

Basic formula of elastic equation in displacement form:

$$\lambda \nabla(\nabla \cdot \mathbf{u}) + G \nabla \cdot [(\nabla \mathbf{u})^T + \nabla \mathbf{u}] + \mathbf{F} = 0 \quad (1)$$

for test function \mathbf{u} , which is zero on part $\Gamma_D \subset \partial\Omega$ weak form:

$$\lambda(\nabla(\nabla \cdot \mathbf{u}), \mathbf{v})_\Omega + G(\nabla \cdot [(\nabla \mathbf{u})^T + \nabla \mathbf{u}], \mathbf{v})_\Omega + (\mathbf{F}, \mathbf{v})_\Omega = 0. \quad (2)$$

The inner product is defined by (d=2 or 3):

$$(\mathbf{u}, \mathbf{v}) = \sum_{i=1}^d \int_{\Omega} \mathbf{u}_i(\mathbf{x}) \mathbf{v}_i(\mathbf{x}) d\mathbf{x} \quad (3)$$

we denote $\boldsymbol{\epsilon}(\mathbf{u}) = \frac{1}{2}[(\nabla \mathbf{u})^T + \nabla \mathbf{u}]$, which is strain tensor. Therefore, the weak form is:

$$\lambda(\nabla(\nabla \cdot \mathbf{u}), \mathbf{v})_\Omega + 2G(\nabla \cdot \boldsymbol{\epsilon}(\mathbf{u}), \mathbf{v})_\Omega + (\mathbf{F}, \mathbf{v})_\Omega = 0. \quad (4)$$

We integrate by parts and use divergence theorem:

$$-\lambda(\nabla \cdot \mathbf{u}, \nabla \cdot \mathbf{v})_\Omega + \lambda(\mathbf{n} \cdot \mathbf{v}, \nabla \cdot \mathbf{u})_{\partial\Omega} - 2G(\boldsymbol{\epsilon}(\mathbf{u}), \nabla \mathbf{v})_\Omega + 2G(\boldsymbol{\epsilon}(\mathbf{u}), \mathbf{n} \otimes \mathbf{v})_{\partial\Omega} + (\mathbf{F}, \mathbf{v})_\Omega = 0. \quad (5)$$

where \mathbf{n} is the normal vector at boundary $\partial\Omega$ For boundary condition, notice the symmetric property of $\boldsymbol{\epsilon}(\mathbf{u})$ we have:

$$\lambda(\mathbf{n} \cdot \mathbf{v}, \nabla \cdot \mathbf{u})_{\partial\Omega} + 2G(\boldsymbol{\epsilon}(\mathbf{u}), \mathbf{n} \otimes \mathbf{v})_{\partial\Omega} = (\mathbf{v}, ((\lambda \nabla \cdot \mathbf{u}) \mathbf{I}_d + 2G \boldsymbol{\epsilon}(\mathbf{u})) \cdot \mathbf{n})_{\Gamma_N} \quad (6)$$

It follows:

$$(\mathbf{v}, ((\lambda \nabla \cdot \mathbf{u}) \mathbf{I}_d + 2G \boldsymbol{\epsilon}(\mathbf{u})) \cdot \mathbf{n})_{\Gamma_N} = (\mathbf{v}, \boldsymbol{\sigma} \cdot \mathbf{n})_{\Gamma_N} \quad (7)$$

Therefore, if we prescribe stress boundary condition:

$$\boldsymbol{\sigma} \cdot \mathbf{n} = \mathbf{T}$$

, then the neumann condition contributes to the right hand side of the weak form integration $(\mathbf{v}, \mathbf{T})_{\Gamma_N}$ In conclusion, the weak formulation follows:

$$\lambda(\nabla \cdot \mathbf{u}, \nabla \cdot \mathbf{v})_\Omega + 2G(\boldsymbol{\epsilon}(\mathbf{u}), \nabla \mathbf{v})_\Omega = (\mathbf{F}, \mathbf{v})_\Omega + (\mathbf{v}, \mathbf{T})_{\Gamma_N}. \quad (8)$$