

Next-to-Leading Order QCD Corrections to Inclusive Heavy-Flavor Production in Polarized Deep-Inelastic Scattering

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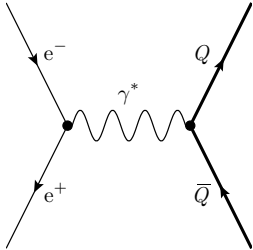
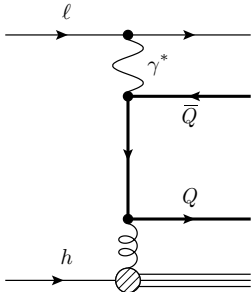
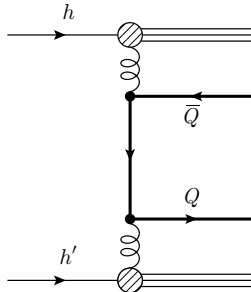
Introduction - Heavy Quarks (HQ)

- Heavy Quarks (HQ): $c(m_c = 1.5 \text{ GeV})$, $b(m_b = 4.75 \text{ GeV})$, $t(m_t = 175 \text{ GeV})$
- EIC will reach region with HQ relevant to structure functions
- compare unpolarized case @HERA: at small $x \sim 30\%$ charm contributions
- dominated by PGF \rightarrow measure Δg
- scheme for massive quarks in polarized PDFs?
- first NLO computation of process
- need improved charm tagging
- full inclusive cross section is complicated to reconstruct
- no hadronization here

- scale of hard process in a perturbative regime
 $m > \Lambda_{QCD}$
- finite mass m^2 ensures full inclusive cross sections
- full m^2 dependence makes computations complicated: phase space + matrix elements
- 2-scale problem: encounter $\ln\left(\frac{s-4m^2}{4m^2}\right)$ and/or $\ln(Q^2/m^2)$
- keep analytic expressions



Introduction - Experimental Setups

e^-e^+ -annihilation (SIA)	deep inelastic scattering (DIS)	Drell-Yan process (DY)
$e^- + e^+ \rightarrow \bar{Q} + X[Q]$	$\ell + h \rightarrow \ell' + \bar{Q} + X[Q]$	$h + h' \rightarrow \bar{Q} + X[Q]$
		
LEP, ILC	HERA, COMPASS, EIC	Tevatron, LHC
gluon	factorization	top, Higgs

cross section (xs):
$$\frac{d^2\sigma}{dxdy} = \frac{2\pi y\alpha^2}{Q^4} L^{\mu\nu} W_{\mu\nu} \quad (1)$$

hadronic tensor:
$$W_{\mu\nu} = \left(-g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) F_1(x, Q^2) + \frac{P_\mu P_\nu}{P \cdot q} F_2(x, Q^2) + i\epsilon_{\mu\nu\alpha\beta} \frac{q^\alpha S^\beta}{P \cdot q} g_1(x, Q^2) \quad (2)$$

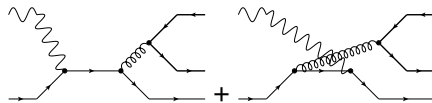
$$F_L(x, Q^2) = F_2(x, Q^2) - 2xF_1(x, Q^2) \quad (3)$$

unpolarized xs:
$$\frac{d^2\sigma}{dxdy} = \frac{2\pi\alpha^2}{xyQ^2} \left(Y_+ F_2(x, Q^2) - y^2 F_L(x, Q^2) \right) \quad (4)$$

polarized xs:
$$\frac{d^2\Delta\sigma}{dxdy} = \frac{4\pi\alpha^2}{xyQ^2} Y_- \cdot 2xg_1(x, Q^2) \quad (5)$$

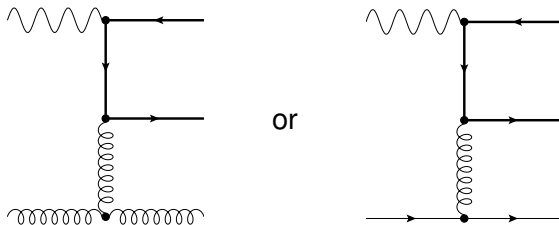
$$Y_\pm = 1 \pm (1 - y)^2 \quad (6)$$

- use factorisation theorem: $s \rightarrow \xi S_h + \text{PDF}$
- $g(k_1) + \gamma^*(q) \rightarrow \bar{Q}(p_2) + Q(p_1)$
- three massive particles: $2 \cdot m^2 > 0, q^2 = -Q^2 < 0$
- compute 2-to-3-phase space: e.g. $dPS_3 \sim dt_1 ds_4 d\Omega_n$



- compute diagrams:
- $\Rightarrow 2xg_1(x) \sim e_u^2 \cdot \xi \Delta u(\xi) \otimes d_{P,q}^{(1)}(\chi, \chi')$
- $d_{P,q}^{(1)}(\chi, \chi') = c_1(\chi, \chi') \ln(\chi) + c_2(\chi, \chi') \text{Li}_2\left(\frac{1+\chi'}{1+\chi}\right) + \dots \checkmark$
- $\frac{m^2}{s} = \frac{\chi}{(1+\chi)^2}$ and $\frac{m^2}{s+Q^2} = \frac{m^2}{s'} = \frac{\chi'}{(1+\chi')^2}$ and $\frac{m^2}{Q^2} = \frac{\chi_q}{(1-\chi_q)^2}$

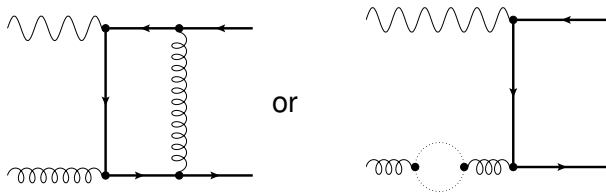
collinear poles appear in, e.g.,



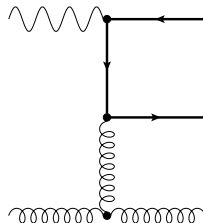
- remove by mass factorization $\rightarrow \overline{\text{MS}}$
- $\Rightarrow 2xg_1(x) \sim e_H^2 \cdot \xi \Delta g(\xi) \otimes \ln(\mu_F^2/m^2) \bar{c}_{P,g}^{F,(1)}(\chi, \chi_q)$
- $\bar{c}_{P,g}^{F,(1)}(\chi, \chi_q) = c_1(\chi, \chi_q) \ln(\chi) + c_2(\chi, \chi_q) \text{Li}_2\left(\frac{1-\chi_q}{1+\chi}\right) + \dots$ (\checkmark for $Q^2 \gg m^2$)

Computation Review - UV and IR Poles

virtual diagrams are, e.g.,



soft poles appear in the limit of a soft gluon $k_2 \rightarrow 0$, e.g.,

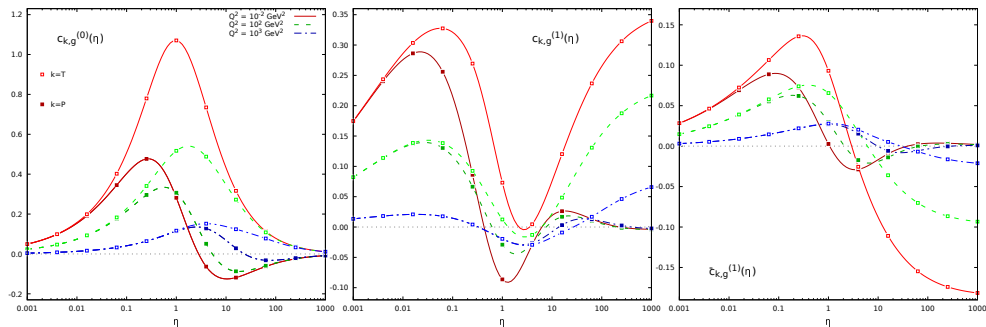


soft + virtual + renormalization ($\overline{\text{MS}}_m$) + factorization is finite!

$$\begin{aligned}
 D_0(m^2, 0, q^2, m^2, t, s, 0, m^2, m^2, m^2) &= \frac{iC_\epsilon}{\beta s t_1} \times \left[-\frac{2}{\epsilon} \ln(\chi) - 2 \ln(\chi) \ln\left(\frac{-t_1}{m^2}\right) \right. \\
 &+ \text{Li}_2(1 - \chi^2) - 4\zeta(2) + \ln^2(\chi_q) + 2 \text{Li}_2(-\chi\chi_q) + 2 \text{Li}_2\left(\frac{-\chi}{\chi_q}\right) \\
 &+ 2 \ln(\chi\chi_q) \ln(1 + \chi\chi_q) + 2 \ln\left(\frac{\chi}{\chi_q}\right) \ln\left(1 + \frac{\chi}{\chi_q}\right) \left. \right] \\
 \int \frac{d\Omega_n}{t' u_7^2} &\sim -\frac{2\pi(m^2 + s_4)(s' + t_1)}{s_4 t_1^2 u_1^2} \left[-2 + \frac{t_1 u_1 (-q^2 s_4 + (2m^2 + s_4)(s' + u_1))}{(s' + t_1) (q^2 s_4 t_1 + m^2 (s' + u_1)^2)} \right. \\
 &+ \frac{2}{\epsilon} + \ln\left(\frac{t_1^2 u_1^2 (m^2 + s_4)}{(s' + t_1)^2 (m^2 (s' + u_1)^2 + q^2 t_1 s_4)}\right) \left. \right]
 \end{aligned}$$

OOO, I'VE THOUGHT OF A NEW ONE!
TWO SQUIGGLES AND A BACKWARDS G!

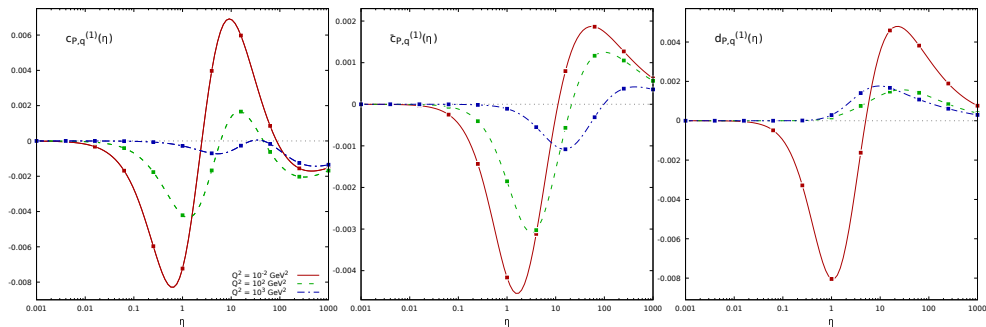
$$2xg_1(x) \sim \xi \Delta g(\xi) \otimes \left(c_{P,g}^{(0)} + 4\pi\alpha_s \left[c_{P,g}^{(1)} + \ln\left(\frac{\mu^2}{m^2}\right) \bar{c}_{P,g}^{(1)} \right] \right) \quad (7)$$



$$\eta = \frac{s-4m^2}{4m^2}, \quad m = m_b = 4.75 \text{ GeV}$$

Partonic Results - Light Quark Channel

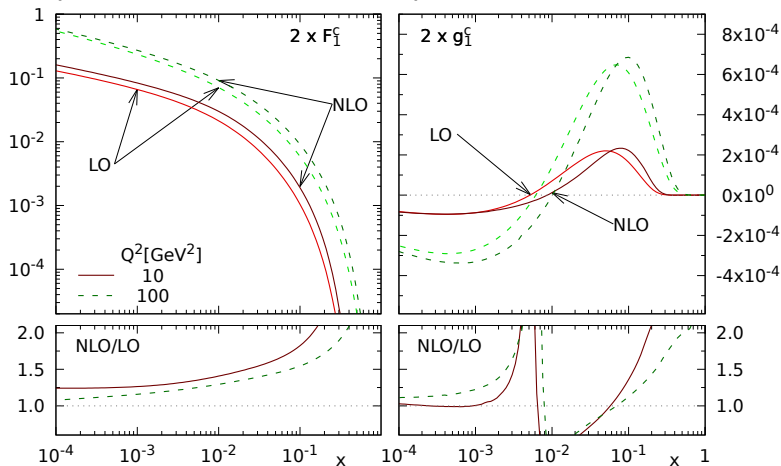
$$2xg_1(x) \sim \sum_{q \in \{u,d,s\}} \xi(\Delta q(\xi) + \Delta \bar{q}(\xi)) \otimes \left(e_H^2 \left[c_{P,q}^{(1)} + \ln \left(\frac{\mu^2}{m^2} \right) \bar{c}_{P,q}^{(1)} \right] + e_q^2 d_{P,q}^{(1)} \right) \quad (8)$$



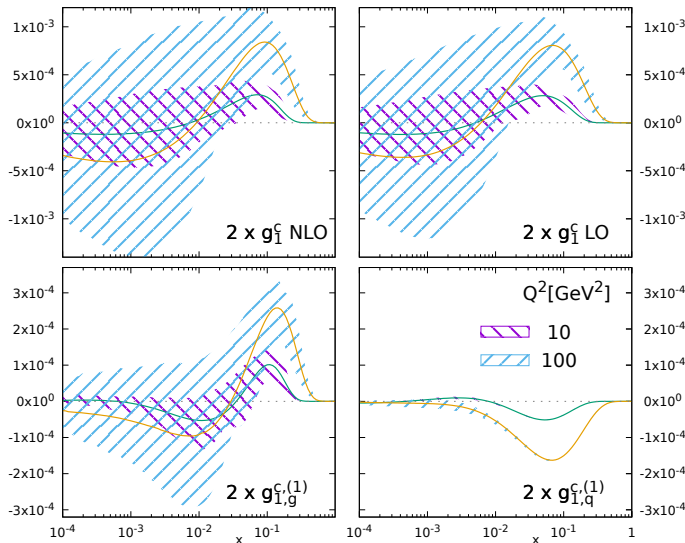
$$\eta = \frac{s-4m^2}{4m^2}, \quad m = m_b = 4.75 \text{ GeV}$$

Hadronic Results - Unpolarized vs. Polarized

unpolarized \sim MSTW2008 \leftrightarrow polarized \sim DSSV2014

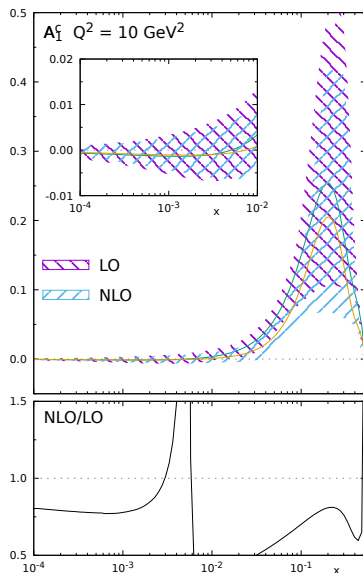


Hadronic Results - PDF Uncertainties DSSV (I)



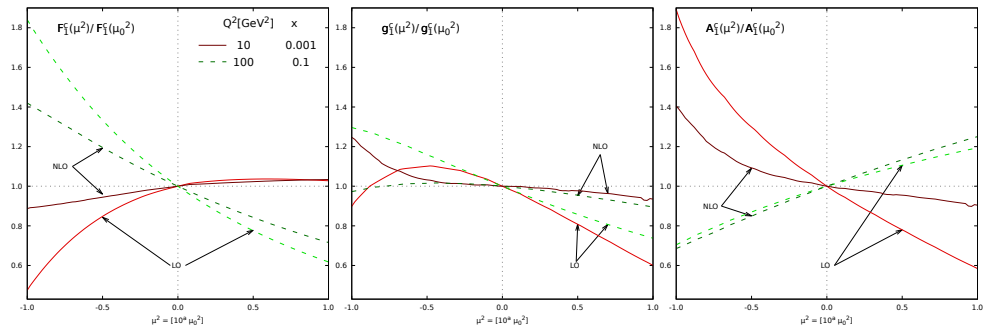
- gluon dominates
- quarks only make small contributions
- sign unclear
- quark band small
- default gluon is small

Hadronic Results - PDF Uncertainties DSSV (II)



- $A_1^c(x, Q^2) = \frac{g_1^c(x, Q^2)}{F_1^c(x, Q^2)}$
- error band are only due to DSSV uncertainties (no correlations!)
- sign unclear
- need measurement of $\mathcal{O}(10^{-3})$
- $NLO \approx LO$

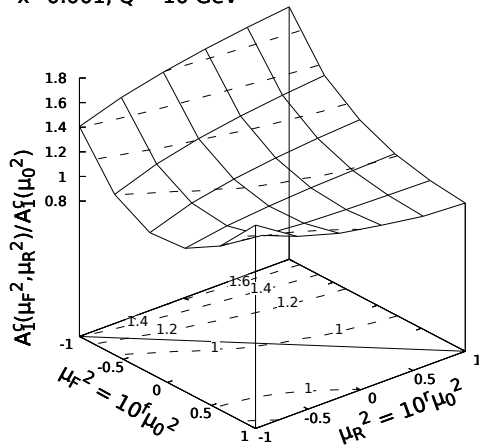
Hadronic Results - Scale Uncertainties (I)



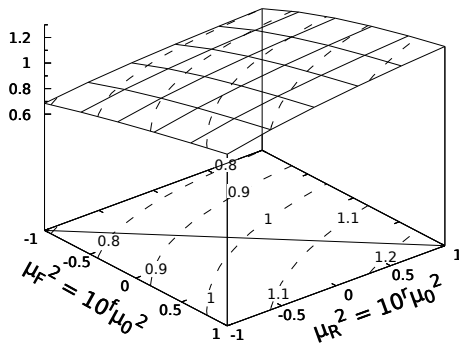
$$\mu^2 = 10^a \mu_0^2 \text{ with } \mu_0^2 = 4m^2 + Q^2 \quad (9)$$

Hadronic Results - Scale Uncertainties (II)

$x=0.001, Q^2=10 \text{ GeV}^2$



$x=0.1, Q^2=100 \text{ GeV}^2$



- inclusive distributions: $\frac{dg_1}{dp_{T,\bar{Q}}}, \frac{dg_1}{dy_{\bar{Q}}}$
- correlated distributions: $\frac{dg_1}{dM_{Q\bar{Q}}^2}, \frac{dg_1}{d\phi_{Q\bar{Q}}}$
- full neutral current (NC) contributions: $F_3^{Z\gamma}, g_4^{Z\gamma}, g_5^{Z\gamma}$ and F_2^Z, F_L^Z, g_1^Z
- distributions of full NC structure functions: $\frac{dg_1^{NC}}{dp_{T,\bar{Q}}}, \frac{dg_1^{NC}}{dM_{Q\bar{Q}}^2}$