$$

Therefore, from (18), (19) and (12), we get

$$\|u\|_{W^1 L^\Phi} \leq C_3(|\bar{u}| + \|u'\|_{L^\Phi})$$

where  $C_3 = C_3(T)$ . Finally, as  $\|u\|_{W^1 L^\Phi} \rightarrow \infty$  we conclude that  $(|\bar{u}| + \|u'\|_{L^\Phi}) \rightarrow \infty$ .  $\square$

**Lemma 3.3.** *Let  $\Phi, \Psi$  be complementary functions. The next statements are equivalent:*

1.  $\Psi \in \Delta_2$  globally.
2. There exists an  $N$ -function  $\Phi_1$  such that

$$\Phi(rs) \geq \Phi_1(r)\Phi(s) \text{ for every } r \geq 1, s \geq 0. \quad (20)$$

*Proof.* 1) $\Rightarrow$ 2) By virtue of the  $\Delta_2$ -condition on  $\Psi$ , [18, Thm. 11.7] and [18, Cor. 11.6] (see also [19, Eq. (2.8)]), we get constants  $K > 0$  and  $\alpha_\Phi > 1$  such that

$$\Phi(rs) \geq Kr^\nu \Phi(s), \quad (21)$$

for any  $1 < \nu < \alpha_\Phi$ ,  $s \geq 0$  and  $r > 1$ . This proves (20) with  $\Phi_1(r) = kr^\nu$ , which is an  $N$ -function.

2) $\Rightarrow$ 1) Next, we follow [7, p. 32, Prop. 13] and [7, p. 29, Prop. 9]. Assume that

$$\Phi_1(r)\Phi(s) \leq \Phi(rs) \text{ } r > 1, s \geq 0.$$

Let  $u = \Phi_1(r) \geq \Phi_1(1)$  and  $v = \Phi(s) \geq 0$ . By a well known inequality [7, p. 13, Prop. 1] and (20), we have for  $u \geq \Phi_1(1)$  and  $v > 0$

$$\frac{uv}{\Psi^{-1}(uv)} \leq \Phi^{-1}(uv) \leq \Phi_1^{-1}(u)\Phi^{-1}(v) \leq \frac{4uv}{\Psi_1^{-1}(u)\Psi^{-1}(v)},$$

then

$$\Psi_1^{-1}(u)\Psi^{-1}(v) \leq 4\Psi^{-1}(uv).$$

If we take  $x = \Psi_1^{-1}(u) \geq \Psi_1^{-1}(\Phi_1(1))$  and  $y = \Psi^{-1}(v) \geq 0$ , then

$$\Psi\left(\frac{xy}{4}\right) \leq \Psi_1(x)\Psi(y).$$

Now, taking  $x \geq \max\{8, \Psi_1^{-1}(\Phi_1(1))\}$  we get that  $\Psi \in \Delta_2$  globally.  $\square$

The following lemma generalizes [15, Lemma 5.2].

**Lemma 3.4.** *Let  $\Phi, \Psi$  be complementary  $N$ -functions with  $\Psi \in \Delta_2$  globally. Let  $\Phi_1$  be any  $N$ -function satisfying (20). Then*

$$\lim_{\|u\|_{L^\Phi} \rightarrow \infty} \frac{\int_0^T \Phi(|u|) dt}{\Phi_0(\|u\|_{L^\Phi})} = \infty, \quad (22)$$

for every  $\Phi_0$  with  $\Phi_0 \ll \Phi_1$ .

If (22) holds for some  $N$ -function  $\Phi_0$ , then  $\Psi \in \Delta_2$  (at  $\infty$ ).

*Proof.* By the assumptions on  $\Phi$  and  $\Phi_1$  and inequality (9), for  $r > 1$  we have

$$\frac{\int_0^T \Phi(|u|) dt}{\Phi_0(\|u\|_{L^\Phi})} \geq \Phi_1(r) \frac{\int_0^T \Phi(r^{-1}|u|) dt}{\Phi_0(\|u\|_{L^\Phi})} \geq \frac{\Phi_1(r)}{\Phi_0(\|u\|_{L^\Phi})} \{r^{-1}\|u\|_{L^\Phi} - 1\}.$$

Now, we choose  $r = \frac{\|u\|_{L^\Phi}}{2}$  and as  $\|u\|_{L^\Phi} \rightarrow \infty$  we can assume  $r > 1$  and by [7, Thm. 2 (b), p. 16].

$$\lim_{\|u\|_{L^\Phi} \rightarrow \infty} \frac{\int_0^T \Phi(|u|) dt}{\Phi_0(\|u\|_{L^\Phi})} \geq \lim_{\|u\|_{L^\Phi} \rightarrow \infty} \frac{\Phi_1\left(\frac{\|u\|_{L^\Phi}}{2}\right)}{\Phi_0(\|u\|_{L^\Phi})} = \infty.$$

Finally, if  $\Phi_0$  is an  $N$ -function, then  $\Phi_0(u) \geq k|u|$  for  $k$  small enough and  $|u| > 1$ . Therefore (22) holds for  $\Phi_0(u) = |u|$ , then [15, Lemma 5.2] implies  $\Psi \in \Delta_2$  at  $\infty$ .  $\square$

*Remark 1.* We point out that this lemma can be applied to more cases than [15, Lemma 5.2]. For example, if  $\Phi(u) = u^2$ ,  $\Phi_1$  and  $\Phi_0$  are  $N$ -functions with principal parts equal to  $u^2/\log u$  and  $u^2/(\log u)^2$  respectively (see [6, p. 16] and [6, Sec. 7] for the definition and properties of principal part), then (22) holds for  $\Phi_0$ . However,  $\Phi_0(u)$  is not dominated for any power function  $|u|^\alpha$  for every  $\alpha < 2$ .

**Definition 3.5.** We define the functionals  $J_{C, \Phi_0} : L^\Phi \rightarrow (-\infty, +\infty]$  and  $H_{C, \Phi_0} : \mathbb{R}^n \rightarrow \mathbb{R}$ , where  $C > 0$  and  $\Phi_0$  is an  $N$ -function, by

$$J_{C, \Phi_0}(u) := \rho_\Phi(u) - C\Phi_0(\|u\|_{L^\Phi}), \quad (23)$$

and

$$H_{C, \Phi_0}(x) := \int_0^T F(t, x) dt - C\Phi_0(|x|), \quad (24)$$

respectively.

In [20] and [4] the authors considered, for the  $p$ -laplacian case, potentials  $F$  satisfying the inequality

$$|\nabla F(t, x)| \leq b_1(t)|x|^\alpha + b_2(t), \quad (25)$$

where  $b_1, b_2 \in L_1^1$  and  $\alpha < p$ . Thus, they called  $F$  a sublinear nonlinearity. In this paper, we will consider bounds on  $\nabla F$  of a more general type.

**Definition 3.6.** Let  $\Phi_0$  be a differentiable  $N$ -function. We say that  $G : [0, T] \times \mathbb{R}^n \rightarrow \mathbb{R}$  satisfies a  $\Phi_0$ -grow condition if

$$|G(t, x)| \leq b_1(t)\Phi'_0(|x|) + b_2(t), \quad (26)$$

with  $b_1, b_2 \in L^1([0, T], \mathbb{R})$ .

**Theorem 3.7.** Let  $\Phi$  be an  $N$ -function whose complementary function  $\Psi$  satisfies the  $\Delta_2$  condition globally. Assume that the  $N$ -function  $\Phi_1$  satisfies (20),  $F$  satisfies (C) and (A), and  $\nabla F$  satisfies a  $\Phi_0$ -grow condition for some  $\Delta_2$ -globally  $N$ -function  $\Phi_0$  such that  $\Phi_0 \ll \Phi_1$ . Furthermore, we suppose that

$$\lim_{|x| \rightarrow \infty} \frac{\int_0^T F(t, x) dt}{\Phi_0(|x|)} = +\infty. \quad (27)$$

Then, the problem (1) has at least a solution which minimizes the action integral  $I$  on  $W^1 E_T^\Phi$ .

*Proof.* By the decomposition  $u = \bar{u} + \tilde{u}$ , Cauchy-Schwarz's inequality and (26), we have

$$\begin{aligned} \left| \int_0^T F(t, u) - F(t, \bar{u}) dt \right| &= \left| \int_0^T \int_0^1 \nabla F(t, \bar{u} + s\tilde{u}(t)) \cdot \tilde{u}(t) ds dt \right| \\ &\leq \int_0^T \int_0^1 b_1(t)\Phi'_0(|\bar{u} + s\tilde{u}(t)|)|\tilde{u}(t)| ds dt + \int_0^T \int_0^1 b_2(t)|\tilde{u}(t)| ds dt \\ &=: I_1 + I_2. \end{aligned} \quad (28)$$

On the one hand, by Hölder's and Sobolev-Wirtinger's inequalities we estimate  $I_2$  as follows

$$I_2 \leq \|b_2\|_{L^1} \|\tilde{u}\|_{L^\infty} \leq C_1 \|u'\|_{L^\Phi}, \quad (29)$$

where  $C_1 = C_1(\|b_2\|_{L^1}, T)$ .

On the other hand, since  $\Phi_0 \in \Delta_2$  globally, then  $\Phi'_0 \in \Delta_2$  globally and consequently  $\Phi'_0$  is a quasi-subadditive function, i.e. there exists  $C(\Phi_0) > 0$  such that  $\Phi'_0(a + b) \leq C(\Phi'_0)(\Phi'_0(a) + \Phi'_0(b))$  for every  $a, b \geq 0$ . In this way, we have

$$\Phi'_0(|\bar{u} + s\tilde{u}(t)|) \leq C(\Phi_0)[\Phi'_0(|\bar{u}|) + \Phi'_0(\|\tilde{u}\|_{L^\infty})], \quad (30)$$

for every  $s \in [0, 1]$ .

Now, inequality (30), Hölder's and Sobolev-Wirtinger's inequalities, the monotonicity, the subadditivity and the  $\Delta_2$ -condition on  $\Phi'_0$ , imply that

$$\begin{aligned} I_1 &\leq C(\Phi'_0) \left\{ \Phi'_0(|\bar{u}|) \|b_1\|_{L^1} \|\tilde{u}\|_{L^\infty} + \|b_1\|_{L^1} \Phi'_0(\|\tilde{u}\|_{L^\infty}) \|\tilde{u}\|_{L^\infty} \right\} \\ &\leq C_2 \left\{ \Phi'_0(|\bar{u}|) \|u'\|_{L^\Phi} + \Phi'_0(\|u'\|_{L^\Phi}) \|u'\|_{L^\Phi} \right\}, \end{aligned} \quad (31)$$

where  $C_2 = C_2(\Phi'_0, T, \|b_1\|_{L^1})$ .



Next, by Young's inequality with complementary functions  $\Phi_0$  and  $\Psi_0$  and the fact that  $\Phi_0 \in \Delta_2$  globally, Young's equality [6, Eq. 2.7-2.8] and [7, Thm. 3-(ii), p. 23], we get

$$\begin{aligned} \Phi'_0(|\bar{u}|)\|u'\|_{L^\Phi} &\leq \Psi_0(\Phi'_0(|\bar{u}|)) + \Phi_0(\|u'\|_{L^\Phi}) \\ &\leq |\bar{u}|\Phi'_0(|\bar{u}|) + \Phi_0(\|u'\|_{L^\Phi}) \\ &\leq C(\Phi_0)\Phi_0(|\bar{u}|) + \Phi_0(\|u'\|_{L^\Phi}) \end{aligned} \quad (32)$$

and

$$\Phi'_0(\|u'\|_{L^\Phi})\|u'\|_{L^\Phi} \leq C(\Phi_0)\Phi_0(\|u'\|_{L^\Phi}), \quad (33)$$

with  $C(\Phi_0)$  the constant that comes from the  $\Delta_2$ -condition on  $\Phi_0$ .

From (31), (32), (33) and (29), we have

$$\begin{aligned} I_1 + I_2 &\leq C_3 \left\{ \Phi_0(|\bar{u}|) + \Phi_0(\|u'\|_{L^\Phi}) + \|u'\|_{L^\Phi} \right\} \\ &\leq C_4 \left\{ \Phi_0(|\bar{u}|) + \Phi_0(\|u'\|_{L^\Phi}) + 1 \right\}, \end{aligned} \quad (34)$$

with  $C_3$  and  $C_4$  depending on  $\Phi_0, T, \|b_1\|_{L^1}$  and  $\|b_2\|_{L^1}$ . The last inequality follows from the fact that  $\Phi_0$  is an  $N$ -function, then there exists  $C > 0$  such that  $\Phi_0(x) \geq Cx$  for every  $x \geq 1$ . Thus  $x \leq C\Phi_0(x) + 1$  for every  $x \geq 0$ .

In the subsequent estimates, we use (28), (34), the fact that  $\Phi_0 \in \Delta_2$  and we get

$$\begin{aligned} I(u) &= \rho_\Phi(u') + \int_0^T F(t, u) dt \\ &= \rho_\Phi(u') + \int_0^T [F(t, u) - F(t, \bar{u})] dt + \int_0^T F(t, \bar{u}) dt \\ &\geq \rho_\Phi(u') - C_4\Phi_0(\|u'\|_{L^\Phi}) + \int_0^T F(t, \bar{u}) dt - C_4\Phi_0(|\bar{u}|) - C_4 \\ &\geq \rho_\Phi(u') - C_4\Phi_0(\|u'\|_{L^\Phi}) + H_{C_4, \Phi_0}(\bar{u}) - C_4 \\ &= J_{C_4, \Phi_0}(u') + H_{C_4, \Phi_0}(\bar{u}) - C_4. \end{aligned} \quad (35)$$

Let  $u_n$  be a sequence in  $\mathcal{E}_d^\Phi(\lambda)$  with  $\|u_n\|_{W^1 L^\Phi} \rightarrow \infty$  and we have to prove that  $I(u_n) \rightarrow \infty$ . On the contrary, suppose that for a subsequence, still denoted by  $u_n$ ,  $I(u_n)$  is upper bounded, i.e. there exists  $M > 0$  such that  $|I(u_n)| \leq M$ . As  $\|u_n\|_{W^1 L^\Phi} \rightarrow \infty$ , from Lemma 3.2, we have  $|\bar{u}_n| + \|u'_n\|_{L^\Phi} \rightarrow \infty$ . Passing to a subsequence is necessary, still denoted  $u_n$ , we can assume that  $|\bar{u}_n| \rightarrow \infty$  or  $\|u'_n\|_{L^\Phi} \rightarrow \infty$ . Now, Lemma 3.4 implies that the functional  $J_{C_4, \Phi_0}(u')$  is coercive; and, by (27), the functional  $H_{C_4, \Phi_0}(\bar{u})$  is also coercive, then  $J_{C_4, \Phi_0}(u'_n) \rightarrow \infty$  or  $H_{C_4, \Phi_0}(\bar{u}_n) \rightarrow \infty$ . From the condition (A) on  $F$ , we have that on a bounded set the functional  $H_{C_4, \Phi_0}(\bar{u}_n)$  is lower bounded and also  $J_{C_4, \Phi_0}(u'_n) \geq 0$ . Therefore,  $I(u_n) \rightarrow \infty$  as  $\|u_n\|_{W^1 L^\Phi} \rightarrow \infty$  which contradicts the initial assumption on the behavior of  $I(u_n)$ .

Let  $\{u_n\} \subset W^1 E_T^\Phi$  be a minimizing sequence for the problem  $\inf\{I(u) | u \in W^1 E_T^\Phi\}$ . Since  $I(u_n)$ ,  $n = 1, 2, \dots$ , is upper bounded, the previous part of the proof

shows that  $\{u_n\}$  is norm bounded in  $W^1 E^\Phi$ . Hence, by virtue of [15, Cor. 2.2], we can assume, taking a subsequence if necessary, that  $u_n$  converges uniformly to a  $T$ -periodic continuous (therefore in  $E_T^\Phi$ ) function  $u$ . As  $u'_n \in E^\Phi$  is a norm bounded sequence in  $L^\Phi$ , there exists a subsequence (again denoted by  $u'_n$ ) such that  $u'_n$  converges to a function  $v \in L^\Phi$  in the weak\* topology of  $L^\Phi$ . Since  $E^\Phi$  is weak\* closed, by Lemma 3.1,  $v \in E^\Phi$ . From this fact and the uniform convergence of  $u_n$  to  $u$ , we obtain that

$$\int_0^T \xi' \cdot u \, dt = \lim_{n \rightarrow \infty} \int_0^T \xi' \cdot u_n \, dt = - \lim_{n \rightarrow \infty} \int_0^T \xi \cdot u'_n \, dt = - \int_0^T \xi \cdot v \, dt$$

for every  $T$ -periodic function  $\xi \in C^\infty([0, T], \mathbb{R}^d) \subset E^\Psi$ . Thus  $v = u'$  a.e.  $t \in [0, T]$  (see [21, p. 6]) and  $u \in W^1 E_T^\Phi$ .

Now, taking into account the relations  $[L^1]^* = L^\infty \subset E^\Psi$  and  $L^\Phi \subset L^1$ , we have that  $u'_n$  converges to  $u'$  in the weak topology of  $L^1$ . Consequently, from the semicontinuity of  $I$  (see [15, Lemma 6.1]) we get

$$I(u) \leq \liminf_{n \rightarrow \infty} I(u_n) = \inf_{v \in W^1 E_T^\Phi} I(v).$$

Hence  $u \in W^1 E_T^\Phi$  is a minimum and, since  $I$  is Gâteaux differentiable on  $W^1 E^\Phi$  (see [15, Thm. 3.2]), therefore  $I'(u) \in (W^1 E_T^\Phi)^\perp$ . Thus,

$$\int_0^T \frac{\Phi'(|u'(t)|)}{|u'(t)|} u'(t) \cdot v'(t) dt = - \int_0^T \nabla F(t, u(t)) \cdot v(t) dt,$$

for every  $v \in W^1 E_T^\Phi$ .

From [15, Lemma 2.4] we have  $u'(t)\Phi'(|u'(t)|)/|u'(t)| \in L^\Psi([0, T], \mathbb{R}^n) \hookrightarrow L^1([0, T], \mathbb{R}^n)$ ; and, from condition (A) and the fact that  $u \in L^\infty$ , it follows that  $\nabla F(t, u(t)) \in L^1([0, T], \mathbb{R}^n)$ . Consequently, from [21, p. 6] we obtain that the differential equations in (1) are verified and  $u'(0)\Phi'(|u'(0)|)/|u'(0)| = u'(T)\Phi'(|u'(T)|)/|u'(T)|$  holds. Thus  $u'(0) = u'(T)$ .  $\square$

**Corollary 3.8.** *Let  $\mathcal{L} : [0, T] \times \mathbb{R}^d \times \mathbb{R}^d \rightarrow \mathbb{R}$  be a Carathéodory, differentiable, strictly convex function such that*

$$\mathcal{L}(t, \mathbf{x}, \mathbf{y}) \geq \Phi(|\mathbf{y}|) + F(t, \mathbf{x}), \quad (36)$$

$$|\mathcal{L}(t, \mathbf{x}, \mathbf{y})| \leq a(|\mathbf{x}|) \left( b(t) + \Phi\left(\frac{|\mathbf{y}|}{\lambda} + f(t)\right) \right), \quad (37)$$

$$|D_{\mathbf{x}} \mathcal{L}(t, \mathbf{x}, \mathbf{y})| \leq a(|\mathbf{x}|) \left( b(t) + \Phi\left(\frac{|\mathbf{y}|}{\lambda} + f(t)\right) \right), \quad (38)$$

$$|D_{\mathbf{y}} \mathcal{L}(t, \mathbf{x}, \mathbf{y})| \leq a(|\mathbf{x}|) \left( c(t) + \varphi\left(\frac{|\mathbf{y}|}{\lambda} + f(t)\right) \right). \quad (39)$$

where  $a \in C(\mathbb{R}^+, \mathbb{R}^+)$ ,  $\lambda > 0$ ,  $\Phi$  is an  $N$ -function,  $\varphi$  is the right continuous derivative of  $\Phi$ ,  $b \in L^1_1([0, T])$ ,  $c \in L^1_1([0, T])$  and  $f \in E^\Phi_1([0, T])$ . In addition, the potential  $F$  satisfies conditions (A) and (C) and **convex** and  $\Psi \in \Delta_2$  ????

Then, the problem

$$\begin{cases} \frac{d}{dt} D_{\mathbf{y}} \mathcal{L}(t, \mathbf{u}(t), \dot{\mathbf{u}}(t)) = D_{\mathbf{x}} \mathcal{L}(t, \mathbf{u}(t), \dot{\mathbf{u}}(t)) & \text{a.e. } t \in (0, T) \\ \mathbf{u}(0) - \mathbf{u}(T) = \dot{\mathbf{u}}(0) - \dot{\mathbf{u}}(T) = 0 \end{cases} \quad (40)$$

has at least one solution  $u : [0, T] \rightarrow \mathbb{R}^d$  absolutely continuous, which minimizes the action integral

$$I(u) = \int_0^T \mathcal{L}(t, u(t), u'(t)) dt. \quad (41)$$

**Quizás hay demasiada información en el enunciado y se puede quitar algo....**

*Proof.* In Theorem 3.7 we have seen that  $\Phi(|y|) + F(t, x)$  is coercive, then the function  $\mathcal{L}$  do so. The differentiability and lower semicontinuity of the  $I$  follow from Thm. 3.2 and Lemma 6.1 of [15], respectively. Last, applying Thm. 4.1 and Thm. 6.2 of [15], we get the existence of  $u$  and the initial conditions on  $u$ .  $\square$

## 4 Examples growing conditions with power functions

The employment of  $N$ -functions instead of power functions in inequalities like (26) will allow us to extend some results of [20] and [4] a  $\Phi$ -laplacian operators with  $N$ -functions  $\Phi$  which grow faster than power functions, for example with an exponential growth. Furthermore, we extend some results of [20] and [4] even in the case of  $p$ -laplacian operator (3), because we get bounds that may be sharper than those in [20, 4]. More precisely, in [4, Th. 2.1] X. Tang and X. Zhang obtained existences of solutions of (3) under the assumption (25) for any  $\alpha \in (0, p - 1)$ . Meanwhile, our Theorem 3.7 implies existence for the potential

$$F(t, x) = |x|^p / \sqrt{\ln(2 + |x|)}.$$

We note that this  $F$  does not satisfy (25) for any  $\alpha < p - 1$ . Next we will show a  $N$ -function  $\Phi$  satisfying the hypothesis of Theorem 3.7 for the potential  $F$ .

As it has been seen in the proof of Lemma 3.3 and in [15, Lemma 5.2], we may assume that the function  $\Phi_0$  in Lemma 3.3 is given by  $\Phi_0(x) = |x|^\mu$  with  $0 < \mu < \alpha_\Phi$  and where  $\alpha_\Phi$  is a Matuszewska-Orlicz index (see [18, Ch. 11]). These indices are defined by

$$\alpha_\Phi := \lim_{t \rightarrow 0^+} \frac{\log \left( \sup_{u > 0} \frac{\Phi(tu)}{\Phi(u)} \right)}{\log(t)}, \quad \beta_\Phi := \lim_{t \rightarrow +\infty} \frac{\log \left( \sup_{u > 0} \frac{\Phi(tu)}{\Phi(u)} \right)}{\log(t)}. \quad (42)$$

Hence, following the same lines as the proof of Theorem 3.7, using [15, Lemma 5.2] instead of our Lemma 3.3, we can assume that  $\Phi_0(x) = |x|^\mu$  with  $0 < \mu < \alpha_\Phi$  in Theorem 3.7. **(Se repite el Thm 3.7 y los dos párrafos están un poco enredados!!!).**

The employment of  $N$ -functions instead of power functions in inequalities like (26) will allow us to extend some results of [20] and [4], not only to the  $\Phi$ -laplacian operator, but even in the case of  $p$ -laplacian operator we get bounds that may be sharper than those in [20, 4]. More precisely, in [4, Thm. 2.1] X. Tang and X. Zhang obtained existences of solutions for the  $p$ -laplacian operator under the assumption

$$|\nabla F| \leq b_1(t)|x|^\alpha + b_2(t)$$

for any  $\alpha \in (0, p-1)$ .

Assuming  $\|b_1\|_{L^1}$  small enough, in [2, 4] coercivity was obtained even for the limit value  $\alpha = p-1$  in inequality (26).

This result leans on the fact that

$$\|u\|_{L^\Phi}^{\alpha_\Phi} = O\left(\int_0^T \Phi(|u|) dt\right) \quad \text{for } \|u\|_{L^\Phi} \rightarrow \infty, \quad (43)$$

when  $\Phi(u) = |u|^p$ . Nevertheless, it is no longer the case for any  $N$ -function  $\Phi$  as the following example shows.

In this section, from now on we will suppose that

$$\Phi(u) = \begin{cases} \frac{p-1}{p} u^p & u \leq e \\ \frac{u^p}{\log u} - \frac{e^p}{p} & u > e \end{cases}$$

with  $p > 1$ . Next, we will establish some properties of this function  $\Phi$ .

**Theorem 4.1.** *If  $p \geq \frac{1+\sqrt{2}}{2}$ , then  $\Phi$  is an  $N$ -function.*

*Proof.* We have

$$\varphi(u) = \Phi'(u) = \begin{cases} (p-1)u^{p-1} & := \varphi_1(u) \quad \text{if } u \leq e \\ \frac{u^{p-1}}{\log u} (p - \frac{1}{\log u}) & := \varphi_2(u) \quad \text{if } u \geq e \end{cases}$$

First let us see that  $\Phi'$  is increasing when  $p \geq \frac{1+\sqrt{2}}{2}$ . For this purpose, since  $\varphi_1(e) = \varphi_2(e)$ , it is enough to see that  $\varphi_1$  is increasing on  $[0, e]$  and  $\varphi_2$  is increasing on  $[e, \infty)$  for every  $p \geq \frac{1+\sqrt{2}}{2}$ . Clearly  $\varphi_1$  is an increasing function for  $p > 1$ . On the other hand, an elementary analysis of the function shows that  $\varphi_2'(u) > 0$  on  $[e, \infty)$  if and only if  $p \notin (\frac{1-\sqrt{2}}{2}, \frac{1+\sqrt{2}}{2})$ . Therefore  $\varphi_2$  is an increasing function when  $p \geq \frac{1+\sqrt{2}}{2}$ .

Besides  $\varphi_2(u) \rightarrow \infty$  and  $\varphi_1(u) \rightarrow 0$  as  $u \rightarrow \infty$  and  $u \rightarrow 0$  respectively, provided that  $p > 1$ . Hence,  $\Phi$  is an  $N$ -function.  $\square$

**Theorem 4.2.** *For every  $\varepsilon > 0$ , there exists a positive constant  $C = C(p, \varepsilon)$  such that*

$$C^{-1}t^{p-\varepsilon}\Phi(u) \leq \Phi(tu) \leq Ct^p\Phi(u) \quad t \geq 1, u > 0, \quad (44)$$

*Proof.* If  $u \leq tu \leq e$ , then  $\Phi(tu) = t^p\Phi(u)$  and (44) holds with  $C = 1$ .

If  $u \leq e \leq tu$ , as  $\frac{e^p}{p} > 0$  and  $\log(tu) \geq 1$ , we have  $\Phi(tu) \leq t^p u^p = \frac{p}{p-1} t^p \Phi(u)$ . Thus, the second inequality of (44) holds with  $C = \frac{p}{p-1}$ . On the other hand, as  $f(t) = \frac{t}{\log t}$  is increasing on  $[e, \infty)$ , then  $f((tu)^p) \geq f(e^p) = e^p/p$ . Now,

$$\begin{aligned} \Phi(tu) &= \frac{p(tu)^p}{\log(tu)^p} - \frac{e^p}{p} \\ &= \frac{(p-1)(tu)^p}{\log(tu)^p} + \frac{(tu)^p}{\log(tu)^p} - \frac{e^p}{p} \\ &\geq \frac{p-1}{p} \frac{(tu)^p}{\log(tu)} \\ &\geq \frac{p-1}{p} \frac{t^\varepsilon}{\log t + 1} t^{p-\varepsilon} u^p. \end{aligned}$$

Since  $\varepsilon e^{1-\varepsilon}$  is the minimum value of  $t \mapsto \frac{t^\varepsilon}{\log t+1}$  on the interval  $[1, +\infty)$  then

$$\Phi(tu) \geq \frac{p-1}{p} \varepsilon e^{1-\varepsilon} t^{p-\varepsilon} u^p,$$

which is the first inequality of (44) with  $C = \frac{p}{p-1} \varepsilon^{-1} e^{-1+\varepsilon}$ .

If  $e \leq u \leq tu$ , then

$$\Phi(tu) \leq \frac{t^p u^p}{\log(tu)} \leq \frac{t^p u^p}{\log(u)} = \frac{pt^p v}{\log v}, \quad (45)$$

where  $v := u^p$  and  $v \geq e^p$ . If  $\alpha > 0$ , the function  $x \mapsto \frac{x}{x-\alpha}$  is decreasing on  $(\alpha, \infty)$  and the function  $v \mapsto \frac{pv}{\log v}$  is increasing on  $[e^p, \infty)$ . Therefore, we have

$$\frac{\frac{pv}{\log v}}{\frac{pv}{\log v} - \frac{e^p}{p}} \leq \frac{e^p}{e^p - \frac{e^p}{p}} = \frac{p}{p-1}$$

for every  $v \geq e^p$ . In this way, from (45), we have

$$\Phi(tu) \leq \frac{pt^p}{p-1} \left( \frac{pv}{\log v} - \frac{e^p}{p} \right) = \frac{pt^p}{p-1} \left( \frac{u^p}{\log u} - \frac{e^p}{p} \right)$$

and the second inequality of (44) holds with  $C = \frac{p}{p-1}$ . For the first inequality we have, as it was proved previously,

$$\Phi(tu) \geq \frac{p-1}{p} \frac{(tu)^p}{\log(tu)} = \frac{p-1}{p} \frac{t^\varepsilon \log u^\varepsilon}{\log(t^\varepsilon u^\varepsilon)} \frac{t^{p-\varepsilon} u^p}{\log u}$$

Let  $f(s) = \frac{sA}{\log s+A}$  with  $s \geq 1$  and  $A \geq \varepsilon$ . If  $A \leq 1$ , the function  $f$  attains a minimum on  $[1, \infty)$  at  $s = e^{1-A}$  and the minimum value is  $f(e^{1-A}) = Ae^{1-A} \geq \varepsilon$ . If  $A > 1$ ,  $f$  is increasing on  $[1, \infty)$  and its minimum value is  $f(1) = 1$ . Then,  $f(s) \geq \varepsilon$  in any case, therefore

$$\Phi(tu) \geq \frac{p-1}{p} \varepsilon \frac{t^{p-\varepsilon} u^p}{\log u} \geq \frac{p-1}{p} \varepsilon t^{p-\varepsilon} \Phi(u).$$

Therefore, (44) holds with  $C = \frac{p}{\varepsilon(p-1)}$ , because this  $C$  is the biggest constant that we have obtained in each case under consideration.  $\square$

*Remark 2.* The inequality

$$\Phi(tu) \geq Ct^p \Phi(u)$$

is false for every  $C$  because for every  $u \geq e$  we have

$$\lim_{t \rightarrow \infty} \frac{\Phi(tu)}{t^p \Phi(u)} = 0$$

**Theorem 4.3.**  $\alpha_\Phi = \beta_\Phi = p$

*Proof.* From (42) and (44), we get

$$\beta_\Phi = \lim_{t \rightarrow \infty} \frac{\log \left[ \sup_{u>0} \frac{\Phi(tu)}{\Phi(u)} \right]}{\log t} \leq \lim_{t \rightarrow \infty} \frac{\log C + p \log t}{\log t} = p.$$

On the other hand, employing (42) and performing some elementary calculations, we obtain

$$\alpha_\Phi = \lim_{t \rightarrow 0^+} \frac{\log \left[ \sup_{u>0} \frac{\Phi(tu)}{\Phi(u)} \right]}{\log t} = \lim_{s \rightarrow \infty} \frac{\log \left[ \sup_{v>0} \frac{\Phi(v)}{\Phi(sv)} \right]^{-1}}{\log s} = \lim_{s \rightarrow \infty} \frac{\log \left[ \inf_{v>0} \frac{\Phi(sv)}{\Phi(v)} \right]}{\log s}$$

where  $v := tu$  and  $s := \frac{1}{t}$ . Then, using (44), for every  $\varepsilon > 0$  we have

$$\alpha_\Phi = \lim_{s \rightarrow \infty} \frac{\log \left[ \inf_{v>0} \frac{\Phi(sv)}{\Phi(v)} \right]}{\log s} \geq \lim_{s \rightarrow \infty} \frac{\log C + (p - \varepsilon) \log s}{\log s} \geq p - \varepsilon,$$

therefore  $\alpha_\Phi \geq p$ .

Finally, as  $\alpha_\Phi \leq \beta_\Phi \leq p$ , we get  $\alpha_\Phi = \beta_\Phi = p$ . □

Now, we are able to see that

$$\rho_\Phi(u) = \int_0^T \Phi(|u|) dx \geq C \|u\|_{L^\Phi}^{\alpha_\Phi} = C \|u\|_{L^\Phi}^p$$

is false.

In fact, if we take  $u \equiv t > 0$ , then  $\|u\|_{L^\Phi}^p = C_1 t^p$  where  $C_1 = \|1\|_{L^\Phi}$  and  $\int_0^T \Phi(|u|) dx = C_2 \Phi(t)$  with  $C_2 = T$ . Then, if  $\rho_\Phi(u) \geq C \|u\|_{L^\Phi}^p$  were true, then  $\Phi(t) \geq C t^p$  would also be true; however, this last inequality is false.

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## References

- [1] X.-P. Wu, C.-L. Tang, Periodic solutions of a class of non-autonomous second-order systems, J. Math. Anal. Appl. 236 (2) (1999) 227–235.
- [2] F. Zhao, X. Wu, Periodic solutions for a class of non-autonomous second order systems, J. Math. Anal. Appl. 296 (2) (2004) 422–434.
- [3] F. Zhao, X. Wu, Existence and multiplicity of periodic solution for non-autonomous second-order systems with linear nonlinearity, Nonlinear Analysis: Theory, Methods & Applications 60 (2) (2005) 325–335.

## References

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- [4] X. Tang, X. Zhang, Periodic solutions for second-order Hamiltonian systems with a  $p$ -Laplacian, *Ann. Univ. Mariae Curie-Skłodowska Sect. A* 64 (1) (2010) 93–113.
- [5] R. Adams, J. Fournier, *Sobolev spaces*, Elsevier/Academic Press, Amsterdam, 2003.
- [6] M. A. Krasnosel'skiĭ, J. B. Rutickiĭ, *Convex functions and Orlicz spaces*, P. Noordhoff Ltd., Groningen, 1961.
- [7] M. M. Rao, Z. D. Ren, *Theory of Orlicz spaces*, Vol. 146, Marcel Dekker, Inc., New York, 1991.
- [8] G. Schappacher, A notion of Orlicz spaces for vector valued functions, *Appl. Math.* 50 (4) (2005) 355–386.
- [9] H. Hudzik, L. Maligranda, Amemiya norm equals Orlicz norm in general, *Indag. Math. (N.S.)* 11 (4) (2000) 573–585.
- [10] A. Cianchi, A fully anisotropic Sobolev inequality, *Pacific J. Math.* 196 (2) (2000) 283–295.
- [11] A. Cianchi, Some results in the theory of Orlicz spaces and applications to variational problems, in: *Nonlinear analysis, function spaces and applications*, Vol. 6 (Prague, 1998), *Acad. Sci. Czech Repub.*, Prague, 1999, pp. 50–92.
- [12] N. Clavero, Optimal Sobolev embeddings and Function Spaces, <http://www.maia.ub.edu/~soria/sobolev1.pdf>, last accessed: 2014-12-22. (2011).
- [13] D. Edmunds, R. Kerman, L. Pick, Optimal Sobolev imbeddings involving rearrangement-invariant quasinorms, *J. Funct. Anal.* 170 (2) (2000) 307–355.
- [14] R. Kerman, L. Pick, Optimal Sobolev imbeddings, *Forum Math.* 18 (4) (2006) 535–570.
- [15] S. Acinas, L. Buri, G. Giubergia, F. Mazzone, E. Schwindt, Some existence results on periodic solutions of Euler-Lagrange equations in an Orlicz-Sobolev space setting, *Nonlinear Analysis, TMA.* 125 (2015) 681 – 698.
- [16] J. B. Conway, *A Course in Functional Analysis*, Springer, USA, 1977.
- [17] B. Xu, C.-L. Tang, Some existence results on periodic solutions of ordinary  $p$ -Laplacian systems, *J. Math. Anal. Appl.* 333 (2) (2007) 1228–1236.
- [18] L. Maligranda, Orlicz spaces and interpolation, Vol. 5 of *Seminários de Matemática [Seminars in Mathematics]*, Universidade Estadual de Campinas, Departamento de Matemática, Campinas, 1989.
- [19] S. Acinas, G. Giubergia, F. Mazzone, E. Schwindt, On estimates for the period of solutions of equations involving the  $\phi$ -Laplace operator, *J. Abstr. Differ. Equ. Appl.* 5 (1) (2014) 21–34.

- [20] C.-L. Tang, Periodic solutions for nonautonomous second order systems with sub-linear nonlinearity, *Proc. Amer. Math. Soc.* 126 (11) (1998) 3263–3270.
- [21] J. Mawhin, M. Willem, *Critical point theory and Hamiltonian systems*, Springer-Verlag, New York, 1989.