

1 Introduction

This work is mainly based on the paper “Complete $O(\alpha_S)$ corrections to heavy-flavour structure functions in electroproduction” by Laenen et. al.[1] - that is, it recalculates all properties and formulas. It extends then the application to the equivalent *polarized* processes. The treating of the polarized processes can for example be found in [2] and we will use many ideas and technics from there. **FiXme Error: more**

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1.1 Motivation

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1.2 Notation

To collect all soft and collinear poles we have to calculate in $n = 4 + \epsilon$ dimension. Unfortunaly the extension for *polarized* processes is nontrivial, because the occuring Levi-Civita tensors $\varepsilon_{\mu\nu\rho\sigma}$ and γ_5 . A common choice to deal with these objects is the HVBM prescription[3] that keeps those two objects four dimensional at the price for splitting the full n dimensional space into a $n - 4$ dimensional space, called “hat-space”, and a four dimensional space (that is actually never used).

In leading order (LO) we have to consider the following processes

$$\gamma^*(q; \sigma_q) + g(k_1; \sigma_{k_1}) \rightarrow Q(p_1) + \bar{Q}(p_2) \quad (1)$$

The corresponding parton structure tensor $W_{\mu\mu'}^{(0)}$, can then be written as **FiXme Error: avoid all order expr?**

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$$\begin{aligned} & W_{\mu\mu'}^{(0)}(k_1, q; s, t_1, u_1, q^2; \sigma_{k_1} \sigma_q) \\ &= \frac{1}{2} E_\epsilon K_{\gamma g} \int \frac{d^{n-1} p_1}{2E_1 (2\pi)^{n-1}} \int \frac{d^{n-1} p_2}{2E_2 (2\pi)^{n-1}} \delta(p_1^2 - m^2) \delta(p_2^2 - m^2) \\ & \quad (2\pi)^n \delta^{(n)}(k_1 + q - p_1 - p_2) \mathcal{M}_\mu^{(0)}(\sigma_{k_1}) \mathcal{M}_{\mu'}^{(0)}(\sigma_q) \end{aligned} \quad (2)$$

where the initial $1/2$ is the initial state spin average, $K_{\gamma g}$ is the color average,

$$E_\epsilon := \begin{cases} 1/(1 + \epsilon/2) & \text{unpolarized} \\ 1 & \text{polarized} \end{cases} \quad (3)$$

accounts for initial freedom in n dimensions for bosons and we defined the following Mandelstam variables:

$$s = (q + k_1)^2, \quad t_1 = t - m^2 = (k_1 - p_2)^2 - m^2, \quad u_1 = u - m^2 = (q - p_2)^2 - m^2 \quad (4)$$

$$s' = s - q^2, \quad u'_1 = u_1 - q^2 \quad (5)$$

FiXme Error: move to LO? The Lorentz indices μ and μ' refer to the virtual photon that is exchanged with the scattering lepton.

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By using Lorentz covariance, hermiticity, parity invariance and current conservation the parton structure tensor can be decomposed into several parts:

$$\begin{aligned} W_{\mu\mu'}^{(0)}(k_1, q; s, t_1, u_1, q^2; \sigma_{k_1}, \sigma_q) = & \left(-g_{\mu\mu'} + \frac{q_\mu q_{\mu'}}{q^2} \right) \frac{d^2 \sigma_T(s, t_1, u_1, q^2)}{dt_1 du_1} \\ & + \left(k_{1,\mu} - \frac{k_1 \cdot q}{q^2} q_\mu \right) \left(k_{1,\mu'} - \frac{k_1 \cdot q}{q^2} q_{\mu'} \right) \left(\frac{-4q^2}{s'^2} \right) \\ & \cdot \left(\frac{d^2 \sigma_T(s, t_1, u_1, q^2)}{dt_1 du_1} + \frac{d^2 \sigma_L(s, t_1, u_1, q^2)}{dt_1 du_1} \right) \end{aligned} \quad (6)$$

FiXme Error: extend We can then define appropriate projection operators[4]:

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$$\mathcal{P}_{G,\mu\mu'} = -g_{\mu\mu'} \quad b_G(\epsilon) = \frac{1}{2(1 + \epsilon/2)} \quad (7)$$

$$\mathcal{P}_{L,\mu\mu'} = -\frac{4q^2}{s'^2} k_{1,\mu} k_{1,\mu'} \quad b_L(\epsilon) = 1 \quad (8)$$

$$\mathcal{P}_{P,\mu\mu'} = i\varepsilon_{\mu\mu'\rho\rho'} \frac{q^\rho k_1^{\rho'}}{s'} \quad b_P(\epsilon) = 1 \quad (9)$$

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$$\frac{d^2 \sigma_k(s, t_1, u_1, q^2)}{dt_1 tu_1} = b_k(\epsilon) \mathcal{P}_{k,\mu\mu'} W^{\mu\mu'} \quad (10)$$

with $k \in \{G, L, P\}$ denoting (here and mostly ever after) the projection type. The transverse partonic cross section $d\sigma_T$ can be reconstructed from the above definitions by using

$$\frac{d^2 \sigma_T}{dt_1 du_1} = \frac{d^2 \sigma_G}{dt_1 du_1} + b_G(\epsilon) \frac{d^2 \sigma_L}{dt_1 du_1} \quad (11)$$

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FiXme Error: explain ghosts?

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2 Leading Order Calculations

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3 Next-To-Leading Order Calculations

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4 Mass Factorization

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5 Partonic Results

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6 Hadronic Results

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7 Summary

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A References

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