An introduction to fractional calculus

Fundamental ideas and numerics

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We know a general way to obtain FLMM methods of the form

$$y^{(n)} = T_{m-1}(t_n) + \tau^{\beta} \sum_{j=0}^{s} w_{n,j} f(t_j, y^{(j)}) + \tau^{\alpha} \sum_{j=0}^{n} \omega_{n-j} f(t_j, y^{(j)}),$$

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- ightharpoons we make the starting values for a multi-step method,
- ightharpoons we still need to discuss how we can efficiently treat the memory term.

To initialize the computation we need the values $y^{(0)}, \ldots, y^{(s)}$, s+1, $s=|\mathcal{A}|=\mathcal{A}_{p-1}\cup\{p-1\}$ with p the order of convergence of the FLMM, and $\mathcal{A}_{p-1}=\{\nu\in\mathbb{R}\mid \nu=i+j\alpha,\quad i,j\in\mathbb{N}, \nu< p-1\}.$

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- To avoid mixing methods we evaluate all the approximations at the same time by solving

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- This will be in general an $s \times \dim(y^{(j)})$ nonlinear system that we need to solve before starting the iteration.
- If the value of α is not very small, viz s is moderate, and the system of ODEs is moderate this is manageable.

If we compute the sum on the coefficients ω_i naively for

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we end up having a $O(N^2)$ cost! If we do not perform this task efficiently the numerical solution degenerates in an unworkable task as we either refine our grid or enlarge our computational domain.

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- We can try to "forget" part of the lag-term,
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- We can try an approach with nested meshes to reduce the load,
- We can exploit the fact that this is a convolution and adopt some FFT tricks.

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• If we now want to compute the next r approximations we can separate the lag term as

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If we call $S_r(n,0,r-1) = \sum_{j=0}^{r-1} c_{n-j} f_j$, $n \in \{r,r+1,\ldots,2r-1\}$, the set of partial sums each of length r we can evaluate them with FFT in $O(2r \log_2(2r))$.

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 We can apply the same process recursively if we double every time-interval under consideration

$$y^{(n)} = \phi_n + \sum_{j=0}^{2r-1} c_{n-j} f_j + \sum_{j=2r}^{n} c_{n-j} f_j, \quad n \in \{2r, 2r+1, \dots, 3r-1\},$$

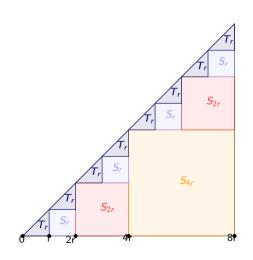
$$y^{(n)} = \phi_n + \sum_{j=0}^{2r-1} c_{n-j} f_j + \sum_{j=2r}^{3r-1} c_{n-j} f_j + \sum_{j=3r}^{n} c_{n-j} f_j, \quad n \in \{3r, 3r+1, \dots, 4r-1\},$$

If we call $S_{2r}(n,0,2r-1)=\sum_{j=0}^{2r-1}c_{n-j}f_j$, and $S_r(n,2r,3r-1)=\sum_{j=2r}^{3r-1}c_{n-j}f_j$ the set of partial sums of lengths 2r and r we can evaluate them with FFT in $O(4r\log_2(4r))$ and $O(2r\log_2(2r))$ respectively.

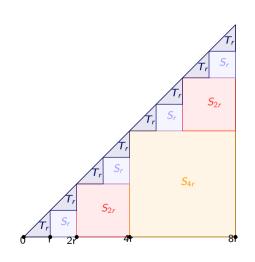
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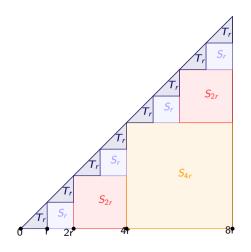


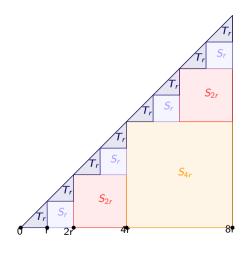
• We can iterate the process for the 4r approximations in the interval $n \in \{4r, \ldots, 8r-1\}$, together with the partial sums $S_{4r}(n, 0, 4r-1)$, $S_{2r}(n, 4r, 6r-1)$, $S_r(n, 6r, 7r-1)$ that can be evaluated in $O(8r \log_2(8r))$, $O(4r \log_2(4r))$ and $O(2r \log_2(2r))$ respectively.



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- At each level we have to complete the recursion by computing

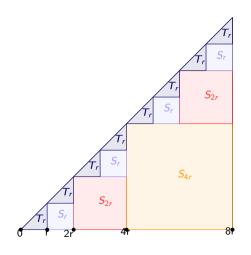
$$T_r(p,n) = \sum_{j=p}^n c_{n-j} f_j, \ p = \ell r,$$
 $n \in \{\ell r, \ell r + 1, \dots, (\ell+1)r - 1\},$ $\ell = 0, 1, 2, \dots$



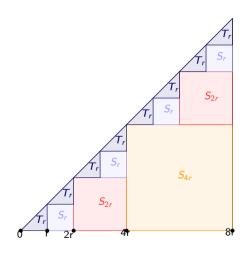


To determine the whole cost we just have to sum the various components

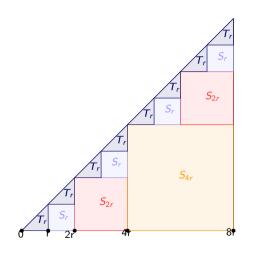
• Assume that $N = 2^{n_t}$



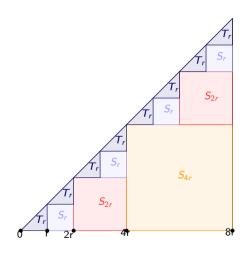
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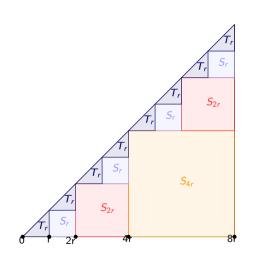
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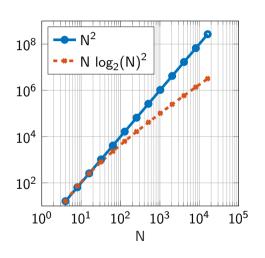


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- In general:

$$N \log_2 N + 2 \frac{N}{2} \log_2 \frac{N}{2} + 4 \frac{N}{4} \log_2 \frac{N}{4} + \cdots + p \frac{N}{p} \log_2 \frac{N}{p} + \frac{N}{r} \frac{r(r+1)}{2}, \quad p = \frac{N}{2r}$$



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- In general:

$$\sum_{j=0}^{\log_2 p} N \log_2 \frac{N}{2^j} + N \frac{r+1}{2} = O(N(\log_2 N)^2).$$

Short-memory principle (Ford and Simpson 2001)

We can try to use a "fixed memory length" to reduce the computational (and memory) load.

$$y(t_{n+1}) = y(t_n) + \frac{1}{\Gamma(\alpha)} \int_{t_n}^{t_{n+1}} (t_{n+1} - \tau)^{\alpha - 1} f(\tau, x(\tau)) d\tau + \frac{1}{\Gamma(\alpha)} \int_{0}^{t_n} ((t_{n+1} - \tau)^{\alpha - 1} - (t_n - \tau)^{\alpha - 1}) f(\tau, y(\tau)) d\tau, \quad \alpha \in (0, 1).$$

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Let us introduce now a fixed window T_M of memory, then

$$E = \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_n - T_M} ((t_{n+1} - \tau)^{\alpha - 1} - (t_n - \tau)^{\alpha - 1}) f(\tau, y(\tau)) d\tau \right|$$

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$$= \frac{M}{\Gamma(\alpha)} \left| (z_{1}^{*})^{\alpha - 1} \tau - (z_{2}^{*})^{\alpha - 1} \tau \right| \quad z_{1}^{*} \in [T_{M}, T_{M} + \tau], \ z_{2}^{*} \in [t_{n}, t_{n+1}]$$

We can try to use a "fixed memory length" to reduce the computational (and memory) load.

$$y(t_{n+1}) = y(t_n) + \frac{1}{\Gamma(\alpha)} \int_{t_n}^{t_{n+1}} (t_{n+1} - \tau)^{\alpha - 1} f(\tau, x(\tau)) d\tau + \frac{1}{\Gamma(\alpha)} \int_{0}^{t_n} ((t_{n+1} - \tau)^{\alpha - 1} - (t_n - \tau)^{\alpha - 1}) f(\tau, y(\tau)) d\tau, \quad \alpha \in (0, 1).$$

$$E = \left| \frac{1}{\Gamma(\alpha)} \int_0^{t_n - T_M} ((t_{n+1} - \tau)^{\alpha - 1} - (t_n - \tau)^{\alpha - 1}) f(\tau, y(\tau)) d\tau \right|$$

$$\leq \frac{M}{\Gamma(\alpha)} \left| \int_0^{t_n - T_M} ((t_{n+1} - \tau)^{\alpha - 1} - (t_n - \tau)^{\alpha - 1}) d\tau \right|$$

$$< \frac{M}{\Gamma(\alpha)} T_M^{\alpha - 1} \tau, \qquad \alpha \in (0, 1).$$

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Let us introduce now a fixed window T_M of memory, then If we have a global error bound $E_{\rm global}$ with step-length τ we just need to choose

$$T_M > \left(\frac{M}{\Gamma(\alpha)E_{\mathsf{global}}}\right)^{1/1-\alpha}, \quad \alpha \in (0,1),$$

while if we have a local error bound E_{local}

$$T_M > \left(\frac{M\tau}{\Gamma(\alpha)E_{local}}\right)^{1/1-\alpha}, \quad \alpha \in (0,1).$$

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The idea can be refined by using nested meshes.

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Scaling properties

$$I_{[0,t]}^{\alpha}f(t)=\int_0^t rac{f(au)}{(t- au)^{1-lpha}}\,\mathrm{d} au$$

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Scaling properties

$$I_{[0,t]}^{\alpha}f(wt) = \int_0^t \frac{f(\tau)}{(wt-\tau)^{1-\alpha}} d\tau$$

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Scaling properties

$$I_{[0,t]}^{\alpha}f(wt)=w^{\alpha}\int_{0}^{t}\frac{f(w au)}{(t- au)^{1-lpha}}\,\mathrm{d} au$$

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Scaling properties

Given $p \in \mathbb{N}$ we then have

$$I_{[0,t]}^{\alpha} f(w^p t) = w^{p\alpha} \int_0^t \frac{f(w^p \tau)}{(t-\tau)^{1-\alpha}} d\tau$$

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• We can use the weight on the mesh

$$\Omega_{\tau}^{\alpha}f(\emph{n}\tau) pprox \emph{I}_{[0,t]}^{\alpha}f(\emph{n}\tau),$$
 step length τ

to compute

$$\Omega^{\alpha}_{w^p\tau}f(nw^p\tau)\approx \textit{I}^{\alpha}_{[0,t]}f(nw^p\tau), \text{ step length } w^p\tau$$

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ullet In summary for any $p\in\mathbb{N}$ we get

$$\Omega^{\alpha}_{\tau}f(n\tau) = \sum_{i=0}^{n} \omega_{n-j}f(j\tau) \iff \Omega^{\alpha}_{w^{p}\tau}f(nw^{p}\tau) = w^{p\alpha}\sum_{i=0}^{n} \omega_{n-j}f(jw^{p}\tau).$$

Nested mesh

Given $\tau \in \mathbb{R}^+$, the mesh $M_{\tau} = \{\tau n, \ n \in \mathbb{N}\}$. Selected $w, r, p \in \mathbb{N}$, w > 0, r > p, we have $M_{w^p\tau} \supset M_{w^r\tau}$ and we decompose the interval as

$$[0,t] = [0,t-w^mT] \cup [t-w^mT,t-w^{m-1}T] \cup \cdots \cup [t-wT,t-T] \cup [t-T,t]$$

for $m \in \mathbb{N}$ the smallest integer such that $t < w^{m+1}T$.

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- We rewrite our integral using the scaling property as

$$I_{[0,t]}^{\alpha}f(t)=I_{[t-T,t]}^{\alpha}f(t)+\sum_{i=0}^{m-1}w^{i\alpha}I_{[t-wT,t-T]}^{\alpha}f(w^{i}t)+w^{m\alpha}I_{[0,t-T]}^{\alpha}f(w^{m}t).$$

In the discrete approximation of

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we approximate

$$\Omega^{\alpha}_{\tau,[t-w^{i+1}T,t-w^{i}T]}f(t)\approx\Omega^{\alpha}_{w^{i}\tau,[t-w^{i+1}T,t-w^{i}T]}f(t)$$

and substitute

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Theorem (Ford and Simpson 2001, Theorem 1)

The nested mesh scheme preserves the order of the underlying quadrature rule on which it is based.

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Proof. For integration over a fixed interval [0, t] the choice of T fixes (independent of h) the number of subranges over which the integral is evaluated, on each of them we have en error $O(h^p)$.

• The first benefit is that we evaluate a fixed number of quadrature coefficients and then re-use them on all successive intervals,

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- This approach cost $O(w^m)$ with respect to $O(w^{2m})$ of the full method,
- We could use linear extrapolation techniques to improve the results.
- Selecting the various parameter may need a bit of tuning.

Available codes

With respect to the ordinary case for which there exists many reliable and high-performance codes, the choices for computing the solution of fractional differential equation is much more *sparse*.

- From (Garrappa 2018)
 - FDE_PI1_Ex.m Explicit Product-Integration of rectanguar type
 - FDE_PI1_Im.m Implicit Product-Integration of rectanguar type
 - FDE_PI2_Im.m Implicit Product-Integration of trapezoidal type
 - FDE_PI12_PC.m Product-Integration with predictor-corrector
- From (Garrappa 2015)
 - FLMM2 Matlab code Three implicit second order Fractional Linear Multistep Methods.

A remark

All these methods use direct-solver for the Newton method inside them, there is space to make improvement on the solution strategies. Furthermore, a challenge that yet remains: can we find a strategy that combines the convolution features and savings on the memory?

What do we have now

We know a general way to obtain FLMM methods of the form

$$y^{(n)} = T_{m-1}(t_n) + \tau^{\beta} \sum_{j=0}^{s} w_{n,j} f(t_j, y^{(j)}) + \tau^{\alpha} \sum_{j=0}^{n} \omega_{n-j} f(t_j, y^{(j)}),$$

- $oldsymbol{arphi}$ starting from the polynomials $(
 ho,\sigma)$ of an implicit order ho method,
- igotimes we have seen how to compute the convolution coefficients ω_n ,
- \bigcirc we have seen how to compute the starting nodes $w_{n,j}$,
- we know how we can compute the starting values for a multi-step method by solving a nonlinear system with Newton,
- we have some hints on how we can efficiently treat the memory term.

Let us write everything for a case, let us start from the 2nd order BDF formula for ODEs

$$y^{(n+2)} - \frac{4}{3}y^{(n+1)} + \frac{1}{3}y^{(n)} = \frac{2}{3}\tau f_{n+2},$$

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• First of all we write down the (ρ, σ) polynomials defining the scheme:

$$\rho(\zeta) = \zeta^2 - \frac{4}{3}\zeta + \frac{1}{3}, \qquad \sigma(\zeta) = \frac{2}{3}\zeta^2.$$

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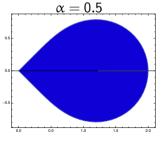
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$$\rho(\zeta)=\zeta^2-\frac{4}{3}\zeta+\frac{1}{3}, \qquad \sigma(\zeta)=\frac{2}{3}\zeta^2.$$

 \bullet Then we compute the generating function $\omega(\zeta)$

$$\omega(\zeta) = \frac{\rho(1/\zeta)}{\sigma(1/\zeta)} = \frac{2}{3\left(1 - 4\zeta/3 + \zeta^2/3\right)}.$$



$$\{1/\omega(\zeta)^{\alpha} : |\zeta| \le 1\}$$

Now we need to expand the convolution coefficients of

$$\omega^{\alpha}(\zeta) = (\omega(\zeta))^{\alpha} = \frac{2^{\alpha}}{3^{\alpha}}(1 - 4\zeta/3 + \zeta^2/3)^{-\alpha}.$$

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Theorem (Henrici 1974, Theorem 1.6c, p. 42)

Let $\phi(\zeta) = 1 + \sum_{n=1}^{+\infty} a_n \zeta^n$ be a formal power series. Then for any $\alpha \in \mathbb{C}$, we have

$$(\phi(\zeta))^{\alpha} = \sum_{n=0}^{+\infty} v_n^{(\alpha)} \zeta^n,$$

where coefficients $v_n^{(\alpha)}$ can be evaluated recursively as

$$v_0^{(\alpha)} = 1,$$
 $v_n^{(\alpha)} = \sum_{j=1}^n \left(\frac{(\alpha+1)j}{n} - 1 \right) a_j v_{n-j}^{(\alpha)}$

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- $a_1 = -4/3$, $a_2 = 1/3$, $a_j = \text{if } j \ge 3$, thus using

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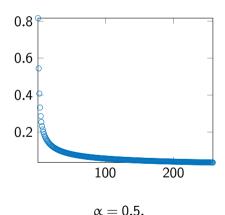
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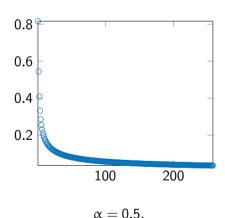
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• we get $\tilde{\omega}_0 = 1$, $\tilde{\omega}_1 = \frac{4}{3}\alpha\tilde{\omega}_0 = \frac{4\alpha}{3}$,

$$\tilde{\omega}_n = \frac{4}{3} \left(1 + \frac{\alpha - 1}{n} \right) \tilde{\omega}_{n-1} + \frac{4}{3} \left(\frac{2(1 - \alpha)}{n} - 1 \right) \tilde{\omega}_{n-2}.$$

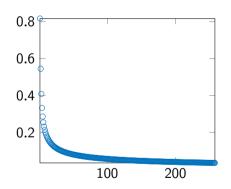


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- Since the a_j are a finite small number, we can compute the coefficients in an O(N) operations,
- We can solve $_{CA}D^{0.5}_{[0,2]}y(t)=-2y(t),\ y(0)=1$

	*	
τ	$ y^{(n)}-y(2) $	order
2^{-6}	1.44e-04	1.61
2^{-7}	4.42e-05	1.71
2^{-8}	1.28e-05	1.79
2^{-9}	3.57e-06	1.84
2^{-10}	9.68e-07	1.88
2^{-11}	2.85e-07	1.76
2^{-12}	8.17e-08	1.80
2^{-13}	2.29e-08	1.84
2^{-14}	6.27e-09	1.87



$$\alpha = 0.5$$
.

- Since the a_j are a finite small number, we can compute the coefficients in an O(N) operations,
- We can solve $_{CA}D^{0.5}_{[0,2]}y(t) = -2y(t)$, y(0) = 1
- For the starting weights we have to solve a 3×3 Vandermonde system:

$$\begin{bmatrix} 1 & 1 & 1 \\ 0 & 1 & \sqrt{2} \\ 0 & 1 & 2 \end{bmatrix} \begin{bmatrix} w_{n,0} \\ w_{n,1} \\ w_{n,2} \end{bmatrix} = \begin{bmatrix} b_1 \\ b_2 \\ b_3 \end{bmatrix}$$

number of time-step times.

Reuse

Since we have a fixed time-grid we can reuse the same factorization for the Vandermonde system and compute all the weights in a single sweep.

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The Brusselator is a model of the autocatalytic chemical reaction, it is described by

$$\begin{cases} \dot{x}_1 = a - (\mu + 1)x_1 + x_1^2 x_2, \\ \dot{x}_2 = \mu x_1 - x_1^2 x_2, \end{cases} \quad a, \mu > 0.$$

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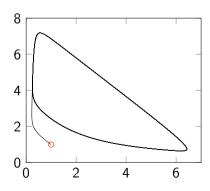
- If $\mu > a^2 + 1$ then a single Brusselator has a unique limit cycle,
- If $(a-1)^2 < \mu \le a^2 + 1$ all the orbits tend to the steady state.

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 $a, \mu > 0.$

```
a = 1 : mu = 4 :
param = [ a , mu ] ;
f_fun = Q(t,y,par) [ ...
par(1) - (par(2)+1)*y(1) + y(1)^2*y(2) ; ...
par(2)*v(1) - v(1)^2*v(2);
t0 = 0 : T = 100 :
y0 = [1 ; 1];
[T,Y] = ode45(@(t,y)
\rightarrow f fun(t,y,param),[t0,T],y0);
figure(1)
plot(Y(:,1),Y(:,2),'k-',y(1,1),y(2,1),'ro')
```

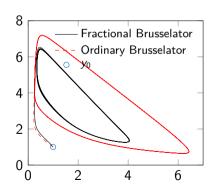


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The cycle of the single fractional Brusselator is contained in the region

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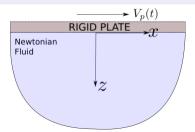
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Of interest (Wang and Li 2007)

Finding the smallest values α_1 , α_2 for which a limit cycle exist is of interest.

Stoke's Second Problem

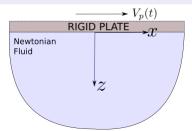
Can we determine the behavior of a half-space of Newtonian, viscous fluid undergoing the motion induced by the prescribed uniform sinusoidal motion of a plate on the surface?



If we write down the equation of motion we find

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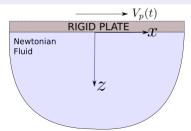
$$\rho \frac{\partial v}{\partial t} = \mu \frac{\partial^2 v}{\partial z^2}$$

• ρ is the *fluid density*, μ is the viscosity, v is the profile of the *transverse fluid velocity*.

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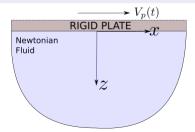


$$\rho[s\tilde{v}(s,z)-v(0,x)]=\mu\frac{d^2\tilde{v}(s,z)}{dz^2},$$

- ρ is the *fluid density*, μ is the viscosity, *v* is the profile of the *transverse fluid velocity*.
- We apply Laplace transform to the equation $\tilde{v} = \mathcal{L}v(s)$,

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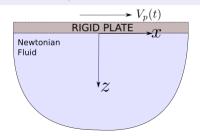


$$\tilde{v}(s,z) = \tilde{v}_{\rho}(s) \exp\left(\sqrt{\frac{\rho s}{\mu}}z\right)$$

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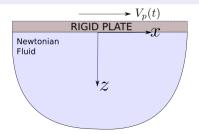


$$\tilde{\sigma}(s,z) = \sqrt{\mu\rho}\sqrt{s}\tilde{v}(s,z) = \sqrt{\mu\rho}\frac{1}{\sqrt{s}}s\tilde{v}(s,z)$$

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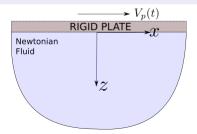


$$ilde{\sigma}(s,z) = \sqrt{\mu
ho} \mathcal{L} \left\{ rac{1}{\Gamma(1/2) t^{1/2}}
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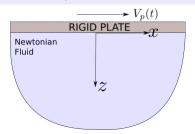


$$\sigma(t,z) = \sqrt{\mu\rho} \frac{1}{\Gamma(1/2)} \int_0^t (t-\tau)^{1/2} v_{\tau}(\tau,z) d\tau.$$

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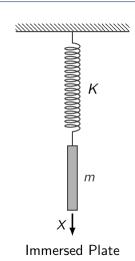
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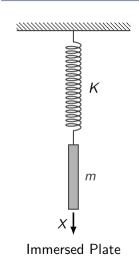
$$\sigma(t,z) = \sqrt{\mu \rho}_{CA} D_{[0,t]}^{1/2} v(t,z).$$

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Assumptions:

- The spring is massless and its oscillations do not disturb the fluid,
- The area A of the plate is sufficiently large as to produce in the fluid adjacent to the plate the velocity field and stresses we just derived,

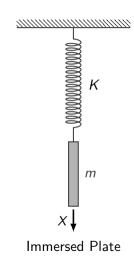


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Using the expression for the strain and $V_p(t,0)=\dot{X}(t)$ we find

$$m\ddot{X} + 2A\sqrt{\mu\rho}_{CA}D_{[0,t]}^{3/2}X + KX = 0.$$

The Bagley-Torvik model is an example of a Linear Multi-Term FDE, that is, something of the form

$$\lambda_{Q\mathit{CA}} D^{\alpha_Q} y(t) + \lambda_{Q-1\mathit{CA}} D^{\alpha_{Q-1}} y(t) + \cdots + \lambda_{2\mathit{CA}} D^{\alpha_2} y(t) + \lambda_{1\mathit{CA}} D^{\alpha_1} y(t) = f(t,y(t)),$$

with

- $\lambda_i \in \mathbb{R} \ \forall i = 1, \ldots, Q$,
- $0 < \alpha_1 < \alpha_2 < \ldots < \alpha_{Q-1} < \alpha_Q$ and $\alpha_Q \neq 0$.

For this problem we have $m_Q = \max m_i$, $m_i = [\alpha_i]$, i = 1, ..., Q initial conditions:

$$y(t_0) = y_0, y'(t_0) = y_0^{(1)}, \dots, y^{(m_Q-1)}(t_0) = y_0^{(m_Q-1)}.$$

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• How can we solve them?

We need to recall one of the properties we have seen of the Caputo derivatives

$$(P_1) \ I_{[t_0,T]}^{\alpha} CAD_{[t_0,T]}^{\alpha} y(t) = y(t) - T_{m-1}[y,t_0](t),$$

$$\begin{array}{ll} (P_2) & I_{[t_0,T]}^{\beta} {}_{CA} D_{[t_0,T]}^{\alpha} y(t) = I_{[t_0,T]}^{\beta} {}_{RL} D_{[t_0,T]}^{\alpha} \left[y(t) - T_{m-1}[y;t_0](t) \right] = \\ & I_{[t_0,T]}^{\beta-\alpha} \left[y(t) - T_{m-1}[y;t_0](t) \right], \; \beta > \alpha. \end{array}$$

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$$(P_2) I_{[t_0,T]}^{\beta} CA D_{[t_0,T]}^{\alpha} y(t) = I_{[t_0,T]}^{\beta} RL D_{[t_0,T]}^{\alpha} [y(t) - T_{m-1}[y;t_0](t)] = I_{[t_0,T]}^{\beta-\alpha} [y(t) - T_{m-1}[y;t_0](t)], \ \beta > \alpha.$$

We start from the multi-term equation

$$\lambda_{QCA}D^{\alpha_Q}y(t) + \lambda_{Q-1CA}D^{\alpha_{Q-1}}y(t) + \cdots + \lambda_{2CA}D^{\alpha_2}y(t) + \lambda_{1CA}D^{\alpha_1}y(t) = f(t,y(t)),$$

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$$\lambda_{Q} I_{[t_{0},T]}^{\alpha_{Q}} [{}_{CA} D^{\alpha_{Q}} y(t)] = -I_{[t_{0},T]}^{\alpha_{Q}} \left[\sum_{i=1}^{Q-1} \lambda_{i} {}_{CA} D^{\alpha_{i}} y(t) + f(t,y(t)) \right],$$

• we multiply both sides by $I_{[t_0,T]}^{\alpha_Q}$,

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$$y(t) - T_{m_Q-1}[y, t_0](t) = -\sum_{i=1}^{Q-1} \frac{\lambda_i}{\lambda_Q} I_{[t_0, t]}^{\alpha_Q - \alpha_i} [y(t) - T_{m_i-1}[y; t_0](t)] + \frac{1}{\lambda_Q} I_{[t_0, T]}^{\alpha_Q} f(t, y(t))$$

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- we multiply both sides by $I_{[t_0,T]}^{\alpha_Q}$,
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- and re-arrange to get an expression for the solution.

First we do a bit of rewriting of

$$y(t) = T_{m_Q-1}[y, t_0](t) - \sum_{i=1}^{Q-1} \frac{\lambda_i}{\lambda_Q} I_{[t_0, t]}^{\alpha_Q - \alpha_i} [y(t) - T_{m_i-1}[y; t_0](t)] + \frac{1}{\lambda_Q} I_{[t_0, T]}^{\alpha_Q} f(t, y(t))$$

we employ the usual fractional integral for polynomials:

$$I_{[t_0,t]}^{\alpha} \mathcal{T}_{m-1}[y;t_0](t) = \sum_{k=0}^{m-1} \frac{(t-t_0)^{k+\alpha}}{\Gamma(k+\alpha)} y^{(k)}(t_0), \quad \substack{\alpha \in \{\alpha_1,\ldots,\alpha_{Q-1}\}, \\ m \in \{m_1,\ldots,m_{Q-1}\}.}$$

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• We use it to simplify the expression

$$\tilde{T}(t) = T_{m_Q-1}[y; t_0](t) + \sum_{i=1}^{Q-1} \frac{\lambda_i}{\lambda_Q} \sum_{k=0}^{m_i-1} \frac{(t-t_0)^{k+\alpha_Q-\alpha_i}}{\Gamma(k+\alpha_Q-\alpha_i+1)} y^{(k)}(t_0).$$

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Now we have an expression that we can treat by adapting one of the Product Integral rules

$$y(t) = \tilde{T}(t) - \sum_{i=1}^{Q-1} \frac{\lambda_i}{\lambda_Q} I_{[t_0,t]}^{\alpha_Q - \alpha_i} y(t) + \frac{1}{\lambda_Q} I_{[t_0,T]}^{\alpha_Q} f(t,y(t)).$$

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We can start from the explicit rectangular product integral rule on a uniform grid

$$y^{(n)} = \tilde{T}(t_n) - \sum_{i=1}^{Q-1} \frac{\lambda_i}{\lambda_Q} \tau^{\alpha_Q - \alpha_i} \sum_{j=0}^{n-1} b_{n-j-1}^{(\alpha_Q - \alpha_i)} y^{(j)} + \frac{1}{\lambda_Q} \sum_{j=0}^{n-1} b_{n-j-1}^{(\alpha_Q)} f(t_j, y^{(j)}).$$

with

$$b_n^{(\alpha)} = [(n+1)^{\alpha} - n^{\alpha}]/\alpha, \quad n = 1, \dots, N.$$

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We can do it similarly for the **Implicit Trapezoidal Rule** and then for the **Predictor-Corrector method** (Diethelm 2003).

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- We don't know how to determine the starting values $w_{n,j}$ for the quadrature. Thus this approach is not viable.

Linear Multi-Term FDEs: generalizing PI rules

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Available codes (Garrappa 2018):

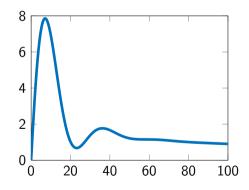
- MT_FDE_PI1_Ex.m Explicit Product-Integration of rectanguar type
- MT_FDE_PI1_Im.m Implicit Product-Integration of rectanguar type
- MT_FDE_PI2_Im.m Implicit Product-Integration of trapezoidal type
- MT_FDE_PI12_PC.m- Product-Integration with predictor-corrector

Linear Multi-Term FDEs: back to Bagley-Torvik

We reached the equation

$$m\ddot{X} + 2A\sqrt{\mu\rho}_{CA}D_{[0,t]}^{3/2}X + KX = 0.$$

```
m = 10; A = 6; K = 3;
mu = 2; rho = 2;
alpha = [2 3/2];
lambda = [m \ 2*A*sqrt(mu*rho)] ;
f_fun = O(t,X) - K*X;
J fun = Q(t,X) - K;
t0 = 0 ; T = 100 ;
XO = [0, 2];
h = 1e-2:
[t, X] = mt fde pi1 ex(alpha, lambda, f fun,
\rightarrow t0, T, X0, h);
```

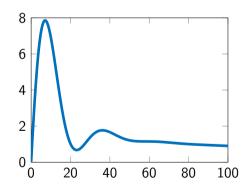


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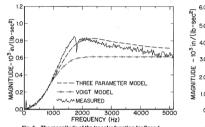
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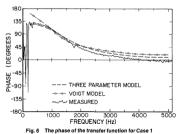


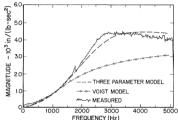
But does it fit the reality?

Linear Multi-Term FDEs: back to Bagley-Torvik

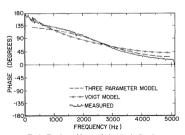


The magnitude of the transfer function for Case 1





The magnitude of the transfer function for Case 5



The phase of the transfer function for Case 5

The model we have derived is a model of the form

$$\begin{split} \sigma(t) &= G_0 \varepsilon(t) + G_1 \dot{\varepsilon}(t), \\ \varepsilon(t) &= \frac{x(t)}{\delta} \\ f(t) &= m \ddot{x}(t) + f_p(t), \\ f_p(t) &= \frac{2A}{\delta} (G_0 + G_{1CA} D^{\alpha} x(t)). \end{split}$$

One can do parameter tuning to find the fractional order from experimental data and compare the results with the integer-order model. The results on the left by Bagley and Torvik 1986 show that the fractional model obtain a hetter fit with the measured data

In the integer-order case we know how to rewrite the equation

$$y^{(n)}(t) = f(t, y^{(n-1)}(t), \dots, y^{(1)}(t), y(t)), \quad y^{(j)}(0) = y_0^{(j)}, \quad j = 0, 1, \dots, n-1,$$

as a system of first-order equations.

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as a system of first-order equations.

Can we do something similar in the fractional case?

In the integer-order case we know how to rewrite the equation

$$y^{(n)}(t) = f(t, y^{(n-1)}(t), \dots, y^{(1)}(t), y(t)), \quad y^{(j)}(0) = y_0^{(j)}, \quad j = 0, 1, \dots, n-1,$$

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(A1) Let us assume that our multi-term equation is of the form

$$_{CA}D^{lpha_k}y(t) = f(t, _{CA}D^{lpha_{k-1}}y(t), \ldots, _{CA}D^{lpha_1}(t), y(t)), \qquad y^{(j)}(0) = y_0^{(j)}, \ j = 0, 1, \ldots, n-1,$$

for
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(A2) Assume also that $\alpha_j \in \mathbb{Q} \ \forall j = 1, 2, ..., k$, and that M is the least common multiple of $\alpha_1, \alpha_2, ..., \alpha_k$.

Theorem (Diethelm 2010, Theorem 8.1)

Under the assumptions (A1) and (A2), set $\gamma=1/M$, and $N=M\alpha_k$, then the IVP is equivalent to

$$\begin{cases} {}_{CA}D^{\gamma}y_{0}(t) = y_{1}(t), \\ {}_{CA}D^{\gamma}y_{1}(t) = y_{2}(t), \\ \vdots \\ {}_{CA}D^{\gamma}y_{N-2}(t) = y_{N-1}(t), \\ {}_{CA}D^{\gamma}y_{N-1}(t) = f(t,y_{0}(t),y_{\alpha_{k-1/M}}(t),\dots,y_{\alpha_{1/M}}(t),y(t)) \end{cases} y_{i}(0) = \begin{cases} y_{0}^{(i/m)}, & \text{if } \frac{j}{M} \in \mathbb{N}_{0}, \\ 0, & \text{otherwise.} \end{cases}$$

 \Rightarrow whenever $\mathbf{y} = (y_0, \dots, y_{N-1})^T$ with $y_0 \in \mathcal{C}^{\lceil \alpha_k \rceil}[0, b]$, for some b > 0, is a solution of the *N*-dimensional system, then $y \equiv y_0$ is a solution of the multi-term FDE.

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 \leftarrow whenever $y \in \mathcal{C}^{\lceil \alpha_k \rceil}([0,b])$ is a solution of the multi-term FDE, then the vector function $\mathbf{y} = (y, c_A D^{\gamma} y, c_A D^{2\gamma} y, \dots, c_A D^{(N-1)\gamma} y)^T$ solves the *N*-dimensional system.

We can relax (A2) from the *rationality requirement* to a requirement on being commensurable².

(A2)' Let $1 \geq \alpha_k > \alpha_{k-1} > \ldots > \alpha_1 > 0$ and assume the equation to be *commensurate*, then we define $\tilde{\alpha}_j = \alpha_j/\alpha_1$ for $j=1,\ldots,k$, let \tilde{M} be the least common multiple of the denominators of the values $\tilde{\alpha}_1,\ldots,\tilde{\alpha}_k$.

Theorem (Diethelm 2010, Theorem 8.2)

Under the assumption (A1) and (A2)', set $\gamma = \alpha_1/\tilde{M}$ and $N = \tilde{M}\alpha_k/\alpha_1$, then the equivalence relation of the *N*-dimensional system and of the multi-term FDE holds as in the previous result.

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- Existence and uniqueness results can be obtained for the single term reformulation,
- See (Ford and Connolly 2009) for other reformulations and comparisons.

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The Method of Lines

Consider a partial differential equations of the form

Find
$$u(\mathbf{x}, t)$$
 s.t. $u_t = \mathcal{L}u$, $\mathbf{x} \in \Omega \subseteq \mathbb{R}^d$, $t \in I \subseteq \mathbb{R}_+$,

where \mathcal{L} is a differential operator, either linear or nonlinear, coupled with the opportune boundary conditions, and given suitable initial conditions.

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A *classical way* of approaching this task is using a **Method Of Lines** (MOL) approach, that is

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2. now we have a (possibly nonlinear, non-autonomous) system of ODEs to which we can apply an integrator.

We can think of using the methods we have seen until now for solving PDEs in which the derivative with respect to time has been substituted by the fractional derivative in the Caputo sense

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Examples:

■ Time-fractional diffusion equation

$$_{CA}D_t^{\alpha}u = \operatorname{div}(p(x)\operatorname{grad}u) - q(x)u + F(x,t), \quad 0 < \alpha \leq 1.$$

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$$(iT_{\rho})^{\alpha}{}_{CA}D_{t}^{\alpha}\psi=-rac{L_{\rho}^{2}}{2N_{m}}\nabla^{2}\psi+N_{\nu}\psi,\quad 0<\alpha\leq1.$$

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Time-fractional Burgers equation equation

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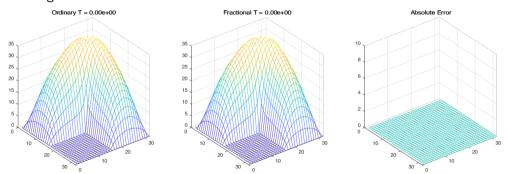
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Time-fractional (incompressible) Navier-Stokes equation

$$\begin{cases} {_{CA}D_t^{\alpha}(u\cdot\nabla)u=\nu\nabla^2u-\frac{1}{\rho}\nabla\rho+f,}\\ \nabla\cdot u=0. \end{cases} \quad 0<\alpha\leq 1.$$

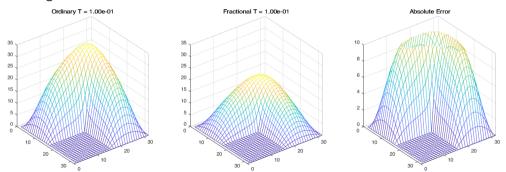
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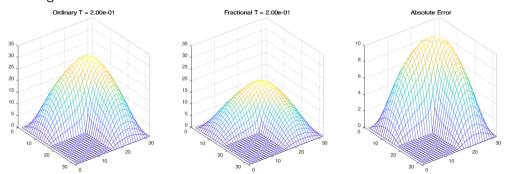
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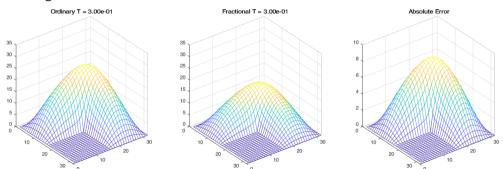
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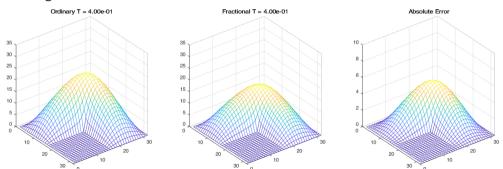
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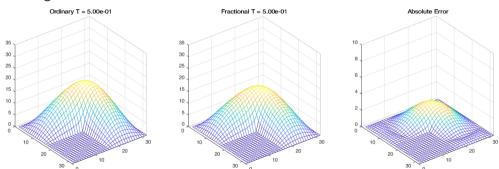
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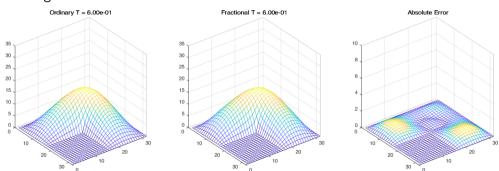
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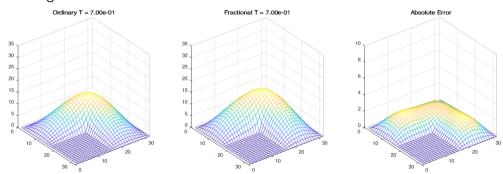
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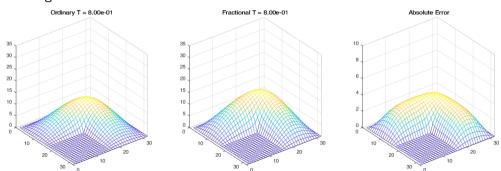
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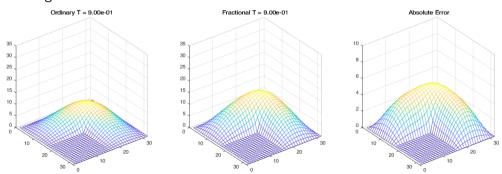
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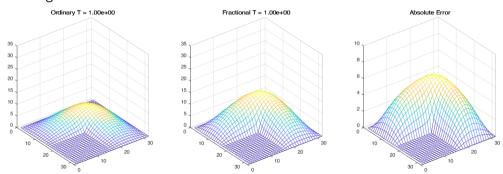
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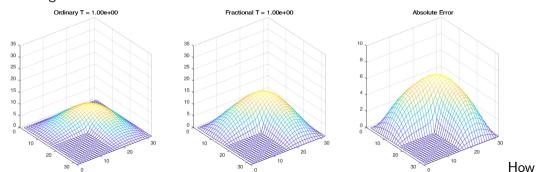
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