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Functional renormalization group for cold atom mixtures

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F. Isaule, I. Morera, A. Polls and B. Juliá-Díaz, PRA **103**, 013318 (2020)

Outline

- 1. Functional renormalization for cold atom systems
- 2. Repulsive Bose-Bose mixtures
- 3. The Bose polaron
- 4. Conclusions

Example: Weakly interacting Bose gas

 In a field theory formulation, a weakly interacting Bose gas is described by the microscopic action

$$S = \int_{x} \left[\varphi^{\dagger} \left(\partial_{\tau} - \frac{\nabla^{2}}{2m} - \mu \right) \varphi + \frac{g}{2} (\varphi^{\dagger} \varphi)^{2} \right]$$

g: repulsive contact potential

It defines the grand-canonical partition function

$$\mathcal{Z}[\varphi] = \int D(\varphi, \varphi^\dagger) e^{-\mathcal{S}[\varphi]} \qquad \rightarrow \qquad \begin{aligned} \Omega &= -T \ln \mathcal{Z} \\ d\Omega &= -PdV - SdT - Nd\mu \end{aligned}$$

$$\Omega = -PdV - SdT - Nd\mu$$

$$\Omega : \text{ grand-canonical potential}$$

Effective action

To obtain Z we need to integrate the different paths

$$\mathcal{Z}[\varphi] = \int D(\varphi, \varphi^{\dagger}) e^{-\mathcal{S}[\varphi]}$$

- Approximations: mean-field, Gaussian
 L. Salasnich and F. Toigo, Phys. Rep. 640, 1 (2016)
- An alternative is to work in terms of an **effective action** Γ that already contains the effect of fluctuations

$$\Gamma[\phi] = -\ln \mathcal{Z}_J[\phi] + \int_x J \cdot \phi \quad \to \quad \Omega = -T \Gamma[\phi_0]$$

$$\phi(x) = \langle \varphi(x) \rangle$$
J: source fields

 The are different ways to compute Γ. For example, in a perturbative expansion:

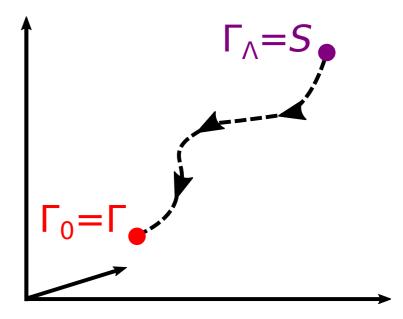
$$\Gamma[\phi] = \mathcal{S}[\phi] + \frac{1}{2} \operatorname{tr} \ln \left(\mathcal{S}^{(2)}[\phi] \right) + \dots$$

Functional renormalization group (FRG)

- The FRG is a modern non perturbative formulation of Wilson's RG.
- A regulator function R_k is added to the theory. It **suppresses** all fluctuations for q < k.
- We work in terms of a k-dependent effective action Γ_k . We follow its flow with a RG equation

$$\partial_k \Gamma_k = \frac{1}{2} \operatorname{tr} \left[\partial_k R_k (\Gamma^{(2)} + R_k)^{-1} \right]$$

C. Wetterich, Phys. Lett. B 301, 90 (1993)



Functional renormalization group (FRG)

- The FRG is used in a variety of fields: high-energy physics, condensed matter, statistical physics, etc.
- It is particularly useful to study **strongly correlated systems** and critical phenomena.

 Recent comprehensive review: N. Dupuis *et al.*, arXiv:2006.04853
- In has been used in different cold atom problems:
 - One-component Bose gases
 S. Floerchinger and C. Wetterich, PRA 77, 053603 (2008). PRA 79, 013601 (2009).
 - Fermi gases: BCS-BEC crossover
 S. Floerchinger et al., PRA 81, 063619 (2010). I. Boettcher et al., PRA 89, 053630 (2014).
 - Bose gases in optical lattices
 A. Rançon and N. Dupuis, PRA 85, 063607 (2012), PRA 86, 043624 (2012).
 - Few bosons: Efimov physics
 S. Floerchinger et al., Few-Body Syst. 51, 153 (2011). R. Schmidt and S. Moroz, PRA 81, 052709 (2010).
 - Fermi polaron
 R. Schmidt and T. Enss, PRA 83, 063620 (2011).

Example: Weakly interacting Bose gas

The microscopic action:

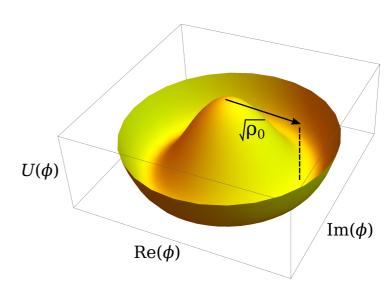
$$S = \int_{x} \left[\varphi^{\dagger} \left(\partial_{\tau} - \frac{\nabla^{2}}{2m} - \mu \right) \varphi + \frac{g}{2} (\varphi^{\dagger} \varphi)^{2} \right]$$

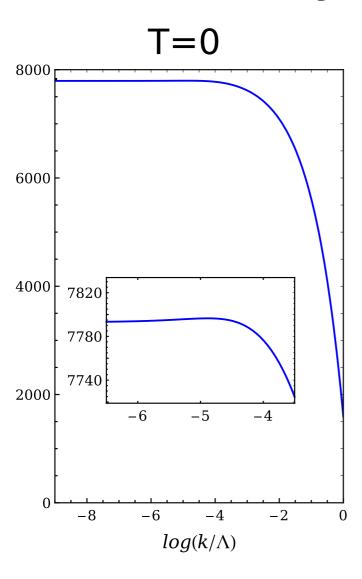
We propose an ansatz based on a derivative expansion:

$$\Gamma_k[\phi] = \int_x \left[\phi^{\dagger} \left(S \partial_{\tau} - \frac{Z}{2m} \nabla^2 - V \partial_{\tau}^2 \right) \phi + U(\rho) \right]$$

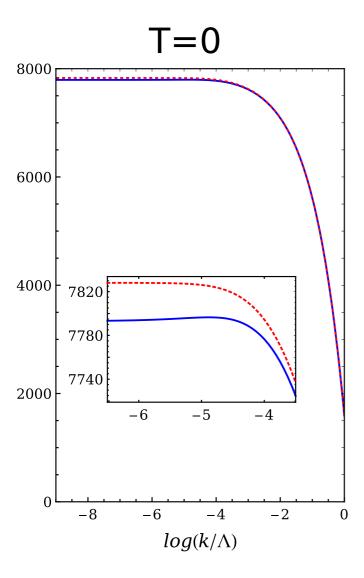
$$U(\rho) = -P + u(\rho - \rho_0) + \frac{\lambda}{2}(\rho - \rho_0)^2, \qquad \rho = \phi^{\dagger}\phi$$

- Couplings flow with k
- $\rho_{\scriptscriptstyle 0}$: condensate density
- Physical inputs: μ, a, T

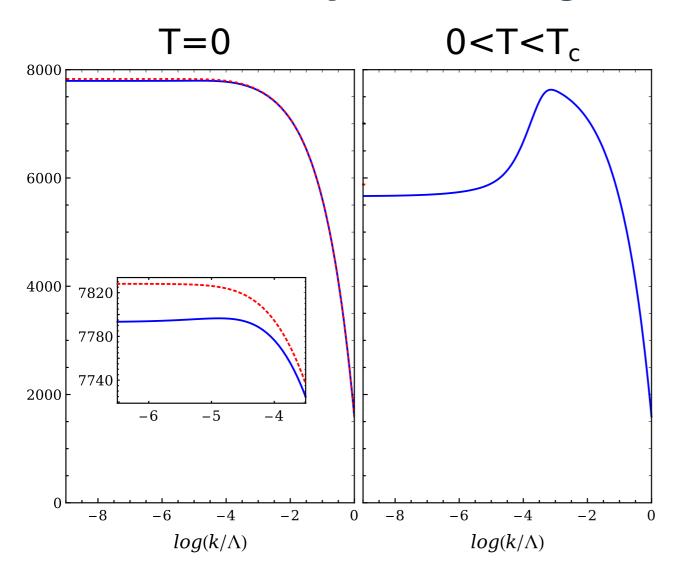




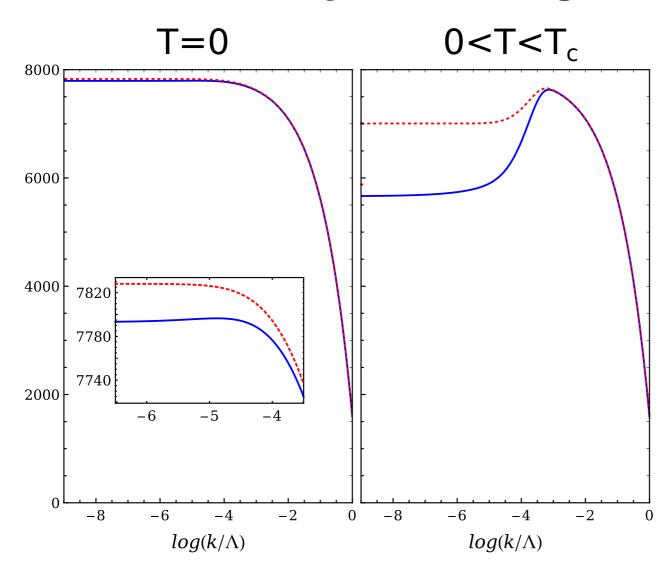
 ρ_0 : condensate density



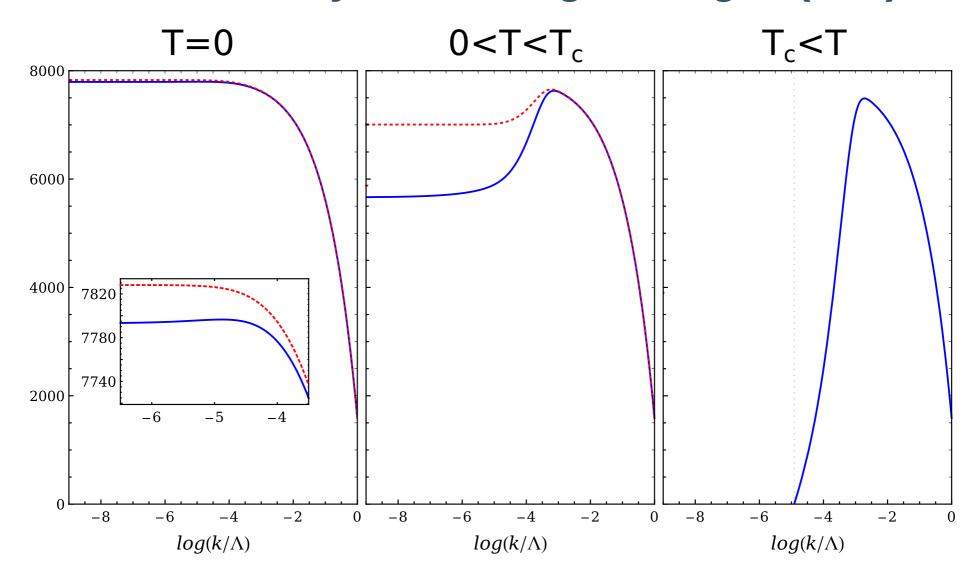
 ρ_0 : condensate density



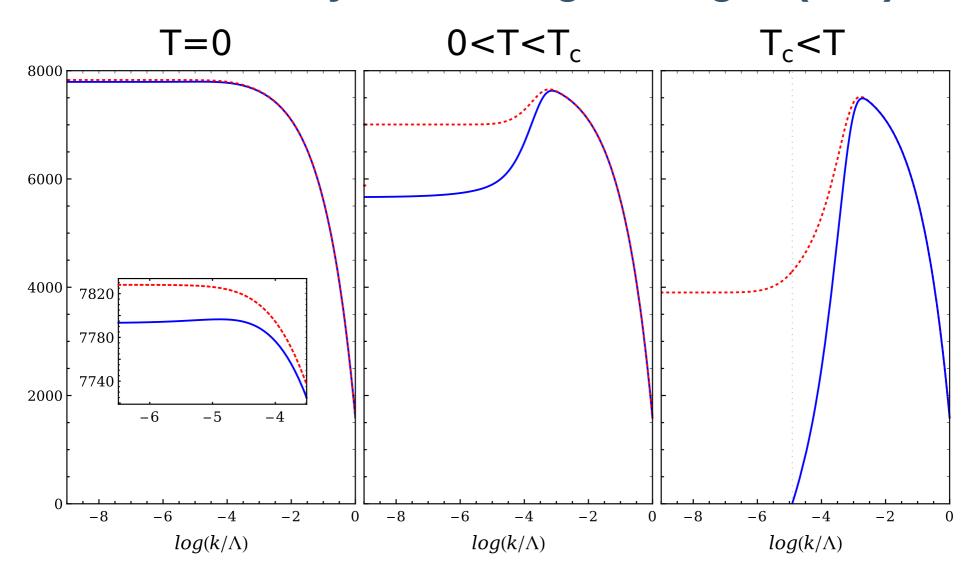
 ρ_0 : condensate density



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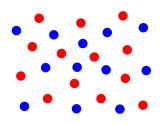
 ρ_0 : condensate density



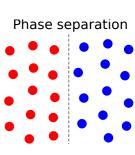
 ρ_0 : condensate density

Bose-Bose mixtures

 Gases with two species of bosons have attracted significant attention in recent years



- The interplay between the two component of the gas leads to rich physics:
 - Spin drag
 - Phase separation (repulsive interspecies potential)
 - Self-bound droplets (attractive interspecies potential)
 D. Petrov, PRL 115, 155302 (2015)
- We study balanced and repulsive Bose-Bose mixtures with the FRG



FRG for repulsive Bose-Bose mixtures

• We consider a balanced Bose-Bose mixture at T=0

$$m_A = m_B, \quad \mu_A = \mu_B, \quad a_{AA} = a_{BB}$$

$$\Gamma_k[\phi] = \int_x \left[\sum_{a=A,B} \psi_a^{\dagger} \left(S \partial_{\tau} - \frac{Z}{2m} \nabla^2 - V \partial_{\tau}^2 \right) \psi_a + U(\rho_A, \rho_B) \right]$$

$$U(\rho_A, \rho_B) = -P + \frac{\lambda}{2}(\rho_A - \rho_0)^2 + \frac{\lambda}{2}(\rho_B - \rho_0)^2 + \lambda_{AB}(\rho_A - \rho_0)(\rho_B - \rho_0)$$

$$\rho_a = \psi_a^{\dagger} \psi_a$$

FRG for repulsive Bose-Bose mixtures

• We consider a balanced Bose-Bose mixture at T=0

$$m_A = m_B, \quad \mu_A = \mu_B, \quad a_{AA} = a_{BB}$$

We propose the following ansatz:

$$\Gamma_{k}[\phi] = \int_{x} \left[\sum_{a=A,B} \psi_{a}^{\dagger} \left(S \partial_{\tau} - \frac{Z}{2m} \nabla^{2} - V \partial_{\tau}^{2} \right) \psi_{a} + U(\rho_{A}, \rho_{B}) \right]$$

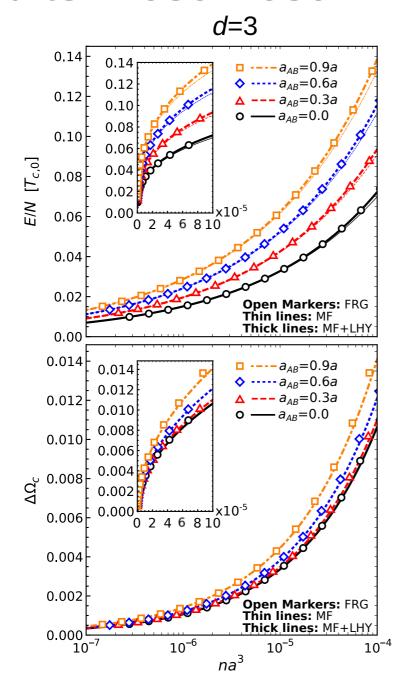
$$U(\rho_{A}, \rho_{B}) = -P + \left(\frac{\lambda}{2} (\rho_{A} - \rho_{0})^{2} + \frac{\lambda}{2} (\rho_{B} - \rho_{0})^{2} \right) + \left(\lambda_{AB} (\rho_{A} - \rho_{0})(\rho_{B} - \rho_{0}) \right)$$

$$\rho_{a} = \psi_{a}^{\dagger} \psi_{a}$$

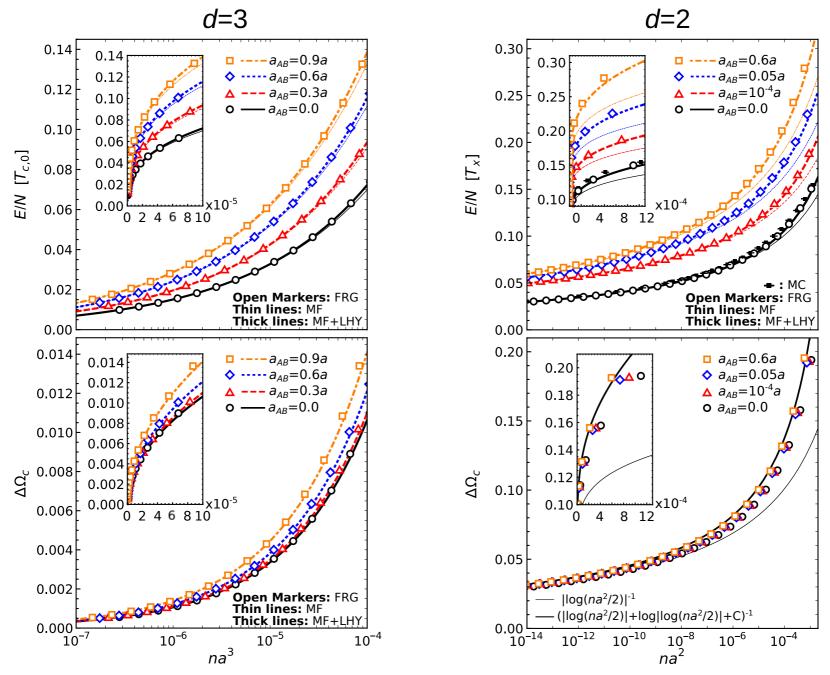
intra-species

inter-species

Results: Bose-Bose mixture



Results: Bose-Bose mixture



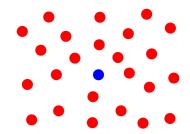
MC: G. Astrakharchik et al., PRA 79, 051602(R) (2009)

FRG for repulsive Bose-Bose mixtures

- FRG calculations compare favourably with known macroscopic results
- Future work:
 - Finite temperature
 - Spin drag
 - Attractive Bose-Bose mixture: liquid phase, dimerization, strongly-interacting regime

The Bose polaron

- Impurity immersed in a weakly-interacting Bose gas
- Several developments in the past few years:



- Experimental realization
 N. Jørgensen et al., PRL 117, 055302 (2016), M. Hu et al., PRL 117, 055301 (2016)
- Theoretical description of strong coupling regime
 J. Levinsen et al., PRL 115, 125302 (2015). N.-E. Guenther et al., PRL 120, 050405 (2018).
 L. Peña Ardila et al., PRA 99, 063607 (2019).
- Bosonic medium: three- and more-body correlations can be important
- Efimov physics

N. T. Zinner, EPL **101**, 60009 (2013). M. Sun et al., PRL **119**, 013401 (2017)

The Bose polaron

Microscopic action

$$\mathcal{S} = \int_x \left[\psi_B^\dagger \left(\partial_\tau - \frac{\nabla^2}{2m_B} - \mu_B \right) \psi_B + \psi_I^\dagger \left(\partial_\tau - \frac{\nabla^2}{2m_I} - \mu_I \right) \psi_I \right. \\ \left. + \frac{g_{BB}}{2} (\psi_B^\dagger \psi_B)^2 + g_{BI} \psi_B^\dagger \psi_I^\dagger \psi_B \psi_I \right] \\ \mu_{\mathrm{B}}: \text{ chemical potential of the medium } \mu_{\mathrm{B}}: \text{ polaron energy}$$

- Boson-impurity interaction can either be repulsive or attractive
- Resonant interaction in attractive branch: $a_{BI} \rightarrow \infty$
- Hubbard-Stratonovich transformation: **Dimer fields** $\phi \sim \psi_B \psi_I$

$$S = \int_{x} \left[\psi_{B}^{\dagger} \left(\partial_{\tau} - \frac{\nabla^{2}}{2m_{B}} - \mu_{B} \right) \psi_{B} + \psi_{I}^{\dagger} \left(\partial_{\tau} - \frac{\nabla^{2}}{2m_{I}} - \mu_{I} \right) \psi_{I} + \nu_{\phi} \phi^{\dagger} \phi \right.$$
$$\left. + \frac{g_{BB}}{2} (\psi_{B}^{\dagger} \psi_{B})^{2} + h_{\Lambda} \left(\phi^{\dagger} \psi_{B} \psi_{I} + \phi \psi_{B}^{\dagger} \psi_{I}^{\dagger} \right) \right]$$

S. P. Rath and R. Schmidt, PRA 88, 053632 (2013)

$$\Gamma_{k} = \int_{x} \left[\psi_{B}^{\dagger} \left(S_{B} \partial_{\tau} - \frac{Z_{B}}{2m_{B}} \nabla^{2} - V_{B} \partial_{\tau}^{2} \right) \psi_{B} + \psi_{I}^{\dagger} \left(S_{I} \partial_{\tau} - \frac{Z_{I}}{2m_{I}} \nabla^{2} + u_{I} \right) \psi_{I} \right.$$

$$\left. + U(\rho_{B}) + \phi^{\dagger} \left(S_{\phi} \partial_{\tau} - \frac{Z_{\phi}}{2m_{\phi}} \nabla^{2} + u_{\phi} \right) \phi + h \left(\phi^{\dagger} \psi_{B} \psi_{I} + \phi \psi_{B}^{\dagger} \psi_{I}^{\dagger} \right) \right.$$

$$\left. + \lambda_{B\phi} \psi_{B}^{\dagger} \phi^{\dagger} \psi_{B} \phi + \dots \right]$$

$$\Gamma_k = \int_x \left[\psi_B^\dagger \left(S_B \partial_\tau - \frac{Z_B}{2m_B} \nabla^2 - V_B \partial_\tau^2 \right) \psi_B + \psi_I^\dagger \left(S_I \partial_\tau - \frac{Z_I}{2m_I} \nabla^2 + u_I \right) \psi_I \right.$$

$$\left. \left(+ U(\rho_B) \right) + \phi^\dagger \left(S_\phi \partial_\tau - \frac{Z_\phi}{2m_\phi} \nabla^2 + u_\phi \right) \phi + h \left(\phi^\dagger \psi_B \psi_I + \phi \psi_B^\dagger \psi_I^\dagger \right) \right.$$

$$\left. + \lambda_{B\phi} \psi_B^\dagger \phi^\dagger \psi_B \phi + \dots \right]$$
Boson-boson interaction

$$\Gamma_k = \int_x \left[\psi_B^\dagger \left(S_B \partial_\tau - \frac{Z_B}{2m_B} \nabla^2 - V_B \partial_\tau^2 \right) \psi_B + \psi_I^\dagger \left(S_I \partial_\tau - \frac{Z_I}{2m_I} \nabla^2 + u_I \right) \psi_I \right.$$

$$\left. + U(\rho_B) + \phi^\dagger \left(S_\phi \partial_\tau - \frac{Z_\phi}{2m_\phi} \nabla^2 + u_\phi \right) \phi + h \left(\phi^\dagger \psi_B \psi_I + \phi \psi_B^\dagger \psi_I^\dagger \right) \right.$$

$$\left. + \lambda_{B\phi} \psi_B^\dagger \phi^\dagger \psi_B \phi + \dots \right]$$
Boson-boson interaction
Boson-impurity interaction

We propose the following ansatz

$$\Gamma_k = \int_x \left[\psi_B^\dagger \left(S_B \partial_\tau - \frac{Z_B}{2m_B} \nabla^2 - V_B \partial_\tau^2 \right) \psi_B + \psi_I^\dagger \left(S_I \partial_\tau - \frac{Z_I}{2m_I} \nabla^2 + u_I \right) \psi_I + \psi_I^\dagger \left(S_I \partial_\tau - \frac{Z_I}{2m_I} \nabla^2 + u_I \right) \psi_I \right]$$
 Boson-boson interaction Boson-impurity interaction
$$\left[+ \lambda_{B\phi} \psi_B^\dagger \phi^\dagger \psi_B \phi + \dots \right]$$

coupling

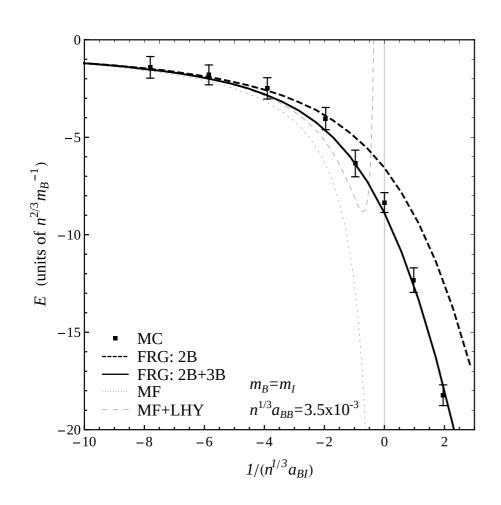
We propose the following ansatz

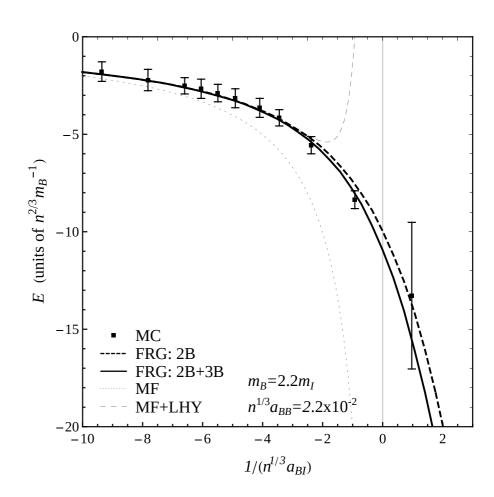
$$\Gamma_k = \int_x \left[\psi_B^\dagger \left(S_B \partial_\tau - \frac{Z_B}{2m_B} \nabla^2 - V_B \partial_\tau^2 \right) \psi_B + \psi_I^\dagger \left(S_I \partial_\tau - \frac{Z_I}{2m_I} \nabla^2 + u_I \right) \psi_I + U(\rho_B) \right] + \phi^\dagger \left(S_\phi \partial_\tau - \frac{Z_\phi}{2m_\phi} \nabla^2 + u_\phi \right) \phi + h \left(\phi^\dagger \psi_B \psi_I + \phi \psi_B^\dagger \psi_I^\dagger \right)$$
Boson-boson interaction

Boson-impurity interaction

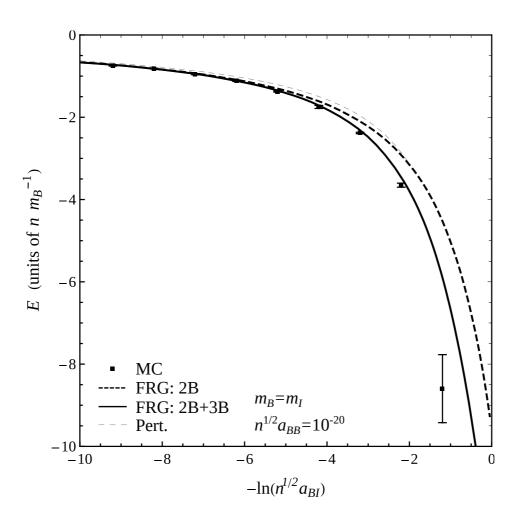
Three-body coupling

Conceptually easy to add three- and more-body correlations





MC: L. Peña Ardila et al., PRA A 99, 063607 (2019)



MC: L. Peña Ardila *et al.*, PRR **2**, 023405 (2020)

Conclusions

- The FRG can provide a successful description of bosonic mixtures.
- It gives a good description of the ground state properties of the Bose polaron within a derivative expansion. Strong coupling regime is reasonably well described.
 - Future work: Finite temperature, four-body correlations, momentum dependence
- Macroscopic properties of repulsive Bose-Bose mixtures are well described.
 - Future work: Attractive Bose-Bose mixtures
- Other related extensions: Bose-Fermi mixtures, SU(N) Fermi gases

Example: Weakly interacting Bose gas

 In a field theory formulation, a weakly interacting Bose gas is described by the microscopic action

$$S = \int_{x} \left[\varphi^{\dagger} \left(\partial_{\tau} - \frac{\nabla^{2}}{2m} - \mu \right) \varphi + \frac{g}{2} (\varphi^{\dagger} \varphi)^{2} \right]$$

g: repulsive contact potential

It defines the grand-canonical partition function

$$\mathcal{Z}[\varphi] = \int D(\varphi, \varphi^{\dagger}) e^{-\mathcal{S}[\varphi]} \longrightarrow \begin{array}{c} \Omega = -T \ln \mathcal{Z} \\ d\Omega = -PdV - SdT - Nd\mu \end{array}$$

 Ω : grand-canonical potential

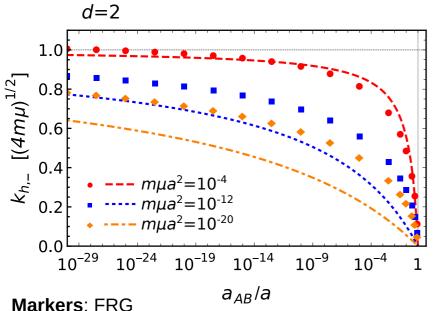
The interaction is related to the two-body T matrix

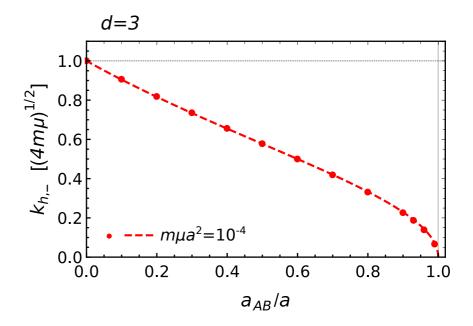
$$T^{2B} = \begin{cases} \frac{4\pi a}{m} & : d = 3\\ \frac{4\pi/m}{\log(2/|\mu|a^2) - 2\gamma_E} & : d = 2 \end{cases}$$

a: s-wave scattering length

Results: Bose-Bose mixture

• A vanishing spin healing scale $k_{h.}$ signals the phase separation





Lines: Bogoliubov

- Within the range of parameters explored, we find that the point of phase separation is $a_{AB} = a$ (mean field result).
- RG studies suggest that the phase separation point in two dimensions occurs for $a_{_{AB}} < a$ at logarithmically small densities

A. K. Kolezhuk, PRA 81, 013601(2010)

$$\Gamma_{k} = \int_{x} \left[\psi_{B}^{\dagger} \left(S_{B} \partial_{\tau} - \frac{Z_{B}}{2m_{B}} \nabla^{2} - V_{B} \partial_{\tau}^{2} \right) \psi_{B} + \psi_{I}^{\dagger} \left(S_{I} \partial_{\tau} - \frac{Z_{I}}{2m_{I}} \nabla^{2} + u_{I} \right) \psi_{I} \right.$$

$$\left. + U(\rho_{B}) + \phi^{\dagger} \left(S_{\phi} \partial_{\tau} - \frac{Z_{\phi}}{2m_{\phi}} \nabla^{2} + u_{\phi} \right) \phi + h \left(\phi^{\dagger} \psi_{B} \psi_{I} + \phi \psi_{B}^{\dagger} \psi_{I}^{\dagger} \right) \right.$$

$$\left. + \lambda_{B\phi} \psi_{B}^{\dagger} \phi^{\dagger} \psi_{B} \phi + \dots \right]$$

- No feedback from the impurity
- <u>Drawbacks of the derivative expansion</u>:
 - Not really suitable to study few-body physics: big mass imbalance? Efimov trimers?
 - Spectral function is not accessible