1 2 to 2 phase space

following [1]:

process:

$$\gamma^*(q) + g(k_1) \to Q(p_1) + \overline{Q}(p_2) \tag{1}$$

kinematics:

$$s = (q + k_1)^2 s' = s - q^2 (2)$$

$$t = (k_1 - p_2)^2 t_1 = t - m^2 (3)$$

$$u = (q - p_2)^2 u_1 = u - m^2 (4)$$

use c.m.s. of incoming particles:

$$q = \left(\frac{s+q^2}{2\sqrt{s}}, 0, 0, \dots, -\frac{s-q^2}{2\sqrt{s}}\right)$$
 (5)

$$k_1 = \frac{s - q^2}{2\sqrt{s}} (1, 0, 0, \dots, 1) \tag{6}$$

such that

$$q + k_1 = (\sqrt{s}, \vec{0}) \qquad k_1^2 = 0$$
 (7)

for the outgoing particles it follows

$$p_1 = \frac{\sqrt{s}}{2} (1, 0, \beta \sin \theta, \dots, \beta \cos \theta)$$
 (8)

$$p_2 = \frac{\sqrt{s}}{2} (1, 0, -\beta \sin \theta, \dots, -\beta \cos \theta)$$
(9)

with $\beta = \sqrt{1 - 4m^2/s}$ such that

$$p_1 + p_2 = (\sqrt{s}, \vec{0})$$
 $p_1^2 = p_2^2 = m^2$ (10)

use n-sphere:

$$d^{D}x = \Omega_{D}x^{D-1}dx = \frac{2\pi^{D/2}}{\Gamma(D/2)}x^{D-1}dx = \frac{\pi^{D/2}}{\Gamma(D/2)}(x^{2})^{(D-2)/2}dx^{2}$$
(11)

compute phase space:

$$PS_2 = \int \frac{d^n p_1}{(2\pi)^{n-1}} \frac{d^n p_1}{(2\pi)^{n-1}} (2\pi)^n \delta^{(n)} (q + k_1 - p_1 - p_2) \delta(p_1^2 - m^2) \delta(p_2^2 - m^2)$$
 (12)

$$= \frac{1}{(2\pi)^{n-2}} \int d^n p_1 \, \delta((q + k_1 - p_1)^2 - m^2) \delta(p_1^2 - m^2)$$
(13)

$$= \frac{1}{(2\pi)^{n-2}} \int dp_{1,0} dp_{1,||} d^2 p_{1,\perp} d^{n-4} \hat{p}_1 \, \delta(s - 2p_{1,0}\sqrt{s}) \delta(p_{1,0}^2 - p_{1,||}^2 - p_{1,\perp}^2 - \hat{p}_1^2 - m^2)$$
(14)

$$= \frac{1}{(2\pi)^{n-2} 2\sqrt{s}} \int dp_{1,||} dp_{1,\perp}^2 d^{n-4} \hat{p}_1 \, \delta(s/4 - p_{1,||}^2 - p_{1,\perp}^2 - \hat{p}_1^2 - m^2)$$
 (15)

$$= \frac{1}{(2\pi)^{n-2} 2\sqrt{s}} \int dp_{1,||} d\hat{p}_1^2 \frac{\pi^{(n-4)/2}}{\Gamma((n-4)/2)} (\hat{p}_1^2)^{(n-6)/2}$$
(16)

$$= \frac{1}{2\sqrt{s}\Gamma((n-4)/2)(4\pi)^{(n-2)/2}} \int dp_{1,||} d\hat{p}_1^2 \, (\hat{p}_1^2)^{(n-6)/2}$$
(17)

Integration borders are

$$p_{1,||} \in \frac{\sqrt{s}}{2}\beta \cdot [-1,1] \qquad \hat{p}_1^2 \in \left(\frac{s\beta^2}{4} - p_{1,||}^2\right) \cdot [0,1]$$
 (18)

if cross section does not depend on hat-space:

$$\int d\hat{p}_1^2 \left(\hat{p}_1^2\right)^{(n-6)/2} = \frac{2}{n-4} \left(\frac{s\beta^2}{4} - p_{1,||}^2\right)^{(n-4)/2} \tag{19}$$

$$\Rightarrow PS_2 = \frac{1}{2\sqrt{s}\Gamma((n-2)/2)(4\pi)^{(n-2)/2}} \int dp_{1,||} \left(\frac{s\beta^2}{4} - p_{1,||}^2\right)^{(n-4)/2} \tag{20}$$

rewrite $p_{1,||}$ to $\cos \theta$:

$$p_{1,||} = \frac{\sqrt{s}}{2}\beta\cos\theta \Rightarrow dp_{1,||} = \frac{\sqrt{s}}{2}\beta\cos\theta, \quad \cos\theta \in [-1,1], \ \hat{p}_1^2 \in \frac{s\beta^2}{4}\left(1-\cos^2\theta\right) \cdot [0,1] \tag{21}$$

rewrite $\cos \theta$ to $t_1 = (k_1 - p_2)^2 - m^2$:

$$\cos \theta = \frac{2t_1/s' + 1}{\beta} \Rightarrow d\cos \theta = \frac{2}{\beta s'} dt_1, \quad t_1 \in \frac{s'}{2} [-\beta - 1, \beta - 1], \ \hat{p}_1^2 \in (-m^2 - \frac{st_1}{s'^2} (s' + t_1)) \cdot [0, 1]$$
(22)

$$p_{1,||} = \sqrt{s} \left(\frac{t_1}{s'} + \frac{1}{2} \right) \Rightarrow dp_{1,||} = \frac{\sqrt{s}}{s'} dt_1$$
 (23)

$$\Rightarrow PS_2 = \frac{1}{2s'\Gamma((n-2)/2)(4\pi)^{(n-2)/2}} \int dt_1 \left(\frac{(t_1(u_1 - q^2) - s'm^2)s' - q^2t_1^2}{s'^2} \right)^{(n-4)/2}$$
(24)

2 2 to 3 phase space

following [2, 3, 1]:

process:

$$\gamma^*(q) + q(k_1) \to Q(p_1) + \overline{Q}(p_2) + q(k_2)$$
 (25)

$$\gamma^*(q) + g(k_1) \to Q(p_1) + \overline{Q}(p_2) + g(k_2)$$
 (26)

2.1 kinematic constraints

definitions of kinematic variables:

$$s = (q + k_1)^2 \qquad \Rightarrow \qquad 2qk_1 = s - q^2 = s' \qquad (27)$$

$$s_3 = (k_2 + p_2)^2 - m^2$$
 \Rightarrow $2k_2p_2 = s_3$ (28)

$$s_4 = (k_2 + p_1)^2 - m^2$$
 \Rightarrow $2k_2p_1 = s_4$ (29)

$$s_5 = (p_1 + p_2)^2 = -u_5$$
 $\Rightarrow 2p_1p_2 = s_5 - 2m^2$ (30)

$$t_1 = (k_1 - p_2)^2 - m^2 = t - m^2$$
 \Rightarrow $2k_1p_2 = -t_1$ (31)

$$t' = (k_1 - k_2)^2$$
 \Rightarrow $2k_1k_2 = -t'$ (32)

$$u_1 = (q - p_2)^2 - m^2 = u - m^2$$
 \Rightarrow $2qp_2 = -u_1 + q^2$ (33)

$$u_6 = (k_1 - p_1)^2 - m^2$$
 \Rightarrow $2k_1 p_1 = -u_6$ (34)

$$u_7 = (q - p_1)^2 - m^2$$
 \Rightarrow $2qp_1 = -u_7 + q^2$ (35)

$$u' = (q - k_2)^2$$
 \Rightarrow $2qk_2 = -u' + q^2$ (36)

impose momentum conservation:

$$q + k_1 = p_1 + p_2 + k_2 \tag{37}$$

contract with 2 times momentum:

$$2q^{2} + s' = -u_{7} + q^{2} - u_{1} + q^{2} - u' + q^{2} \Leftrightarrow 0 = s' + u_{1} + u_{7} + u' - q^{2}$$

$$(38)$$

$$-q^{2} + 0 = -u_{3} - t_{4} - t' \Leftrightarrow 0 = s' + t_{5} + t' + u_{5}$$

$$s - q^{2} + 0 = -u_{6} \qquad -t_{1} \qquad \Leftrightarrow \quad 0 = s' + t_{1} + t' + u_{6} \qquad (39)$$

$$-u_{7} + q^{2} - u_{6} = 2m^{2} \qquad +s_{5} - 2m^{2} + s_{4} \qquad \Leftrightarrow \quad 0 = s_{4} + s_{5} + u_{6} + u_{7} - q^{2} \qquad (40)$$

$$-u_1 + q^2 - t_1 = s_5 - 2m^2 + 2m^2 + s_3 \qquad \Leftrightarrow 0 = s_3 + s_5 + t_1 + u_1 - q^2$$
(41)

$$-u' + q^2$$
 $-t' = s_4$ $+s_3$ $+0$ $\Leftrightarrow 0 = s_3 + s_4 + t' + u' - q^2$
(42)

$$\frac{1}{2}\left((38) + (39) + (41) - (40) - (42)\right) = 0 = s' + t_1 + u_1 - s_4 \tag{43}$$

$$\frac{1}{2}((38) + (39) + (41) - (40) - (42)) = 0 = s' + t_1 + u_1 - s_4$$

$$\frac{1}{2}((38) + (39) + (42) - (40) - (41)) = 0 = s' + t' + u' - s_5$$
(43)

$$\frac{1}{2}((41) + (42) - (38) - (39) - (40)) = 0 = s_3 - s' - u_6 - u_7$$

$$\frac{1}{2}((40) + (41) + (42) - (38) - (39)) = 0 = s_3 + s_4 + s_5 - s$$
(45)

$$\frac{1}{2}\left((40) + (41) + (42) - (38) - (39)\right) = 0 = s_3 + s_4 + s_5 - s \tag{46}$$

$$\frac{1}{2}\left((38) + (40) - (39) - (41) - (42)\right) = 0 = -s_3 - t_1 - t' + u_7 \tag{47}$$

$$\frac{1}{2}\left((39) + (40) + (42) - (38) - (41)\right) = 0 = s_4 + t' - u_1 + u_6 \tag{48}$$

2.2 choose framework

use c.m.s. of recoiling heavy and light quark $(Q(p_1))$ and $q(k_2)$:

$$k_2 = (\omega_2, k_{2,x}, \omega_2 \sin \theta_1 \cos \theta_2, \omega_2 \cos \theta_1, \hat{k}_2) \tag{49}$$

$$p_1 = (E_1, -k_{2,x}, -\omega_2 \sin \theta_1 \cos \theta_2, -\omega_2 \cos \theta_1, -\hat{k}_2)$$
 (50)

where $k_{2,x}=k_{2,x}(\omega_2,\theta_1,\theta_2,\hat{k}_2)$ is such, that $k_2^2=0$.

The only remainding choice is then the alignment of the z-axis: either along k_1 , called "Set I"(2.2.1), along q, called "Set II"(2.2.2) or along p_1 , called "Set III"(2.2.3).

2.2.1 Set I

align k_1 to z-axis:

$$k_1 = (\omega_1, 0, 0, \omega_1, \hat{0}) \tag{51}$$

$$q = (q_0, 0, |\vec{p}_2| \sin \psi, |\vec{p}_2| \cos \psi - \omega_1, \hat{0})$$
(52)

$$p_2 = (E_2, 0, |\vec{p}_2| \sin \psi, |\vec{p}_2| \cos \psi, \hat{0}) \tag{53}$$

constraints, from energy conservation, masses (light quark masses are already fixed: $k_1^2 =$

 $0 = k_2^2$) and external Mandelstam variables:

$$q_0 + \omega_1 = E_1 + E_2 + \omega_2 \tag{54}$$

$$m^2 = p_1^2 = E_1^2 - \omega_2^2 (55)$$

$$m^2 = p_2^2 \qquad = E_2^2 - |\vec{p_2}|^2 \tag{56}$$

$$m^{2} = p_{2}^{2} = E_{2}^{2} - |\vec{p}_{2}|^{2}$$

$$= q_{0}^{2} - |\vec{p}_{2}|^{2} + 2|\vec{p}_{2}|\omega_{1}\cos\psi - \omega_{1}^{2}$$

$$(56)$$

$$= q_{0}^{2} - |\vec{p}_{2}|^{2} + 2|\vec{p}_{2}|\omega_{1}\cos\psi - \omega_{1}^{2}$$

$$(57)$$

$$s = (q + k_1)^2 \qquad = (q_0 + \omega_1)^2 - |\vec{p}_2|^2 \qquad (58)$$

$$t = (k_1 - p_2)^2 \qquad = (\omega_1 - E_2)^2 - |\vec{p}_2|^2 + 2|\vec{p}_2|\omega_1\cos\psi - \omega_1^2 \qquad (59)$$

$$t = (k_1 - p_2)^2 \qquad = (\omega_1 - E_2)^2 - |\vec{p_2}|^2 + 2|\vec{p_2}|\omega_1\cos\psi - \omega_1^2 \qquad (59)$$

$$u = (q - p_2)^2 = (q_0 - E_2)^2 - \omega_1^2$$
(60)

solve:

$$(58) - (57) + (59) - (56) + (60) = s - q^{2} + t - m^{2} + u$$

$$(61)$$

$$= s_4 + m^2 = (E_1 + \omega_2)^2 \tag{62}$$

$$(59) + (60) - (57) = t + u - q^2 = -2(E_1 + \omega_2)E_2$$
(63)

$$\Rightarrow E_2 = -\frac{t + u - q^2}{2\sqrt{s_4 + m^2}} = \frac{s - s_4 - 2m^2}{2\sqrt{s_4 + m^2}} \tag{64}$$

$$(62) \wedge (55) \Rightarrow \omega_2 = \frac{s_4}{2\sqrt{s_4 + m^2}} \tag{65}$$

$$(62) \Rightarrow E_1 = \frac{s_4 + 2m^2}{2\sqrt{s_4 + m^2}} \tag{66}$$

$$(58) + (60) - (56) = s + u - m^2 = 2q_0(E_1 + \omega_2)$$
(67)

$$\Rightarrow q_0 = \frac{s + u_1}{2\sqrt{s_4 + m^2}} \tag{68}$$

$$(59) - (57) = t - q^2 = (\omega_1 - E_2)^2 - q_0^2 \tag{69}$$

$$\Rightarrow \omega_1 = \frac{s' + t_1}{2\sqrt{s_4 + m^2}} = \frac{s_4 - u_1}{2\sqrt{s_4 + m^2}} \tag{70}$$

$$(56) \Rightarrow |\vec{p}_2| = \sqrt{E_2^2 - m^2} = \frac{\sqrt{(s - s_4)^2 - 4sm^2}}{2\sqrt{s_4 + m^2}}$$
 (71)

$$(57) \Rightarrow \cos \psi = \frac{q^2 - q_0^2 + |\vec{p}_2|^2 + \omega_1^2}{2|\vec{p}_2|\omega_1}$$
 (72)

$$= \frac{(u_1 + m^2)(t_1 - s') - (m^2 - q^2 - t_1)(s' + t_1)}{(s' + t_1)\sqrt{(s - s_4)^2 - 4sm^2}}$$
(73)

$$\Rightarrow \sin \psi = 2 \frac{\sqrt{s_4 + m^2} \sqrt{m^2 s'^2 + q^2 t_1 (s' + t_1) - s' t_1 u_1}}{(s' + t_1) \sqrt{(s - s_4)^2 - 4sm}}$$
 (74)

$$t' = -2k_1k_2 = -2\omega_1\omega_2(1 - \cos\theta_1) \tag{75}$$

$$u_6 = -2k_1 p_1 = -2\omega_1 (E_1 + \omega_2 \cos \theta_1) \tag{76}$$

$$(39): 0 = s + t_1 + t' + u_6 - q^2 \quad \checkmark (77)$$

t' is the only variable that can get collinear (for $-q^2 > 0$).

$$s_3 = 2k_2p_2 = 2\omega_2(E_2 - |\vec{p}_2|(\cos\psi\cos\theta_1 + \sin\psi\sin\theta_1\cos\theta_2))$$
 (78)

$$s_5 = (p_1 + p_2)^2 = 2m^2 + 2p_1p_2 (79)$$

$$= 2(m^{2} + E_{1}E_{2} + \omega_{2} |\vec{p}_{2}| (\cos \psi \cos \theta_{1} + \sin \psi \sin \theta_{1} \cos \theta_{2}))$$
 (80)

(41):
$$0 = s_3 + s_5 + t_1 + u_1 - q^2 \quad \checkmark$$
 (81)

$$u' = (q - k_2)^2 = q^2 - 2qk_2 (82)$$

$$= q^{2} - 2(q_{0}\omega_{2} - \omega_{2}(|\vec{p}_{2}|\cos\psi - \omega_{1})\cos\theta_{1} - \omega_{2}|\vec{p}_{2}|\sin\psi\sin\theta_{1}\cos\theta_{2})$$
 (83)

$$u_7 = q^2 - 2qp_1 \tag{84}$$

$$= q^{2} - 2(q_{0}E_{1} + \omega_{2}(|\vec{p}_{2}|\cos\psi - \omega_{1})\cos\theta_{1} + \omega_{2}|\vec{p}_{2}|\sin\psi\sin\theta_{1}\cos\theta_{2})$$
 (85)

$$(38): 0 = s + u_1 + u_7 + u' - 2q^2 \checkmark (86)$$

2.2.2 Set II

align q to z-axis:

$$q = (q_0, 0, 0, q_z, \hat{0}) \tag{87}$$

$$k_1 = (\omega_1, 0, |\vec{p}_2| \sin \psi, |\vec{p}_2| \cos \psi - q_z, \hat{0})$$
 (88)

$$p_2 = (E_2, 0, |\vec{p}_2| \sin \psi, |\vec{p}_2| \cos \psi, \hat{0})$$
(89)

constraints, from energy conservation, masses $(k_2^2 = 0)$ is already fixed) and external

Mandelstam variables:

$$q_0 + \omega_1 = E_1 + E_2 + \omega_2 \tag{90}$$

$$m^2 = p_1^2 \qquad = E_1^2 - \omega_2^2 \tag{91}$$

$$m^2 = p_2^2 \qquad = E_2^2 - |\vec{p}_2|^2 \tag{92}$$

$$q^2 = q_0^2 - q_z^2 (93)$$

$$= \omega_1^2 - |\vec{p}_2|^2 + 2|\vec{p}_2| q_z \cos \psi - q_z^2$$
 (94)

$$s = (q + k_1)^2 \qquad = (q_0 + \omega_1)^2 - |\vec{p_2}|^2 \tag{95}$$

$$t = (k_1 - p_2)^2 \qquad = (\omega_1 - E_2)^2 - q_z^2 \tag{96}$$

$$u = (q - p_2)^2 = (q_0 - E_2)^2 - |\vec{p}_2|^2 + 2|\vec{p}_2| q_z \cos \psi - q_z^2$$
 (97)

solve:

$$(97) - (94) = u = (q_0 - E_2)^2 - \omega_1^2 \tag{98}$$

$$(95) - (93) + (96) - (92) + (98) = s - q^2 + t - m^2 + u$$

$$(99)$$

$$= s_4 + m^2 = (\omega_1 + q_0 - E_2)^2 \tag{100}$$

$$(96) + (97) - (93) = t + u - q^{2} = 2(E_{2} - \omega_{1} - q_{0})E_{2}$$
(101)

$$\Rightarrow E_2 = -\frac{t + u - q^2}{2\sqrt{s_4 + m^2}} = \frac{s - s_4 - 2m^2}{2\sqrt{s_4 + m^2}} = (64)$$
 (102)

$$(95) - (93) + (96) - (92) = s - q^2 + t - m^2 = -2(E_2 - \omega_1 - q_0)\omega_1$$
(103)

$$\Rightarrow \omega_1 = \frac{s' + t_1}{2\sqrt{s_4 + m^2}} = \frac{s_4 - u_1}{2\sqrt{s_4 + m^2}} = (70)$$

$$(100) \wedge (102) \wedge (104) \Rightarrow q_0 = \frac{s + u_1}{2\sqrt{s_4 + m^2}} = (68)$$

$$(95) \wedge (105) \wedge (104) \Rightarrow |\vec{p}_2| = \frac{\sqrt{(s - s_4)^2 - 4m^2 s}}{2\sqrt{s_4 + m^2}} = (71)$$

$$(106)$$

$$(93) \wedge (105) \Rightarrow q_z = \frac{\sqrt{(s' + u_1')^2 - 4q^2t}}{2\sqrt{s_4 + m^2}}$$

$$(107)$$

$$(91) \Rightarrow \omega_2 = \frac{s_4}{2\sqrt{s_4 + m^2}} = (65) \tag{108}$$

$$(90) \Rightarrow E_1 = \frac{s_4 + 2m^2}{2\sqrt{s_4 + m^2}} = (66) \tag{109}$$

$$(93) \Rightarrow \cos \psi = \frac{s(q^2 - t - m^2) - 2q^2(s_4 + m^2) + s_4 u_1}{\sqrt{((s - s_4)^2 - 4m^2 s)((s' + u_1')^2 - 4q^2 t)}} \neq (73)$$
(110)

$$u' = (q - k_2)^2 = q^2 - 2qk_2 \qquad = q^2 - 2(q_0\omega_2 - q_z\omega_2\cos\theta_1$$
 (111)

$$u' = (q - k_2)^2 = q^2 - 2qk_2 = q^2 - 2(q_0\omega_2 - q_z\omega_2\cos\theta_1)$$
(111)

$$u_7 = q^2 - 2qp_1 = q^2 - 2(q_0E_1 + q_z\omega_2\cos\theta_1)$$
(112)

$$0 = s' + u'_1 + u_7 + u'$$
(113)

$$s_3 = 2k_2p_2 = 2\omega_2(E_2 - |\vec{p}_2|(\cos\psi\cos\theta_1 + \sin\psi\sin\theta_1\cos\theta_2)) = (78)$$
 (114)

$$s_5 = (p_1 + p_2)^2 = 2m^2 + 2p_1p_2 \tag{115}$$

$$= 2(m^2 + E_1 E_2 + \omega_2 |\vec{p}_2| (\cos \psi \cos \theta_1 + \sin \psi \sin \theta_1 \cos \theta_2)) = (80)$$
 (116)

$$t' = -2k_1k_2 (118)$$

$$= -2\omega_2(\omega_1 - (|\vec{p}_2|\cos\psi - q_z)\cos\theta_1 - |\vec{p}_2|\sin\psi\sin\theta_1\cos\theta_2)$$
 (119)

$$u_6 = -2k_1 p_1 \tag{120}$$

$$= -2(\omega_1 E_1 + \omega_2(|\vec{p}_2|\cos\psi - q_z)\cos\theta_1 + \omega_2|\vec{p}_2|\sin\psi\sin\theta_1\cos\theta_2)$$
 (121)

same as before: t' is the only variable that can get collinear (for $-q^2 > 0$) - now collinear as [ABC] variable.

2.2.3 Set III

align p_2 to z-axis:

$$q = (q_0, 0, |\vec{q}| \sin \psi, |\vec{q}| \cos \psi, \hat{0})$$
(123)

$$k_1 = (\omega_1, 0, -|\vec{q}| \sin \psi, -|\vec{q}| \cos \psi + p_{2,z}, \hat{0})$$
(124)

$$p_2 = (E_2, 0, 0, p_{2,z}, \hat{0}) \tag{125}$$

constraints, from energy conservation, masses $(k_2^2 = 0)$ is already fixed) and external

Mandelstam variables:

$$q_0 + \omega_1$$
 = $E_1 + E_2 + \omega_2$ (126)
 $m^2 = p_1^2$ = $E_1^2 - \omega_2^2$ (127)

$$m = p_{1} \qquad = E_{1} - \omega_{2} \qquad (127)$$

$$m^{2} = p_{2}^{2} \qquad = E_{2}^{2} - p_{2,z}^{2} \qquad (128)$$

$$q^{2} \qquad = q_{0}^{2} - |\vec{q}|^{2} \qquad (129)$$

$$0 \qquad = \omega_{1}^{2} - |\vec{q}|^{2} + 2|\vec{q}| p_{2,z} \cos \psi - p_{2,z}^{2} \qquad (130)$$

$$s = (q + k_{1})^{2} \qquad = (q_{0} + \omega_{1})^{2} - p_{2,z}^{2} \qquad (131)$$

$$t = (k_{1} - p_{2})^{2} \qquad = (\omega_{1} - E_{2})^{2} - |\vec{q}|^{2} \qquad (132)$$

$$u = (q - p_{2})^{2} \qquad = (q_{0} - E_{2})^{2} - |\vec{q}|^{2} + 2|\vec{q}| p_{2,z} \cos \psi - p_{2,z}^{2} \qquad (133)$$

$$q^2 = q_0^2 - |\vec{q}|^2 \tag{129}$$

$$0 = \omega_1^2 - |\vec{q}|^2 + 2|\vec{q}| p_{2,z} \cos \psi - p_{2,z}^2$$
 (130)

$$s = (q + k_1)^2 \qquad = (q_0 + \omega_1)^2 - p_{2,z}^2 \tag{131}$$

$$t = (k_1 - p_2)^2 \qquad = (\omega_1 - E_2)^2 - |\vec{q}|^2 \tag{132}$$

$$u = (q - p_2)^2 = (q_0 - E_2)^2 - |\vec{q}|^2 + 2|\vec{q}|p_{2,z}\cos\psi - p_{2,z}^2$$
 (133)

solve:

$$(133) - (130) = u = (q_0 - E_2)^2 - \omega_1^2 \tag{134}$$

$$(131) - (129) + (132) - (128) + (134) = s - q^2 + t - m^2 + u$$
(135)

$$= s_4 + m^2 = (\omega_1 + q_0 - E_2)^2 = (100)$$
 (136)

$$(132) + (133) - (129) = t + u - q^2 = 2(E_2 - \omega_1 - q_0)E_2$$
(137)

$$\Rightarrow E_2 = -\frac{t + u - q^2}{2\sqrt{s_4 + m^2}} = \frac{s - s_4 - 2m^2}{2\sqrt{s_4 + m^2}} = (64) = (102)$$

(138)

$$(131) - (129) + (132) - (128) = s - q^2 + t - m^2 = -2(E_2 - \omega_1 - q_0)\omega_1$$
 (139)

$$\Rightarrow \omega_1 = \frac{s' + t_1}{2\sqrt{s_4 + m^2}} = \frac{s_4 - u_1}{2\sqrt{s_4 + m^2}} = (70) = (104)$$

(140)

$$(136) \wedge (138) \wedge (140) \Rightarrow q_0 = \frac{s + u_1}{2\sqrt{s_4 + m^2}} = (68) = (105)$$

$$(141)$$

$$(129) \wedge (141) \Rightarrow |\vec{q}| = \frac{\sqrt{(s' + u_1')^2 - 4q^2t}}{2\sqrt{s_4 + m^2}} = (107)$$

$$(142)$$

$$(127) \wedge (126) \Rightarrow E_1 = \frac{s_4 + 2m^2}{2\sqrt{s_4 + m^2}} = (66) = (109)$$
(143)

$$(126) \Rightarrow \omega_2 = \frac{s_4}{2\sqrt{s_4 + m^2}} = (65) = (108) \tag{144}$$

$$(128) \wedge (138) \Rightarrow p_{2,z} = \frac{\sqrt{(s - s_4)^2 - 4m^2 s}}{2\sqrt{s_4 + m^2}} = (71) = (106)$$
 (145)

$$(129) \Rightarrow \cos \psi = \frac{s(q^2 - t - m^2) - 2q^2(s_4 + m^2) + s_4 u_1}{\sqrt{((s - s_4)^2 - 4m^2 s)((s' + u_1')^2 - 4q^2 t)}}$$

$$= (110) \neq (73)$$

$$s_3 = 2k_2p_2 = 2\omega_2(E_2 - p_{2,z}\cos\theta_1) \tag{147}$$

$$s_5 = (p_1 + p_2)^2 = 2m^2 + 2p_1p_2 \tag{148}$$

$$= 2(m^2 + E_1 E_2 + \omega_2 p_{2,z} \cos \theta_1) \tag{149}$$

$$u' = (q - k_2)^2 = q^2 - 2qk_2 (151)$$

$$= q^2 - 2\omega_2(q_0 - |\vec{q}|(\cos\psi\cos\theta_1 + \sin\psi\sin\theta_1\cos\theta_2))$$
 (152)

$$u_7 = q^2 - 2qp_1 \tag{153}$$

$$= q^{2} - 2(q_{0}E_{1} + |\vec{q}| \omega_{2}(\cos\psi\cos\theta_{1} + \sin\psi\sin\theta_{1}\cos\theta_{2})))$$
 (154)

$$t' = -2k_1k_2 (156)$$

$$= -2\omega_2(\omega_1 + (|\vec{q}|\cos\psi - p_{2,z})\cos\theta_1 + |\vec{q}|\sin\psi\sin\theta_1\cos\theta_2)$$
 (157)

$$u_6 = -2k_1 p_1 \tag{158}$$

$$= -2(\omega_1 E_1 - \omega_2(|\vec{q}|\cos\psi - p_{2,z})\cos\theta_1 - \omega_2|\vec{q}|\sin\psi\sin\theta_1\cos\theta_2)$$
 (159)

$$(39): 0 = s' + t_1 + t' + u_6 \checkmark (160)$$

same as before: t' is the only variable that can get collinear (for $-q^2 > 0$) - now collinear as [ABC] variable.

2.3 phase space integrals

at phase space integration there occur integrations over propagators[4, 2, 3]; the propagators can be decomposed in 2 types: [ab] and [ABC]; the needed integrals then reduce to the master formula:

$$I_n^{(k,l)} = \int_0^{\pi} d\theta_1 \sin^{n-3}(\theta_1) \int_0^{\pi} d\theta_2 \sin^{n-4}(\theta_2) (a + b\cos(\theta_1))^{-k} (A + B\cos(\theta_1) + C\sin(\theta_1)\cos(\theta_2))^{-l}$$

(161)

$$= \int d\Omega_n \left(a + b \cos(\theta_1) \right)^{-k} \left(A + B \cos(\theta_1) + C \sin(\theta_1) \cos(\theta_2) \right)^{-l}$$
(162)

the integrals can be further destinguished by the range of k, l and the type of collinearity (following the notation in [4]):

- "non collinear": $a^2 \neq b^2 \wedge A^2 \neq B^2 + C^2 \rightarrow I_{0,n}^{(k,l)}$
- "single collinear a": $a = -b \wedge A^2 \neq B^2 + C^2 \rightarrow I_{a,n}^{(k,l)}$
- "single collinear A": $a^2 \neq b^2 \wedge A^2 = B^2 + C^2 \rightarrow I_{A,n}^{(k,l)}$

• "double collinear": $a = -b \wedge A = -\sqrt{B^2 + C^2} \rightarrow I_{aA,n}^{(k,l)}$

Use $n = 4 + \epsilon$.

2.3.1 integral helper

define helper integral

$$\hat{I}^{(q)}(\nu) := \int_{0}^{\pi} dt \, \sin^{\nu-3}(t) \cos^{q}(t) \tag{163}$$

It is [5, eq. 5.12.6]:

$$\int_0^{\pi} (\sin t)^{\alpha - 1} e^{i\beta t} dt = \frac{\pi}{2^{\alpha - 1}} \frac{e^{i\pi\beta/2}}{\alpha B ((\alpha + \beta + 1)/2, (\alpha - \beta + 1)/2)} \quad \text{if } \Re(\alpha) > 0 \quad (164)$$

$$\Rightarrow \hat{I}^{(0)}(n) = \frac{\pi}{2^{n-3}(n-2)} \frac{1}{B((n-1)/2, (n-1)/2)}$$
(165)

$$\Rightarrow \hat{I}^{(0)}(n-1) = \frac{\pi}{2^{n-4}(n-3)} \frac{1}{B((n-2)/2, (n-2)/2)} = B((n-3)/2, 1/2)$$
 (166)

If q is odd: $\hat{I}^{(q)} = 0$, due to symetry of kernel; if q is even: q = 2p with $p \in \mathbb{N}$:

$$\hat{I}^{(2p)}(\nu) = \frac{1}{2^{2p}} \sum_{k=0}^{2p} {2p \choose k} \int_{0}^{\pi} \sin^{\nu-3}(t) \exp(2i(k-p)t) dt$$
(167)

$$= \frac{\pi}{2^{2p+\nu-3}(\nu-2)} \sum_{k=0}^{2p} {2p \choose k} \frac{\exp(i\pi(k-p))}{B((\nu-1)/2 + (k-p), (\nu-1)/2 - (k-p))}$$
(168)

$$= \frac{\pi}{2^{2p+\nu-3}(\nu-2)} \sum_{l=-p}^{p} {2p \choose p+l} \frac{(-1)^l}{B((\nu-1)/2+l,(\nu-1)/2-l)}$$
(169)

$$= \frac{\pi\Gamma(\nu-1)(2p)!}{2^{2p+\nu-3}(\nu-2)\Gamma(\frac{n-1}{2}+p)\Gamma(\frac{n-1}{2}+p)} \left(\frac{1}{(p!)^2} \frac{\Gamma(\frac{\nu-1}{2}+p)}{\Gamma(\frac{\nu-1}{2})} \frac{\Gamma(\frac{\nu-1}{2}-p)}{\Gamma(\frac{\nu-1}{2})} + 2\sum_{l=1}^p \frac{(-1)^l}{(p+l)!(p-l)!} \frac{\Gamma(\frac{\nu-1}{2}+p)}{\Gamma(\frac{\nu-1}{2}+l)} \frac{\Gamma(\frac{\nu-1}{2}-p)}{\Gamma(\frac{\nu-1}{2}-l)}\right)$$
(170)

$$= \frac{2^{3-\nu}\pi\Gamma(\nu-1)}{(\nu-2)\Gamma(\frac{n-1}{2}+p)\Gamma(\frac{n-1}{2}+p)} \cdot \frac{\Gamma(\frac{\nu-1}{2}-p)}{2^{p}\Gamma(\frac{\nu-1}{2})} \cdot \frac{(2p)!}{2^{p}p!} \cdot p! \left(\frac{1}{(p!)^{2}} \frac{\Gamma(\frac{\nu-1}{2}+p)}{\Gamma(\frac{\nu-1}{2})} + 2\sum_{l=1}^{p} \frac{(-1)^{l}}{(p+l)!(p-l)!} \frac{\Gamma(\frac{\nu-1}{2}+p)}{\Gamma(\frac{\nu-1}{2}+l)} \frac{\Gamma(\frac{\nu-1}{2})}{\Gamma(\frac{\nu-1}{2}-l)}\right)$$
(171)

TODO: prove

FiXme Error:

$$p! \left(\frac{1}{(p!)^2} \frac{\Gamma(\frac{\nu-1}{2} + p)}{\Gamma(\frac{\nu-1}{2})} + 2 \sum_{l=1}^{p} \frac{(-1)^l}{(p+l)!(p-l)!} \frac{\Gamma(\frac{\nu-1}{2} + p)}{\Gamma(\frac{\nu-1}{2} + l)} \frac{\Gamma(\frac{\nu-1}{2})}{\Gamma(\frac{\nu-1}{2} - l)} \right)$$
(172)

$$= \frac{1}{p!} \frac{\Gamma(-\frac{1}{2} + p)}{\Gamma(-\frac{1}{2})} + 2 \sum_{l=1}^{p} \frac{(-1)^{l} p!}{(p+l)! (p-l)!} \frac{\Gamma(-\frac{1}{2} + p)}{\Gamma(-\frac{1}{2} + l)} \frac{\Gamma(-\frac{1}{2})}{\Gamma(-\frac{1}{2} - l)}$$
(173)

$$=1\tag{174}$$

$$\Rightarrow \hat{I}^{(2p)}(\nu) = \frac{2^{3-\nu}\pi\Gamma(\nu-1)}{(\nu-2)\Gamma(\frac{n-1}{2}+p)\Gamma(\frac{n-1}{2}-p)} \cdot \frac{\Gamma(\frac{\nu-1}{2}-p)}{2^p\Gamma(\frac{\nu-1}{2})} \cdot \frac{(2p!)}{2^pp!}$$
(175)

$$= \frac{\sqrt{\pi}(2p)!}{2^{2p}p!} \frac{\Gamma((\nu-2)/2)}{\Gamma(\frac{\nu-1}{2}+p)}$$
 (176)

2.3.2 any collinearity and $-k, -l \in \mathbb{N}_0$

If $-k, -l \in \mathbb{N}_0$ $I_n^{(k,l)}$ can always be reduced in a straight forward manner to combinations of $\hat{I}^{(q)}(n)$ and this way one finds[4, Ch. 5][2, App. C]:

$$I_n^{(0,0)} = \hat{I}^{(0)}(n-1) \cdot \hat{I}^{(0)}(n) = \frac{2\pi}{n-3}$$
(177)

$$I_4^{(0,0)} = 2\pi (178)$$

$$I_n^{(-1,0)} = \hat{I}^{(0)}(n-1) \cdot (a\hat{I}^{(0)}(n) + b\hat{I}^{(1)}(n)) = \frac{2\pi a}{n-3}$$
(179)

$$I_4^{(-1,0)} = 2\pi a \tag{180}$$

$$I_n^{(0,-1)} = \hat{I}^{(0)}(n-1) \cdot (A\hat{I}^{(0)}(n) + B\hat{I}^{(1)}(n)) + C\hat{I}^{(1)}(n-1)\hat{I}^{(0)}(n)$$
(181)

$$=\frac{2\pi A}{n-3}\tag{182}$$

$$I_4^{(0,-1)} = 2\pi A \tag{183}$$

$$I_n^{(-2,0)} = \hat{I}^{(0)}(n-1) \cdot (a^2 \hat{I}^{(0)}(n) + 2ab\hat{I}^{(1)}(n) + b^2 \hat{I}^{(2)}(n))$$
(184)

$$=2\pi \left(\frac{a^2(n-1)+b^2}{(n-1)(n-3)}\right)$$
 (185)

$$I_4^{(-2,0)} = 2\pi(a^2 + b^2/3) \tag{186}$$

$$I_n^{(0,-2)} = \hat{I}^{(0)}(n-1) \cdot (A^2 \hat{I}^{(0)}(n) + B^2 \hat{I}^{(2)}(n)) + C^2 \hat{I}^{(2)}(n-1) \hat{I}^{(0)}(n+2)$$
(187)

$$=2\pi \left(\frac{A^2(n-1)+B^2+C^2}{(n-1)(n-3)}\right)$$
(188)

$$I_4^{(0,-2)} = 2\pi (A^2 + (B^2 + C^2)/3)$$
(189)

$$I_n^{(-1,-1)} = \hat{I}^{(0)}(n-1) \cdot (aA\hat{I}^{(0)}(n) + bB\hat{I}^{(2)}(n)) = 2\pi \left(\frac{aA(n-1) + bB}{(n-1)(n-3)}\right)$$
(190)

$$I_4^{(-1,-1)} = 2\pi(aA + bB/3) \tag{191}$$

2.3.3 single collinear a

If $-l \in \mathcal{N}$ one finds:

$$\hat{I}_a^{(k,q)}(\nu) = \int_0^\pi \frac{\sin^{\nu-3} t}{(1 - \cos(t))^k} \cos^q(t) dt$$
 (192)

$$= \int_{0}^{\pi} \frac{\sin^{\nu-3}(t)}{(1-\cos^{2}(t))^{k}} \cos^{q}(t) (1+\cos(t))^{k} dt$$
 (193)

$$= \int_{0}^{\pi} \sin^{\nu - 3 - 2k}(t) \cos^{q}(t) (1 + \cos(t))^{k} dt$$
 (194)

$$= \sum_{l=0}^{k} {k \choose l} \hat{I}^{(q+l)} (\nu - 2k)$$
 (195)

this way one finds[4, Ch. 5][2, App. C]:

$$I_{a,n}^{(1,0)} = \frac{1}{a}\hat{I}^{(0)}(n-1)\cdot\hat{I}^{(0)}(n-2)$$
(196)

$$=\frac{2\pi}{a(n-4)}\tag{197}$$

$$I_{a,n}^{(1,-1)} = \frac{1}{a}\hat{I}^{(0)}(n-1) \cdot \left(A\hat{I}^{(0)}(n-2) + B\hat{I}^{(2)}(n-2)\right)$$
(198)

$$= \frac{2\pi}{a} \frac{(A(n-3)+B)}{(n-3)(n-4)} \approx \frac{2\pi}{a} \left(\frac{A+B}{\epsilon} - 2B + O(\epsilon) \right)$$
(199)

$$I_{a,n}^{(1,-2)} = \frac{1}{a} \left(\hat{I}^{(0)}(n-1) \cdot \left(A^2 \hat{I}^{(0)}(n-2) + (B^2 + 2AB) \hat{I}^{(2)}(n-2) \right) + C^2 \hat{I}^{(2)}(n-1) \hat{I}^{(0)}(n) \right)$$
(200)

$$= \frac{2\pi}{a} \left(\frac{A^2}{n-4} + \frac{2AB + B^2}{(n-4)(n-3)} + \frac{C^2}{(n-3)(n-2)} \right)$$
 (201)

$$\approx \frac{2\pi}{a} \left(\frac{(A+B)^2}{\epsilon} + \frac{C^2}{2} - 2AB - B^2 + O(\epsilon) \right)$$
 (202)

$$I_{a,n}^{(2,0)} = \frac{1}{a^2} \hat{I}^{(0)}(n-1) \cdot \left(\hat{I}^{(0)}(n-4) + \hat{I}^{(2)}(n-4)\right)$$
(203)

$$=\frac{2\pi}{a^2(n-6)} \approx -\frac{\pi}{a^2} + O(\epsilon) \tag{204}$$

$$I_{a,n}^{(2,-1)} = \frac{1}{a^2} \hat{I}^{(0)}(n-1) \cdot \left(A \left(\hat{I}^{(0)}(n-4) + \hat{I}^{(2)}(n-4) \right) + 2B \hat{I}^{(2)}(n-4) \right)$$
 (205)

$$= \frac{2\pi}{a^2} \left(\frac{A}{n-6} + \frac{2B}{(n-6)(n-4)} \right) \approx -\frac{2\pi}{a^2} \left(\frac{B}{\epsilon} + \frac{A+B}{2} \right) + O(\epsilon)$$
 (206)

$$I_{a,n}^{(2,-2)} = \frac{1}{a^2} \left(\hat{I}^{(0)}(n-1) \cdot \left(A^2 (\hat{I}^{(0)}(n-4) + \hat{I}^{(2)}(n-4)) + 4AB\hat{I}^{(2)}(n-4) \right) \right)$$

$$+B^{2}(\hat{I}^{(2)}(n-4)+\hat{I}^{(4)}(n-4)))+C^{2}\hat{I}^{(2)}(n-1)(\hat{I}^{(0)}(n-2)+\hat{I}^{(2)}(n-2)))$$
(207)

$$= \frac{2\pi}{a^2} \left(\frac{A^2}{n-6} + \frac{4AB}{(n-6)(n-4)} + \frac{B^2n}{(n-6)(n-4)(n-3)} + \frac{C^2}{(n-4)(n-3)} \right)$$
(208)

$$\approx \frac{2\pi}{a^2} \left(\frac{-2AB - 2B^2 + C^2}{\epsilon} + \frac{B^2 - A^2}{2} - AB - C^2 + O(\epsilon) \right)$$
 (209)

(210)

It is [4, Ch. 5]:

$$I_{a,n}^{(1,1)} = \frac{\pi}{a(A+B)} \left(\frac{2}{\epsilon} + \ln\left(\frac{(A+B)^2}{A^2 - B^2 - C^2}\right) \right) + O(\epsilon)$$
 (211)

$$I_{a,n}^{(2,1)} = \frac{\pi}{a^2(A+B)} \left(\frac{B^2 + AB + C^2}{(A+B)^2} \left(\frac{2}{\epsilon} + \ln\left(\frac{(A+B)^2}{A^2 - B^2 - C^2}\right) \right) \right)$$

$$-\frac{2C^2}{\left(A+B\right)^2} - 1 + O(\epsilon) \tag{212}$$

From this are the following integrals derived[4, Ch. 5]:

$$I_{a,n}^{(1,2)} = -\frac{\partial}{\partial A} I_{a,n}^{(1,1)} \tag{213}$$

$$= \frac{\pi}{a(A+B)^2} \left(\frac{2}{\epsilon} + \ln\left(\frac{(A+B)^2}{A^2 - B^2 - C^2}\right) + \frac{2A(A+B)}{A^2 - B^2 - C^2} - 2 \right) + O(\epsilon)$$
 (214)

$$I_{a,n}^{(2,2)} = -\frac{\partial}{\partial A} I_{a,n}^{(2,1)} \tag{215}$$

$$= \frac{\pi}{a^{2}(A+B)^{2}} \left(\frac{2B^{2} + 2AB + 3C^{2}}{(A+B)^{2}} \left(\frac{2}{\epsilon} + \ln\left(\frac{(A+B)^{2}}{A^{2} - B^{2} - C^{2}}\right) \right) + \frac{2A^{2}}{A^{2} - B^{2} - C^{2}} - \frac{8C^{2}}{(A+B)^{2}} - 3 \right) + O(\epsilon)$$
(216)

If $-k \in \mathcal{N}$ use I_0 with b = -a.

2.3.4 double collinear

as said in [4, Ch. 5]: if $0 \le -\frac{C}{A}, \frac{B}{A} \le 1$ use [3, eq. A11] with $\cos \kappa = -\frac{B}{A}$:

$$I_{aA,n}^{(k,l)} = \frac{2\pi 2^{-(k+l)}}{a^k A^l} \frac{\Gamma(1+\epsilon)}{\Gamma^2(1+\epsilon/2)} B(1+\frac{\epsilon}{2}-k, 1+\frac{\epsilon}{2}-l)_2 F_1\left(k, l; 1+\frac{\epsilon}{2}; \frac{A-B}{2A}\right)$$
(217)

but we will not need it here.

2.3.5 non collinear

If $-l \in \mathcal{N}$ the θ_2 integration can be performed using the integral helper and the problem reduces then to the following integral:

$$\hat{I}_{0}^{(k,q,p)}(\epsilon) = \int_{0}^{\pi} d\theta_{1} \frac{\sin^{1+\epsilon}(\theta_{1}) \sin^{q}(\theta_{1}) \cos^{p}(\theta_{1})}{(a+b \cos(\theta_{1}))^{k}}$$

$$= \frac{1}{2a^{k}} \left((1+(-1)^{p})B\left(\frac{2+q+\epsilon}{2}, \frac{1+p}{2}\right) {}_{3}F_{2}\left(\frac{1+k}{2}, \frac{k}{2}, \frac{1+p}{2}; \frac{1}{2}, \frac{3+q+p+\epsilon}{2}; \frac{b^{2}}{a^{2}}\right) \right)$$

$$\frac{b}{a}k(-1+(-1)^{p})B\left(\frac{2+q+\epsilon}{2}, \frac{2+p}{2}\right) {}_{3}F_{2}\left(\frac{1+k}{2}, \frac{2+k}{2}, \frac{2+p}{2}; \frac{3}{2}, \frac{4+q+p+\epsilon}{2}; \frac{b^{2}}{a^{2}}\right) \right)$$
(219)

for k = 1 this simplifies to

$$\hat{I}_{0}^{(1,q,p)}(\epsilon) = \int_{0}^{\pi} d\theta_{1} \frac{\sin^{1+\epsilon}(\theta_{1}) \sin^{q}(\theta_{1}) \cos^{p}(\theta_{1})}{(a+b \cos(\theta_{1}))}$$

$$= \frac{1}{2a} \left((1+(-1)^{p}) B\left(\frac{2+q+\epsilon}{2}, \frac{1+p}{2}\right) {}_{2}F_{1}\left(1, \frac{1+p}{2}; \frac{3+q+p+\epsilon}{2}; \frac{b^{2}}{a^{2}}\right) \right)$$

$$\frac{b}{a} (-1+(-1)^{p}) B\left(\frac{2+q+\epsilon}{2}, \frac{2+p}{2}\right) {}_{2}F_{1}\left(1, \frac{2+p}{2}; \frac{4+q+p+\epsilon}{2}; \frac{b^{2}}{a^{2}}\right)$$
(221)

TODO: structure does NOT match [4, Ch. 5] - neither A,B,C nor ϵ

FiXme Error: fix to Bojak For $I_{0,n}^{(1,1)}$ one finds[4, Ch. 5]:

$$I_{0,n}^{(1,1)} = \frac{\pi}{\sqrt{X}} \ln \left(\frac{aA - bB + \sqrt{X}}{aA - bC - \sqrt{X}} \right)$$
 (222)

with
$$X = (aB - bA)^2 - (a^2 - b^2)C^2$$
 (223)

 $I_{0,n}^{(1,-3)}$ is given by [4, Ch. 5] and from those two all other integrals can be deduced using the techiques described in [4, Ch. 5]: increase k or l by differentiation or interchange k and l by a rotation:

$$I_n^{(k,l)} = I_n^{(l,k)} \left(a \leftrightarrow A, b \to -\sqrt{B^2 + C^2}, B \to \frac{-bB}{\sqrt{B^2 + C^2}}, C \to \frac{-bC}{\sqrt{B^2 + C^2}} \right)$$
 (224)

2.4 needed set in matrix element

define a shortcut:

$$\mathcal{V}_{a,b}(x,y) = \left(x^k y^l\right)_{k=0..a,l=0..b}$$
(225)

It is

$$A_{G,1} = \sum_{k,l=0}^{3} \left(\mathcal{C}_{A_{G,1}} \right)_{(k,l)} t'^{-2+k} u_7^{-2+l} = \operatorname{tr} \left(\mathcal{C}_{A_{G,1}} \frac{\mathcal{V}_{3,3}(t', u_7)^t}{t'^2 u_7^2} \right)$$
(226)

$$A_{L,1} = \sum_{k=0}^{4} \sum_{l=0}^{2} \left(\mathcal{C}_{A_{L,1}} \right)_{(k,l)} t'^{-2+k} u_{7}^{-2+l} \qquad = \operatorname{tr} \left(\mathcal{C}_{A_{L,1}} \frac{\mathcal{V}_{4,2}(t', u_{7})^{t}}{t'^{2} u_{7}^{2}} \right)$$
(227)

and we will thus need the integrals

$$\left(\mathcal{I}_{A_{G,1}} \right)_{(k,l)} = \int d\Omega_n \, \frac{1}{t'^2 u_7^2} \left(\mathcal{V}_{3,3}(t', u_7) \right)_{(k,l)}$$
 (228)

$$\left(\mathcal{I}_{A_{L,1}}\right)_{(k,l)} = \int d\Omega_n \, \frac{1}{t'^2 u_7^2} \left(\mathcal{V}_{4,2}(t', u_7)\right)_{(k,l)} \tag{229}$$

with

$$a(t') = -b(t') = -2\omega_1 \omega_2 \qquad = -\frac{s_4(s' + t_1)}{2(s_4 + m^2)}$$
(230)

$$A(u_7) = q^2 - 2q_0 E_1$$

$$2(s_4 + m^2)$$

$$= q^2 - \frac{(s_4 + 2m^2)(s + u_1)}{2(s_4 + m^2)}$$
(231)

$$B(u_7) = -2\omega_2(|\vec{p}_2|\cos\psi - \omega_1) \qquad = \frac{s_4}{2} \left(1 - \frac{s + u_1}{s_4 + m^2} + \frac{s' - t_1}{s' + t_1} \right)$$
(232)

$$C(u_7) = -2\omega_2 |\vec{p}_2| \sin \psi \tag{233}$$

that is $I_{a,n}^{(-2...2,-1...2)}$. With this we find

$$A + B = -\frac{t_1 u_1}{s' + t_1} \tag{234}$$

$$\frac{(A+B)^2}{A^2 - B^2 - C^2} = \frac{(s_4 + m^2)t_1^2 u_1^2}{(s'+t_1)^2 (s_4 q^2 t_1 + m^2 (s'+u_1)^2)}$$
(235)

$$2B(A+B) + 3C^{2} = -\frac{s_{4}(m^{2}s'(3s's_{4} + 2t_{1}u_{1}) + t_{1}(q^{2}(s_{4} - u_{1})(3s_{4} - u_{1}) - u_{1}(s's_{4} + t_{1}u_{1})))}{(s_{4} + m^{2})(s' + t_{1})^{2}}$$
(236)

With this we get

$$\int d\Omega_n A_{j,1} = \operatorname{tr}(\mathcal{C}_{A_{j,1}} \left(\mathcal{I}_{A_{j,1}} \right)^t) \qquad j = G, L$$
(237)

It is

$$A_{G,2} = \sum_{k,l=0}^{3} \left(\mathcal{C}_{A_{G,2}} \right)_{(k,l)} s_5^{-2+k} u'^{-2+l} = \operatorname{tr} \left(\mathcal{C}_{A_{G,2}} \frac{\mathcal{V}_{3,3}(s_5, u')^t}{s_5^2 u'^2} \right)$$
(238)

$$A_{L,2} = \sum_{k=0}^{4} \sum_{l=0}^{3} \left(\mathcal{C}_{A_{G,2}} \right)_{(k,l)} s_5^{-2+k} u'^{-2+l} \qquad = \operatorname{tr} \left(\mathcal{C}_{A_{L,2}} \frac{\mathcal{V}_{4,3}(s_5, u')^t}{s_5^2 u'^2} \right)$$
(239)

and we will thus need the integrals

$$\left(\mathcal{I}_{A_{G,2}}\right)_{(k,l)} = \int d\Omega_n \, \frac{1}{s_5^2 u'^2} \left(\mathcal{V}_{3,3}(s_5, u')\right)_{(k,l)} \tag{240}$$

$$\left(\mathcal{I}_{A_{L,2}}\right)_{(k,l)} = \int d\Omega_n \, \frac{1}{s_5^2 u'^2} \left(\mathcal{V}_{4,3}(s_5, u')\right)_{(k,l)} \tag{241}$$

but as both s_5 and u' are of [ABC] type we have to apply partial fractioning, following the ideas of [4, Ch. 5]. It is

$$(44): s_5 = s - q^2 + t' + u' (242)$$

so end up with a form of $\mathcal{V}(p+q,q)$ where p is [ab] and q and p+q are [ABC]. The aim is then to get to a form with fractions of $\frac{p}{q}$ and/or $\frac{p+q}{p}$ and indeed this can be achieved. [TODO]

FiXme Error: rewrite

A References

- [1] M. Stratmann. Private communications, 2015.
- [2] W. Beenakker, H. Kuijf, W. L. van Neerven, and J. Smith, "QCD Corrections to Heavy Quark Production in p anti-p Collisions," Phys. Rev. **D40** (1989) 54–82.
- [3] W. L. Van Neerven, "Dimensional regularization of mass and infrared singularities in two-loop on-shell vertex functions," <u>Nuclear Physics B</u> **268** no. 2, (May, 1986) 453–488. http://www.sciencedirect.com/science/article/pii/0550321386901653.
- [4] I. Bojak,

 NLO QCD corrections to the polarized photoproduction and hadroproduction of heavy quarks.

 PhD thesis, Dortmund U., 2000. arXiv:hep-ph/0005120 [hep-ph].
- [5] "NIST Digital Library of Mathematical Functions." Http://dlmf.nist.gov/, release 1.0.10 of 2015-08-07. http://dlmf.nist.gov/. Online companion to [6].
- [6] F. W. J. Olver, D. W. Lozier, R. F. Boisvert, and C. W. Clark, eds., NIST Handbook of Mathematical Functions. Cambridge University Press, New York, NY, 2010. Print companion to [5].

List of Corrections

Error:	prove			 															1	3
Error:	fix to Bojak			 															1	6
Error:	rewrite	_		 							_			_			_		1	۶