

# Next-to-Leading Order QCD Corrections to Inclusive Heavy-Flavor Production in Polarized Deep-Inelastic Scattering

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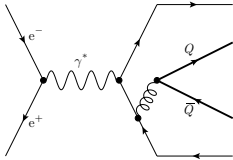
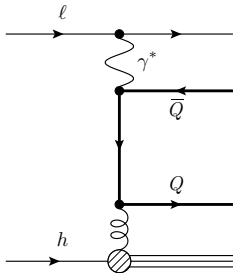
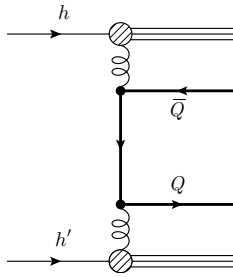
# Outline

- 1 Introduction
- 2 Computation Review
- 3 Partonic Results
- 4 Hadronic Results
- 5 Outlook

# Introduction - Heavy Quarks

HQ are good

# Introduction - Experimental Setups

$e^-e^+$ -annihilation (SIA)	deep inelastic scattering (DIS)	Drell-Yan process (DY)
$e^- + e^+ \rightarrow \bar{Q} + X[Q]$	$\ell + h \rightarrow \bar{Q} + X[Q]$	$h + h' \rightarrow \bar{Q} + X[Q]$
		
LEP, ILC	HERA, COMPASS, EIC	Tevatron, LHC
gluon	factorization	top, Higgs

# Introduction - Structure Functions

cross section: 
$$\frac{d^2\sigma}{dx dy} = \frac{2\pi y \alpha^2}{Q^4} L^{\mu\nu} W_{\mu\nu} \quad (1)$$

hard. tensor: 
$$W_{\mu\nu} = \left( -g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) F_1(x, Q^2) + \frac{P_\mu P_\nu}{P \cdot q} F_2(x, Q^2) + i\epsilon_{\mu\nu\alpha\beta} \frac{q^\alpha S^\beta}{P \cdot q} g_1(x, Q^2) \quad (2)$$

$$F_L(x, Q^2) = F_2(x, Q^2) - 2x F_1(x, Q^2) \quad (3)$$

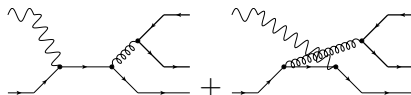
unpol. xs: 
$$\frac{d^2\sigma}{dx dy} = \frac{2\pi\alpha^2}{xyQ^2} \left( Y_+ F_2(x, Q^2) - y^2 F_L(x, Q^2) \right) \quad (4)$$

pol. xs: 
$$\frac{d^2\Delta\sigma}{dx dy} = \frac{4\pi\alpha^2}{xyQ^2} Y_- \cdot 2xg_1(x, Q^2) \quad (5)$$

$$Y_\pm = 1 \pm (1-y)^2 \quad (6)$$

# Computation Review

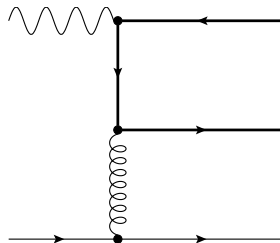
- use factorisation theorem:  $s \rightarrow \xi S_h + \text{PDF}$
- $g(k_1) + \gamma^*(q) \rightarrow \bar{Q}(p_2) + Q(p_1)$
- three massive particles:  $2 \cdot m^2 > 0, q^2 = -Q^2 < 0$
- compute 2-to-3-phase space: e.g.  $dPS_3 \sim dt_1 ds_4 d\theta_1 d\theta_2$



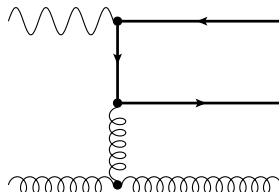
- compute diagrams:  $\Rightarrow 2 \times g_1(x) \sim e_u^2 \cdot \xi \Delta f_u(\xi) \otimes d_{P,q}^{(1)}(\chi, \chi')$
- $d_{P,q}^{(1)}(\chi, \chi') = c_1(\chi, \chi') \ln(\chi) + c_2(\chi, \chi') \text{Li}_2\left(\frac{1+\chi'}{1+\chi}\right) + \dots \checkmark$
- $\frac{m^2}{s} = \frac{\chi}{(1+\chi)^2}$  and  $\frac{m^2}{s+Q^2} = \frac{m^2}{s'} = \frac{\chi'}{(1+\chi')^2}$  and  $\frac{m^2}{Q^2} = \frac{\chi_q}{(1-\chi_q)^2}$

# Computation Review - Collinear Poles

collinear poles appear as  $1/\epsilon$  in



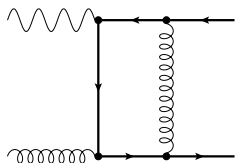
or



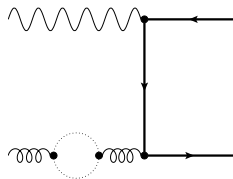
- remove by mass factorization  $\rightarrow \overline{\text{MS}}_m$
- $\Rightarrow 2xg_1(x) \sim e_H^2 \cdot \xi \Delta g(\xi) \otimes \ln(\mu_F^2/m^2) \bar{c}_{P,g}^{F,(1)}(\chi, \chi_q)$
- $\bar{c}_{P,g}^{F,(1)}(\chi, \chi_q) = c_1(\chi, \chi_q) \ln(\chi) + c_2(\chi, \chi_q) \text{Li}_2\left(\frac{1-\chi_q}{1+\chi}\right) + \dots$  ( $\checkmark$  for  $Q^2 \gg m^2$ )

# Computation Review - Virtual and Soft Poles

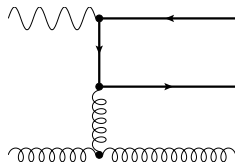
virtual diagrams are, e.g.,



or



soft poles appear in the limit of a soft gluon  $k_2 \rightarrow 0$ , e.g.,



soft + virtual + renormalization + factorization is finite!



# Partonic Results

cg, cq, dq

ALL g1

- inclusive distributions:  $\frac{dg_1}{dp_T}, \frac{dg_1}{dy}$
- correlated distributions:  $\frac{dg_1}{dM_{Q\bar{Q}}^2}, \frac{dg_1}{d\phi}$
- full neutral current (NC) contributions:  $F_3^{Z\gamma}, g_4^{Z\gamma}, g_5^{Z\gamma}$  and  $F_2^Z, F_L^Z, g_1^Z$
- distributions of full NC structure functions:  $\frac{dg_1^{NC}}{dp_T}, \frac{dg_1^{NC}}{dM_{Q\bar{Q}}^2}$