

Implementation of a Fluid-Solid interaction solver

Numerical methods for partial differential equations

Giulia Cavoletti, Matteo Parimbelli, Federico Pizzolato, Paolo Giuseppe Potì

23/02/2026



POLITECNICO
MILANO 1863

NMPDE

Giulia Cavoletti, Matteo Parimbelli, Federico Pizzolato, Paolo Giuseppe Potì

Strong Formulation

The coupled FSI problem is defined on $\Omega = \Omega_f \cup \Omega_s$:

Governing equations

$$\begin{cases} -\nu \Delta \mathbf{u} + \nabla p = \mathbf{0} & \text{in } \Omega_f, \\ \nabla \cdot \mathbf{u} = 0 & \text{in } \Omega_f, \\ -\nabla \cdot \sigma(\mathbf{d}) = \mathbf{0} & \text{in } \Omega_s, \end{cases}$$

with

$$\sigma(\mathbf{d}) = \mu(\nabla \mathbf{d} + (\nabla \mathbf{d})^T) + \lambda(\nabla \cdot \mathbf{d})\mathbf{I}.$$

Boundary & interface conditions

$$\begin{cases} \mathbf{u} = \varphi & \text{on } \partial\Omega_{f_1}, \\ \nu \frac{\partial \mathbf{u}}{\partial \mathbf{n}} - p \mathbf{n} = \mathbf{0} & \text{on } \partial\Omega_{f_2}, \\ \mathbf{u} = \mathbf{0}, \quad \mathbf{d} = \mathbf{0} & \text{on } \Sigma, \partial\Omega_s \setminus \Sigma, \\ \sigma(\mathbf{d})\mathbf{n} = \nu \frac{\partial \mathbf{u}}{\partial \mathbf{n}} - p \mathbf{n} & \text{on } \Sigma. \end{cases}$$

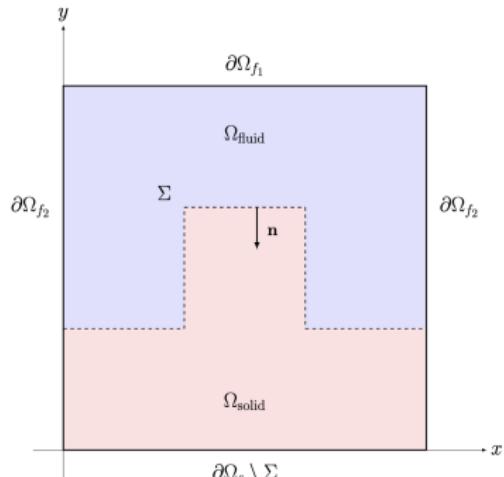


Figure 1: Domain Ω and boundary partition.



Weak Formulation

$$V_0 = \{\mathbf{v} \in [H^1(\Omega_f)]^2 : \\ \mathbf{v}|_{\Sigma} = \mathbf{0}, \mathbf{v}|_{\partial\Omega_{f_1}} = \mathbf{0}\}$$

$$\mathbb{Q} = L^2(\Omega_f)$$

$$\mathbb{B} = [H^1_{\Gamma_s}(\Omega_s)]^2$$

Find $\mathbf{u}_0 \in V$, $p \in Q$, $\mathbf{d} \in B$ such that:

$$a(\mathbf{u}_0, \mathbf{v}) + b(\mathbf{v}, p) + b(\mathbf{u}_0, q) + c(\mathbf{d}, \mathbf{b}) + m(p, \mathbf{b}) + n(\mathbf{u}_0, \mathbf{b}) = F(\mathbf{v}) + G(q)$$

for all $\{\mathbf{v}, q, \mathbf{b}\} \in Y$, where the bilinear forms are:

$$V = \{\mathbf{v} \in [L^2(\Omega)]^2 : \\ \mathbf{v}|_{\Omega_f} \in V_0, \mathbf{v}|_{\Omega_s} = \mathbf{0}\}$$

$$Q = \{q \in L^2(\Omega) : \\ q|_{\Omega_f} \in \mathbb{Q}, q|_{\Omega_s} = 0\}$$

$$B = \{\mathbf{w} \in [L^2(\Omega)]^2 : \\ \mathbf{w}|_{\Omega_s} \in \mathbb{B}, \mathbf{w}|_{\Omega_f} = \mathbf{0}\}$$



Weak Formulation

Fluid (Stokes)

$$a(\mathbf{u}, \mathbf{v}) = \int_{\Omega_f} \nu \nabla \mathbf{u} : \nabla \mathbf{v} \, d\mathbf{x},$$

$$b(\mathbf{v}, p) = - \int_{\Omega_f} p (\nabla \cdot \mathbf{v}) \, d\mathbf{x}.$$

Lifting functionals

$$F(\mathbf{v}) = -a(R_\varphi, \mathbf{v}),$$

$$G(q) = -b(R_\varphi, q).$$

Solid (Elasticity)

$$\begin{aligned} c(\mathbf{d}, \mathbf{b}) = \int_{\Omega_s} & [\mu(\nabla \mathbf{d} + (\nabla \mathbf{d})^T) : \nabla \mathbf{b} \\ & + \lambda(\nabla \cdot \mathbf{d})(\nabla \cdot \mathbf{b})] \, d\mathbf{x}, \end{aligned}$$

Interface coupling

$$m(p, \mathbf{b}) = - \int_{\Sigma} p \mathbf{n} \cdot \mathbf{b} \, d\Gamma,$$

$$n(\mathbf{u}, \mathbf{b}) = \int_{\Sigma} \nu \frac{\partial \mathbf{u}}{\partial \mathbf{n}} \cdot \mathbf{b} \, d\Gamma.$$



Discretization

We introduce a triangulation \mathcal{T}_h of Ω and finite-dimensional subspaces $V_h \subset V_0$, $Q_h \subset Q$, $B_h \subset B$. Using **Taylor-Hood** elements to satisfy the LBB (inf-sup) condition:

Finite element spaces

V_h : vector velocity, Q_{p+1}^d on Ω_f ,

Q_h : scalar pressure, Q_p on Ω_f ,

B_h : vector displacement, Q_r^d on Ω_s ,

with fields extended by zero outside their subdomain. Here Q_r^d denotes d -dimensional polynomials of degree r on each element $K \in \mathcal{T}_h$.

Discrete expansions

$$\mathbf{u}_h = \sum_{i=1}^{N_h^u} u_i \boldsymbol{\varphi}_i, \quad p_h = \sum_{i=1}^{N_h^p} p_i \boldsymbol{\varepsilon}_i$$

$$\mathbf{d}_h = \sum_{i=1}^{N_h^d} d_i \boldsymbol{\psi}_i,$$

where $\{\boldsymbol{\varphi}_i\}$, $\{\boldsymbol{\varepsilon}_i\}$, $\{\boldsymbol{\psi}_i\}$ are the vector/scalar Lagrangian bases for velocity, pressure and displacement respectively.



Algebraic system

Assembling the bilinear forms yields:

$$\underbrace{\begin{bmatrix} A & B^T & 0 \\ B & 0 & 0 \\ N & M & C \end{bmatrix}}_A \begin{bmatrix} \mathbf{U} \\ \mathbf{P} \\ \mathbf{D} \end{bmatrix} = \begin{bmatrix} \mathbf{F} \\ \mathbf{G} \\ \mathbf{0} \end{bmatrix},$$

where the matrix blocks are:

$$A_{ij} = a(\varphi_j, \varphi_i),$$

$$B_{ij} = b(\varphi_j, \varepsilon_i),$$

$$C_{ij} = c(\psi_j, \psi_i),$$

$$M_{ij} = m(\varepsilon_j, \psi_i),$$

$$N_{ij} = n(\varphi_j, \psi_i).$$

$\mathbf{U}, \mathbf{P}, \mathbf{D}$ collect the DoF coefficients for velocity, pressure and displacement; \mathbf{F}, \mathbf{G} arise from the lifting R_φ of the inlet data φ .



Preconditioner

Reordering DoFs as $(\mathbf{U}, \mathbf{P}, \mathbf{D})$, the global matrix has a **block-triangular** structure:

$$A = \begin{pmatrix} A_{uu} & A_{up} & 0 \\ A_{pu} & 0 & 0 \\ R_{du} & R_{dp} & A_{dd} \end{pmatrix} = \begin{pmatrix} A_{\text{fluid}} & 0 \\ R & A_{\text{solid}} \end{pmatrix},$$

Outer (global) preconditioner

Exploiting the block-triangular form, the exact inverse satisfies:

$$\widetilde{A^{-1}} = \begin{pmatrix} \widetilde{A_{\text{fluid}}^{-1}} & 0 \\ -\widetilde{A_{\text{solid}}^{-1}} R \widetilde{A_{\text{fluid}}^{-1}} & \widetilde{A_{\text{solid}}^{-1}} \end{pmatrix},$$

requiring one approximate solve each with A_{fluid} and A_{solid} .

$$A_{uu} \tilde{\mathbf{u}} = \mathbf{x}_u \rightarrow \tilde{\mathbf{p}} = \nu M_p^{-1} (\mathbf{x}_p - A_{pu} \tilde{\mathbf{u}}) \rightarrow A_{dd} \tilde{\mathbf{d}} = \mathbf{x}_d - R_{du} \tilde{\mathbf{u}} - R_{dp} \tilde{\mathbf{p}}.$$

Inner (Stokes) preconditioner

The fluid block is approximated via:

$$P = \begin{bmatrix} A_{uu} & 0 \\ A_{pu} & \frac{1}{\nu} M_p \end{bmatrix},$$

where ν is the fluid viscosity and $(M_p)_{ij} = \int_{\Omega_f} \varepsilon_i \varepsilon_j \, d\mathbf{x}$ is the **pressure mass matrix**.



Adaptive Mesh Refinement

After each solve, the mesh is locally refined to concentrate resolution where it is most needed.

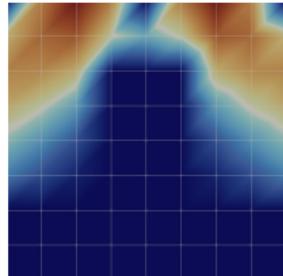
Kelly error estimator

For each cell $K \in \mathcal{T}_h$, the local indicator is:

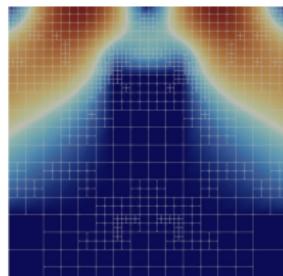
$$\eta_K^2 = \sum_{F \in \partial K} c_F \int_{\partial K_F} \left[\left[a \frac{\partial u_h}{\partial n} \right] \right]^2 d\Gamma,$$

where $\llbracket \cdot \rrbracket$ denotes the jump across face F , c_F is a face-geometry factor and a a coefficient depending on the specific pde. For each subdomain, the global estimate is:

$$\eta^2 = \sum_{K \in \mathcal{T}_h} \eta_K^2.$$



Initial mesh



Final refined mesh



Tests

Two test cases were run, sharing the following parameters:

Parameter	Symbol	Value
Kinematic viscosity	ν	2.0
Lamé parameter (shear)	μ	1.0 Pa
Lamé parameter (bulk)	λ	10.0 Pa
Kelly estimator weights	$w_{\text{Stokes}} / w_{\text{elas}}$	4.0 / 1.0
Refinement target		top 30% of cells
Refinement cycles		5 (2D) - 3 (3D)

Test 1 — 2D

Inlet velocity on $\partial\Omega_{f_1}$:

$$v_{\text{in}} = \sin(\pi x)$$

Test 2 — 3D

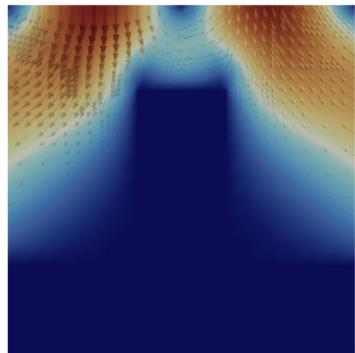
Inlet velocity on $\partial\Omega_{f_1}$:

$$v_{\text{in}} = \sin(\pi x) \sin(\pi y)$$

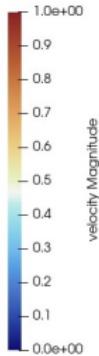
In both cases, zero-displacement conditions are imposed on $\partial\Omega_s \setminus \Sigma$, and Gauss quadrature of order $p + 2$ is used.



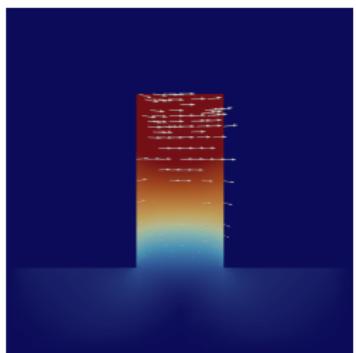
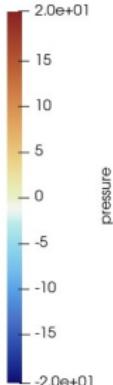
Simulation Results, 2D



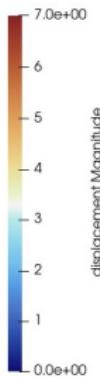
Velocity vectors



Pressure field



Displacement field



Simulation results, 3D

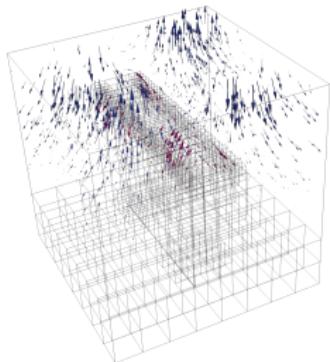
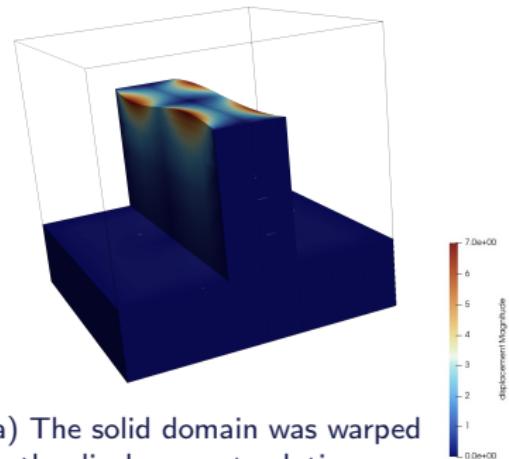


Figure 2: Vector plot for velocity (in blue) and for solid displacement (in red).



(a) The solid domain was warped by the displacement solution



Strong Scalability

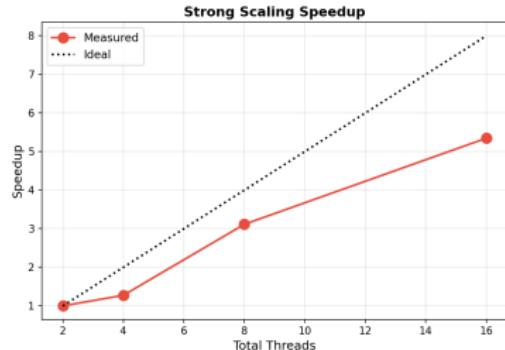


Figure 4: Speedup on the unrefined mesh: nearly linear trend confirms good parallel efficiency.

Within the same refinement cycle, a $2\times$ increase in cores yields a $\approx 2\times$ reduction in wall time.

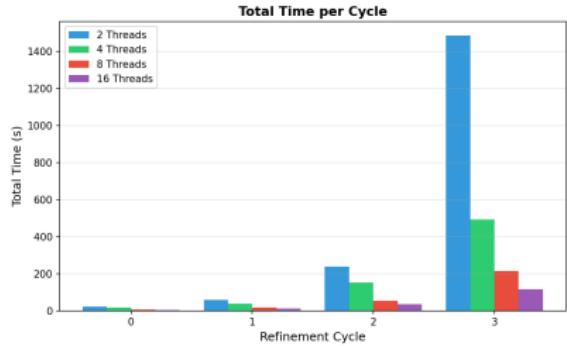


Figure 5: Total time per cycle: doubling resources approximately halves execution time within each cycle.



Weak Scalability & Preconditioner Effectiveness

Weak Scalability

Problem size grows proportionally with the thread count.

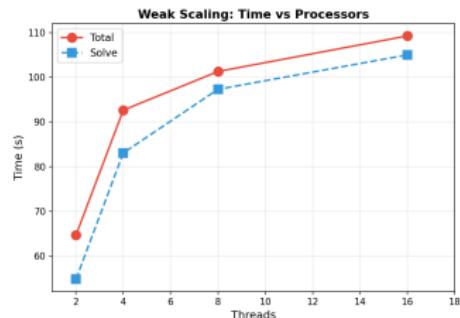


Figure 6: Execution time vs. threads. Sub-linear growth that flattens confirms solid weak scalability.

Preconditioner vs. Direct Solver

GMRES + block preconditioner compared against SparseDirectUMFPACK.

Cycle	Iters	Prec [s]	UMFPACK [s]
0	48	20.9	4.1
1	51	46.7	19.7
2	51	196.3	96.7
3	73	853.1	1198.2

- GMRES iteration count stays stable, proving the preconditioner controls the condition number effectively.

