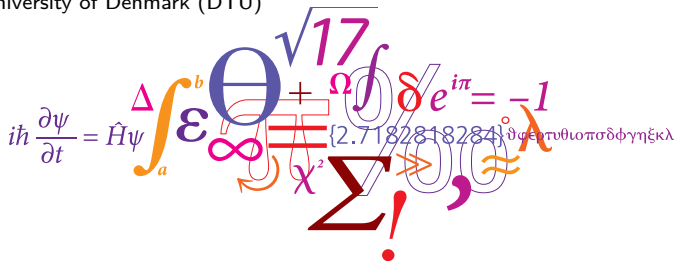


Introduction to discontinuous Galerkin methods

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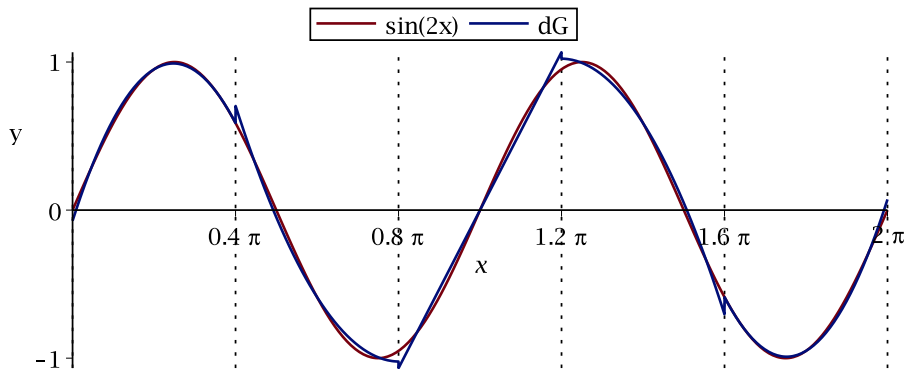
Outline

- Introduction
- Numerical methods
 - Integration
 - Derivatives
 - Local discontinuous Galerkin
- Numerical Experiments
- Conclusions

Discontinuous Galerkin (dG) methods

- are a whole class of methods e.g. Runge Kutta dG (RKDG), local dG (LDG), interior penalty methods, or hybridized dG (\rightarrow G. Giorgiani's talk)
- combine advantages of finite element methods (ease of handling difficult geometry), finite volume (flux conservation) and finite differences (absence of global mass matrix, easily parallelisable)
- B. Cockburn and C. W. Shu. Runge–Kutta discontinuous Galerkin methods for convection-dominated problems. J. Sci. Comput., 16(3):173–261, 2001. doi:10.1023/a:1012873910884
Review article, covers RKDG schemes for advection type problems including e.g shock capturing schemes (slope limiters), as well as LDG for parabolic and elliptic equations
- Notation unfortunately horrible (overwhelming generality evades practicality)
- This presentation: https://feltor-dev.github.io/doc/dg/html/dg_introduction.pdf

A polynomial in each cell



$$f_h(x) = \sum_{n=1}^N \sum_{k=0}^{P-1} \bar{f}^{nk} p_{nk}(x) \quad \text{Discontinuous !}$$

A polynomial in each cell

$$p_{nk}(x) := \begin{cases} p_k\left(\frac{2}{h}(x - x_n)\right), & \text{for } x - x_n \in \left[-\frac{h}{2}, \frac{h}{2}\right] \\ 0, & \text{else.} \end{cases} \quad (1)$$

- p_k are Legendre polynomials (in this work at least)
- orthogonal and complete
- P is the number of base polynomials per cell (typically 3 or 4)
- In each cell we have X-space and L-space (x_j^a and w_j are abscissas and weights of Gauss-Legendre quadrature)

$$\bar{f}^k = \sum_{j=0}^{P-1} F^{kj} f_j \quad F^{kj} := \frac{2k+1}{2} w_j p_k(x_j^a) \quad (2a)$$

$$f_j = \sum_{k=0}^{P-1} B_{jk} \bar{f}^k \quad B_{jk} := p_k(x_j^a) \quad (2b)$$

We directly have an order P interpolation formula (just evaluate the polynomials at x)

$$f_h(x) = \sum_{n=1}^N \sum_{k=0}^{P-1} \bar{f}^{nk} p_{nk}(x)$$

Integrating polynomials yields

$$\langle f_h, g_h \rangle := \int_a^b f_h(x) g_h(x) \, dx = \sum_{n=1}^N \sum_{j=0}^{P-1} \frac{hw_j}{2} f_{nj} g_{nj} = \sum_{n=1}^N \sum_{k=0}^{P-1} \frac{h}{2k+1} \bar{f}^{nk} \bar{g}^{nk}$$

equals Gauss-Legendre integration (of order $2P - 1$). Note that its implementation is a reduction.

A first try (imply summation over repeated indices)

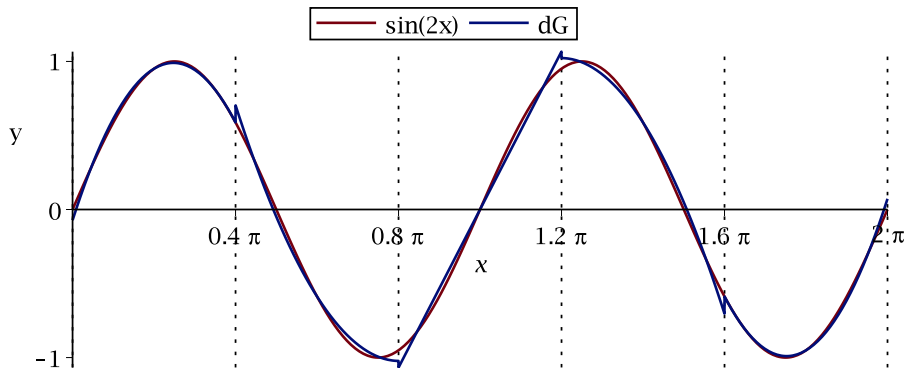
$$\partial_x f_h(x) = \bar{f}^{ni} \partial_x p_{ni}(x) \quad ??$$

Loose one order and completely wrong for $P = 1$

Weak formulation (project with base functions, integrate by part)

$$\int_{C_n} \partial_x f_h(x) p_{ni}(x) dx = (f_h p_{ni})|_{x_{n-1/2}}^{x_{n+1/2}} - \int_{C_n} f_h(x) \partial_x p_{ni}(x) dx$$

Remember $f_h(x)$ is double valued on cell-boundaries !?



$$\int_{C_n} \partial_x f_h(x) p_{ni}(x) dx = (f_h p_{ni})|_{x_{n-1/2}}^{x_{n+1/2}} - \int_{C_n} \bar{f}^{nk} p_{nk} \partial_x p_{ni} dx$$

Freely choose the flux $f_h \rightarrow \hat{f}$ on cell boundaries.

The flux \hat{f} is what defines a dG method!

$$\hat{f}_C(x) = \frac{1}{2} \left(\lim_{\varepsilon \rightarrow 0, \varepsilon > 0} f_h(x + \varepsilon) + \lim_{\varepsilon \rightarrow 0, \varepsilon > 0} f_h(x - \varepsilon) \right), \text{ centered} \quad (3a)$$

$$\hat{f}_F(x) = \lim_{\varepsilon \rightarrow 0, \varepsilon > 0} f_h(x + \varepsilon), \text{ forward} \quad (3b)$$

$$\hat{f}_B(x) = \lim_{\varepsilon \rightarrow 0, \varepsilon > 0} f_h(x - \varepsilon), \text{ backward} \quad (3c)$$

There are many more (e.g. slope limiters, Godunov flux, Lax-Friedrich ...)

$$\bar{D}_{x,per}^C = \frac{1}{2} \begin{pmatrix} (M - M^T) & RL & & & -RL^T \\ -RL^T & (M - M^T) & RL & & \\ & -RL^T & \dots & & \\ & & \dots & RL & \\ RL & & -RL^T & (M - M^T) & \end{pmatrix}$$

- M and RL are $P \times P$ block matrices.
- Leads to familiar finite difference formulas for $P = 1$!

$$-\frac{\partial}{\partial x} \left(\chi(x) \frac{\partial \phi(x)}{\partial x} \right) = \rho(x) \quad (4)$$

Split into first order equations

$$j' = \partial_x \phi, \quad (5a)$$

$$j = \chi j', \quad (5b)$$

$$\rho = -\partial_x j. \quad (5c)$$

- Apply same procedure on each equation
- Choose fluxes such that $\hat{\phi}$ is independent of j_h , but \hat{j} might depend on ϕ_h to penalize the jumps between cells (local dG) $\rightarrow \rho = D_x^T \chi D_x + J$
- Self-adjoint (symmetric) discretization easily invertible with a (preconditioned) conjugate gradient method

On structured grids the discretization is immediately extensible to higher dimensions:

$$-\frac{\partial}{\partial x} \left(\chi(x, y) \frac{\partial \phi(x, y)}{\partial x} \right) - \frac{\partial}{\partial y} \left(\chi(x, y) \frac{\partial \phi(x, y)}{\partial y} \right) = \rho(x, y)$$

Solve on $[0, \pi] \times [0, \pi]$ with Dirichlet boundary condition

$$\chi(x, y) = 1 + \sin(x) \sin(y)$$

$$\rho(x, y) = 2 \sin(x) \sin(y) [\sin(x) \sin(y) + 1] - \sin^2(x) \cos^2(y) - \cos^2(x) \sin^2(y)$$

# of cells	ε_{res}	Forward		Backward		Centered			
		iterations	L^2 error	iterations	L^2 error	Order	iterations	L^2 error	Order
$P = 3$									
17^2	1.0E-06	181	4.77E-05	181	4.77E-05	-	113	5.37E-06	-
34^2	1.0E-07	403	5.22E-06	403	5.22E-06	3.19	259	3.67E-07	3.87
68^2	1.0E-08	893	5.93E-07	892	5.93E-07	3.14	583	2.64E-08	3.80
136^2	1.0E-09	1946	6.97E-08	1946	6.97E-08	3.09	1277	1.92E-09	3.78
$P = 4$									
17^2	1.0E-08	357	4.62E-07	357	4.62E-07	-	221	7.60E-07	-
34^2	1.0E-09	793	2.47E-08	795	2.47E-08	4.22	498	5.54E-08	3.78
68^2	1.0E-09	1637	1.48E-09	1637	1.48E-09	4.06	1035	3.80E-09	3.87
136^2	1.0E-10	3505	9.13E-11	3505	9.13E-11	4.02	2223	2.49E-10	3.93
$P = 5$									
17^2	1.0E-09	581	1.57E-08	580	1.57E-08	-	354	2.16E-09	-
34^2	1.0E-10	1277	3.62E-10	1277	3.62E-10	5.44	782	3.51E-11	5.95
68^2	1.0E-11	2751	8.39E-12	2752	8.39E-12	5.43	1697	6.68E-13	5.71
136^2	1.0E-12	5816	2.03E-13	5816	2.03E-13	5.37	3597	4.01E-14	4.06

- In practice it is easier to work in configuration space than with the polynomial coefficients
- Well parallelizable (we have implementations for OpenMP, Cuda, MPI, MPI+OpenMP as well as MPI+Cuda)
- A higher order method does not automatically lead to a more accurate solution (in practice $P = 3$ or $P = 4$ is absolutely sufficient).
- Beware of supraconvergence (inversion of $\partial_x^2 \phi = \rho$ converges, but $\partial_x^2 \phi$ does not necessarily converge on its own)
- LDG quite robust; works well for elliptic equations of the form $\sum_{ij} \partial_i (\chi^{ij} \partial_j \phi) = \rho$ with χ positive (semi-)definite and also $(1 + \xi \Delta) \phi = \rho$