## An introduction to fractional calculus

**Fundamental ideas and numerics** 

#### Fabio Durastante

Università di Pisa

■ fabio.durastante@unipi.it

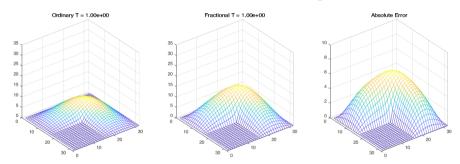




May, 2022

## **Subdiffusion equations**

At the end of the last lecture we had observed the following behavior:



for the solution of:

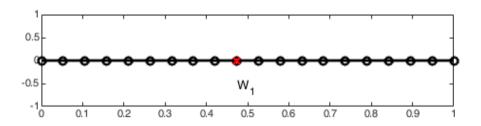
$$_{CA}D_t^{\alpha}u = 0.05\nabla^2u, \quad \alpha = 0.3, 1.$$

The visual effect seemed to be a slowing down of the diffusion.

- Consider a 1D lattice with cell size  $\Delta x$ .
- In discrete time steps of span  $\Delta t$  a test particle jumps to one of its neighbour sites,

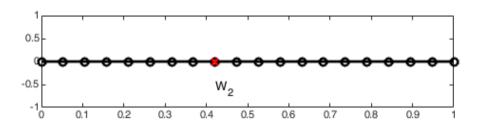
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- The process can be modelled by the master equation

$$W_j(t + \Delta t) = \frac{1}{2}W_{j-1}(t) + \frac{1}{2}W_{j+1}(t)$$



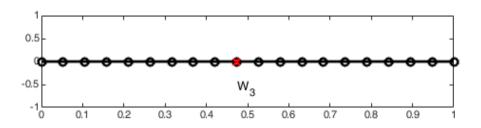
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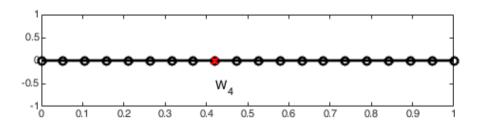
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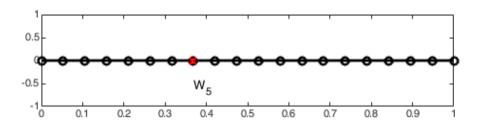
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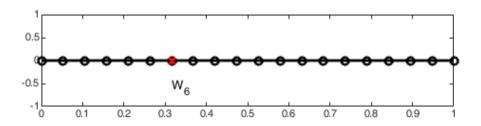
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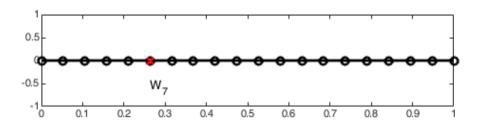
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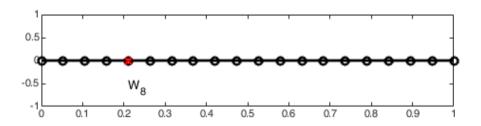
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- The prefactor 1/2 tells us that the **process is isotropic** with respect to the left/right direction.
- If we let  $\Delta t \to 0$ ,  $\Delta x \to 0$  and do a Taylor expansion in both  $\Delta$  and  $\Delta x$  we get

$$egin{aligned} W_j(t+\Delta t) = &W_j(t) + \Delta t rac{\partial W_j}{\partial t} + O([\Delta t]^2), & ext{for } \Delta t 
ightarrow 0, \ W_{j\pm 1}(t) = &W(x,t) \pm \Delta x rac{\partial W}{\partial x} + rac{(\Delta x)^2}{2} rac{\partial^2 W}{\partial x^2} + O([\Delta x]^3), & ext{for } \Delta x 
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$$W(x,t) + \Delta t \frac{\partial W}{\partial t} + O(\Delta t^2) = W(x,t) + \frac{1}{2} \Delta x^2 \frac{\partial^2 W}{\partial x^2} + O(\Delta x^3)$$

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obtaining

$$\frac{\partial W}{\partial t} = \frac{\Delta x^2}{2\Delta t} \frac{\partial^2 W}{\partial x^2} + O\left(\Delta x^3 + \Delta t\right)$$

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obtaining

$$\frac{\partial W}{\partial t} = K_1 \frac{\partial^2 W}{\partial x^2}, \qquad K_1 = \lim_{\substack{\Delta x \to 0 \\ \Delta t \to 0}} \frac{\Delta x^2}{2\Delta t} < \infty.$$

### **Brownian motion**

$$\frac{\partial W}{\partial t} = K_1 \frac{\partial^2 W}{\partial x^2}$$

Let us call X the random variable measuring the distance covered in two consecutive jumps

• Assume that the pdf of X (appropriately normalised) has existing moments

$$\overline{X} = \sum_{i} X_{i}, \quad \overline{X^{2}},$$

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• Then the **central limit theorem** assures that exists

$$V=rac{\overline{X}}{\Delta t}$$
 (Mean velocity)  $K=rac{\overline{X^2}-\overline{X}^2}{2\Delta}$  (Diffusion coefficient)

and that

$$W(x,t) = \frac{1}{2\sqrt{\pi K_1 t}} \exp\left(-\frac{x^2}{4K_1 t}\right).$$

### Brownian motion: the Fourier domain

We can rewrite

$$W(x,t) = \frac{1}{2\sqrt{\pi K_1 t}} \exp\left(-\frac{x^2}{4K_1 t}\right).$$

in the Fourier domain as

$$W(k,t) = \exp(-K_1k^2t), \qquad W_0(x) = \lim_{t \to 0^+} W(x,t) = \delta(x),$$

that solve the Fourier transformed diffusion equation

$$\frac{\partial W}{\partial t} = -K_1 k^2 W(k, t),$$

that is a **relaxation equation**, for a fixed wavenumber k.

The Continuous Time Random Walk model (CTRW):

**?** Both the **length of a given jump**, and the **waiting time** elapsing between two successive jumps are drawn from a pdf  $\psi(x,t)$ 

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- ★ The jump length pdf

$$\lambda(x) = \int_0^{+\infty} \psi(x, y) \, \mathrm{d}t,$$

#### Jump length

 $\lambda(x)dx$  produces the probability for a jump length in the interval (x, x + dx).

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The waiting time pdf

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#### Waiting time

w(t)dt produces the probability for a waiting time in the interval (t, t + dt).

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The waiting time pdf

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• If the jump length and waiting time are **independent random variables** then:

$$\psi(x,t) = w(t)\lambda(x)$$

To categorise different CTRW one can look at the quantities

$$T = \int_0^{+\infty} tw(t) dt$$
, (Characteristic waiting time),

and

$$\Sigma^2 = \int_{-\infty}^{+\infty} x^2 \lambda(x) \, \mathrm{d}x$$
 (Jump length variance),

specifically, are they finite? Do they diverge?

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The master (Langevin) equation for this process is then given by

$$\eta(x,t) = \int_{-\infty}^{+\infty} \mathrm{d}x' \int_{0}^{+\infty} \mathrm{d}t' \eta(x',t') \psi(x-x',t-t') + \delta(x)\delta(t),$$

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Pdf of having arrived at position x at time  $t - \eta(x, t)$  – having just arrived at x' at time  $t' - \eta(x', t')$  – with initial condition  $\delta(x)$ .

Then if we use

$$\eta(x,t) = \int_{-\infty}^{+\infty} \mathrm{d}x' \int_{0}^{+\infty} \mathrm{d}t' \, \eta(x',t') \psi(x-x',t-t') + \delta(x) \delta(t),$$

we can write the pdf of being in x at time t as

$$W(x,t) = \int_0^t \eta(x,t') \Psi(t-t'), \mathrm{d}t, \qquad \Psi(t) = 1 - \int_0^t w(t') \, \mathrm{d}t',$$

where the latter is the cumulative probability assigned to the probability of no jump event during the time interval t-t'.

#### Fact

If both T and  $\Sigma^2$  are finite the long-time limit corresponds to Brownian motion, e.g.,  $w(t)=\tau^{-1}exp(-t/\tau),\ T=\tau,\ \lambda(x)=(4\pi\sigma^2)^{-1/2}\exp(-x^2/4\sigma^2),\ \Sigma^2=2\sigma^2$ , we recover the standard diffusion equation.

## The CTRW in the Fourier-Laplace domain

We take

$$W(x,t) = \int_0^t \eta(x,t') \Psi(t-t'), \mathrm{d}t, \qquad \Psi(t) = 1 - \int_0^t w(t') \, \mathrm{d}t',$$

and rewrite it again in the **Fourier-Laplace domain** (Fourier for the space variable, Laplace for the time one) as

$$W(k,u) = \frac{1 - w(u)}{u} \frac{W_0(k)}{1 - \psi(k,u)}, \qquad W_0(k) = \int_{-\infty}^{+\infty} W_0(x) e^{-i2\pi kx} dx.$$

In the Brownian case

$$w(u) \sim 1 - u\tau + O(\tau^2), \quad \lambda(k) \sim 1 - \sigma^2 k^2 + O(k^4), \quad W_0(x) = \delta(x)$$

then

$$W(k, u) = \frac{1}{u + K_1 k^2}, \quad K_1 = \sigma^2 / \tau.$$

### Long rests

The characteristic waiting time  $T = \int_0^{+\infty} tw(t) dt$  diverges, but the jump length variance  $\Sigma^2 = \int_{-\infty}^{+\infty} x^2 \lambda(x) dx$  is finite.

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• and then obtain the expression for W(k, u) in the Fourier-Laplace space

$$W(k, u) = \frac{W_0(k)/u}{(1+K_\alpha u^{-\alpha} k^2)}.$$

To get an expression of the equation we use the Laplace transform for fractional integrals:

$$\mathcal{L}\left\{I_{[0,t]}^{-\alpha}W(x,t)\right\}=u^{-\alpha}W(x,u),\qquad \alpha\geq 0,$$

and together with

$$W(k,u) = \frac{W_0(k)/u}{(1 + K_\alpha u^{-\alpha} k^2)}.$$

we infer the fractional integral equation

$$W(x,t) - W_0(x) = I_{[0,t]} K_{\alpha} \frac{\partial^2}{\partial x^2} W(x,t).$$

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we infer the fractional integral equation, and apply derivative w.r.t. to time

$$\frac{\partial}{\partial t} \left( W(x,t) - W_0(x) \right) = \frac{\partial}{\partial t} \left( I_{[0,t]} K_{\alpha} \frac{\partial^2}{\partial x^2} W(x,t) \right).$$

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We can compute also the mean squared displacement

$$\langle x^2(t)\rangle = \mathcal{L}^{-1}\left\{\lim_{k\to 0} -\frac{d^2}{dk^2}W(k,u)\right\} = \frac{2K_{\alpha}}{\Gamma(1+\alpha)}t^{\alpha}.$$

We have obtained a Fractional Differential Equation:

$$rac{\partial \mathcal{W}}{\partial t} = {}_{\mathit{RL}}D^{lpha}_{[0,t]}\mathcal{K}_{lpha}rac{\partial^2}{\partial x^2}\mathcal{W}(x,t), \qquad 0 < lpha < 1$$

but this is not the model we started looking at, that was

$$_{CA}D_{[0,t]}^{lpha}rac{\partial \mathit{W}}{\partial \mathit{t}}=\mathit{K}_{lpha}rac{\partial^{2}}{\partial x^{2}}\mathit{W}(\mathit{x},\mathit{t}), \qquad 0$$

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We now have an *interpretation* of what a Fractional Derivative with respect to time is. We will come back to this when we will speak about fractional derivative with respect to space.

We start from the FDE

$$\begin{cases} {}_{CA}D^{\alpha}_{[t_0,t]}u(t)+\lambda y(t)=f(t),\\ u(0)=u_0, \end{cases} \quad \alpha\in\mathbb{R}_{>0}, \quad \lambda\in\mathbb{R}, \ u(t):[t_0,T]\to\mathbb{R}.$$

Then we rewrite the solution as

$$u(t)=e_{lpha,1}(t-t_0;\lambda)u_0+\int_{t_0}^t e_{lpha,lpha}(t-s;\lambda)f(s)\,\mathrm{d} s,\quad e_{lpha,eta}=t^{eta-1}\mathcal{E}_{lpha,eta}(-\lambda t^lpha),$$

for  $E_{\alpha,\beta}(z)$  the Mittag-Leffler (ML) function with two parameters.

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$$\begin{cases} {}_{CA}D^{\alpha}_{[t_0,t]}u(t)+\lambda y(t)=f(t),\\ u(0)=u_0, \end{cases} \quad \alpha\in\mathbb{R}_{>0}, \quad \lambda\in\mathbb{R}, \ u(t):[t_0,T]\to\mathbb{R}.$$

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- We can use this formulation to build different PI rules,
- We can use it to address the problem

$$_{CA}D^{\alpha}_{[t_0,t]}U(t)+Ay(t)=F(U(t)), \quad U(0)=U_0.$$

For both the approaches we need reliable ways for **computing** the **ML function** on both the **real line** and with **matrix argument**.

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In general, we expect to mostly need matrix function-times-vector operations:

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We postpone it to after we have discussed the actual necessities we have.

We start from the formula

$$u(t) = e_{\alpha,1}(t - t_0; \lambda)u_0 + \int_{t_0}^t e_{\alpha,\alpha}(t - s; \lambda)f(s) ds, \quad e_{\alpha,\beta} = t^{\beta-1}E_{\alpha,\beta}(-\lambda t^{\alpha}),$$

and select a grid  $\{t_i\}_{i=0}^N$ , then

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Then a PI rule for

$$u(t_n) = e_{\alpha,1}(t_n - t_0; \lambda)u_0 + \tau^{\alpha} \sum_{j=0}^{n-1} \int_0^1 e_{\alpha,\alpha}((t-t_j)/\tau - r; \tau^{\alpha}\lambda) f(t_j + r\tau) dr.$$

is obtained by selecting q+1 distinct nodes  $0 \le c_0 < c_1 < \dots < c_q \le 1$  and replacing  $f(t_j+r\tau)$  with

$$ho_j^{[q]}(t_j+r au)=\sum_{\ell=0}^q L_\ell^{[q]}(r)f(t_j+c_\ell au), \quad r\in[0,1], \quad L_\ell^{[q]} ext{ Lagrange basis element of degree } q.$$

Then the PI rule is

$$u^{(n)} = e_{\alpha,1}(t_n - t_0; \lambda)y_0 + \tau^{\alpha} \sum_{j=0}^{n-1} \sum_{\ell=0}^{q} \omega_{\ell}^{[q;\alpha]}(n-j; \tau^{\alpha}\lambda)f(t_j + c_{\ell}\tau).$$

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And selecting the weights

$$\omega_{\ell}^{[q;\alpha]}(n,z) = \int_0^1 e_{\alpha,\alpha}(n-j-r;z) L_{\ell}^{[q]}(r) dr.$$

#### Theorem (Garrappa and Popolizio 2011, Theorem 4.2)

Let  $\alpha > 0$  and  $f(t) \in \mathcal{C}^{q+2}([t_0, T])$ . The error of a q-step exponential PI rule is given by

$$u(t_n) - u^{(n)} = \tau^{q+1} \frac{C_0^{[q]}}{(q+1)!} \int_{t_0}^{t_n} e_{\alpha,\alpha}(t_n - s; \lambda) f^{(q+1)}(s) ds + O(\tau^{q+1+\alpha}),$$

where the constant  $C_0^{[q]}$  depends only on the nodes  $c_\ell$ .

- For q=2,  $c_0=0$ ,  $c_1=1/2$   $c_2=1$ , one finds  $C_0^{[2]}=0$ , thus an interpolatory formula of order  $O(\tau^{q+1+\alpha})$ .
- **?** The **general idea** is to select nodes  $c_{\ell}$  in such way that

$$C_{\mathbf{v}}^{[q]} = \int_{0}^{1} \omega_{q}(r) \xi(1-\mathbf{v}, 1-r) \, \mathrm{d}r, \quad \mathbf{v} \in \mathbb{R},$$

for  $\xi$  the *Hurwitz zeta function*, are zeroed out in the error expansion for the method.

## The MOL/Matrix case

Let us go back to the case that sparked our interest in going "exponential", that was the  $\mathsf{MOL}$  problem

$$\begin{cases} {}_{CA}D^{\alpha}_{[0,t]}\mathbf{u}(t)+A\mathbf{u}(t)=\mathbf{g}(t), & t>0, \\ \mathbf{u}(0)=\mathbf{u}_0. \end{cases}$$

By the variation of constant formula, we have seen that we can express the solution as

$$\mathbf{u}(t) = E_{\alpha,1}(-t^{\alpha}A)\mathbf{u}_0 + \int_0^t (t-s)^{\alpha-1}E_{\alpha,\alpha}(-A(t-s)^{\alpha})\mathbf{g}(s)\,\mathrm{d}s.$$

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- In the general case we then have to apply one of the PI rules to compute the integral term.
- If  $g(s) = \sum_{k=0}^{q} s^k \mathbf{v}_k$  for some vectors, we can compute the integral on the right-hand side in *closed form* and obtain

$$\mathbf{u}(t) = E_{lpha,1}(-t^lpha A)\mathbf{y}_0 + \sum_{k=0}^q \Gamma(k+1)t^{lpha+k} E_{lpha,lpha+k+1}(-t^lpha A)\mathbf{v}_k, \qquad t>0.$$

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#### Matrix functions: the normal case

If A is a normal matrix, and f is a function existing on the spectrum of A, then

$$f(A) = Uf(\Lambda)U^H$$
,  $U^HU = I$ ,  $\Lambda = \operatorname{diag}(\lambda_1, \dots, \lambda_n)$ ,  $A\mathbf{u}_i = \lambda_i \mathbf{u}_i$ ,  $U = [\mathbf{u}_1, \dots, \mathbf{u}_n]$ .

This is, e.g., sufficient for the cases in which

- A is the discretization of a self-adjoint operator,
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What about the *non-normal* and *nond-diagonalizable* case? For diagonalizable matrices, we can use the eigendecomposition at the same way.

#### Matrix functions: the Jordan Canonical Form

#### Jordan Canonical Form

We recall that any matrix  $A \in \mathbb{C}^{n \times n}$  can be expressed in Jordan canonical form as

$$Z^{-1}AZ = J = \operatorname{diag}(J_1, \ldots, J_p), \quad ext{ for } J_k = J_k(\lambda_k) = egin{bmatrix} \lambda_k & 1 & & & \\ & \lambda_k & \ddots & & \\ & & \ddots & 1 & \\ & & & \lambda_k \end{bmatrix} \in \mathbb{C}^{m_k imes m_k},$$

where Z is nonsingular and  $m_1 + m_2 + ... + m_p = n$ . If each block in which the eigenvalue  $\lambda_k$  appears is of size 1 then  $\lambda_k$  is said to be a *semisimple* eigenvalue.

• This is a *theoretical object*, it is useful to prove and define *things*, not to implement *things*.

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- This is a *theoretical object*, it is useful to prove and define *things*, not to implement *things*.
- Now that we have a decomposition of the matrix, we need to introduce a suitable definition of **being defined on the spectrum**.

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### Matrix functions: the general case

Let us denote by  $\lambda_1, \ldots, \lambda_s$  the distinct eigenvalues of A, and by  $n_i$  the order of the largest Jordan block in which the  $\lambda_i$  appears, i.e., the *index* of the eigenvalue  $\lambda_i$ .

#### Defined on the spectrum

The function f is defined on the spectrum of A if the values

$$f^{(j)}(\lambda_i), \qquad j = 0, 1, \dots, n_i - 1, \quad i = 1, \dots s,$$

exist, where  $f^{(j)}$  denotes the jth derivative of f, with  $f^{(0)} = f$ .

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 $oldsymbol{\Lambda}$  Again for the ML function and  $\alpha > 0$  we have no problem with this.

#### Matrix functions: the general case

#### Matrix function

Lef f be defined on the spectrum of  $A \in \mathbb{C}^{n \times n}$ , which is represented in Jordan canonical form as  $Z^{-1}AZ = J$ ,

$$f(A) = Zf(J)Z^{-1} = Z \operatorname{diag}(f(J_1), \dots, f(J_p))Z^{-1},$$

where

$$f(J_k) = egin{bmatrix} f(\lambda_k) & f'(\lambda_k) & \dots & rac{f^{(m_k-1)}(\lambda_k)}{(m_k-1)!} \ & f(\lambda_k) & \ddots & dots \ & & \ddots & f'(\lambda_k) \ & & & f(\lambda_k) \end{bmatrix}.$$

Moreover, let f be a multivalued function and suppose some eigenvalues occur in more than one Jordan block. If the same choice of branch of f is made in each block, then we say that f(A) is a *primary matrix function*.

# Matrix functions: computing f(A) and f(A)v

To march our scheme for

$$\mathbf{u}(t) = E_{\alpha,1}(-t^{\alpha}A)\mathbf{u}_0 + \int_0^t (t-s)^{\alpha-1}E_{\alpha,\alpha}(-A(t-s)^{\alpha})\mathbf{g}(s)\,\mathrm{d}s.$$

we need to compute operations of the form  $f(A)\mathbf{v}$ , nevertheless, we will have to compute  $f(\cdot)$  at least on some small matrix.

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#### Schur decomposition and matrix functions

Given a matrix A there exist always a matrix Q such that  $Q^*Q = I$ , and a upper triangular matrix T such that  $A = QTQ^*$ . Then, if f is defined on the spectrum of A we can compute f(A) as  $f(A) = Qf(T)Q^*$ .

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But how do we compute the matrix function of an upper triangular matrix?

## Matrix functions: the upper triangular case

Assumption we assume that T is such that each block  $T_{i,j}$  has clustered eigenvalues, and distinct diagonal blocks have far enough eigenvalues.

• If the assumption doesn't hold we look for a block permutation.

$\begin{bmatrix} (T_{1,1})_{1,1} & (T_{1,1})_{1,2} \\ 0 & (T_{1,1})_{2,2} \end{bmatrix}$	$T_{1,2}$
0	$\begin{bmatrix} (T_{2,2})_{1,1} & (T_{2,2})_{1,2} \\ 0 & (T_{2,2})_{2,2} \end{bmatrix}$

⚠ Close eigenvalues may lead to severe accuracy loss, even far apert eigenvalues can produce more inaccurate answers than expected, see (Davies and Higham 2003).

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• To evaluate  $f(T_{ii})$  we use the Taylor series in  $\sigma$ 

$$f(T_{i,i}) = \sum_{k=0}^{+\infty} \frac{f^{(k)}}{k!} M^k,$$

for 
$$\sigma = \frac{\operatorname{trace}(T_{i,i})}{m}$$
,  $m = \dim(T_{i,i})$ , and  $M = T_{i,i} - \sigma I$ .

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0	$(T_{2,2})_{1,1}$	$(T_{2,2})_{1,2}$ $(T_{2,2})_{2,2}$

For the off-diagonal blocks we apply the block-Parlett recurrence  $F_{i,i} = f(T_{i,i}), i = 1, \ldots, n;$ for  $j = 2, \ldots, n$  do for i = j - 1, j - 2, ..., 1 do Solve Sylvester equation for  $F_{i,i}$ :  $T_{i,i}F_{j,j} - F_{i,j}T_{j,j} = F_{i,i}T_{i,j} - T_{i,j}F_{j,j} + \sum_{k=0}^{j-1} (F_{i,k} - T_{k,j} - T_{i,k}F_{k,j}).$ end end

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#### What we need

To use the algorithm we have sketched out, we need to be able to compute the derivatives of the ML function sufficiently accurately.

### Derivatives of the ML function

The key observation for this task is

$$\frac{d^k}{dz^k}E_{\alpha,\beta}(z)=\sum_{i=0}^{+\infty}\frac{(j+k)_kz^j}{\Gamma(\alpha j+\alpha k+\beta)}=\frac{k!}{\Gamma(k+1)}\sum_{i=0}^{+\infty}\frac{\Gamma(j+k+1)z^j}{j!\Gamma(\alpha j+\alpha k+\beta)}=k!E_{\alpha,\alpha k+\beta}^{k+1}(z),$$

where

$$E_{\alpha,\beta}^{\gamma}(z) = \frac{1}{\Gamma(\gamma)} \sum_{i=0}^{+\infty} \frac{\Gamma(1+\gamma)z^{i}}{j!\Gamma(\alpha j + \beta)},$$

is called the Prabhakar function.

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is called the Prabhakar function.

Its **efficient computation** can be obtained, similarly to the ML function, by means of a *Laplace transform inversion* 

$$\mathcal{L}\left\{t^{\beta-1}E_{\alpha,\beta}^{\gamma}(t^{\alpha}z)\right\}(s)=\frac{s^{\alpha\gamma-\beta}}{(s^{\alpha}-t^{\alpha}z)^{\gamma}},\quad \Re(s)>0,\quad |t^{\alpha}zs^{-\alpha}|<1.$$

We select t = 1 in

$$\mathcal{L}\left\{t^{eta-1} \mathcal{E}_{lpha,eta}^{\gamma}(t^{lpha}z)
ight\}(s) = rac{s^{lpha\gamma-eta}}{(s^{lpha}-t^{lpha}z)^{\gamma}}, \quad \mathfrak{R}(s) > 0, \quad |t^{lpha}zs^{-lpha}| < 1.$$

Having selected t = 1 we have

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Since

$$\frac{d^k}{dz^k}E_{\alpha,\beta}(z)=k!E_{\alpha,\alpha k+\beta}^{k+1}(z)=\frac{k!}{2\pi i}\int_{\mathcal{C}}e^sH_k(s;z)\mathrm{d}s\equiv I_k(z),$$

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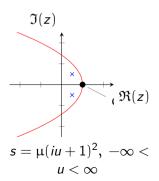
$$\mathcal{L}\left\{E_{\alpha,\beta}^{\gamma}(z)\right\}(s) = \frac{s^{\alpha\gamma-\beta}}{(s^{\alpha}-z)^{\gamma}}, \quad \Re(s) > 0, \quad |zs^{-\alpha}| < 1, \ H_k(z;z) = \frac{s^{\alpha-\beta}}{(s^{\alpha}-z)^{k+1}}.$$

Since

$$\frac{d^k}{dz^k}E_{\alpha,\beta}(z)=k!E_{\alpha,\alpha k+\beta}^{k+1}(z)=\frac{k!}{2\pi i}\int_{\mathcal{C}}e^sH_k(s;z)\mathrm{d}s\equiv I_k(z),$$

• we use the *Optimal Parabolic Contour* we have already discussed in **Lecture 2** to determine the deformation of the Bromwich line to evaluate

$$I_k^{[N]} = \frac{k!h}{2\pi i} \sum_{j=-N}^N e^{\sigma(u_j)} H_k(\sigma(u_j); z) \sigma'(u_j).$$



We needed the ML derivatives to apply Schur-Parlett to non-diagonalizable matrices.



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### Diagonalization by perturbation

Let A be nonnormal

$$\tilde{A} = A + E$$

for E a suitable perturbation is likely to be diagonalizable. Diagonalizable matrices are **dense in**  $\mathbb{C}^{n\times n}$ , for a given A and machine precision  $\epsilon$  then the best approximate diagonalization can be measured in terms of

$$\sigma(A, \varepsilon) = \inf_{E, V} \sigma(A, V, E, \varepsilon) = \inf_{E, V} \{ \kappa_2(V) \varepsilon + ||E||_2 \}.$$



We needed the ML derivatives to apply Schur-Parlett to non-diagonalizable matrices.

### Diagonalization by perturbation

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We can expect to measure on f(A) by estimating

$$||f(A+E)-f(A)|| \lesssim ||L_f(A,E)|| \leq ||L_f(A)|||E||,$$

for  $L_f(A, E)$  the Fréchet derivative of f at A in direction E,  $||L_f(A)|| = \max_{||E||=1} \{||L_f(A, E)||\}$ .



#### Fréchet derivative

The **Fréchet derivative** of a matrix function  $f: \mathbb{C}^{n \times n} \to \mathbb{C}^{n \times n}$  at a point  $X \in \mathbb{C}^{n \times n}$  is a linear mapping  $L: \mathbb{C}^{n\times n} \to \mathbb{C}^{n\times n} \to \mathbb{C}^{n\times n} \to L_f(X, E)$  such that for all  $E \in \mathbb{C}^{n\times n}$  we find

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The idea from (Higham and Liu 2021) is to use a structured perturbation: "take E to be upper triangular standard Gaussian matrix."



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#### What precision do we need?

To have  $\kappa_1(V)u_h \lesssim u$  we select for  $c_m u \approx \min_i |\operatorname{diag}(\tilde{t}_{1,1}I - \tilde{T}_{2,2})|$ 

$$u_h \lesssim \frac{c_m u^2}{\max_{i < i} |\tilde{t}_{i,i}| \left(\max_{i < j} |\tilde{t}_{i,j}|/c_m u + 1\right)^{k-2}}, \quad k = \text{"size of the Jordan block"} \geq 2.$$

### From small to large matrices

We now know how to compute  $E_{\alpha,\beta}(A)$  for a *small matrix* A, either with

- Classical Schur-Parlett algorithm with Laplace inversion technique for the needed derivative of the ML function (Garrappa and Popolizio 2018),
  - https://it.mathworks.com/matlabcentral/fileexchange/66272-mittag-leffler-function-with-matrix-arguments
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What about large matrices?

### Projection methods for matrix functions

We can exploit the subspace projection idea, take  $V \in \mathbb{R}^{n imes k}$  spanning a given subspace  $\mathcal{W}_k$ 

$$f(A)\mathbf{v} \approx V f(V^T A V) V^T \mathbf{v}$$
  $V^T A V \in \mathbb{R}^{k \times k}$ ,  $k \ll n$ .

## **Krylov Projection Methods**

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### A general framework

Given a set of scalars  $\{\sigma_1,\ldots,\sigma_{k-1}\}\subset \overline{\mathbb{C}}$  (the extended complex plane), that are not eigenvalues of A, let

$$q_{k-1}(z) = \prod_{j=1}^{k-1} (\sigma_j - z).$$

The rational Krylov subspace of order k associated with A,  $\mathbf{v}$  and  $q_{k-1}$  is defined by

$$Q_k(A, \mathbf{v}) = [q_{k-1}(A)]^{-1} \mathcal{K}_k(A, \mathbf{v}), \qquad \mathcal{K}_k(A, \mathbf{v}) = \operatorname{Span}\{\mathbf{v}, A\mathbf{v}, \dots, A^{k-1}\mathbf{v}\}.$$

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$$C_i = (\mu_i \sigma_i A - I) (\sigma_i I - A)^{-1}$$
, and  $Q_k(A, \mathbf{v}) = \operatorname{Span}\{\mathbf{v}, C_1 \mathbf{v}, \dots, C_{k-1} \dots C_2 C_1 \mathbf{v}\}$ .

## Krylov Projection Methods: special cases

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$$(\mu_j, \sigma_j) = \left\{ egin{array}{ll} (1, \infty), & ext{for } j ext{ even,} \\ (0, 0), & ext{for } j ext{ odd.} \end{array} 
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Shift-And-Invert  $W_k(A, \mathbf{v}) = \operatorname{Span}\{\mathbf{v}, (\sigma I - A)^{-1}\mathbf{v}, \dots, (\sigma I - A)^{-(k-1)}\mathbf{v}\}$ , take  $\mu_j = 0$  and  $\sigma_j = \sigma$  for each j,

To estimate the convergence behavior of general projection methods in the non-normal we need the concept of **field of values** (or *numerical range*.)

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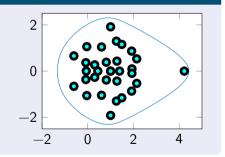
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### Field of Values/Numerical Range

Given  $A \in \mathbb{C}^{N \times N}$  we denote its **field of values** as

$$W(A) = \left\{ rac{\langle \mathbf{x}, A\mathbf{x} 
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where  $\langle \cdot, \cdot \rangle$  represents the Euclidean inner product.



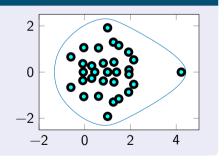
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It has many **properties**, e.g.,  $W(A) \subseteq D(0, ||A||)$  (disk centered on 0 with radius ||A||), is compact, sub-additive  $W(A+B) \subseteq W(A)+W(B)$ , unitarily invariant  $W(UAU^H)=UW(A)U^H$ , etc. see (Benzi 2021).

#### **Assumptions:**

(A1) We assume that  $\exists a > 0, \ \theta \in [0, \pi/2)$  such that

$$W(A) \subset \Sigma_{\theta,a} = \{\lambda \in \mathbb{C} : |\arg(\lambda) - a| \le \theta\}.$$

(A2)  $\beta > 0$ ,  $\alpha \in (0,2)$  be such that  $\alpha \pi/2 < \pi - \theta$ ,  $\epsilon > 0$  and

$$\frac{\alpha\pi}{2} < \mu \le \min\{\pi, \alpha\pi\}, \quad \mu < \pi - \theta.$$

**Method of choice:** we use polynomial Krylov method  $\mathcal{K}_m(A, \mathbf{v})$ :

$$AV_m = V_m H_m + h_{m+1,m} v_{m+1} \mathbf{e}_m^T, \quad \operatorname{Span} V_m = \operatorname{Span}\{\mathbf{v}_i\}_{i=1}^m = \mathcal{K}_m(A, \mathbf{v}), \quad H_m = V_m^H A V_m.$$

#### We want to bound:

$$R_m = E_{\alpha,\beta}(-A)\mathbf{v} - V_m E_{\alpha,\beta}(-H_m)\mathbf{e}_1, \quad m \ge 1.$$

We first express the error in integral form, starting from (Podlubny 1999, Theorem 1.1)

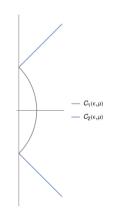
$$E_{\alpha,\beta}(z) = \frac{1}{2\alpha\pi i} \int_{C(\varepsilon,\mu)} \frac{\exp(\lambda^{1/\alpha})\lambda^{1-\beta/\alpha}}{\lambda - z} \, \mathrm{d}\lambda, \quad z \in G^{-}(\varepsilon,\mu),$$

where

•  $\forall \, \varepsilon > 0, \, 0 < \mu < \pi$ 

$$C(\varepsilon,\mu) = \bigcup \ \begin{cases} C_1(\varepsilon,\mu) = \{\lambda \ : \ \lambda = \varepsilon \exp(i\varphi), & -\mu \le \varphi \le \mu\}, \\ C_2(\varepsilon,\mu) = \{\lambda \ : \ \lambda = r \exp(\pm i\mu), & r \ge \varepsilon\}. \end{cases}$$

• The contour  $C(\varepsilon, \mu)$  divides the complex plane into two domains,  $G^-(\varepsilon, \mu)$  and  $G^+(\varepsilon, \mu)$  lying respectively on the left and on the right of  $C(\varepsilon, \mu)$ .



## An Expression for the Error

From the previous we find

$$E_{\alpha,\beta}(-A) = \frac{1}{2\alpha\pi i} \int_{C(\varepsilon,\mu)} \exp(\lambda^{1/\alpha}) \lambda^{1-\beta/\alpha} (\lambda I + A)^{-1} \, \mathrm{d}\lambda, \quad \sigma(-A) \in G^{-}(\varepsilon,\mu),$$

and together with

$$R_m = E_{\alpha,\beta}(-A)\mathbf{v} - V_m E_{\alpha,\beta}(-H_m)\mathbf{e}_1, \quad m \ge 1,$$

we write

$$R_m = rac{1}{2 lpha \pi i} \int_{C(\varepsilon, \Pi)} \exp(\lambda^{1/lpha}) \lambda^{1-eta/lpha} \delta_m(\lambda), \mathrm{d}\lambda,$$

for

$$\delta_m(\lambda) = (\lambda I + A)^{-1} \mathbf{v} - V_m(\lambda I + H_m)^{-1} \mathbf{e}_1$$
$$= (\lambda I + A)^{-1} \mathbf{v} - V_m(\lambda I + H_m)^{-1} V_m^H \mathbf{v}.$$

Observe now that

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For an arbitrary  $\mathbf{y} \in \mathbb{C}^m$  we have then

$$(\lambda I + A)^{-1}\mathbf{v} - V_m(\lambda I + H_m)^{-1}V_m^H\mathbf{v} = \Delta_m(\mathbf{v} - (\lambda I + A)V_m\mathbf{y}) = \Delta_m p_m(A)\mathbf{v},$$

where  $p_m(z)$  is a polynomial of degree  $\leq m$  with  $p_m(-\lambda) = 1$ .

We have therefore proved that

$$\|\delta_m(A)\| \leq \|(\lambda I + A)^{-1} - V_m(\lambda I + H_m)^{-1}V_m^H\|\|p_m(A)\mathbf{v}\|, \forall p_m \in \mathbb{P}_{\leq m}[z] \text{ with } p_m(-\lambda) = 1.$$

By using (Diele, Moret, and Ragni 2008/09, Lemma 2) we also have the following expression

$$\|\delta_m(\lambda)\| = \frac{\prod_{j=1}^m h_{j+1,j}}{|\det(\lambda I + H_m)|} \|(\lambda I + A)^{-1} \mathbf{v}_{m+1}\|.$$

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To obtain the first bound we call then

$$D(\lambda) = \operatorname{dist}(\lambda, W(-A)) \quad \forall \lambda \in C(\varepsilon, \mu).$$

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#### Representation function

Using (A1) and (A2) we can find a function  $v(\phi)$  such that

$$\forall \lambda = |\lambda| \exp(\pm i\varphi) \in C(\varepsilon, \mu) \quad D(\lambda) > \nu(\varphi)|\lambda|, \quad \nu(\varphi) > \nu > 0.$$

#### Theorem (Moret and Novati 2011, Theorem 3.2)

Let assumptions (A1) and (A2) hold, then for  $m \ge 1$  and for every M > 0 we have

$$||R_m|| \leq \frac{\exp(M) \prod_{j=1}^m h_{j+1,j}}{\pi v^{m+1} M^{m\alpha+\beta-1}} \left( \frac{\mu}{\alpha} + \frac{\exp(-M(|\cos(\mu/\alpha)|+1))}{m\alpha-1+\beta} \right).$$

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$$||R_m|| = \left| \left| \frac{1}{2\alpha\pi i} \int_{C(\varepsilon,\mu)} \exp(\lambda^{1/\alpha}) \lambda^{1-\beta/\alpha} \delta_m(\lambda), d\lambda \right| \right|$$

$$\leq \frac{\prod_{j=1}^m h_{j+1,j}}{2\pi\alpha} \int_{C(\varepsilon,\mu)} \frac{\left| \exp(\lambda^{1/\alpha}) \lambda^{1-\beta/\alpha} \right|}{D(\lambda)^{m+1}} |d\lambda|.$$

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$$||R_m|| \leq \frac{\prod_{j=1}^m h_{j+1,j}}{2\pi\alpha} (I_1 + I_2),$$

with

$$\mathit{I}_1 = \int_{\mathit{C}_1(\epsilon,\mu)} \frac{\left|\exp(\lambda^{1/\alpha})\lambda^{1-\beta/\alpha}\right|}{\mathit{D}(\lambda)^{m+1}} \left|\mathrm{d}\lambda\right| \leq 2\epsilon^{\frac{1-\beta}{\alpha}-m} \int_0^\mu \frac{\exp(\epsilon^{1/\alpha}\cos(\phi/\alpha))}{\nu(\phi)^{m+1}} \,\mathrm{d}\phi,$$

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$$||R_m|| \leq \frac{\exp(M) \prod_{j=1}^m h_{j+1,j}}{\pi \gamma^{m+1} M^{m\alpha+\beta-1}} \left( \frac{\mu}{\alpha} + \frac{\exp(-M(|\cos(\mu/\alpha)|+1))}{m\alpha-1+\beta} \right).$$

**Proof.** We use  $\|(\lambda I + A)^{-1}\| \le D(\lambda)^{-1}$  and  $W(H_m) \subseteq W(A)$  in the error expression  $R_m$ 

$$||R_m|| \leq \frac{\prod_{j=1}^m h_{j+1,j}}{2\pi\alpha} \left( 2\varepsilon^{\frac{1-\beta}{\alpha}-m} \int_0^\mu \frac{\exp(\varepsilon^{1/\alpha}\cos(\varphi/\alpha))}{\nu(\varphi)^{m+1}} d\varphi + I_2 \right),$$

with

$$\begin{split} I_2 = & \int_{C_2(\varepsilon,\mu)} \frac{\left| \exp(\lambda^{1/\alpha}) \lambda^{1-\beta/\alpha} \right|}{D(\lambda)^{m+1}} \left| \mathrm{d} \lambda \right| \leq \frac{2}{\nu^{m+1}} \int_{\varepsilon}^{+\infty} \frac{r^{\frac{1-\beta}{\alpha}} \exp(-r^{\frac{1}{\alpha}} |\cos(\mu/\alpha)|)}{r^{m+1}} \, \mathrm{d} r \\ = & \frac{2}{\nu^{m+1}} \int_{\varepsilon^{1/\alpha}}^{+\infty} \frac{\exp(-s |\cos(\mu/\alpha)|)}{s^{m\alpha+\beta}} \, \mathrm{d} s \leq \frac{2\alpha \exp(-\varepsilon^{1/\alpha} |\cos(\mu/\alpha)|)}{(m\alpha+\beta-1)\nu^{m+1} \varepsilon^{\frac{m\alpha+\beta-1}{\alpha}}}. \end{split}$$

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The result follows then by setting  $\varepsilon=M^{lpha}$  and simplifying the expression.

 $oldsymbol{\Lambda}$  With the same proof another bound for the case of small lpha can be obtained.

### A First Error Bound: small $\alpha$ s

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#### Corollary (Moret and Novati 2011, Corollary 3.3)

Let assumptions (A1) and (A2) hold. Let  $m \ge 1$  be such that  $m\alpha + \beta > 0$ , then for every M > 0, we have

$$||R_m|| \leq \frac{\exp(M) \prod_{j=1}^m h_{j+1,j}}{4\nu^{m+1} M^{m\alpha}} \frac{4M^{1-\beta}}{\pi} \left( \frac{\mu}{\alpha} + \frac{\exp(-M(1+|\cos(\mu/\alpha)|))}{M|\cos(\mu/\alpha)|} \right).$$

### A First Error Bound: some observations

**‡** The ML function is entire for  $\alpha > 0 \Rightarrow$  superlinear convergence for large enough m:

$$M = m lpha + eta - 1 \ \Rightarrow \ \|R_m\| \propto \left(rac{\exp(1)}{M}
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lacktriangle To better understand this, we use that for every monic polynomial of degree m we find

$$\prod_{j=1}^{m} h_{j+1,j} \leq ||q_{m}(A)v||,$$

Therefore, if we take  $q_m$  as the monic Faber polynomial associated to a closed convex subset  $\Omega \supset W(-A)$  we get the bound in terms of the logarithmic capacity  $\gamma$  of  $\Omega$ .

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Therefore, if we take  $q_m$  as the **monic Faber polynomial** associated to a closed convex subset  $\Omega \supset W(-A)$  we get the bound in terms of the **logarithmic capacity**  $\gamma$  of  $\Omega$ .

 $\Rightarrow$  we have discovered:

$$||R_m|| \propto \left(\frac{\exp(1)}{m\alpha}\right)^{m\alpha} \left(\frac{\gamma}{\gamma}\right)^m.$$

# **Specialized bounds**

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#### Theorem (Moret and Novati 2011, Theorem 3.5)

Assume that A is Hermitian with  $\sigma(A)\subseteq [a,b]\subset [0,+\infty)$ . Assume that  $0<\alpha<1,\ \beta\geq\alpha$ . Let  $\mu\leq\pi/2,\ \alpha\pi/2<\mu<\alpha\pi$ . Then for every index  $m\geq1$  and for every M>0 we have

$$||R_m|| \leq \frac{4M^{1-\beta}}{\pi} \left( \frac{\mu}{\alpha} + \frac{\exp(-M(1+|\cos(\mu/\alpha)|))}{M|\cos(\mu/\alpha)|} \right) \exp(M) \Phi(u(M^{\alpha}\exp(i\mu)))^{-m}.$$

for 
$$\Phi(u) = u + \sqrt{u^2 - 1}$$
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#### Limiting relation

If  $\alpha \to 0$ ,  $\beta = 1$ , we have  $E_{0,1}(-z) = (1+z)^{-1}$ , |z| < 1. Then setting  $\mu = \alpha \pi$  and letting M = 1, we find

$$||R_m|| \le \frac{4(\pi \exp(1) - \exp(-1))}{\pi \Phi(u(1))^m}$$

We remain under the assumptions (A1) and (A2) and consider the matrix

$$Z = (I + hA)^{-1}, \qquad h > 0,$$

together with the space  $\mathcal{K}_m(Z, \mathbf{v})$ .

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We can write the analogous Arnoldi relation for  $U_m = [\mathbf{u}_1, \dots, \mathbf{u}_m]$  spanning  $\mathcal{K}_m(Z, \mathbf{v})$ :

$$ZU_m = U_m S_m + s_{m+1,m} u_{m+1} \mathbf{e}_m^T, \qquad S_m = U_m^H Z U_m.$$

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The **approximation** is then given by

$$\mathbf{y} = f(A)\mathbf{v} \approx \mathbf{y}_m = V_m f(B_m)\mathbf{e}_1$$
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We can repeat the general error analysis using

$$R_m = E_{\alpha,\beta}(-A)\mathbf{v} - U_m E_{\alpha,\beta}(-B_m)\mathbf{e}_1 = \frac{1}{2\pi\alpha i} \int_{C(\varepsilon,\mu)} \exp(\lambda^{1/\alpha}) \lambda^{(1-\beta)/\alpha} b_m(\lambda) d\lambda,$$

for 
$$b_m(\lambda) = (\lambda I + A)^{-1}\mathbf{v} - U_m(\lambda I + B_m)^{-1}\mathbf{e}_1$$
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# **Error bound (Moret and Novati 2011)**

#### Theorem (Moret and Novati 2011, Theorem 4.3)

For every matrix A satisfying (A1) and (A2), assume  $0 < \alpha < 1$  and  $\beta \ge \alpha$ . Then, there exists a function g(h), continuous in any bounded interval  $0 < h_1 \le h \le h_2$ , such that for m > 2.

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$$||R_m|| \leq \frac{g(h)}{m-1}.$$

#### Theorem (Moret and Novati 2011, Theorem 4.5)

Assume that A is Hermitian with  $\sigma(A)\subseteq [a,+\infty)$ ,  $a\geq 0$ . Assume  $0<\alpha\leq 2/3$  and  $\beta\geq \alpha$ . Then, for every  $m\geq 1$  we have

$$||R_m|| \le \frac{K_1 Q_m h^{\frac{\beta-1}{\alpha}}}{(1+\sqrt{2})^{m-1}} + \frac{K_2 h^{\beta/\alpha}}{(m-1)^2} \exp\left(-\frac{h^{-1/\alpha}}{\sqrt{2}}\right),$$

where  $Q_m = \max_{0 \le |\phi| \le 3\alpha\pi/4} \exp\left(h^{-1/\alpha}\cos\phi/\alpha\right) (1-\cos\phi)^{\frac{m-1}{2}}$ , with  $K_1$ ,  $K_2$  constants.

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$$\mathbf{u}(t) = E_{\alpha,1}(-t^{\alpha}A)\mathbf{u}_0 + \int_0^t (t-s)^{\alpha-1}E_{\alpha,\alpha}(-A(t-s)^{\alpha})\mathbf{g}(s)\,\mathrm{d}s.$$

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**Research ideas:** finding better rational approximations/poles/expansions together with error analysis for the ML function.

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#### Other extensions

A variant with *restart* is discussed in (Moret and Popolizio 2014), the combination with other matrix-functions in (Moret and Novati 2019).

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