# Ref of Fu's 1993 paper

September 18, 2015

**The linearized drift kinetic equation is given by**

$$\left(\partial_t + \mathbf{v}_d \cdot \nabla + v_\parallel \hat{\mathbf{b}} \cdot \nabla\right) g = ie\frac{\partial F}{\partial \epsilon}\left(\omega - \omega_\star\right)\frac{i}{\omega}\mathbf{v}_d \cdot \delta\mathbf{E}_\perp$$

where,$\mathbf{v}_d = \frac{\hat{\mathbf{b}}}{\Omega} \times \left(\mu \nabla B + \kappa v_\parallel^2\right) \approx \frac{v_\perp^2/2+v_\parallel^2}{\Omega}\hat{\mathbf{b}} \times \kappa$, $\mu = \frac{v_\perp^2}{2B}$, $\omega_\star = \frac{i\hat{\mathbf{b}}\times\nabla F\cdot\nabla}{\Omega \partial F/\partial\epsilon}$, $\delta\mathbf{E}_\perp = i\omega\vec{\xi} \times \mathbf{B}$.

**The term $\mathbf{v}_d \cdot \delta\mathbf{E}_\perp$ can be expressed by the following form.**

$$\mathbf{v}_d \cdot \delta\mathbf{E}_\perp = \frac{v_\perp^2/2+v_\parallel^2}{\Omega}\hat{\mathbf{b}} \times \kappa \cdot \delta\mathbf{E}_\perp = i\omega\frac{v_\perp^2/2+v_\parallel^2}{\Omega}\hat{\mathbf{b}} \times \vec{\kappa} \cdot \left(\vec{\xi} \times \mathbf{B}\right) \tag{1}$$

$$= -i\omega B\frac{v_\perp^2/2+v_\parallel^2}{\Omega}\vec{\kappa}\cdot\vec{\xi} = -i\omega B\frac{v_\perp^2/2+v_\parallel^2}{\Omega}\left(\nabla\theta\kappa_\theta + \nabla r\kappa_r\right)\cdot\left(\xi_\theta\nabla\theta + \xi_r\nabla r\right)$$

$$= -i\omega B\frac{v_\perp^2/2+v_\parallel^2}{\Omega}\left(\nabla\theta\cdot\nabla\theta\xi_\theta\kappa_\theta + \nabla r\cdot\nabla r\xi_r\kappa_r + \nabla\theta\cdot\nabla r\xi_r\kappa_\theta + \nabla r\cdot\nabla\theta\xi_\theta\kappa_r\right)$$

$$= -i\omega B\frac{\epsilon}{\Omega}\left(\frac{\Lambda}{b} + 2\left(1 - \frac{\Lambda}{b}\right)\right)\left(\left(g^{\theta\theta}\kappa_\theta + g^{r\theta}\kappa_r\right)\xi_\theta + \left(g^{rr}\kappa_r + g^{r\theta}\kappa_\theta\right)\xi_r\right)$$

with $g^{rr} = \nabla r\cdot\nabla r$, $g^{\theta r} = \nabla\theta\cdot\nabla r$, $g^{\theta\theta} = \nabla\theta\cdot\nabla\theta$. $\xi_r\nabla r = \xi_r\mathbf{e}_r$, $\xi_\theta\nabla\theta = \frac{1}{r}\xi_\theta\mathbf{e}_\theta$, $\kappa_\theta = -\frac{1}{R}\frac{\partial R}{\partial\theta}$, $\kappa_r \approx -\frac{1}{R}\frac{\partial R}{\partial r} - \frac{r}{q^2R^2}$ [G. Y. Fu, PHYSICS OF PLASMAS 13 2006]. $\Lambda = \frac{\mu B_0}{\epsilon}$, $b = B_0/B \approx 1 + (r/R_0)\cos\theta$, $\epsilon = \frac{1}{2}v^2$.

**The linearized drift kinetic equation is rewritten as**

$$\frac{d}{dt}g = ie\frac{\partial F}{\partial\epsilon}\left(\omega - \omega_\star\right)B\frac{\epsilon}{\Omega}\left(\frac{\Lambda}{b} + 2\left(1 - \frac{\Lambda}{b}\right)\right)\left(\left(g^{\theta\theta}\kappa_\theta + g^{r\theta}\kappa_r\right)\xi_\theta + \left(g^{rr}\kappa_r + g^{r\theta}\kappa_\theta\right)\xi_r\right) \tag{2}$$

$$\frac{d}{dt}g = H\left(r, \theta, \phi, t\right)$$

**The solution of perturbed distribution function $g$ is obtained in the followings.** At equilibrium, the projection of the orbit on the polodial cross section is a closed curve. For either mirror-trapped or passing orbit, we define the bounce time [F. Porcelli, R. Stankiewicz, and W. Kerner, Phys. Plasmas 1 1994]

$$\tau_b = \oint d\tau = \oint \frac{d\psi}{\dot{\psi}} = \oint \frac{d\theta}{\dot{\theta}} \tag{3}$$

as the time it takes to close an equilibrium orbit on the poloidal plane. We assume that perturbations have the form

$$X^{(1)} = \hat{X}^{(1)}\left(r,\theta\right)\exp\left(-i\omega t + in\phi\right) \tag{4}$$

The formal solution of the nonadibatic distribution $g$ is

$$g = \int_{-\infty}^{t} ie\frac{\partial F}{\partial \epsilon}\left(\omega - \omega_{\star}\right) B\frac{\epsilon}{\Omega} G\left(\tau\right) d\tau \tag{5}$$

with

$$G = \left(\frac{\Lambda}{b} + 2\left(1 - \frac{\Lambda}{b}\right)\right)\left(\left(g^{\theta\theta}\kappa_{\theta} + g^{r\theta}\kappa_{r}\right)\xi_{\theta} + \left(g^{rr}\kappa_{r} + g^{r\theta}\kappa_{\theta}\right)\xi_{r}\right)$$

where $G\left(\tau\right) = \hat{G}\left[r\left(\tau\right),\theta\left(\tau\right)\right]\exp\left(-i\omega\tau + in\phi\left(\tau\right)\right)$ and the $\tau$ dependence is through the following equations:

$$\dot{r} = \dot{\mathbf{R}}\cdot\nabla r, \dot{\theta} = \dot{\mathbf{R}}\cdot\nabla\theta, \dot{\phi} = \dot{\mathbf{R}}\cdot\nabla\phi \tag{6}$$

Let us separate $\phi\left(\tau\right)$ into its secular and osillating parts:

$$\phi\left(\tau\right) = \left\langle\dot{\phi}\right\rangle\tau + \widetilde{\phi}\left(\tau\right) \tag{7}$$

where the brackets indicate bounce averaging.

The quantity $\tilde{G}\left[r\left(\tau\right),\theta\left(\tau\right)\right] = \hat{G}\left[r\left(\tau\right),\theta\left(\tau\right)\right]\exp\left(in\tilde{\phi}\left(\tau\right)\right)$ is a periodic function of $\tau$, which can be expanded in Fourier series,

$$\tilde{G}\left(\tau\right) = \sum_{-\infty}^{\infty} Y_p\left(\Lambda,\epsilon,\bar{r};\sigma\right)\exp\left(ip\omega_b\tau\right) \tag{8}$$

where,

$$Y_p\left(\Lambda,\epsilon,\bar{r};\sigma\right) = \frac{1}{\tau_b}\oint d\tau\tilde{G}\left(\tau\right)\exp\left(-ip\omega_b\tau\right) \tag{9}$$

with $r\left(\tau\right) = \bar{r} + \rho\cos\theta\left(\tau\right)$, $\rho$ repesents the finite orbit width.

When $\tilde{G}$ is expressed by $\cos k\theta$, $\sin k\theta$ instead of $\exp ik\theta$, for $p = 0$,

$$Y_p\left(\Lambda,\epsilon,\bar{r};\sigma\right) = \frac{1}{\tau_b}\oint d\tau\tilde{G}\left(\tau\right)\exp\left(-ip\omega_b\tau\right) \tag{10}$$

for $p \neq 0$,

$$Y_p\left(\Lambda,\epsilon,\,\bar{r};\sigma\right)=\frac{2}{\tau_b}\oint d\tau\tilde{G}\left(\tau\right)\exp\left(-ip\omega_b\tau\right) \tag{11}$$

Carrying out the time integration, the solution of $g$ is obtained

$$g=e\frac{\partial F}{\partial\epsilon}\left(\omega-\omega_\star\right)B\frac{\epsilon}{\Omega}\sum_{-\infty}^{\infty}Y_p\left(\Lambda,\epsilon,\bar{r};\sigma\right)\frac{\exp\left[i\left(n\left\langle\dot{\phi}\right\rangle+p\omega_b-\omega\right)t\right]}{n\left\langle\dot{\phi}\right\rangle+p\omega_b-\omega} \tag{12}$$

**The formula of $\delta W_k$ is derived as follows.**

$$\delta W_k=\int d^3x\vec{\xi}^\star\cdot\nabla\cdot\delta\mathbf{P}_k=-e\int d^3x\int d^3v\left(\frac{i}{\omega}\mathbf{v}_d\cdot\delta\mathbf{E}_\perp\right)^\star g$$

$$=-e\int d^3x\int d^3vgB\frac{\epsilon}{\Omega}\left(\frac{\Lambda}{b}+2\left(1-\frac{\Lambda}{b}\right)\right)\left(\left(g^{\theta\theta}\kappa_\theta+g^{r\theta}\kappa_r\right)\xi_\theta^\star+\left(g^{rr}\kappa_r+g^{r\theta}\kappa_\theta\right)\xi_r^\star\right)$$

$$=-e\int d^3x\int d^3vgB\frac{\epsilon}{\Omega}G^\star$$

where $G^\star=\hat{G}^\star\left[r\left(\tau\right),\theta\left(\tau\right)\right]\exp\left(i\omega\tau-in\phi\left(\tau\right)\right)$ . Let $\tilde{G}^\star\left[r\left(\tau\right),\theta\left(\tau\right)\right]=\hat{G}^\star\left[r\left(\tau\right),\theta\left(\tau\right)\right]\exp\left(-in\tilde{\phi}\left(\tau\right)\right)$, which is a periodic function of $\tau$, which can be expanded in Fourier series,

$$\tilde{G}^\star\left(\tau\right)=\sum_{-\infty}^{\infty}Y_p^\star\left(\Lambda,\epsilon,\bar{r};\sigma\right)\exp\left(-ip\omega_b\tau\right) \tag{13}$$

where,

$$Y_p^\star\left(\Lambda,\bar{r};\sigma\right)=\frac{1}{\tau_b}\oint d\tau\tilde{G}^\star\left(\tau\right)\exp\left(ip\omega_b\tau\right) \tag{14}$$

with $r\left(\tau\right)=\bar{r}+\rho\cos\theta\left(\tau\right)$.

$$\delta W_k=-e^2\int d^3x\int d^3v\frac{\partial F}{\partial\epsilon}\left(\omega-\omega_\star\right)B^2\frac{\epsilon^2}{\Omega^2}\sum_{-\infty}^{\infty}Y_p\left(\Lambda,\bar{r};\sigma\right)$$

$$\cdot\frac{\exp\left[i\left(n\left\langle\dot{\phi}\right\rangle+p\omega_b-\omega\right)\tau\right]}{n\left\langle\dot{\phi}\right\rangle+p\omega_b-\omega}\sum_{-\infty}^{\infty}Y_{p'}^\star\left(\Lambda,\epsilon,\bar{r};\sigma\right)\exp\left(i\omega\tau-in\left\langle\dot{\phi}\right\rangle\tau-ip'\omega_b\tau\right) \tag{15}$$

$$\delta W_k = -e^2 \int d^3x \int d^3v \frac{\partial F}{\partial \epsilon} \left(\omega - \omega_\star\right) B^2 \frac{\epsilon^2}{\Omega^2}$$

$$\cdot \sum_{-\infty}^{\infty} Y_p \left(\Lambda, \epsilon, \bar{r}; \sigma\right) \frac{\exp\left[ip\omega_b\tau\right]}{n\left\langle \dot{\phi} \right\rangle + p\omega_b - \omega} \sum_{-\infty}^{\infty} Y_{p'}^{\star} \left(\Lambda, \epsilon, \bar{r}; \sigma\right) \exp\left(-ip'\omega_b\tau\right) \tag{16}$$

Using $d^3v = \sqrt{2}\pi \frac{1}{b\sqrt{1-\frac{\Lambda}{b}}} d\Lambda \epsilon^{1/2} d\epsilon$, $d^3x = 2\pi J dr d\theta$, yields

$$\delta W_k = -e^2 \int 2\pi J dr d\theta \int \sqrt{2}\pi \frac{1}{b\sqrt{1-\frac{\Lambda}{b}}} d\Lambda \epsilon^{1/2} d\epsilon \frac{\partial F}{\partial \epsilon} \left(\omega - \omega_\star\right) B^2 \frac{\epsilon^2}{\Omega^2}$$

$$\cdot \sum_{-\infty}^{\infty} Y_p \left(\Lambda, \epsilon, \bar{r}; \sigma\right) \frac{\exp\left(ip\omega_b\tau\right)}{n\left\langle \dot{\phi} \right\rangle + p\omega_b - \omega} \sum_{-\infty}^{\infty} Y_{p'}^{\star} \left(\Lambda, \epsilon, \bar{r}; \sigma\right) \exp\left(-ip'\omega_b\tau\right) \tag{17}$$

Applying $d\tau = \frac{qR_0}{\sigma\sqrt{2\epsilon}b\sqrt{1-\frac{\Lambda}{b}}} d\theta$, $\sigma = \pm 1$ for the direction of $v_\parallel$, one finially obtains

$$\delta W_k = -4\pi^2 \frac{e^2 B^2}{\Omega^2} \frac{1}{R_0} \int \frac{J}{q} dr \int d\Lambda \epsilon^3 d\epsilon \frac{\partial F}{\partial \epsilon} \tau_b \left(\omega - \omega_\star\right)$$

$$\cdot \sum_{-\infty}^{\infty} \frac{\left|Y_p\right|^2}{n\left\langle \dot{\phi} \right\rangle + p\omega_b - \omega} \tag{18}$$

Note that $\tilde{\phi} \cong 0$, $\left\langle \dot{\phi} \right\rangle \cong \omega_D^0 + q\omega_b, \omega_D^0 \approx 0$ for passing particles. For internal kink mode $\nabla \cdot \vec{\xi} = 0$, the forms of the perturbation are $\xi_r = \xi_0 \cos\theta$, $\xi_\theta = -\xi_0 r \sin\theta$ within the region $q = 1$ rational surface $r_s = 1$, and $\xi_r = \xi_0 \left(\frac{\Delta r - r + (r_s - \Delta r/2)}{\Delta r}\right) \cos\theta$, $\xi_\theta = -\xi_0 \left(\frac{\Delta r - r + (r_s - \Delta r/2)}{\Delta r}\right) r \sin\theta + \xi_0 \left(\frac{r}{\Delta r}\right) r \sin\theta$ in the inertial region $r_s - \frac{\Delta r}{2} \leq r \leq r_s + \frac{\Delta r}{2}$.

In angle-action coordinate,

$$J_b = \frac{1}{2\pi} \int p_\parallel ds \cong p_{\parallel e} R_\parallel \int_{-\theta_b}^{\theta_b} \sqrt{1 - \kappa^{-1} \sin^2 \frac{\theta}{2}} \frac{d\theta}{\pi} \tag{19}$$

$$J_t = \frac{1}{2\pi} \int p_\parallel ds \cong p_{\parallel e} R_\parallel \int_{-\pi}^{\pi} \sqrt{1 - \kappa^{-1} \sin^2 \frac{\theta}{2}} \frac{d\theta}{\pi} \tag{20}$$

The formulas of bounce/transit frequency is given by [Alain J. Brizard, PHYSICS OF PLASMAS 18 2011]

$$\omega_b = \frac{\partial H}{\partial J_b} = \left(\frac{\partial J_b}{\partial E}\right)^{-1} = \frac{\pi\omega_\parallel}{2K(\kappa)}, \kappa < 1 \tag{21}$$

$$\omega_t = \frac{\partial H}{\partial J_t} = \left(\frac{\partial J_t}{\partial E}\right)^{-1} = \frac{\pi\sqrt{\kappa}\omega_\parallel}{K(\kappa^{-1})}, \kappa > 1 \tag{22}$$

where $\omega_\parallel = \frac{1}{qR}\sqrt{\varepsilon\mu B_0} = \frac{\sqrt{\epsilon}}{qR}\sqrt{\varepsilon\Lambda}$, $\kappa = \frac{1-\Lambda(1-\varepsilon)}{2\varepsilon\Lambda}$, $\varepsilon = \frac{r}{R_0}$. $K$ denotes the complete elliptic integral of the first kind.

**The normalized relations of the quantities are** $F = \frac{n}{v_h^3}\bar{F}_0$**,** $v_h = \sqrt{\frac{2T_h}{M}}$**,** $\epsilon = \frac{T_h}{M}\bar{\epsilon}$**,** $r = a\bar{x}$**,** $J = aR_0\bar{J}$**,** $R = R_0\bar{R}$**,** $\omega_t = \frac{v_h}{R_0}\bar{\omega}_t$, $\frac{1}{\tau_t} = \frac{v_h}{2\pi R_0}\bar{\omega}_t = \frac{v_h}{R_0}\frac{\bar{\omega}_t}{2\pi} = \frac{v_h}{R_0}\frac{1}{\bar{\tau}_t}$**,** $\omega = \frac{v_h}{R_0}\bar{\omega}$, $\omega_\phi = \frac{v_h}{R_0}\bar{\omega}_\phi$, $\omega_\star = \frac{v_h}{R_0}\bar{\omega}_\star$ **.**

$$\delta W_k = -\pi^2\frac{e^2B^2}{\Omega^2}a^2R_0n_0\frac{T_h}{M}\int\frac{\bar{J}}{q}dx\int d\Lambda\bar{\epsilon}^3d\bar{\epsilon}\frac{\partial\bar{F}}{\partial\bar{\epsilon}}\bar{\tau}_b\left(\bar{\omega}-\bar{\omega}_\star\right)$$

$$\cdot\sum_{-\infty}^{\infty}\frac{\left|\bar{Y}_p\right|^2}{n\left\langle\bar{\dot{\phi}}\right\rangle + p\bar{\omega}_b - \bar{\omega}} \tag{23}$$

For passing particles,

$$\bar{\omega}_b = \frac{\pi\sqrt{\kappa}}{K(\kappa^{-1})}\frac{\sqrt{\varepsilon\Lambda/2}}{q}\sqrt{\bar{\epsilon}} \tag{24}$$

$$\omega_b t = \bar{\omega}_b\frac{v_h}{R}\int_0^\theta\frac{qR_0}{\sqrt{2(T/M)\bar{\epsilon}}b\sqrt{1-\frac{\Lambda}{b}}}d\theta \tag{25}$$

$$= \int_0^\theta\frac{\pi\sqrt{\kappa}}{K(\kappa^{-1})}\sqrt{\varepsilon\Lambda/2}\frac{1}{b\sqrt{1-\frac{\Lambda}{b}}}d\theta$$

$$Y_p(\Lambda,\bar{r};\sigma) = \frac{1}{\tau_b}\oint d\tau\tilde{G}(\tau)\exp(-ip\omega_b\tau) \tag{26}$$

$$= \frac{\omega_b}{2\pi}\oint d\tau\tilde{G}\left[r(\tau),\theta(\tau)\right]\exp(-ip\omega_b\tau)$$

$$= \frac{1}{2\pi}\oint d(\omega_b\tau)\tilde{G}\left[r(\tau),\theta(\tau)\right]\exp(-ip\omega_b\tau)$$

$$= \frac{1}{2\pi}\int_0^{2\pi}d\theta\frac{\sigma\pi\sqrt{\kappa}}{K(\kappa^{-1})}\frac{\sqrt{\varepsilon\Lambda/2}}{b\sqrt{1-\frac{\Lambda}{b}}}\tilde{G}\left[r(\tau),\theta(\tau)\right]\exp\left(-ip\int_0^\theta d\theta'\frac{\sigma\pi\sqrt{\kappa}}{K(\kappa^{-1})}\frac{\sqrt{\varepsilon\Lambda/2}}{b\sqrt{1-\frac{\Lambda}{b}}}\right)$$

where

$$\tilde{G}\left[r\left(\tau\right),\theta\left(\tau\right)\right]=\left(\frac{\Lambda}{b\left(r,\theta\right)}+2\left(1-\frac{\Lambda}{b\left(r,\theta\right)}\right)\right)$$

$$\cdot\left(\left(g^{\theta\theta}\kappa_{\theta}+g^{r\theta}\kappa_{r}\right)\hat{\xi}_{\theta}\left(\theta,\bar{r}+\rho\cos\theta\right)+\left(g^{rr}\kappa_{r}+g^{r\theta}\kappa_{\theta}\right)\hat{\xi}_{r}\left(\theta,\bar{r}+\rho\cos\theta\right)\right)$$

Note that the effect of finite orbit width is induced radially in the perturbed $\xi\left(r\right)$ for simply. Furthermore, $\tilde{G}$ is a normalized quantity, so is $Y_p$,

$$\tilde{G}\left[x\left(\tau\right),\theta\left(\tau\right)\right]=\left(\frac{\Lambda}{b\left(x,\theta\right)}+2\left(1-\frac{\Lambda}{b\left(x,\theta\right)}\right)\right)$$

$$\cdot\left(\left(\bar{g}^{\theta\theta}\bar{\kappa}_{\theta}+\bar{g}^{r\theta}\bar{\kappa}_{r}\right)\bar{\hat{\xi}}_{\theta}\left(\theta,\bar{x}+\frac{\rho}{a}\cos\theta\right)+\left(\bar{g}^{rr}\bar{\kappa}_{r}+\bar{g}^{r\theta}\bar{\kappa}_{\theta}\right)\bar{\hat{\xi}}_{r}\left(\theta,\bar{x}+\frac{\rho}{a}\cos\theta\right)\right)$$

where, the normalized displacements are $\bar{\hat{\xi}}_{\theta m}=\hat{\xi}_{\theta m}/a^2$, $\bar{\hat{\xi}}_{rm}=\hat{\xi}_{rm}/a$.

The slowing down distribution function of fast ions is given by[M. Schneller 2013]

$$F\left(x,\bar{\epsilon},\Lambda\right)=\frac{n_0}{v_h^3}\frac{1}{\bar{\epsilon}^{3/2}+\bar{\epsilon}_c^{3/2}}Erfc\left(\frac{\bar{\epsilon}-\bar{\epsilon}_0}{\Delta\bar{\epsilon}}\right)\exp\left[-\left(\frac{x-x_0}{\Delta x}\right)^2\right]\exp\left(\frac{\Lambda-\Lambda_0}{\Delta\Lambda^{0.2}}\right)^2 \tag{27}$$

The normalized metric tensors are

$$\bar{g}^{rr}=1+2\Delta'\cos\theta \tag{28}$$

with $\Delta'=\left(\varepsilon+\alpha\right)/4$, $\varepsilon=\frac{r}{R_0}$, $\alpha=-R_0q^2d\beta/dr$, $\beta=\frac{2\mu_0P}{B^2}$ set $\alpha=0$ if $\beta=0$, or assume $\bar{g}^{rr}=1$ without toroidal effect, $\theta$ independent. Specially, in low beta limit,

$$\bar{g}^{rr}=1+\frac{1}{2}\varepsilon\cos\theta \tag{29}$$

$$\bar{g}^{\theta\theta}=\frac{1}{x^2}\left[1-2\left(\varepsilon+\Delta'\right)\cos\theta\right] \tag{30}$$

assume $\bar{g}^{\theta\theta}=\frac{1}{x^2}$ without toroidal effect. Specially, in low beta limit,

$$\bar{g}^{\theta\theta}=\frac{1}{x^2}\left[1-\frac{5}{2}\varepsilon\cos\theta\right] \tag{31}$$

$$\bar{g}^{r\theta}=-\frac{1}{x}\left[\varepsilon+\left(r\Delta'\right)'\right]\sin\theta \tag{32}$$

specially,

$$\bar{g}^{r\theta}=-\frac{1}{x}\frac{3}{2}\varepsilon\sin\theta \tag{33}$$

for low beta limit. and $\bar{g}^{r\theta} = 0$ without toroidal effect.
The normalized curvature are in low beta limit

$$\bar{\kappa}_r = -\frac{a}{R}\cos\theta + \frac{a}{R}\frac{\varepsilon}{4} - \frac{a}{R}\frac{5}{4}\varepsilon\left(\cos 2\theta - 1\right) - \left(\frac{a}{R}\right)^2\frac{x}{q} \tag{34}$$

$$\bar{\kappa}_\theta = \varepsilon\sin\theta + \frac{5}{4}\varepsilon^2\sin 2\theta \tag{35}$$

with $R = R_0 + r\cos\theta - \Delta\left(r\right) + r\eta\left(r\right)\left(\cos 2\theta - 1\right).\eta\left(r\right) = \left(\varepsilon + \Delta'\right)/2$.
The normalized $\omega_\star$ is

$$\bar{\omega}_\star = \frac{1}{2}\frac{m\partial\bar{F}/\partial x}{\partial\bar{F}/\partial\epsilon}\frac{1}{x}\frac{R}{a}\frac{\rho_h}{a} \tag{36}$$

where, $m$ is poloidal mode number, $\rho_h = v_h/\Omega$, $v_h = \sqrt{2T_h/M}, \Omega = Be/M$.
The normalized $\xi$ are

$$\bar{\xi}_r\left(\theta, x\right) = \bar{\xi}_0\cos\theta$$

$$\bar{\xi}_\theta\left(\theta, x\right) = -\bar{\xi}_0 x\sin\theta$$

within $q = 1$ surface. In the inertial region $r_s - \frac{\Delta r}{2} \leq r \leq r_s + \frac{\Delta r}{2}$,

$$\bar{\xi}_r\left(\theta, x\right) = \bar{\xi}_0\left(\frac{\bar{\Delta} r - x + \left(\bar{r}_s - \overline{\Delta r}/2\right)}{\overline{\Delta r}}\right)\cos\theta$$

$$\bar{\xi}_\theta\left(\theta, x\right) = -\bar{\xi}_0\left(\frac{\overline{\Delta r} - x + \left(\bar{r}_s - \overline{\Delta r}/2\right)}{\overline{\Delta r}}\right)x\sin\theta + \bar{\xi}_0\left(\frac{x}{\overline{\Delta r}}\right)x\sin\theta$$

with $\bar{\xi}_0 = \xi_0/a$, $\bar{r}_s = r_s/a$, $\overline{\Delta r} = \Delta r/a$, $x = r/a$.