## Periodic solutions of Euler-Lagrange equations in an anisotropic Orlicz-Sobolev space setting

Sonia Acinas \*

Dpto. de Matemática, Facultad de Ciencias Exactas y Naturales Universidad Nacional de La Pampa (L6300CLB) Santa Rosa, La Pampa, Argentina sonia.acinas@gmail.com

Fernando D. Mazzone †

Dpto. de Matemática, Facultad de Ciencias Exactas, Físico-Químicas y Naturales
Universidad Nacional de Río Cuarto
(5800) Río Cuarto, Córdoba, Argentina,
fmazzone@exa.unrc.edu.ar

### **Abstract**

### 1 Introduction

In this paper we obtain existence of solutions for systems of equations of the type:

$$\begin{cases} \frac{d}{dt} \nabla_y \mathcal{L}(t, u(t), u'(t)) = D_x \mathcal{L}(t, u(t), u'(t)) & \text{a.e. } t \in (0, T), \\ u(0) - u(T) = u'(0) - u'(T) = 0, \end{cases}$$
 (P)

where the function  $\mathcal{L}:[0,T]\times\mathbb{R}^d\times\mathbb{R}^d\to\mathbb{R},\ d\geqslant 1$  (called the *Lagrange function* or *lagrangian*) satisfies that it is measurable in t for each  $(x,y)\in\mathbb{R}^d\times\mathbb{R}^d$  and continuously differentiable in (x,y) for almost every  $t\in[0,T]$ . The unknown function  $u:[0,T]\to\mathbb{R}^d$  is assumed absolutely continuous.

Our approach involves the direct method of the calculus of variations in the framework of *anisotropic Orlicz-Sobolev spaces*. We suggest the article [18] for definitions and main results on anisotropic Orlicz spaces, see also [3]. These spaces allow us to unify and extend previous results on existence of solutions for systems like (P).

2010 AMS Subject Classification. Primary: . Secondary: .

Keywords and phrases. .

<sup>\*</sup>SECyT-UNRC and FCEyN-UNLPam

<sup>†</sup>SECyT-UNRC, FCEyN-UNLPam and CONICET

Through this article we say that a function  $\Phi : \mathbb{R}^d \to [0, +\infty)$  belongs to  $N_\infty$  class if  $\Phi$  is convex,  $\Phi(0) = 0$ ,  $\Phi(y) > 0$  if  $y \neq 0$  and  $\Phi(-y) = \Phi(y)$ , and

$$\lim_{|y| \to \infty} \frac{\Phi(y)}{|y|} = +\infty. \tag{1}$$

where  $|\cdot|$  denotes the euclidean norm on  $\mathbb{R}^d$ . From [4, Cor. 2.35] an  $N_{\infty}$  function is continuous.

Associated to  $\Phi$  we have the *complementary function*  $\Psi$  which is defined in  $\xi \in \mathbb{R}^d$  as

$$\Psi(\xi) = \sup_{y \in \mathbb{R}^d} y \cdot \xi - \Phi(y), \tag{2}$$

then, from the continuity of  $\Phi$  and (1), we have that  $\Psi : \mathbb{R}^d \to [0, \infty)$ . It is easy to see that  $\Psi$  is a convex function such that  $\Psi(0) = 0$ ,  $\Psi(-\xi) = \Psi(\xi)$  [9, Ch. 2]. Moreover,  $\Psi$  satisfies (1) (see [18, Thm. 2.2]). i.e.  $\Psi$  is an  $N_{\infty}$  function.

Some examples of  $N_{\infty}$  functions are the following.

Example 1.1.  $\Phi_p(y) := |y|^p/p$ , for  $1 . In this case <math>\Psi(\xi) = |\xi|^q/q$ , q = p/(p-1). Example 1.2. If  $\Phi : \mathbb{R} \to [0, +\infty)$  is a  $N_\infty$  function on  $\mathbb{R}$  then  $\overline{\Phi}(y) = \Phi(|y|)$  is a  $N_\infty$  function on  $\mathbb{R}^d$ . In this example, as in the previous one, the function  $\Phi$  is *radial*, i.e. the value of  $\Phi(y)$  depends on the norm of y and not on its direction. These cases are not authentically anisotropic.

Example 1.3. An anisotropic function  $\Phi(y)$  depends on the direction of y. For example, if  $1 < p_1, p_2 < \infty$ , we define  $\Phi_{p_1, p_2} : \mathbb{R}^d \times \mathbb{R}^d \to [0, +\infty)$  by

$$\Phi_{p_1,p_2}(y_1,y_2) \coloneqq \frac{|y_1|^{p_1}}{p_1} + \frac{|y_2|^{p_2}}{p_2}.$$

Then  $\Phi_{p_1,p_2}$  is an  $N_{\infty}$  function. In this case the complementary function is  $\Phi_{q_1,q_2}$  with  $q_i = p_i/(p_i - 1)$ .

More generally, if  $\Phi_k : \mathbb{R}^d \to [0, +\infty)$ , k = 1, ..., n, are  $N_\infty$  functions, then  $\Phi : \mathbb{R}^d \times \cdots \times \mathbb{R}^d \to [0, +\infty)$  defined by  $\Phi(y_1, ..., y_n) = \Phi_1(y_1) + \cdots + \Phi_n(y_n)$  is an  $N_\infty$  function. These functions are truly anisotropic, i.e. |x| = |y| does not imply that  $\Phi(x) = \Phi(y)$ .

Example 1.4. If  $\Phi : \mathbb{R} \to [0, +\infty)$  is an  $N_{\infty}$  function and  $O \in GL(d, \mathbb{R})$ , then  $\Phi(y) = \Phi(Oy)$  is an  $N_{\infty}$  function.

Example 1.5. An anisotropic  $N_{\infty}$  function is not necessarily controlled by power functions if it does not satisfy the  $\Delta_2$  condition (see xxxxx). For example  $\Phi: \mathbb{R}^d : \to [0, +\infty)$  defined by  $\Phi(y) = \exp(|y|) - 1$  is an  $N_{\infty}$  function.

The appearance of Orlicz Spaces in this paper is due to the fact that we will consider the following structure condition on the lagrangian:

$$|\mathcal{L}| + |\nabla_x \mathcal{L}| + \Psi(\nabla_y \mathcal{L}) \le a(x) \left\{ b(t) + \Phi\left(\frac{y}{\lambda}\right) \right\},$$
 (S)

for a.e.  $t \in [0,T]$ , where  $a \in C(\mathbb{R}^d, [0,+\infty)), b \in L^1([0,T], [0,+\infty))$ .

Our condition (S) includes structure conditions that have previously been considered in the literature. For example, it is easy to see that, when  $\Phi(x)$  is as in Example

1.1, then the condition (S) is equivalent to the structure condition in [9, Th. 1.4]. If  $\Phi$  is a radial  $N_{\infty}$  function such that  $\Psi$  satisfies that  $\Delta_2$  function then (S) is essentially equivalent?????? to conditions [1, Eq. (2)-(4)] (see xxxx mas abajo). If  $\Phi$  is as in Example 1.3 and  $\mathcal{L} = \mathcal{L}(t, x_1, x_2, y_1, y_2)$  is a lagrangian with  $\mathcal{L} : [0, T] \times \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}^d \times \mathbb{R}^d \to \mathbb{R}$  then inequality (S) is related to estructure conditions like [24, Lemma 3.1, Eq. (3.1)]. As can be seen, condition (S) is a more compact expression than [24, Lemma 3.1, Eq. (3.1)] and moreover weaker, because (S) does not imply a control of  $|D_{y_1}L|$  independent of  $y_2$ . We will return to this point later.

An important example of lagrangian is giving by:

$$\mathcal{L}_{\Phi,F}(t,x,y) \coloneqq \Phi(y) + F(t,x). \tag{3}$$

Here the function F(t, x), which is often referred to a potential, be differentiable with respect to x for a.e.  $t \in [0, T]$ . Moreover F satisfies the following conditions:

- (C) F and its gradient  $\nabla_x F$ , with respect to  $x \in \mathbb{R}^d$ , are Carathéodory functions, i.e. they are measurable functions with respect to  $t \in [0,T]$ , for every  $x \in \mathbb{R}^d$ , and they are continuous functions with respect to  $x \in \mathbb{R}^d$  for a.e.  $t \in [0,T]$ .
- (A) For a.e.  $t \in [0,T]$ , it holds that

$$|F(t,x)| + |\nabla_x F(t,x)| \le a(x)b(t). \tag{4}$$

where 
$$a \in C(\mathbb{R}^d, [0, +\infty))$$
 and  $0 \le b \in L^1([0, T], \mathbb{R})$ .

The lagrangian  $\mathcal{L}_{\Phi,F}$  satisfies condition (S). In order to prove this, the only non trivial fact that we should establish is that  $\Psi(\nabla_y \mathcal{L}) \leq a(x) \{b(t) + \Phi(y/\lambda)\}$ . But, from inequality xxxx below,  $\Psi(\nabla_y \mathcal{L}) = \Psi(\nabla \Phi(y)) \leq \Phi(2y)$ .

The laplacian  $\mathcal{L}_{\Phi,F}$  leads to the system

$$\begin{cases} \frac{d}{dt} \nabla \Phi(u'(t)) = \nabla_x F(t, u(t)) & \text{a.e. } t \in (0, T), \\ u(0) - u(T) = u'(0) - u'(T) = 0, \end{cases}$$
  $(\mathbf{P}_{\Phi})$ 

Problem  $(P_{\Phi})$  contains, as a particular case, many problems that are usually considered in the literature. For example, the classic book [9] deals mainly with problem (P) for the lagrangian  $\mathcal{L}_{\Phi,F}$  with  $\Phi(x)=|x|^2/2$ , through various methods: direct, dual action, minimax, etc. The results in [9] were extended and improved in several articles, see [22, 20, 25, 21, 28] to cite some examples. The case  $\Phi(y)=|y|^p/p$ , for arbitrary  $1 were considered in [24, 23], among other papers. In this case, <math>(P_{\Phi})$  is reduced to the p-laplacian system

$$\begin{cases} \frac{d}{dt} \left( u'(t) |u'|^{p-2} \right) = \nabla F(t, u(t)) & \text{a.e. } t \in (0, T) \\ u(0) - u(T) = u'(0) - u'(T) = 0. \end{cases}$$
 ( $P_p$ )

If  $\Phi$  is as in Example 1.3 and  $F:[0,T]\times\mathbb{R}^d\times\mathbb{R}^d\to\mathbb{R}$  is a Carathéodory function, then the equations  $(P_{\Phi})$  become

$$\begin{cases} \frac{d}{dt} \left( |u_1'|^{p_1 - 2} u_1' \right) = F_{x_1}(t, u) & \text{a.e. } t \in (0, T) \\ \frac{d}{dt} \left( |u_2'|^{p_2 - 2} u_2' \right) = F_{x_2}(t, u) & \text{a.e. } t \in (0, T) \\ u(0) - u(T) = u'(0) - u'(T) = 0, \end{cases}$$
  $(\mathbf{P_{p_1, p_2}})$ 

where  $x = (x_1, x_2) \in \mathbb{R}^d \times \mathbb{R}^d$  and  $u(t) = (u_1(t), u_2(t)) \in \mathbb{R}^d \times \mathbb{R}^d$ . In the literature, these equations are known as  $(p_1, p_2)$ -Laplacian system, see [27, 13, 26, 10, 11, 12, 7].

In conclusion, the problem (P) subject to conditions (S) contains several problems that have been considered by many authors in the past.

### 2 Anisotropic Orlicz and Orlicz-Sobolev spaces

In this section, we give a short introduction to Orlicz and Orlicz-Sobolev spaces of vector valued functions associated to anisotropic  $N_{\infty}$  functions  $\Phi: \mathbb{R}^n \to [0, +\infty)$ . References for these topics are [5, 18, 19].

We say that  $\Phi: \mathbb{R}^d \to [0, +\infty)$  satisfies the  $\Delta_2^{\infty}$ -condition, denoted by  $\Phi \in \Delta_2^{\infty}$ , if there exist constants K > 0 and  $M \geqslant 0$  such that

$$\Phi(2x) \leqslant K\Phi(x),\tag{5}$$

for every  $|x| \ge M$ .

We denote by  $\mathcal{M} := \mathcal{M}([0,T],\mathbb{R}^d)$ , with  $d \ge 1$ , the set of all measurable functions (i.e. functions which are limits of simple functions) defined on [0,T] with values on  $\mathbb{R}^d$  and we write  $u = (u_1, \dots, u_d)$  for  $u \in \mathcal{M}$ .

Given an  $N_{\infty}$  function  $\Phi$  we define the modular function  $\rho_{\Phi}: \mathcal{M} \to \mathbb{R}^+ \cup \{+\infty\}$  by

$$\rho_{\Phi}(u) \coloneqq \int_0^T \Phi(u) \ dt.$$

Now, we introduce the *Orlicz class*  $C^{\Phi} = C^{\Phi}([0,T],\mathbb{R}^d)$  by setting

$$C^{\Phi} := \{ u \in \mathcal{M} | \rho_{\Phi}(u) < \infty \}. \tag{6}$$

The Orlicz space  $L^{\Phi}$  =  $L^{\Phi}$  ([0, T],  $\mathbb{R}^d$ ) is the linear hull of  $C^{\Phi}$ ; equivalently,

$$L^{\Phi} := \{ u \in \mathcal{M} | \exists \lambda > 0 : \rho_{\Phi}(\lambda u) < \infty \}. \tag{7}$$

The Orlicz space  $L^{\Phi}$  equipped with the Luxemburg norm

$$\|u\|_{L^\Phi}\coloneqq\inf\left\{\lambda\left|\rho_\Phi\left(\frac{v}{\lambda}\right)dt\leqslant1\right\},$$

is a Banach space.

The subspace  $E^{\Phi} = E^{\Phi}\left([0,T],\mathbb{R}^d\right)$  is defined as the closure in  $L^{\Phi}$  of the subspace  $L^{\infty}\left([0,T],\mathbb{R}^d\right)$  of all  $\mathbb{R}^d$ -valued essentially bounded functions. It is shown that (see [18, Thm. 5.1])  $u \in E^{\Phi}$  if and only if  $\rho_{\Phi}(\lambda u) < \infty$  for any  $\lambda > 0$ . The equality  $L^{\Phi} = E^{\Phi}$  is true if and only if  $\Phi \in \Delta_2^{\infty}$  (see [18, Thm. 5.2]).

A generalized version of *Hölder's inequality* holds in Orlicz spaces (see [18, Thm. 7.2]). Namely, if  $u \in L^{\Phi}$  and  $v \in L^{\Psi}$  then  $u \cdot v \in L^{1}$  and

$$\int_{0}^{T} v \cdot u \, dt \le 2 \|u\|_{L^{\Phi}} \|v\|_{L^{\Psi}}. \tag{8}$$

By  $u \cdot v$  we denote the usual dot product in  $\mathbb{R}^d$  between u and v.

We consider the subset  $\Pi(E^{\Phi}, r)$  of  $L^{\Phi}$  given by

$$\Pi(E^{\Phi}, r) := \{ u \in L^{\Phi} | d(u, E^{\Phi}) < r \}.$$

This set is related to the Orlicz class  $C^{\Phi}$  by means of inclusions, namely,

$$\Pi(E^{\Phi}, r) \subset rC^{\Phi} \subset \overline{\Pi(E^{\Phi}, r)} \tag{9}$$

for any positive r. This relation is a trivial generalization of [18, Thm. 5.6]. If  $\Phi \in \Delta_2^{\infty}$ , then the sets  $L^{\Phi}$ ,  $E^{\Phi}$ ,  $\Pi(E^{\Phi},r)$  and  $C^{\Phi}$  are equal.

As usual, if  $(X, \|\cdot\|_X)$  is a normed space and  $(Y, \|\cdot\|_Y)$  is a linear subspace of X, we write  $Y \hookrightarrow X$  and we say that Y is *embedded* in X when there exists C > 0 such that  $\|y\|_X \leqslant C\|y\|_Y$  for any  $y \in Y$ . With this notation, Hölder's inequality states that  $L^{\Phi} \hookrightarrow [L^{\Psi}]^*$ , where a function  $v \in L^{\Phi}$  is associated to  $\xi_v \in [L^{\Psi}]^*$  being

$$\xi_v(u) = \langle \xi_v, u \rangle = \int_0^T v \cdot u \, dt, \tag{10}$$

We highlight the following result that is a consequence of Theorems 7.1 and 7.3 in [18].

**Proposition 2.1.** If  $\Psi$  satisfies the  $\Delta_2^{\infty}$ -condition then  $L^{\Phi}([0,T],\mathbb{R}^d) = [L^{\Psi}([0,T],\mathbb{R}^d)]^*$ .

We define the Sobolev-Orlicz space  $W^1L^{\Phi}$  by

$$W^1L^\Phi\left([0,T],\mathbb{R}^d\right)\coloneqq\left\{u|u\in AC\left([0,T],\mathbb{R}^d\right)\text{ and }u'\in L^\Phi\left([0,T],\mathbb{R}^d\right)\right\},$$

where  $AC\left([0,T],\mathbb{R}^d\right)$  denotes the space of all  $\mathbb{R}^d$  valued absolutely continuous functions defined on [0,T]. The space  $W^1L^\Phi\left([0,T],\mathbb{R}^d\right)$  is a Banach space when equipped with the norm

$$||u||_{W^1L^{\Phi}} = ||u||_{L^{\Phi}} + ||u'||_{L^{\Phi}}.$$
(11)

We introduce the following subspaces of  $W^1L^\Phi$ 

$$W^{1}E^{\Phi} = \{u \in W^{1}L^{\Phi} | u' \in E^{\Phi}\},\$$

$$W^{1}E^{\Phi}_{T} = \{u \in W^{1}E^{\Phi} | u(0) = u(T)\}.$$
(12)

In order to find a modulus of continuity for functions in  $W^1L^{\Phi}$ , and from there, to obtain compact embedding of  $W^1L^{\Phi}$ , we define the function  $A_{\Phi}: \mathbb{R}^+ \to \mathbb{R}^+$  by

$$A_{\Phi}(s) = \min\left\{\Phi(x) \,\middle|\, |x| = s\right\},\tag{13}$$

Let us establish some elementary properties of  $A_{\Phi}$  that we will use in this article.

**Proposition 2.2.** The function  $A_{\Phi}$  has the following properties:

- 1.  $A_{\Phi}$  is continuous,
- 2.  $A_{\Phi}(s)/s$  is increasing,

- 3.  $A_{\Phi}(|x|)$  is the greatest radial minorant of  $\Phi(x)$ ,
- 4.  $\Phi$  is  $N_{\infty}$  if and only if  $\lim_{s\to+\infty} A_{\Phi}(s)/s = +\infty$ .

*Proof.* It is well known that finite and convex functions defined on finite dimensional vector spaces are locally Lipschitz functions (see [4]). This fact implies item 1 immediately.

In order to prove item 2, suppose 0 < r < s and  $x \in \mathbb{R}^d$  with  $A_{\Phi}(s) = \Phi(x)$ . Then, from the definition of  $A_{\Phi}$  and the convexity of  $\Phi$ ,

$$\frac{A_{\Phi}(r)}{r} \leqslant \frac{\Phi\left(\frac{r}{s}x\right)}{r} \leqslant \frac{\Phi\left(x\right)}{s} = \frac{A_{\Phi}(s)}{s}.$$

Property in items 3 and 4 are obtained easily.

Example 2.1. We compute  $A_{\Phi}$  for the function  $\Phi = \Phi_{p_1,p_2}$  given in Example (1.3). We apply the method of Lagrange multipliers (see [8, Ch. 11]) to solve the next minimization problem subject to constraints

$$\begin{cases} \text{ minimize } \Phi_{p_1, p_2}(y_1, y_2) \\ \text{ subject to } |y_1|^2 + |y_2|^2 = r^2 \end{cases}.$$

The first order conditions are

$$\begin{cases} |y_1|^{p_1-2}y_1 + \lambda y_1 &= 0\\ |y_2|^{p_2-2}y_2 + \lambda y_2 &= 0\\ |y_1|^2 + |y_2|^2 &= r^2 \end{cases}$$
(14)

These equations are solved, among others, by the following two sets of critical points: a) |x| = r, y = 0 and  $\lambda = -r^{p_1-2}$  and b) x = 0, |y| = r and  $\lambda = -r^{p_2-2}$ . These sets are infinite when d > 1. Associated with these critical points we have the following critical values: a)  $r^{p_1}/p_1$  and b)  $r^{p_2}/p_2$ .

We deal with  $p_1 \le 2$  and  $p_2 \le 2$  being one of them (suppose  $p_2$ ) different from 2. The remaining cases can be treated with similar techniques.

If  $(y_1,y_2)$  solve (14) with  $y_1 \neq 0$  and  $y_2 \neq 0$  then  $|y_2| = |y_1|^{\frac{p_1-2}{p_2-2}}$  and  $\lambda = -|y_1|^{p_1-2}$ . We use second order conditions for constrained problems. We have to consider the tangent plane at the point  $(y_1,y_2) \in \mathbb{R}^{2n}$ , i.e.  $M = \{(\xi,\eta) \in \mathbb{R}^{2n} : \xi y_1^t + \eta y_2^T = 0\}$ . Let L be the Lagrangian associated to the constrained problem:  $L(y_1,y_2,\lambda) = \Phi(y_1,y_2) + \lambda H(y_1,y_2)$  being H = 0 the constraint. We must analyze the positivity of the quadratic form associated to the matrix of second partial derivatives  $\mathcal{H} = D^2 \Phi + \lambda D^2 H$  on the subspace M. By elementary computations we have for  $(\xi,\eta) \in M$ 

$$(\xi,\eta)^t \mathcal{H}(\xi,\eta) = |\lambda|(\xi^t x)^2 [|y_1|^{-2}(p_1-2) + (p_2-2)|y_2|^{-2}],$$

on the subspace M. We note that  $(-y_2, y_1) \in M$  and  $(-y_2, y_1)^t \mathcal{H}(-y_2, y_1) < 0$ . Then, by second order necessary conditions [8, p.333], at  $(y_1, y_2)$  there cannot be a minimum. Therefore, the minima occur at  $y_1 = 0$  or  $y_2 = 0$ , then

$$A_{\Phi}(x,y) = \min\{r^{p_1}/p_1, r^{p_2}/p_2\}.$$

More generally, it holds that

$$K_1 \min\{r^{p_1}, r^{p_2}\} \le A_{\Phi} \le K_2 \min\{r^{p_1}, r^{p_2}\}$$

with  $K_1, K_2 > 0$ , for every  $1 < p_1, p_2 < \infty$ .

As it is customary, we will use the decomposition  $u = \overline{u} + \widetilde{u}$  for a function  $u \in L^1([0,T])$  where  $\overline{u} = \frac{1}{T} \int_0^T u(t) \ dt$  and  $\widetilde{u} = u - \overline{u}$ .

The following lemma is an elementary generalization to anisotropic Sobolev-Orlicz spaces of known results of Sobolev spaces.

**Lemma 2.3.** Let  $\Phi : \mathbb{R}^d \to [0, +\infty)$  be a Young's function and let  $u \in W^1L^{\Phi}([0, T], \mathbb{R}^d)$ . Let  $A_{\Phi} : \mathbb{R}^+ \to \mathbb{R}^+$  be the function defined by (13). Then

1. For every  $s, t \in [0, T]$ ,  $s \neq t$ ,

$$|u(t) - u(s)| \le ||u'||_{L^{\Phi}} |s - t| A_{\Phi}^{-1} \left( \frac{1}{|s - t|} \right)$$
 (Morrey's inequality)

$$||u||_{L^{\infty}} \leqslant A_{\Phi}^{-1}\left(\frac{1}{T}\right) \max\{1, T\} ||u||_{W^1L^{\Phi}}$$
 (Sobolev's inequality)

2. We have  $\widetilde{u} \in L^{\infty}([0,T],\mathbb{R}^d)$  and

$$\|\widetilde{u}\|_{L^{\infty}} \leqslant TA_{\Phi}^{-1}\left(\frac{1}{T}\right)\|u'\|_{L^{\Phi}}$$
 (Sobolev-Wirtinger's inequality)

3. If  $\Phi$  is  $N_{\infty}$  then the space  $W^1L^{\Phi}([0,T],\mathbb{R}^d)$  is compactly embedded in the space of continuous functions  $C([0,T],\mathbb{R}^d)$ .

*Proof.* By the absolutely continuity of u, Jensen's inequality and the definition of the Luxemburg norm, we have

$$\Phi\left(\frac{u(t) - u(s)}{\|u'\|_{L^{\Phi}}|s - t|}\right) \leqslant \Phi\left(\frac{1}{|s - t|} \int_{s}^{t} \frac{u'(r)}{\|u'\|_{L^{\Phi}}} dr\right) 
\leqslant \frac{1}{|s - t|} \int_{s}^{t} \Phi\left(\frac{u'(r)}{\|u'\|_{L^{\Phi}}}\right) dr \leqslant \frac{1}{|s - t|}.$$

By Proposition 2.2(3) we have  $A_{\Phi}^{-1}\Phi(x) \ge |x|$ , therefore we get

$$\frac{|u(t) - u(s)|}{\|u'\|_{L^{\Phi}}|s - t|} \leqslant A_{\Phi}^{-1} \left(\frac{1}{|s - t|}\right),$$

then 1 holds.

Now, we use Morrey's inequality and Proposition 2.2 (2) and we have

$$|u(t) - \overline{u}| = \left| \frac{1}{T} \int_0^T u(t) - u(s) \, ds \right|$$

$$\leq \frac{1}{T} \int_0^T |u(t) - u(s)| \, ds$$

$$\leq ||u'||_{L^{\Phi}} T A_{\Phi}^{-1} \left( \frac{1}{T} \right)$$

In order to prove the Sobolev's inequality, we note that, using Jensen's inequality and the definition of  $\|u\|_{L^{\Phi}}$ , we obtain

$$\Phi\left(\frac{\overline{u}}{\|u\|_{L^{\Phi}}}\right) \leqslant \frac{1}{T} \int_{0}^{T} \Phi\left(\frac{u(s)}{\|u\|_{L^{\Phi}}}\right) ds \leqslant \frac{1}{T}$$

Then by By Proposition 2.2(3)

$$|\overline{u}| \leqslant A_{\Phi}^{-1} \left(\frac{1}{T}\right) \|u\|_{L^{\Phi}}.$$

Therefore, from this and (Sobolev-Wirtinger's inequality) we get

$$\begin{split} \|u\|_{L^{\infty}} & \leq |\overline{u}| + \|\tilde{u}\|_{L^{\infty}} \\ & \leq A_{\Phi}^{-1} \left(\frac{1}{T}\right) \|u\|_{L^{\Phi}} + T A_{\Phi}^{-1} \left(\frac{1}{T}\right) \|u'\|_{L^{\Phi}} \\ & \leq A_{\Phi}^{-1} \left(\frac{1}{T}\right) \max\{1, T\} \|u\|_{W^{1} L^{\Phi}} \end{split}$$

In order to prove item 3, we take a bounded sequence  $u_n$  in  $W^1L^{\Phi}\left([0,T],\mathbb{R}^d\right)$ . Since  $\Phi$  is  $N_{\infty}$ , from Proposition 2.2(4) we obtain  $sA_{\Phi}^{-1}(1/s) \to 0$  when  $s \to 0$ . Therefore (Morrey's inequality) implies that  $u_n$  are equicontinuous. Furthermore (??) implies that  $u_n$  is bounded in  $C\left([0,T],\mathbb{R}^d\right)$ . Therefore by the Arzela-Ascoli Theorem we obtain a subsequence  $n_k$  and  $u \in C\left([0,T],\mathbb{R}^d\right)$  with  $u_{n_k} \to u$  in  $C\left([0,T],\mathbb{R}^d\right)$ .

QUIZAS LO QUE VIENE TENDRIA QUE IR CON LAS OTRAS DESIGUAL-DADES... O MODIFICAR EL LEMA ANTERIOR....

We get an anisotropic version of Sobolev-Wirtinger inequality, as follows.

**Proposition 2.4.** If  $u \in W^1L^{\Phi}$ , then

$$\Phi(\tilde{u}(t)) \leqslant \int_0^T \Phi(\max\{1, T\}u'(r)) dr. \tag{15}$$

L

*Proof.* Writing  $\tilde{u}(t) = u(t) - \overline{u}$  and applying Jensen's inequality, we get

$$\Phi(\tilde{u}(t)) = \Phi\left(\frac{1}{T} \int_0^T u(t) - u(s) \, ds\right)$$

$$\leqslant \frac{1}{T} \int_0^T \Phi(u(t) - u(s)) \, ds \leqslant \frac{1}{T} \int_0^T \Phi\left(\int_s^t |t - s| u'(r) \frac{dr}{|t - s|}\right) \, ds$$

$$\leqslant \frac{1}{T} \int_0^T \frac{1}{|t - s|} \int_s^t \Phi\left(|t - s| u'(r)\right) \, dr \, ds$$

As  $\Phi$  is convex, we have

$$\frac{1}{|t-s|}\Phi(|t-s|u'(r)) \leqslant \begin{cases} \Phi(u'(r)) & if \quad |t-s| \leqslant 1\\ \Phi(Tu'(r)) & if \quad 1 < |t-s| < T \end{cases}$$

Then  $\frac{1}{|t-s|}\Phi\left(|t-s|u'(r)\right) \leqslant \Phi(\max\{1,T\}u'(r))$  and (15) follows.

**Lemma 2.5.** Let  $\{u_n\}_{n\in\mathbb{N}}$  be a sequence of functions in  $\Pi(E^{\Phi},1)$  converging to  $u\in\Pi(E^{\Phi},1)$  in the  $L^{\Phi}$ -norm. Then, there exist a subsequence  $u_{n_k}$  and a real valued function  $h\in L^1([0,T],\mathbb{R})$  such that  $u_{n_k}\to u$  a.e. and  $\Phi(u_{n_k})\leqslant h$  a.e.

*Proof.* Since  $d(u, E^{\Phi}) < 1$  and  $u_n$  converges to u, there exists  $u_0 \in E^{\Phi}$ , a subsequence of  $u_n$  (again denoted  $u_n$ ) and 0 < r < 1 such that  $d(u_n, u_0) < r$ . Let  $\lambda_0 \in (r, 1)$ . By extracting more subsequences, if necessary, we can assume that  $u_n \to u$  a.e. and

$$\lambda_n := \|u_{n+1} - u_n\|_{L^{\Phi}} < \frac{1 - \lambda_0}{2^n}, \quad \text{for } n \geqslant 1.$$

We can assume  $\lambda_n > 0$  for every  $n = 0, \ldots$ 

Let  $\lambda := 1 - \sum_{n=0}^{\infty} \lambda_n$  and define  $h : [0,T] \to \mathbb{R}$  by

$$h(x) = \lambda \Phi\left(\frac{u_0}{\lambda}\right) + \sum_{n=0}^{\infty} \lambda_n \Phi\left(\frac{u_{n+1} - u_n}{\lambda_n}\right). \tag{16}$$

Note that  $\sum_{n=0}^{\infty} \lambda_n + \lambda = 1$ , therefore for any  $n = 1, \dots$ 

$$\begin{split} \Phi(u_n) &= \Phi\left(\lambda \frac{u_0}{\lambda} + \sum_{j=0}^{n-1} \lambda_j \frac{u_{j+1} - u_j}{\lambda_j}\right) \\ &\leq \lambda \Phi\left(\frac{u_0}{\lambda}\right) + \sum_{j=0}^{n-1} \lambda_j \Phi\left(\frac{u_{j+1} - u_j}{\lambda_j}\right) \leq h \end{split}$$

Since  $u_0 \in E^{\Phi} \subset C^{\Phi}$  and  $E^{\Phi}$  is a subspace we have that  $\Phi(u_0/\lambda) \in L^1([0,T],\mathbb{R})$ . On the other hand  $||u_{n+1} - u_n||_{L^{\Phi}} \leq \lambda_n$ , therefore

$$\int_0^T \Phi\left(\frac{u_{j+1} - u_j}{\lambda_j}\right) dt \le 1.$$

Then  $h \in L^1([0,T],\mathbb{R})$ .

### DEMOSTRACION ALTERNATIVA PARA EL LEMA y con $\lambda$

**Lemma 2.6.** Let  $\{u_n\}_{n\in\mathbb{N}}$  be a sequence of functions in  $\Pi(E^{\Phi},\lambda)$  converging to  $u\in\Pi(E^{\Phi},\lambda)$  in the  $L^{\Phi}$ -norm. Then, there exist a subsequence  $\Phi(\frac{u_{n_k}}{\lambda})$  and a function  $h\in L^1([0,T],\mathbb{R})$  such that  $\Phi(\frac{u_{n_k}}{\lambda})\to\Phi(\frac{u}{\lambda})$  a.e. and  $\Phi(\frac{u_{n_k}}{\lambda})\leqslant h$  a.e.

*Proof.* As  $u \in \Pi(E^{\Phi}, \lambda)$ , we consider  $\Lambda \in (0, \lambda)$ . In this way,  $d(u, E^{\Phi}) < \Lambda < \lambda$  and, taking into account (9),  $\Phi(\frac{u}{\lambda}) \in L^1([0, T], \mathbb{R})$ .

Applying [3, Lemma 3.1] with  $x+y=\frac{u_n}{\lambda}$ ,  $x=\frac{u}{\lambda}$ ,  $k=\frac{\lambda}{\Lambda}$ ,  $0<\epsilon<\frac{\Lambda}{\lambda}$  and  $C_\epsilon=\frac{\Lambda}{\epsilon(\lambda-\Lambda)}$ , we have

$$\int_{0}^{T} \left| \Phi\left(\frac{u_{n}}{\lambda}\right) - \Phi\left(\frac{u}{\lambda}\right) \right| dt$$

$$\leq \epsilon \int_{0}^{T} \left| \Phi\left(\frac{u}{\lambda}\right) - \frac{\lambda}{\Lambda} \Phi\left(\frac{u}{\lambda}\right) \right| dt + 2 \int_{0}^{T} \Phi\left(\frac{\Lambda}{\epsilon(\lambda - \Lambda)} \frac{u_{n} - u}{\lambda}\right) dt.$$
(17)

Let  $\eta > 0$ . Since  $\Phi(\frac{u}{\lambda}) \in L^1([0,T,\mathbb{R}])$ , we can choose  $\epsilon$  such that

$$\epsilon \int_0^T \left| \Phi\left(\frac{u}{\lambda}\right) - \frac{\lambda}{\Lambda} \Phi\left(\frac{u}{\lambda}\right) \right| dt < \eta. \tag{18}$$

From the fact that  $u_n \to u$  in the  $L^{\Phi}$ -norm, there exists  $n_0$  such that  $\|u_n - u\|_{L^{\Phi}} < \epsilon^2(\lambda - \Lambda)$  for every  $n \geqslant n_0$ . Then, by the convexity of  $\Phi$  and the definition of Orlicz norm, we get

$$\int_{0}^{T} \Phi\left(\frac{\Lambda}{\epsilon(\lambda - \Lambda)} \frac{u_{n} - u}{\lambda}\right) \leq \epsilon \frac{\Lambda}{\lambda} \int_{0}^{T} \Phi\left(\frac{u_{n} - u}{\epsilon^{2}(\lambda - \Lambda)}\right) < \epsilon$$
(19)

Thus, from (17), (18) and (19), we obtain that  $\Phi(\frac{u_n}{\lambda})$  converges to  $\Phi(\frac{u}{\lambda})$  in the  $L^1$ -norm. Now, [2, Thm. 4.9] implies that there exist a subsequence  $\Phi(\frac{u_{n_k}}{\lambda})$  and a function  $h \in L^1([0,T],\mathbb{R})$  such that  $\Phi(\frac{u_{n_k}}{\lambda}) \to \Phi(\frac{u}{\lambda})$  a.e. and  $\Phi(\frac{u_{n_k}}{\lambda}) \leqslant h$  a.e.

# 3 Differentiability Gateâux of action integrals in anisotropic Orlicz spaces

In this section we give a brief introduction to superposition operators between anisotropic Orlicz Spaces. We apply these results to obtain Gateâux differentiability of action integrals associated to lagrangian functions defined on Sobolev-Orlicz spaces.

Henceforth we assume that  $f:[0,T]\times\mathbb{R}^d\to\mathbb{R}^d$  is a Carathéodory function, i.e.

(C) f is measurable with respect to  $t \in [0,T]$  for every  $x \in \mathbb{R}^d$ , and f is a continuous function with respect to  $x \in \mathbb{R}^d$  for a.e.  $t \in [0,T]$ .

**Definition 3.1.** For  $f:[0,T]\times\mathbb{R}^d\to\mathbb{R}^d$  we denote by f the Nemytskii (o superposition) operator defined for functions  $u:[0,T]\to\mathbb{R}^d$  by

$$fu(t) = f(t, u(t))$$

In the following Theorem we enumerate some known properties for superposition operators defined on anisotropic Orlicz spaces of vector functions. For the proofs see [6] for scalar functions and [16, 15, 14] for the generalization to  $\mathbb{R}^d$ -valued (moreover Banach spaces valued) functions in a anisotropic Orlicz Spaces (moreover modular anisotropic spaces).

**Theorem 3.2.** We assume that f satisfies condition ((C)) and that  $\Phi_1, \Phi_2 : \mathbb{R}^d \to [0, +\infty)$  are anisotropic Young functions. Then

- Measurability. The operator f maps measurable function into measurable functions
- 2. Extensibility. If the operator  ${\bf f}$  acts from the ball  $B_{L^{\Phi_1}}(r)\coloneqq\{u\in L^{\Phi_1}|\|u\|_{L^{\Phi_1}}< r\}$  into the space  $L^{\Phi_2}$  or the space  $E^{\Phi_2}$  then  ${\bf f}$  can be extended from  $\Pi(E^{\Phi_1},r)$  into space  $L^{\Phi_2}$  or  $E^{\Phi_2}$ , respectively.
- 3. Continuity. If the operator f acts from  $\Pi(E^{\Phi_1}, r)$  into space  $E^{\Phi_2}$ , then f is continuous.

Given a continuous function  $a \in C(\mathbb{R}^n, \mathbb{R}^+)$ , we define the composition operator  $a : \mathcal{M}_d \to \mathcal{M}_d$  by  $\mathbf{a}(u)(x) = a(u(x))$ .

We will often use the following result whose proof can be performed as that of Corollary 2.3 in [1].

**Lemma 3.3.** If  $a \in C(\mathbb{R}^d, \mathbb{R}^+)$  then  $a : W^1L^{\Phi} \to L^{\infty}([0,T])$  is bounded. More concretely, there exists a non decreasing function  $A : \mathbb{R}^+ \to \mathbb{R}^+$  such that  $\|a(u)\|_{L^{\infty}([0,T])} \le A(\|u\|_{W^1L^{\Phi}})$ .

*Proof.* Let  $A \in C(\mathbb{R}^+, \mathbb{R}^+)$  be a non decreasing, continuous function defined by  $\alpha(s) := \sup_{\|x\| \leq s, x \in \mathbb{R}^d} |a(x)|$ . If  $u \in W^1L_d^{\Phi}$  then, by Sobolev's inequality, for a.e.  $t \in [0, T]$ 

$$a(u(t)) \leqslant \alpha(\|u\|_{L^{\infty}}) \leqslant \alpha\left(A_{\Phi}^{-1}\left(\frac{1}{T}\right) \max\{1, T\}\|u\|_{W^{1}L^{\Phi}}\right) =: A(\|u\|_{W^{1}L^{\Phi}}).$$

HABRÍA QUE VER DÓNDE SE UBICA LA CONDICIÓN DE ESTRUCTURA...QUIZÁS EN LA INTRODUCCIÓN?....

Next, we deal with the differentiability of the action integral

$$I(u) = \int_0^T \mathcal{L}(t, u(t), \dot{u}(t)) dt.$$
 (20)

**Theorem 3.4.** Let  $\mathcal{L}$  be a differentiable Carathéodory function satisfying (S). Then the following statements hold:

1. The action integral given by (20) is finitely defined on  $\mathcal{E}^{\Phi} := W^1 L^{\Phi} \cap \{u | \dot{u} \in \Pi(E^{\Phi}, 1)\}.$ 

2. The function I is Gâteaux differentiable on  $\mathcal{E}^{\Phi}$  and its derivative I' is demicontinuous from  $\mathcal{E}^{\Phi}$  into  $\left[W^1L^{\Phi}\right]^*$ . Moreover, I' is given by the following expression

$$\langle I'(u), v \rangle = \int_0^T \left\{ D_x \mathcal{L}(t, u, \dot{u}) \cdot v + \nabla_y \mathcal{L}(t, u, \dot{u}) \cdot \dot{v} \right\} dt. \tag{21}$$

3. If  $\Psi \in \Delta_2$  then I' is continuous from  $\mathcal{E}^{\Phi}$  into  $\left[W^1L^{\Phi}\right]^*$  when both spaces are equipped with the strong topology.

*Proof.* Let  $u \in \mathcal{E}^{\Phi}$ . As

$$\dot{u} \in \Pi(E^{\Phi}, 1) \subset C_1^{\Phi} \tag{22}$$

and (9), then  $\Phi(\dot{u}(t)) \in L^1$ . Now,

$$|\mathcal{L}(\cdot, u, u)| + |\nabla_x \mathcal{L}(\cdot, u, u)| + \Psi(\nabla_y \mathcal{L}(\cdot, u, u)) \le A(\|u\|_{W^1 L^{\Phi}})(b + \Phi(u)) \in L^1, \quad (23)$$

by (S) and Lemma 3.3. Thus item (1) is proved.

We split up the proof of item 2 into four steps.

Step 1. The non linear operator  $u \mapsto \nabla_x \mathcal{L}(t, u, u)$  is continuous from  $\mathcal{E}^{\Phi}$  into  $L^1([0, T])$  with the strong topology on both sets.

Let  $\{u_n\}_{n\in\mathbb{N}}$  be a sequence of functions in  $\mathcal{E}^{\Phi}$  and let  $u\in\mathcal{E}^{\Phi}$  such that  $u_n\to u$  in  $W^1L^{\Phi}$ . By (Sobolev's inequality), we have

$$|u_n(t) - u(t)| \le T A_{\Phi}^{-1} \left(\frac{1}{T}\right) ||u_n - u||_{L^{\Phi}}$$

then  $u_n \to u$  uniformly. As  $u_n \to u \in \mathcal{E}^{\Phi}$ , by Lemma 2.5, there exist a subsequence of  $u_{n_k}$  (again denoted  $u_{n_k}$ ) and a function  $h \in L^1([0,T],\mathbb{R})$  such that  $u_{n_k} \to u$  a.e. and  $\Phi(u_{n_k}) \leq h$  a.e.

Since  $u_{n_k}$ ,  $k=1,2,\ldots$ , is a strong convergent sequence in  $W^1L^{\Phi}$ , it is a bounded sequence in  $W^1L^{\Phi}$ . According to item (3) of Lemma 2.3, there exists M>0 such that  $\|\boldsymbol{a}(u_{n_k})\|_{L^{\infty}} \leq M$ ,  $k=1,2,\ldots$  From the previous facts and (23), we get

$$|\nabla_x \mathcal{L}(\cdot, u_{n_k}, u_{n_k})| \leq a(|u_{n_k}|)(b + \Phi(u_{n_k})) \leq M(b + h) \in L^1.$$

On the other hand, by the continuous differentiability of  $\mathcal{L}$ , we have

$$\nabla_x \mathcal{L}(t, u_{n_k}(t), u_{n_k}(t)) \to \nabla_x \mathcal{L}(t, u(t), u(t))$$
 for a.e.  $t \in [0, T]$ .

Applying the Dominated Convergence Theorem we conclude the proof of step 1. Step 2. The non linear operator  $u \mapsto \nabla_y \mathcal{L}(t,u,\dot{u})$  is continuous from  $\mathcal{E}^{\Phi}$  with the strong topology into  $\left[L^{\Phi}\right]^*$  with the weak\* topology.

Let  $u \in \mathcal{E}^{\Phi}$ . From (23) it follows that

$$\nabla_y \mathcal{L}(\cdot, u, \dot{u}) \in C^{\Psi}. \tag{24}$$

Note that (23), (24) and the imbeddings  $W^1L^\Phi \to L^\infty$  and  $L^\Psi \to \left[L^\Phi\right]^*$  imply that the second member of (21) defines an element of  $\left[W^1L^\Phi\right]^*$ .

Let  $u_n, u \in \mathcal{E}^{\Phi}$  such that  $u_n \to u$  in the norm of  $W^1L^{\Phi}$ . We must prove that  $\nabla_y \mathcal{L}(\cdot, u_n, u_n) \stackrel{w^*}{\rightharpoonup} \nabla_y \mathcal{L}(\cdot, u, u)$ . On the contrary, there exist  $v \in L^{\Phi}$ ,  $\epsilon > 0$  and a subsequence of  $\{u_n\}$  (denoted  $\{u_n\}$  for simplicity) such that

$$|\langle \nabla_{y} \mathcal{L}(\cdot, u_n, \dot{u}_n), v \rangle - \langle \nabla_{y} \mathcal{L}(\cdot, u, \dot{u}), v \rangle| \ge \epsilon. \tag{25}$$

We have  $u_n \to u$  in  $L^\Phi$  and  $u_n \to u$  in  $L^\Phi$ . By Lemma 2.5, there exist a subsequence of  $\{u_n\}$  (again denoted  $\{u_n\}$  for simplicity) and a function  $h \in L^1([0,T],\mathbb{R})$  such that  $u_n \to u$  uniformly,  $u_n \to u$  a.e. and  $\Phi(u_n) \leqslant h$  a.e. As in the previous step, since  $u_n$  is a convergent sequence, Lemma 3.3 implies that  $a(|u_n(t)|)$  is uniformly bounded by a certain constant M>0. Therefore, from inequality (23) with  $u_n$  instead of u, we have

$$\Psi(\nabla_{\nu}\mathcal{L}(\cdot, u_n, \dot{u}_n)) \leq M(b+h) \in L^1.$$
(26)

As  $v \in L^{\Phi}$  there exists  $\lambda > 0$  such that  $\Phi(\frac{v}{\lambda}) \in L^1$ . Now, by Young inequality and (26), we have

$$\lambda \nabla_{y} \mathcal{L}(\cdot, u_{n_{k}}, \dot{u}_{n_{k}}) \cdot \frac{v(t)}{\lambda}$$

$$\leq \lambda \left[ \Psi(\nabla_{y} \mathcal{L}(\cdot, u_{n_{k}}, \dot{u}_{n_{k}})) + \Phi\left(\frac{v}{\lambda}\right) \right]$$

$$\leq \lambda M(b+h) + \lambda \Phi\left(\frac{v}{\lambda}\right) \in L^{1}$$
(27)

Finally, from the Lebesgue Dominated Convergence Theorem, we deduce

$$\int_{0}^{T} \nabla_{y} \mathcal{L}(t, u_{n_{k}}, \dot{u}_{n_{k}}) \cdot v \, dt \to \int_{0}^{T} \nabla_{y} \mathcal{L}(t, u, \dot{u}) \cdot v \, dt \tag{28}$$

which contradicts the inequality (25). This completes the proof of step 2.

Step 3. We will prove (21). For  $u \in \mathcal{E}^{\Phi}$  and  $0 \neq v \in W^1 L^{\Phi}$ , we define the function

$$H(s,t) \coloneqq \mathcal{L}(t,u(t) + sv(t),u(t) + sv(t)).$$

For  $|s| \leq s_0 := \min\{\left(1 - d(u, E^{\Phi})\right) / \|v\|_{W^1L^{\Phi}}, 1 - d(u, E^{\Phi})\}$ , using triangle inequality we get  $d\left(u + sv, E^{\Phi}\right) < 1$  and thus  $u + sv \in \Pi(E^{\Phi}, 1)$ . These facts imply, in virtue of Theorem 3.4 item 1, that I(u + sv) is well defined and finite for  $|s| \leq s_0$ .

We also have  $\|u+sv\|_{W^1L^\Phi}\leqslant \|u\|_{W^1L^\Phi}+s_0\|v\|_{W^1L^\Phi}$ ; then, by Lemma 3.3, there exists M>0 such that  $\|a(u+sv)\|_{L^\infty}\leqslant M$ .

Let  $\lambda > 0$  such that  $\Phi(\frac{\dot{v}}{\lambda}) \in L^{1}$ . On the other hand, if  $\dot{v} \in L^{\Phi}$  and  $|s| \leq s_0 \lambda^{-1}$ , from the convexity and the parity of  $\Phi$ , we get

$$\begin{split} & \Phi(\dot{u} + s\dot{v}) = \Phi\left((1 - s_0)\frac{\dot{u}}{1 - s_0} + s_0\frac{s}{s_0}\dot{v}\right) \leqslant (1 - s_0)\Phi\left(\frac{\dot{u}}{1 - s_0}\right) + s_0\Phi\left(\frac{s}{s_0}\dot{v}\right) \\ & \leqslant (1 - s_0)\Phi\left(\frac{\dot{u}}{1 - s_0}\right) + s_0\Phi\left(\frac{\dot{v}}{\lambda}\right) \in L^1 \end{split}$$

As  $\dot{u} \in \Pi(E^{\Phi}, 1)$  then

$$d\left(\frac{\dot{u}}{1-s_0}, E^{\Phi}\right) = \frac{1}{1-s_0}d(\dot{u}, E^{\Phi}) < 1$$

and therefore  $\frac{u}{1-s_0} \in C^{\Phi}$ .

Now, applying (23), (27), the fact that  $v \in L^{\infty}$  and  $\dot{v} \in L^{\Phi}$ , we get

$$|D_{s}H(s,t)| = \left| \nabla_{x}\mathcal{L}(t,u+sv,\dot{u}+s\dot{v}) \cdot v + \lambda \nabla_{y}\mathcal{L}(t,u+sv,\dot{u}+s\dot{v}) \cdot \frac{\dot{v}}{\lambda} \right|$$

$$\leq M \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] |v|$$

$$+ \lambda \left[ \Psi(\nabla_{y}\mathcal{L}(t,u+sv,\dot{u}+s\dot{v})) + \Phi\left(\frac{\dot{v}}{\lambda}\right) \right]$$

$$\leq M \left\{ \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] |v| \right\} + \lambda M \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] + \lambda \Phi\left(\frac{\dot{v}}{\lambda}\right)$$

$$= M \left[ b(t) + \Phi(\dot{u}+s\dot{v}) \right] (|v| + \lambda) + \lambda \Phi\left(\frac{\dot{v}}{\lambda}\right) \in L^{1}.$$
(29)

Consequently, I has a directional derivative and

$$\langle I'(u), v \rangle = \frac{d}{ds} I(u + sv) \Big|_{s=0} = \int_0^T \left\{ \nabla_x \mathcal{L}(t, u, \dot{u}) \cdot v + \nabla_y \mathcal{L}(t, u, \dot{u}) \cdot \dot{v} \right\} dt.$$

Moreover, from the previous formula, (23), (24), and Lemma 2.3, we obtain

$$|\langle I'(u), v \rangle| \le \|\nabla_x \mathcal{L}\|_{L^1} \|v\|_{L^{\infty}} + \|\nabla_y \mathcal{L}\|_{L^{\Psi}} \|\dot{v}\|_{L^{\Phi}} \le C \|v\|_{W^1 L^{\Phi}}$$

with a appropriate constant C.

This completes the proof of the Gâteaux differentiability of I.

Step 4. The operator  $I': \mathcal{E}^{\Phi} \to \left[W^1L_d^{\Phi}\right]^*$  is demicontinuous. This is a consequence of the continuity of the mappings  $u \mapsto \nabla_x \mathcal{L}(t, u, u)$  and  $u \mapsto \nabla_y \mathcal{L}(t, u, u)$ . Indeed, if  $u_n, u \in \mathcal{E}^{\Phi}$  with  $u_n \to u$  in the norm of  $W^1L^{\Phi}$  and  $v \in W^1L^{\Phi}$ , then

$$\langle I'(u_n), v \rangle = \int_0^T \left\{ \nabla_x \mathcal{L}(t, u_n, \dot{u}_n) \cdot v + \nabla_y \mathcal{L}(t, u_n, \dot{u}_n) \cdot \dot{v} \right\} dt$$

$$\to \int_0^T \left\{ \nabla_x \mathcal{L}(t, u, \dot{u}) \cdot v + \nabla_y \mathcal{L}(t, u, \dot{u}) \cdot \dot{v} \right\} dt$$

$$= \langle I'(u), v \rangle.$$

In order to prove item 3, it is necessary to see that the maps  $u \mapsto \nabla_x \mathcal{L}(t, u, u)$  and  $u \mapsto \nabla_y \mathcal{L}(t, u, u)$  are norm continuous from  $\mathcal{E}^{\Phi}$  into  $L^1$  and  $L^{\Psi}$ , respectively.

The continuity of the first map has already been proved in step 1.

Let  $u_n, u \in \mathcal{E}^{\Phi}$  with  $||u_n - u||_{W^1L^{\Phi}} \to 0$ .

Applying Lemma 2.5 to  $\dot{u}_n$ , there exists a subsequence (denoted  $\dot{u}_n$  for simplicity) such that  $u_n \in L^{\Phi}$  and a function  $h \in L^1$  such that  $\Psi(u_n) \leq h$  and  $u_n \to u$  a.e.

Then, by (27) we have  $\Psi(v_n) \leq m(t) \in L^1$  being  $v_n := \nabla_y \mathcal{L}(\cdot, u_n, u_n)$  and m(t) :=M(b+h). In addition, from the continuous differentiability of  $\mathcal{L}$ , we have that  $v_n \to v$ a.e. where  $\nabla_{u} \mathcal{L}(\cdot, u, u)$ .

As  $\Psi \in \Delta_2$ , there exists  $c : \mathbb{R}^+ \to \mathbb{R}^+$  such that  $\Psi(\lambda x) \leq c(|\lambda|)\Psi(x)$ . Then,

$$\begin{split} &\Psi(\frac{v_n-v}{\lambda})\leqslant c(|\lambda|^{-1})\Psi(v_n-v) \text{ for every } \lambda\in\mathbb{R}.\\ &\text{Therefore, } \Psi(\frac{v_n-v}{\lambda})\to 0 \text{ a.e. as } n\to\infty \text{ and } \Psi(\frac{v_n-v}{\lambda})\leqslant c(|\lambda|^{-1})K\Psi(v_n)+\\ &\Psi(v))\leqslant c(|\lambda|^{-1})K\big[m(t)+\Psi(v)\big])\in L^1. \end{split}$$

Now, by Dominated Convergence Theorem, we get  $\int \Psi(\frac{v_n - v}{\lambda}) dt \to 0$  for every  $\lambda > 0$ . Thus,  $v_n \to v$  in  $L^{\Psi}$ .

The continuity of I' follows from the continuity of  $\nabla_x \mathcal{L}$  and  $\nabla_y \mathcal{L}$  using the formula (21).

# ALGO DE ESTO ES NECESARIO SI VAMOS A USAR LAS RELACIONES < y $\ll$ , QUIZAS EN LOS PRELIMINARES

There exist several orders and equivalence relations between N-functions (see [17, Sec. 2.2]). Following [17, Def. 1, pp. 15-16] we say that the N-function  $\Phi_2$  is *stronger* than the N-function  $\Phi_1$ , in symbols  $\Phi_1 \prec \Phi_2$ , if there exist a > 0 and  $x_0 \ge 0$  such that

$$\Phi_1(x) \leqslant \Phi_2(ax), \quad x \geqslant x_0. \tag{30}$$

The N-functions  $\Phi_1$  and  $\Phi_2$  are equivalent  $(\Phi_1 \sim \Phi_2)$  when  $\Phi_1 < \Phi_2$  and  $\Phi_2 < \Phi_1$ . We say that  $\Phi_2$  is essentially stronger than  $\Phi_1$   $(\Phi_1 \ll \Phi_2)$  if and only if for every a > 0 there exists  $x_0 = x_0(a) \geqslant 0$  such that (30) holds. Finally, we say that  $\Phi_2$  is completely stronger than  $\Phi_1$   $(\Phi_1 \ll \Phi_2)$  if and only if for every a > 0 there exist K = K(a) > 0 and K = K(a) > 0 and K = K(a) > 0 such that

We assume that there exist an  $N_{\infty}$  function  $\Phi_0$  such that  $\Phi_0 \ll \Phi$ , a function  $a \in L^1([0,T],\mathbb{R})$  such that  $a \geqslant 1$  and a constant M > 0 such that

$$\Psi_0(a^{-1}(t)\nabla_x F) \leqslant \Phi_0(x), \ |x| \geqslant M$$
 (Sub)

**Proposition 3.5.** Let 1 and suppose that <math>F satisfies (A). Then,

$$|\nabla_x F| \le a(t)|x|^{p_0'-1} + b(t), \text{ for } p_0' < p,$$
 (31)

if and only if F satisfies (Sub) with  $\Phi_0(x) = |x|^{p'}$  for all  $p' \in (1, p)$ .

*Proof.* Suppose that (Sub) holds with  $\Phi_0(x) = |x|^{p'}$  for 1 < p' < p - 1. Then,

$$(a^{-1}(t)|\nabla_x F|)^{q'_0} \leq |x|^{p'_0}, \ for \ |x| \geq M$$

As F sastisfies (A), we have  $|\nabla_x F| \le b(t) \in L^1$  for  $|x| \le M$ . PONEMOS SOLO b O Kb?? No confunde? Then,

$$|\nabla_x F| \le a|x|^{\frac{p'_0}{q'_0}} + b(t) \le a|x|^{p'_0-1} + b(t)$$

Now, we assume that F satisfies (31). If  $|x| \ge 1$ , we have

$$|\nabla_x F| \le a(t)|x|^{p_0'-1} + b(t) \le (a(t) + b(t))|x|^{\frac{p_0'}{q_0'}}$$

which is condition (Sub) with a(t) + b(t) instead of a(t).

**Theorem 3.6.** Let  $\Phi$  be an N-function whose complementary function  $\Psi$  satisfies the  $\Delta_2^{\infty}$ -condition.

Let F be a potential that satisfies (C), (A) and the following conditions:

1. (Sub) for some  $N_{\infty}$  function  $\Phi_0$  such that  $\Phi_0 \ll \Phi$ 

2.

$$\lim_{|x| \to \infty} \frac{\int_0^T F(t, x) dt}{\Psi_0(2|x|)} = +\infty, \tag{32}$$

for some N-function  $\Psi_0$  complementary function of  $\Phi_0$ .

EN LO OUE SIGUE, PONEMOS LA NUEVA CONDICION DE ESCTRUCTURA O DEFINIMOS UN CONJUNTO DONDE SE VERIFICAN ESAS CONDICIONES???

Now, if the lagrangian  $\mathcal{L}(t,x,y)$  is strictly convex with respect to  $y \in \mathbb{R}^d$ ,  $\mathcal{L} \in$  $\mathfrak{A}(a,b,c,\lambda,f,\Phi)$ ,  $D_y\mathcal{L}(0,x,y)=D_y\mathcal{L}(T,x,y)$  and (??) holds, then the problem (P) has at least a solution which minimizes the action integral I on  $W^1E_T^{\Phi}$ .

### FALTA UNA BUENA REVISION DE LO ESCRITO EN LA PRUEBA

*Proof.* By the decomposition  $u = \overline{u} + \tilde{u}$ , Young's inequality, (Sub), the convexity of  $\Phi_0$ and (15), we obtain

$$\left| \int_{0}^{T} F(t,u) - F(t,\overline{u}) dt \right| \leqslant \int_{0}^{T} \int_{0}^{1} \left| \nabla_{x} F(t,\overline{u} + s\widetilde{u}) \widetilde{u} \right| ds dt$$

$$= \int_{0}^{T} a(t) \int_{0}^{1} a^{-1}(t) \left| \nabla_{x} F(t,\overline{u} + s\widetilde{u}) \widetilde{u} \right| ds dt$$

$$\leqslant \int_{0}^{T} a(t) \int_{0}^{1} \Psi_{0}(a^{-1}(t) \nabla_{x} F(t,\overline{u} + s\widetilde{u})) + \Phi_{0}(\widetilde{u}(t)) ds dt$$

$$\leqslant \int_{0}^{T} a(t) \left[ \int_{0}^{1} \Phi_{0}(\overline{u} + s\widetilde{u}) ds + \Phi_{0}(\widetilde{u}(t)) \right] dt$$

$$\leqslant \int_{0}^{T} a(t) (\Phi_{0}(2\overline{u}) + 2\Phi_{0}(2\widetilde{u})) dt$$

$$\leqslant C_{1}\Phi_{0}(2\overline{u}) + \int_{0}^{T} a(t) 2\Phi_{0}(2\widetilde{u}) dt$$

$$\leqslant C_{1}\Phi_{0}(2\overline{u}) + C_{2} \int_{0}^{T} \Phi_{0}(C_{T}u'(t)) dt$$

$$(33)$$

Then, by (33) we get

$$\int_{0}^{T} \Phi(u') + F(t, u) dt = \int_{0}^{T} \{\Phi(u') + [F(t, u) - F(t, \overline{u})] + F(t, \overline{u})\} dt$$

$$\geq \int_{0}^{T} [\Phi(u') - C_{2}\Phi_{0}(C_{T}u'(t))] dt - C_{1}\Phi_{0}(2\overline{u}) + \int_{0}^{T} F(t, \overline{u}) dt$$
(34)

If  $\|u\|_{W^1L^\Phi} \to \infty$ , then  $\|u'\|_{L^\Phi} \to \infty$  or  $|\overline{u}| \to \infty$ . Now, if  $|\overline{u}| \to \infty$ , by hypothesis we have  $\frac{1}{\Psi_0(2\overline{u})} \int_0^\infty F(t,\overline{u}) \, dt \to \infty$ .

If  $\|u_n\|_{L^\Phi} \to \infty$  for a sequence  $u_n$ . Let M>0 such that  $\frac{\Phi(x)}{C_2\Phi_0(C_Tx)}\geqslant 2$  for  $|x|\geqslant M$ . Set  $A_n:=\{t\in[0,T]:|u_n'(t)|\geqslant M\}$ . Now, if  $\|u_n'\|_{L^\Phi}\to\infty$  then  $\|u_n'\chi_{A_n}\|_{L^\Phi}\to\infty$  or  $\|u_n'\chi_{A_n^C}\|_{L^\Phi}\to\infty$ .

As  $\Phi$  is continuous, then  $\Phi$  is bounded on  $\overline{B}_r(0) = \{x \in \mathbb{R}^n : |x| \le r\}$ Let  $M_r \coloneqq \max_{\overline{B}_r(0)} \Phi(x)$ .

As  $M_r \to 0$  when  $r \to 0$ , we can choose r such that  $M_r T \leqslant 1$ . Then  $\int_0^T \Phi(\frac{u}{\frac{1}{r} \|u\|_{L^{\infty}}}) dt \leqslant M_r T \leqslant 1$  and consequently  $\|u\|_{L^{\Phi}} \leqslant r^{-1} \|u\|_{L^{\infty}}$ , i.e.

 $L^{\infty} \hookrightarrow L^{\Phi}$ . Thus,  $\|u_n'\chi_{A_n^C}\|_{L^{\Phi}}$  is bounded and  $\|u_n'\chi_{A_n}\|_{L^{\Phi}} \to \infty$ . By Ameniya norm, we have  $\|u_n'\|_{L^{\Phi}} \leqslant 1 + \int_{A_n} \Phi(u_n') \, dt$ , then

$$\int_{A_n} \Phi(u_n') dt \to \infty. \tag{35}$$

Now, as  $\frac{\Phi(x)}{2} \ge C_2 \Phi_0$  for  $|x| \ge M$ , we have

$$\int_{0}^{T} \left[ \Phi(u'_{n}) - C_{2} \Phi_{0}(C_{T} u'_{n}) \right] dt 
\geqslant \int_{A_{n}} \left[ \Phi(u'_{n}) - C_{2} \Phi_{0}(C_{T} u'_{n}) \right] dt - C_{2} \int_{[0,T]/A_{n}} \Phi_{0}(C_{T} u'_{n}) dt 
\geqslant \frac{1}{2} \int_{A_{n}} \Phi(u'_{n}) dt - C_{2} T \max \{ \Phi_{0}(x) : ||x|| \leqslant C_{T} M \}.$$
(36)

Finally, by (35), we conclude that  $\int_0^T \left[ \Phi(u_n') - C_2 \Phi_0(C_T u_n') \right] dt \to \infty$ 

## Acknowledgments

The authors are partially supported by a UNRC grant number 18/C417. The first author is partially supported by a UNSL grant number 22/F223.

### References

- [1] S. Acinas, L. Buri, G. Giubergia, F. Mazzone, and E. Schwindt. Some existence results on periodic solutions of Euler-Lagrange equations in an Orlicz-Sobolev space setting. *Nonlinear Analysis, TMA.*, 125:681 698, 2015.
- [2] Haim Brezis. Functional analysis, Sobolev spaces and partial differential equations. Springer Science & Business Media, 2010.
- [3] M Chamra and J Maksymiuk. Anisotropic orlicz-sobolev spaces of vector valued functions and lagrange equations. *arXiv* preprint arXiv:1702.08683, 2017.
- [4] F. Clarke. Functional Analysis, Calculus of Variations and Optimal Control. Graduate Texts in Mathematics. 2013.
- [5] W. Desch and R. Grimmer. On the well-posedness of constitutive laws involving dissipation potentials. *Trans. Amer. Math. Soc*, (353):5095–5120, 2001.

- [6] M.A. Krasnosel'skii, P.P. Zabreyko, E.I. Pustylnik, and P.E. Sobolevski. *Integral operators in spaces of summable functions*. Mechanics: Analysis. Springer Netherlands, 2011.
- [7] Chun Li, Zeng-Qi Ou, and Chun-Lei Tang. Periodic solutions for non-autonomous second-order differential systems with (q, p)-laplacian. *Electronic Journal of Differential Equations*, 2014(64):1–13, 2014.
- [8] David G Luenberger and Yinyu Ye. *Linear and nonlinear programming*, volume 228. Springer, 2015.
- [9] J. Mawhin and M. Willem. *Critical point theory and Hamiltonian systems*. Springer-Verlag, New York, 1989.
- [10] Daniel Pasca. Periodic solutions of a class of nonautonomous second order differential systems with (q, p)-laplacian. *Bulletin of the Belgian Mathematical Society-Simon Stevin*, 17(5):841–851, 2010.
- [11] Daniel Paşca and Chun-Lei Tang. Some existence results on periodic solutions of nonautonomous second-order differential systems with (q, p)-laplacian. *Applied Mathematics Letters*, 23(3):246–251, 2010.
- [12] Daniel Pasca and Chun-Lei Tang. Some existence results on periodic solutions of ordinary (q, p)-laplacian systems. *Journal of applied mathematics & informatics*, 29(1.2):39–48, 2011.
- [13] Daniel Pasca and Zhiyong Wang. On periodic solutions of nonautonomous second order hamiltonian systems with (q, p)-laplacian. *Electronic Journal of Qualitative Theory of Differential Equations*, 2016(106):1–9, 2016.
- [14] Ryszard Płuciennik. Boundedness of the superposition operator in generalized Orlicz spaces of vector-valued functions. *Bull. Pol. Acad. Sci., Math.*, 33:531â540, 1985.
- [15] Ryszard Płuciennik. On some properties of the superposition operator in generalized Orlicz spaces of vector-valued functions. *Ann. Soc. Math. Pol., Ser. I, Commentat. Math.*, 25:321â337, 1985.
- [16] Ryszard Płuciennik. The superposition operator in Musielak-Orlicz spaces of vector-valfued functions. Abstract analysis, Proc. 14th Winter Sch., Srní/Czech. 1986, Suppl. Rend. Circ. Mat. Palermo, II. Ser. 14, 411-417 (1987)., 1987.
- [17] M. M. Rao and Z. D. Ren. *Theory of Orlicz spaces*, volume 146. Marcel Dekker, Inc., New York, 1991.
- [18] G. Schappacher. A notion of Orlicz spaces for vector valued functions. *Appl. Math.*, 50(4):355–386, 2005.
- [19] M. S. Skaff. Vector valued Orlicz spaces. II. *Pacific J. Math.*, 28(2):413–430, 1969.

- [20] C.-L. Tang. Periodic solutions for nonautonomous second order systems with sublinear nonlinearity. *Proc. Amer. Math. Soc.*, 126(11):3263–3270, 1998.
- [21] C. L. Tang and X.-P. Wu. Periodic solutions for second order systems with not uniformly coercive potential. *J. Math. Anal. Appl.*, 259(2):386–397, 2001.
- [22] Chun-Lei Tang. Periodic solutions of non-autonomous second-order systems with  $\gamma$ -quasisubadditive potential. *Journal of Mathematical Analysis and Applications*, 189(3):671–675, 1995.
- [23] X. Tang and X. Zhang. Periodic solutions for second-order Hamiltonian systems with a *p*-Laplacian. *Ann. Univ. Mariae Curie-Skłodowska Sect. A*, 64(1):93–113, 2010.
- [24] Y. Tian and W. Ge. Periodic solutions of non-autonomous second-order systems with a *p*-Laplacian. *Nonlinear Anal.*, 66(1):192–203, 2007.
- [25] X.-P. Wu and C.-L. Tang. Periodic solutions of a class of non-autonomous second-order systems. *J. Math. Anal. Appl.*, 236(2):227–235, 1999.
- [26] Xiaoxia Yang and Haibo Chen. Periodic solutions for a nonlinear (q, p)-laplacian dynamical system with impulsive effects. *Journal of Applied Mathematics and Computing*, 40(1-2):607–625, 2012.
- [27] Xiaoxia Yang and Haibo Chen. Existence of periodic solutions for sublinear second order dynamical system with (q, p)-laplacian. *Mathematica Slovaca*, 63(4):799–816, 2013.
- [28] F. Zhao and X. Wu. Periodic solutions for a class of non-autonomous second order systems. *J. Math. Anal. Appl.*, 296(2):422–434, 2004.