

Onde piane in un mezzo con perdite I

$$\nabla \times \hat{\mathbf{H}} = \hat{\mathbf{j}} + j\omega\varepsilon \hat{\mathbf{E}}$$

$$\nabla \times \hat{\mathbf{E}} = -j\omega\mu \hat{\mathbf{H}}$$

Legge di Ohm:

$$\hat{\mathbf{j}} = \sigma \hat{\mathbf{E}}$$

Onde piane in un mezzo con perdite II

$$\nabla \times \hat{\mathbf{E}} = -j\omega\mu \hat{\mathbf{H}}$$

$$\nabla \times \hat{\mathbf{H}} = \sigma \hat{\mathbf{E}} + j\omega\varepsilon \hat{\mathbf{E}}$$

$$\nabla \times \hat{\mathbf{E}} = -j\omega\mu \hat{\mathbf{H}}$$

$$\nabla \times \hat{\mathbf{H}} = [\sigma + j\omega\varepsilon] \hat{\mathbf{E}}$$

Onde piane in un mezzo con perdite III

$$\nabla \times \nabla \times \hat{\mathbf{E}} = -j\omega\mu \nabla \times \hat{\mathbf{H}} = -j\omega\mu [\sigma + j\omega\varepsilon] \hat{\mathbf{E}}$$

$$\nabla \times \nabla \times \hat{\mathbf{E}} = \nabla \nabla \cdot \hat{\mathbf{E}} - \nabla^2 \hat{\mathbf{E}}$$

$$\nabla \cdot \hat{\mathbf{E}} = 0$$

$$\nabla^2 \hat{\mathbf{E}} - \gamma^2 \hat{\mathbf{E}} = 0 \qquad \gamma^2 = j\omega\mu [\sigma + j\omega\varepsilon]$$

Onde piane in un mezzo con perdite IV

$$\gamma = j\omega \sqrt{\epsilon\mu} \left[1 + \frac{\sigma}{j\omega\epsilon} \right]^{1/2}$$

$$\gamma = j\omega \sqrt{(\epsilon' - j\epsilon'')\mu} = j\omega \sqrt{\mu\epsilon' (1 - j \tan \delta)}$$

γ è la costante di fase

*Se assumiamo che l'onda sia uniforme
nel piano trasversale alla propagazione:*

$$\partial_z^2 \hat{E}_x - \gamma^2 \hat{E}_x = 0$$

$$\hat{E}_x = E_x^+ \exp(-\gamma z) + E_x^- \exp(+\gamma z)$$

Onda piana progressiva con perdite : E

$$\gamma = \alpha + j\beta \quad \hat{E}_x = E_x^+ \exp(-(\alpha + j\beta) z) =$$

$$\hat{E}_x = E_x^+ \exp(-\alpha z) \exp(-j\beta z)$$

$$E_x(z, t) = \text{Re} \left\{ \hat{E}_x \exp(j\omega t) \right\} =$$

$$= E_x^+ \exp(-\alpha z) \cos(\omega t - \beta z)$$

Onda piana progressiva con perdite : H

$$\hat{H}_y = \frac{-j}{\omega\mu} \partial_z \hat{E}_x = \frac{-j\gamma}{\omega\mu} E_x^+ \exp(-\gamma z)$$

$$H_y(z, t) = -\operatorname{Re} \left\{ \frac{j(\alpha + j\beta)}{\omega\mu} E_x^+ \exp(-(\alpha + j\beta)z) \exp(j\omega t) \right\} =$$

$$H_y(z, t) = -\operatorname{Re} \left\{ \frac{j\alpha - \beta}{\omega\mu} E_x^+ \exp(-(\alpha + j\beta)z) \exp(j\omega t) \right\} =$$

$$H_y(z, t) = -E_x^+ \frac{\exp(-\alpha z)}{\omega\mu} \operatorname{Re} \left\{ (j\alpha - \beta) \exp(j(\omega t - \beta z)) \right\} =$$

$$= +E_x^+ \frac{\exp(-\alpha z)}{\omega\mu} [\beta \cos(\omega t - \beta z) + \alpha \sin(\omega t - \beta z)]$$

Propagazione buon conduttore

σ molto maggiore $\omega\varepsilon$ $\sigma = 5.81 \cdot 10^7 \text{ S/m (Cu)}$

a 1 GHz $\omega\varepsilon = 2\pi \cdot 10^9 \cdot 8.854 \cdot 10^{-12} = 0.055 \text{ S/m}$

$$\gamma = j\omega \sqrt{\varepsilon\mu} \left[1 + \frac{\sigma}{j\omega\varepsilon} \right]^{1/2} \approx \sqrt{j\omega\mu\sigma}$$

$$= \frac{1+j}{\sqrt{2}} \sqrt{\omega\mu\sigma} = (1+j) \sqrt{\pi f\mu\sigma}$$

Profondità di penetrazione

$$\delta = \frac{1}{\sqrt{\pi f \mu \sigma}}$$

$$\hat{E}_x = E_x^+ \exp(-z / \delta) \exp(-jz / \delta)$$

Profondità di penetrazione di Al, Cu, Au, Ag

$$f = 10 \text{ GHz} \qquad \delta = \frac{1}{\sqrt{\pi f \mu \sigma}}$$

$$\delta_{Al} = 8.14 \cdot 10^{-7} \text{ m} \qquad \delta_{Au} = 7.86 \cdot 10^{-7} \text{ m}$$

$$\delta_{Cu} = 6.6 \cdot 10^{-7} \text{ m} \qquad \delta_{Ag} = 6.4 \cdot 10^{-7} \text{ m}$$

A 100 Hz...

$$\delta_{Al} = 8.14 \text{ mm}$$

$$\delta_{Au} = 7.86 \text{ mm}$$

$$\delta_{Cu} = 6.6 \text{ mm}$$

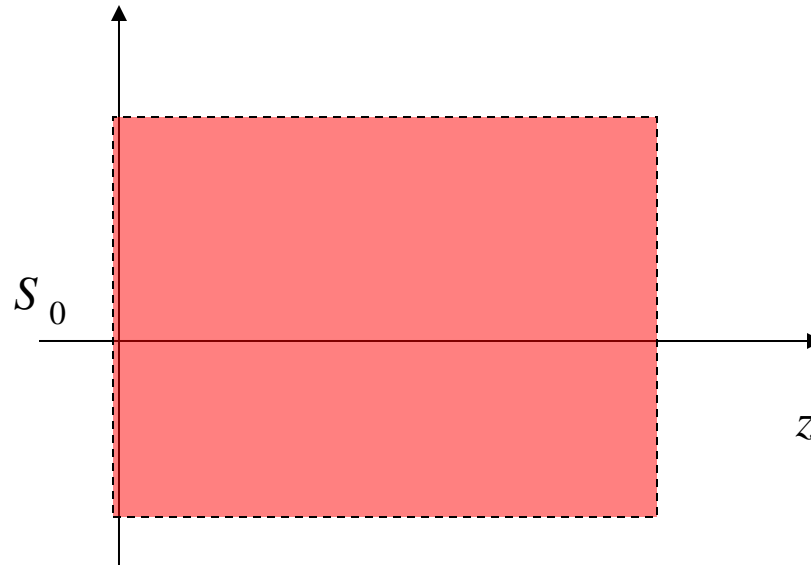
$$\delta_{Ag} = 6.4 \text{ mm}$$

Impedenza d'onda in un buon conduttore

$$\frac{j\omega\mu_0}{jk} = \frac{j\omega\mu_0}{\gamma} \Rightarrow \eta \approx (1 + j) \sqrt{\frac{\omega\mu_0}{2\sigma}} = (1 + j) \frac{1}{\sigma\delta}$$

L'angolo di fase dell'impedenza d'onda vale 0° per un Dielettrico senza perdite, mentre è un valore compreso tra -45° e $+45^\circ$ per un mezzo qualsiasi.

Potenza assorbita da un buon conduttore



$$P_{av} = \frac{1}{2} \operatorname{Re} \int_{\partial V} \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* \cdot \mathbf{n} dS = \frac{1}{2} \operatorname{Re} \int_{S_0} \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* \cdot \mathbf{u}_z dS =$$

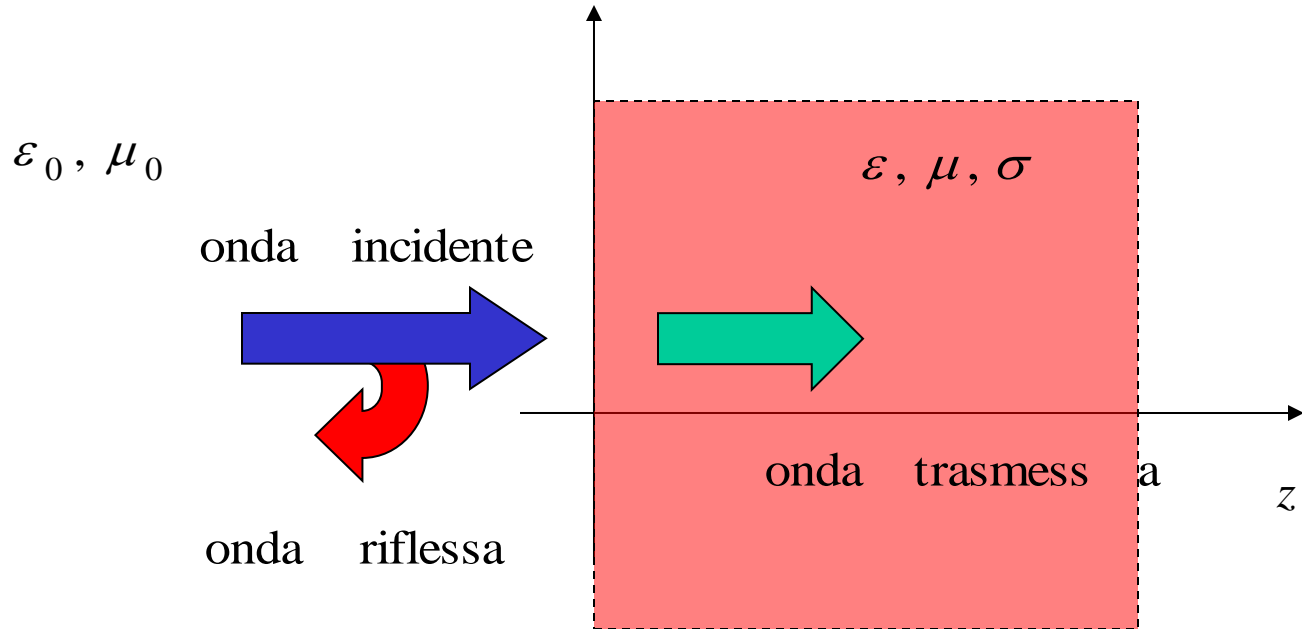
$$P_{av} = \frac{1}{2} \operatorname{Re} \left\{ \eta \int_{S_0} \hat{\mathbf{H}} \cdot \hat{\mathbf{H}}^* dS \right\} = \frac{1}{2} \operatorname{Re} \{ \eta \} \int_{S_0} |\hat{\mathbf{H}}|^2 dS = \frac{1}{2} R_s \int_{S_0} |\hat{\mathbf{H}}|^2 dS =$$

$$= \frac{1}{2} R_s \int_{S_0} |\hat{\mathbf{J}}_s|^2 dS$$

Resistenza superficiale [Ω]

$$R_s = \operatorname{Re} \left\{ (1 + j) \sqrt{\frac{\omega \mu}{2 \sigma}} \right\} = \sqrt{\frac{\omega \mu}{2 \sigma}} = \frac{1}{\sigma \delta}$$

Riflessione di un'onda piana



$$\hat{\mathbf{E}}_i = \mathbf{u}_x E_0 \exp(-jk_0 z)$$

$$\hat{\mathbf{H}}_i = \mathbf{u}_y \frac{E_0}{\eta_0} \exp(-jk_0 z)$$

$$\hat{\mathbf{E}}_r = \mathbf{u}_x \Gamma E_0 \exp(jk_0 z)$$

$$\hat{\mathbf{H}}_r = -\mathbf{u}_y \Gamma \frac{E_0}{\eta_0} \exp(jk_0 z)$$

Continuità dei campi tangenziali

Onda trasmessa:

$$\hat{\mathbf{E}}_t = \mathbf{u}_x TE_0 \exp(-\gamma z)$$

$$\eta = \frac{j\omega\mu}{\gamma}$$

$$\hat{\mathbf{H}}_t = \mathbf{u}_y T \frac{E_0}{\eta} \exp(-\gamma z)$$

$$\gamma = j\omega \sqrt{\mu\epsilon} \sqrt{1 - j\sigma/(\omega\epsilon)}$$

Continuità delle componenti tangenziali:

$$\hat{\mathbf{E}}_i(z=0) + \hat{\mathbf{E}}_r(z=0) = \hat{\mathbf{E}}_t(z=0)$$

$$\Leftrightarrow 1 + \Gamma = T$$

$$\hat{\mathbf{H}}_i(z=0) + \hat{\mathbf{H}}_r(z=0) = \hat{\mathbf{H}}_t(z=0)$$

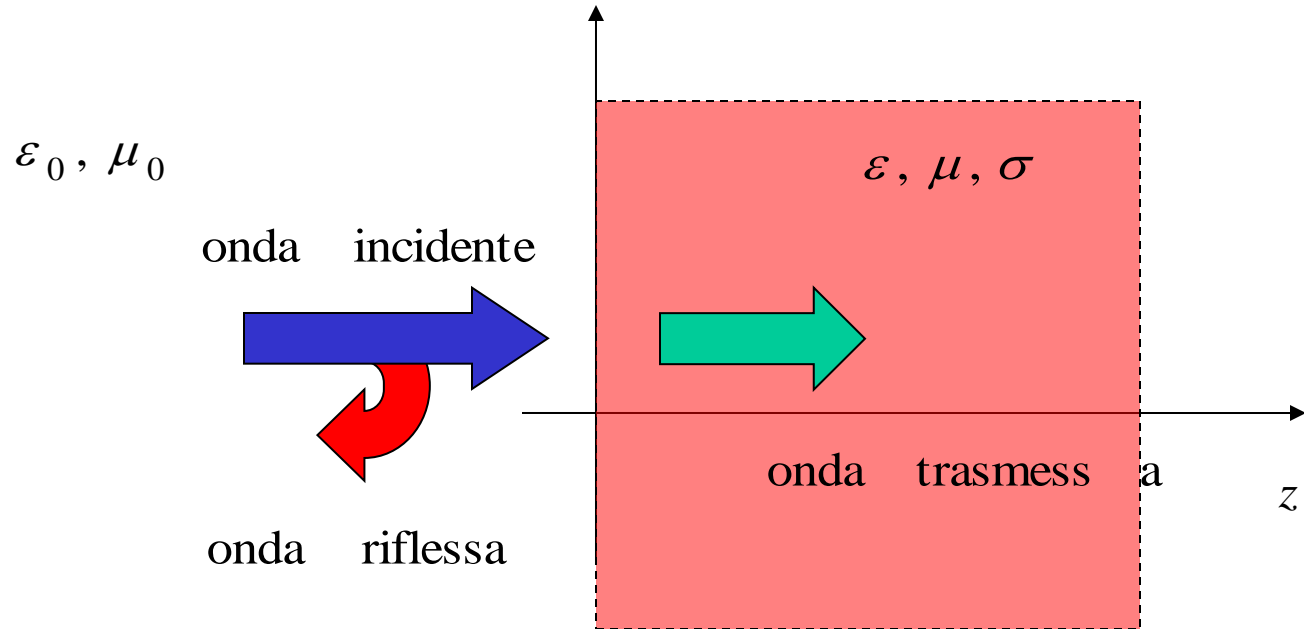
$$\Leftrightarrow \frac{1 - \Gamma}{\eta_0} = \frac{T}{\eta}$$

Coefficienti di riflessione e trasmissione

$$\Gamma = \frac{\eta - \eta_0}{\eta + \eta_0}$$

$$T = 1 + \Gamma = \frac{2\eta}{\eta + \eta_0}$$

Riflessione di un'onda piana



$$\hat{\mathbf{E}} = \hat{\mathbf{E}}_i + \hat{\mathbf{E}}_r = \mathbf{u}_x E_0 [\exp(-jk_0 z) + \Gamma \exp(jk_0 z)]$$

$$\hat{\mathbf{H}} = \hat{\mathbf{H}}_i + \hat{\mathbf{H}}_r = \mathbf{u}_y \frac{E_0}{\eta_0} [\exp(-jk_0 z) - \Gamma \exp(+jk_0 z)]$$

Vettore di Poynting ($z < 0$)

$$\hat{\mathbf{E}} \times \hat{\mathbf{H}}^* = \mathbf{u}_x \times \mathbf{u}_y E_0 [\exp(-jk_0 z) + \Gamma \exp(jk_0 z)] \cdot \frac{E_0^*}{\eta_0} [\exp(jk_0 z) - \Gamma^* \exp(-jk_0 z)]$$

$$= \mathbf{u}_z \frac{|E_0|^2}{\eta_0} \left\{ 1 - |\Gamma|^2 + 2j \operatorname{Im} [\Gamma \exp(+2jk_0 z)] \right\}$$

Vettore di Poynting ($z > 0$)

$$\hat{\mathbf{E}} \times \hat{\mathbf{H}}^* = \mathbf{u}_x \times \mathbf{u}_y E_0 T \exp(-\gamma z) \cdot \frac{E_0^* T^*}{\eta^*} \exp(-\gamma^* z) =$$

$$\Gamma = \frac{\eta - \eta_0}{\eta + \eta_0}$$

$$T = 1 + \Gamma = \frac{2\eta}{\eta + \eta_0}$$

$$= \mathbf{u}_z \frac{|E_0|^2}{\eta^*} |T|^2 \exp(-2\alpha z)$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}}; \quad \eta = \frac{j\omega\mu_0}{\gamma} \quad \gamma = j\omega\sqrt{\mu\epsilon}\sqrt{1 - j\sigma/(\omega\epsilon)}$$

Assenza di perdite $\sigma=0$, $\alpha=0$

$$\gamma = j\omega \sqrt{\mu\epsilon}$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}}; \quad \eta = \frac{j\omega\mu}{j\omega\sqrt{\mu\epsilon}} = \sqrt{\frac{\mu}{\epsilon}}$$

$$\Gamma = \frac{\eta - \eta_0}{\eta + \eta_0}$$

$$T = 1 + \Gamma = \frac{2\eta}{\eta + \eta_0}$$

Entrambi reali

$$(z < 0) \quad \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* = \mathbf{u}_z \frac{|E_0|^2}{\eta_0} \left\{ 1 - |\Gamma|^2 + 2j\Gamma \sin(2k_0 z) \right\}$$

$$(z > 0) \quad \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* = \mathbf{u}_z \frac{|E_0|^2}{\eta} |T|^2$$

$$(z > 0) \quad \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* = \mathbf{u}_z \frac{|E_0|^2}{\eta} \left| \frac{2\eta}{\eta + \eta_0} \right|^2 = \mathbf{u}_z \frac{|E_0|^2}{(\eta + \eta_0)^2} 4\eta$$

$$(z < 0) \quad \text{Re} \left\{ \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* \right\} = \mathbf{u}_z \frac{|E_0|^2}{\eta_0} \left(1 - \frac{(\eta - \eta_0)^2}{(\eta + \eta_0)^2} \right) =$$

$$= \mathbf{u}_z \frac{|E_0|^2}{\eta_0} \frac{4\eta\eta_0}{(\eta + \eta_0)^2}$$

Buon conduttore

$$\eta = (1 + j) \sqrt{\frac{\omega \mu_0}{2 \sigma}} = \frac{1 + j}{\sigma \delta}$$

$$\gamma = \alpha + j\beta = \frac{1 + j}{\delta}$$

$$\begin{aligned} (z > 0) \quad \hat{\mathbf{E}} \times \hat{\mathbf{H}}^* &= \mathbf{u}_z \frac{\left| E_0 \right|^2}{\eta^*} \left| T \right|^2 \exp(-2\alpha z) = \\ &= \mathbf{u}_z \frac{\left| E_0 \right|^2}{2} \sigma \delta (1 + j) \left| T \right|^2 \exp(-2\alpha z) = \end{aligned}$$

$$|\Gamma|^2 = \left| \frac{(1 + j) / (\sigma \delta) - \eta_0}{(1 + j) / (\sigma \delta) + \eta_0} \right|^2 = \frac{(1 / (\sigma \delta) - \eta_0)^2 + 1 / (\sigma \delta)^2}{(1 / (\sigma \delta) + \eta_0)^2 + 1 / (\sigma \delta)^2}$$

$$1 - |\Gamma|^2 = \frac{(1 / (\sigma \delta) + \eta_0)^2 - (1 / (\sigma \delta) - \eta_0)^2}{(1 / (\sigma \delta) + \eta_0)^2 + 1 / (\sigma \delta)^2} =$$

$$= \frac{4 \eta_0 / (\sigma \delta)}{(1 / (\sigma \delta) + \eta_0)^2 + 1 / (\sigma \delta)^2}$$

$$\left| T \right|^2 = \left| 1 + \Gamma \right|^2 = \left| 1 + \frac{(1 + j) / (\sigma \delta) - \eta_0}{(1 + j) / (\sigma \delta) + \eta_0} \right|^2 =$$

$$= \left| \frac{(1 + j) / (\sigma \delta) + \eta_0 + (1 + j) / (\sigma \delta) - \eta_0}{(1 + j) / (\sigma \delta) + \eta_0} \right|^2 =$$

$$= \left| \frac{2(1 + j) / (\sigma \delta)}{(1 + j) / (\sigma \delta) + \eta_0} \right|^2 = \frac{8 / (\sigma \delta)^2}{\left(1 / (\sigma \delta) + \eta_0 \right)^2 + 1 / (\sigma \delta)^2}$$

($z = 0^+$)

$$\hat{\mathbf{E}} \times \hat{\mathbf{H}}^* = \mathbf{u}_z \frac{|E_0|^2}{2} \sigma \delta (1 + j) \frac{8 / (\sigma \delta)^2}{(1 / (\sigma \delta) + \eta_0)^2 + 1 / (\sigma \delta)^2} =$$

$$\Rightarrow \operatorname{Re} (\hat{\mathbf{E}} \times \hat{\mathbf{H}}^*) = \mathbf{u}_z |E_0|^2 \frac{4 / (\sigma \delta)}{(1 / (\sigma \delta) + \eta_0)^2 + 1 / (\sigma \delta)^2}$$

($z = 0^-$)

$$\operatorname{Re} (\hat{\mathbf{E}} \times \hat{\mathbf{H}}^*) = \mathbf{u}_z \frac{|E_0|^2}{\eta_0} \frac{4 \eta_0 / (\sigma \delta)}{(1 / (\sigma \delta) + \eta_0)^2 + 1 / (\sigma \delta)^2}$$

Onda piana riflessa da un conduttore

Un'onda piana incide normalmente su un semispazio riempito di rame.

Se $f=1\text{GHz}$, si calcoli la costante di propagazione, l'impedenza, la profondità di penetrazione. Si calcolino inoltre i coefficienti di riflessione e di trasmissione.

$$\sigma = 5.83 \cdot 10^7 \text{ S/m} \quad \delta = \frac{1}{\sqrt{\pi f \mu \sigma}} = 2.088 \cdot 10^{-6} \text{ m}$$

$$\gamma = \frac{1 + j}{\delta} = (4.789 + j4.789) \cdot 10^5 \text{ m}^{-1}$$

$$\eta = \frac{1 + j}{\sigma \delta} = (8.239 + j8.239) \cdot 10^{-3} \Omega$$

Onda piana riflessa da un conduttore II

$$\Gamma = \frac{\eta - \eta_0}{\eta + \eta_0} = 1 \angle 179.99^\circ$$

$$T = \frac{2\eta}{\eta + \eta_0} = 6.181 \times 10^{-5} \angle 45^\circ$$

Onda piana incidente su uno slab dielettrico

