Next-to-Leading Order QCD Corrections to Inclusive Heavy-Flavor Production in Polarized Deep-Inelastic Scattering

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Outline

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- 3 Partonic Results
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Introduction - Heavy Quarks

HQ are good





Introduction - Experimental Setups

e^- - e^+ -annihilation (SIA)	deep inelastic scattering (DIS)	Drell-Yan process (DY)
$e^- + e^+ ightarrow \overline{Q} + X[Q]$	$\ell + h \rightarrow \ell' + \overline{Q} + X[Q]$	$h+h' o \overline{Q} + X[Q]$
	$\begin{array}{c} \ell \\ \hline \\ \gamma^* \\ \hline \overline{Q} \\ \\ \end{matrix}$	$\begin{array}{c} h \\ \hline Q \\ \hline Q \\ \hline \\ h' \\ \hline \end{array}$
LEP, ILC	HERA, COMPASS, EIC	Tevatron, LHC
gluon	factorization	top, Higgs

Introduction - Structure Functions

cross section:
$$\frac{d^2\sigma}{dxdy} = \frac{2\pi y \alpha^2}{O^4} L^{\mu\nu} W_{\mu\nu}$$
 (1)

hard. tensor:

$$W_{\mu\nu} = \left(-g_{\mu\nu} + \frac{q_{\mu}q_{\nu}}{q^2}\right)F_1(x,Q^2) + \frac{P_{\mu}P_{\nu}}{P \cdot q}F_2(x,Q^2)$$

$$+ i\epsilon_{\mu\nu\alpha\beta} \frac{q^{\alpha}S^{\beta}}{P \cdot q} g_1(X, Q^2)$$
 (2)

$$F_L(x, Q^2) = F_2(x, Q^2) - 2xF_1(x, Q^2)$$
 (3)

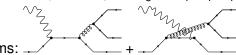
$$\frac{d^{2}\sigma}{dxdy} = \frac{2\pi\alpha^{2}}{xvQ^{2}} \left(Y_{+}F_{2}(x,Q^{2}) - y^{2}F_{L}(x,Q^{2}) \right)$$
(4)

$$\frac{d^2\Delta\sigma}{dxdy} = \frac{4\pi\alpha^2}{x_V Q^2} Y_{-} \cdot 2xg_1(x, Q^2)$$
 (5)

$$Y_{\pm} = 1 \pm (1 - y)^2 \tag{6}$$

Computation Review

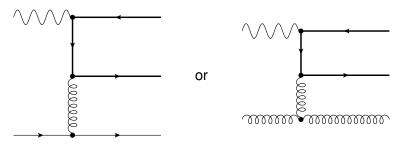
- use factorisation theorem: $s \rightarrow \xi S_h + PDF$
- \blacksquare g(k_1) + $\gamma^*(q) \rightarrow \overline{Q}(p_2) + Q(p_1)$
- three massive particles: $2 \cdot m^2 > 0$, $q^2 = -Q^2 < 0$
- \blacksquare compute 2-to-3-phase space: e.g. $\textit{dPS}_3 \sim \textit{dt}_1 \textit{ds}_4 \textit{d}\theta_1 \textit{d}\theta_2$



- compute diagrams:
- $\blacksquare \Rightarrow 2xg_1(x) \sim e_u^2 \cdot \xi \Delta u(\xi) \otimes d_{P,q}^{(1)}(\chi,\chi')$
- $d_{P,q}^{(1)}(\chi,\chi') = c_1(\chi,\chi')\ln(\chi) + c_2(\chi,\chi')\operatorname{Li}_2\left(\frac{1+\chi'}{1+\chi}\right) + \dots \checkmark$

Computation Review - Collinear Poles

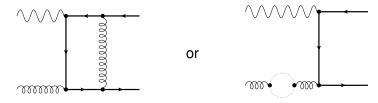
collinear poles appear in, e.g.,



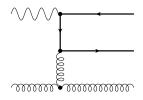
- remove by mass factorization $\rightarrow MS_m$
- $\blacksquare \Rightarrow 2xg_1(x) \sim e_H^2 \cdot \xi \Delta g(\xi) \otimes \ln(\mu_F^2/m^2) \bar{c}_{P,g}^{F,(1)}(\chi,\chi_q)$
- $\bar{\boldsymbol{c}}_{P,g}^{F,(1)}(\chi,\chi_q) = \boldsymbol{c}_1(\chi,\chi_q) \ln(\chi) + \boldsymbol{c}_2(\chi,\chi_q) \operatorname{Li}_2\left(\frac{1-\chi_q}{1+\chi}\right) + \dots \left(\checkmark \operatorname{for} \right)$ $Q^2 \gg m^2$)

Computation Review - UV,IR and Virtual Poles

virtual diagrams are, e.g.,



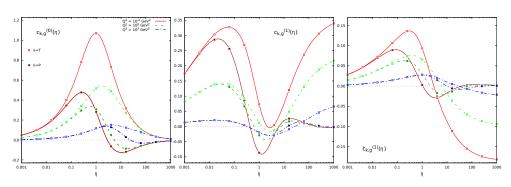
soft poles appear in the limit of a soft gluon $\textit{k}_2 \rightarrow 0, \text{ e.g.},$



soft + virtual + renormalization + factorization is finite!

Partonic Results - Gluon Channel

$$2xg_1(x) \sim \xi \Delta g(\xi) \otimes \left(c_{P,g}^{(0)} + 4\pi\alpha_s \left[c_{P,g}^{(1)} + \ln\left(\frac{\mu^2}{m^2}\right)\bar{c}_{P,g}^{(1)}\right]\right)$$
(7)



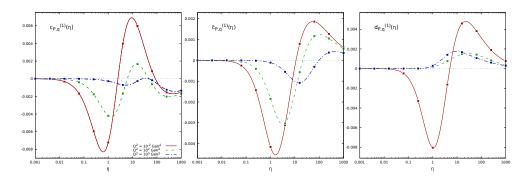
$$\eta = \frac{s-4m^2}{4m^2}, \quad m = m_b = 4.75\, \text{GeV}$$



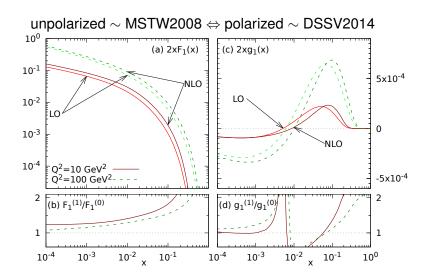


Partonic Results - Light Quark Channel

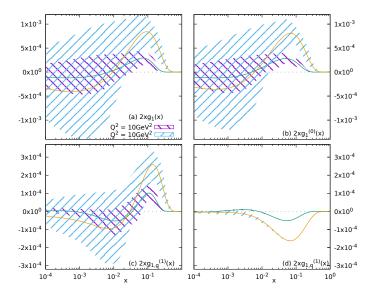
$$2xg_1(x) \sim \sum_{q \in \{u,d,s\}} \xi \Delta q(\xi) \otimes \left(e_H^2 \left[c_{P,q}^{(1)} + \ln \left(\frac{\mu^2}{m^2} \right) \bar{c}_{P,q}^{(1)} \right] + e_q^2 d_{P,q}^{(1)} \right)$$
(8)



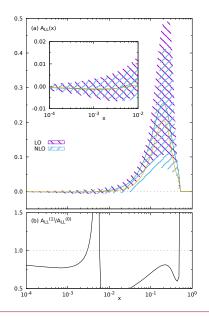
Hadronic Results - Unpolarized vs. Polarized



Hadronic Results - PDF Uncertainties DSSV (I)

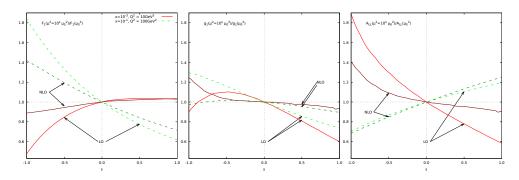


Hadronic Results - PDF Uncertainties DSSV (II)

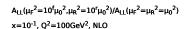


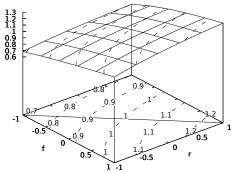
Hadronic Results - Scale Uncertainties (I)

$$\mu_0^2 = 4m^2 + Q^2$$

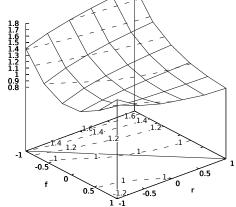


Hadronic Results - Scale Uncertainties (II)









Outlook

- inclusive distributions: $\frac{dg_1}{dp_{T,\bar{Q}}}$, $\frac{dg_1}{dy_{\bar{Q}}}$
- correlated distributions: $\frac{dg_1}{dM_{Q\bar{Q}}^2}$, $\frac{dg_1}{d\phi_{Q\bar{Q}}}$
- full neutral current (NC) contributions: $F_3^{Z\gamma}, g_4^{Z\gamma}, g_5^{Z\gamma}$ and F_2^Z, F_L^Z, g_1^Z
- distributions of full NC structure functions: $\frac{dg_1^{NC}}{dp_{T,\bar{Q}}}$, $\frac{dg_1^{NC}}{dM_{Q\bar{Q}}^2}$