

# dwk++ Documentation

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January 26, 2016

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## 1 Introduction

dwk++ A small code to get  $\delta W_k$  and solver fishbone dispersion relation. This c++ code is used to calculate delta\_W\_k in tokamak plasma by 3D integration (r,Lambda,E). A slowing down distribution function and a kink like mode structure is used. The goal of this code is to find the Eigenvalues with a mode structure input.

## 2 The defintion of $\delta W_k$

### 2.1 The normalized $\delta W_k$ in the code:

$$\delta W_k = \sum_p \int_0^1 dr \int \frac{d\Lambda}{B} \int E^3 dE \tau_b(\omega - \omega_*) \frac{\partial F}{\partial E} \frac{|Y_p|^2}{n\omega_\phi + p\omega_b - \omega} \quad (1)$$

$\sigma = \pm 1$  is the direction of  $v_{\parallel}$ ,  $m$  and  $n$  is poloidal and toroidal mode number,  $\Lambda = \frac{\mu}{E}$ ,  $E$  is the fast ion energy,  $\tau_b$  is the particle bounce time. The slowing down distribution function of fast ions is:

$$F = \frac{1}{C_f} F(r, \epsilon, \Lambda) = \frac{1}{C_f} \frac{1}{E^{3/2} + E_c^{3/2}} \text{erfc} \left( \frac{E - E_0}{\Delta E} \right) \exp \left[ -\left( \frac{r - r_0}{\Delta r} \right)^2 \right] \exp \left[ -\left( \frac{\Lambda - \Lambda_0}{\Delta \Lambda} \right)^2 \right] \quad (2)$$

$$C_f = \int dE d\Lambda \frac{1}{E^{3/2} + E_c^{3/2}} \text{erfc} \left( \frac{E - E_0}{\Delta E} \right) \exp \left[ -\left( \frac{\Lambda - \Lambda_0}{\Delta \Lambda} \right)^2 \right] \quad (3)$$

$$\frac{\partial F}{\partial E} = \frac{-1}{C_f} \left[ \frac{2 \exp(-(\frac{E-E_0}{\Delta E})^2)}{\sqrt{\pi} \Delta E (E^{3/2} + E_c^{3/2})} + \frac{3\sqrt{E} \text{erfc}(\frac{E-E_0}{\Delta E})}{2(E^{3/2} + E_c^{3/2})^2} - \frac{2\Lambda(\Lambda - \Lambda_0) \text{erfc}(\frac{E-E_0}{\Delta E})}{E \Delta \Lambda^2 (E^{3/2} + E_c^{3/2})} \right] \quad (4)$$

$$* \exp \left[ -\left( \frac{r - r_0}{\Delta r} \right)^2 \right] \exp \left[ -\left( \frac{\Lambda - \Lambda_0}{\Delta \Lambda} \right)^2 \right] \quad (5)$$

$$\frac{\partial F}{\partial r} = \frac{1}{C_f} \frac{\text{erfc}(\frac{E-E_0}{\Delta E})}{E^{3/2} + E_c^{3/2}} \frac{2(r_0 - r)}{\Delta r^2} \exp \left[ -\left( \frac{r - r_0}{\Delta r} \right)^2 \right] \exp \left[ -\left( \frac{\Lambda - \Lambda_0}{\Delta \Lambda} \right)^2 \right] \quad (6)$$

The diamagnetic frequency:

$$\omega_* = \frac{m}{2r} \frac{\rho_0}{\varepsilon_0} \frac{\partial F / \partial r}{\partial F / \partial E} \quad (7)$$

where  $\rho_0$  is the gyro radius with injection energy.  $\varepsilon_0$  is the inverse aspect-ratio, and  $\varepsilon = \frac{r}{R_0}$ . The transit frequency for passing particle is given below:

$$\omega_b = \frac{\pi \sqrt{\kappa}}{K(\kappa^{-1})} \frac{\sqrt{\varepsilon \Lambda / 2}}{q} \sqrt{E} \quad (8)$$

The transit frequency in toroidal:  $\omega_\phi = q\omega_b$ . Particle bounce time:  $\tau_b = \frac{2\pi}{\omega_b}$ . The integral along the particle orbits:

$$Y(r, \Lambda) = \frac{1}{2\pi} \int_0^{2\pi} \chi d\theta B_\Lambda (\Lambda_b + 2(1 - \Lambda_b)) G(r, \theta) \exp(-i\chi p \Theta) \quad (9)$$

where  $\chi(r, \Lambda) = \frac{\sigma \pi \sqrt{\kappa} \sqrt{\varepsilon \Lambda / 2}}{K(\kappa^{-1})}$ ,  $\kappa = \frac{1 - \Lambda(1 - \varepsilon)}{2\varepsilon \Lambda}$ .  $K$  denotes the complete elliptic integral of the first kind.  $\Lambda_b = \frac{\Lambda}{b}$ ,  $B_\Lambda(r, \Lambda, \theta) = \frac{1}{b\sqrt{(1 - \Lambda_b)}}$ ,  $b = B/B_0 = 1 + (r/R_0)\cos\theta$ .

$$G = (g^{\theta\theta} \kappa_\theta + g^{r\theta} \kappa_r) \xi_\theta(\hat{r}(\bar{r}, \rho_d, \theta), \theta) + (g^{rr} \kappa_r + g^{r\theta} \kappa_\theta) \xi_r(\hat{r}(\bar{r}, \rho_d, \theta), \theta) \quad (10)$$

where  $\hat{r} = \bar{r} + \rho_d \cos\theta$ .

$$\rho_d = \frac{q}{2} \rho_0 \sqrt{\frac{E}{1 - \Lambda/b}} \left[ \frac{\Lambda}{b} + 2(1 - \frac{\Lambda}{b}) \right] = \frac{q}{2} \rho_0 \sqrt{\frac{E}{1 - \Lambda_b}} [\Lambda_b + 2(1 - \Lambda_b)] \quad (11)$$

To simplify the code, we use  $\Lambda_0$  instead of  $\Lambda$  in  $\rho_d$ .  $g^{rr} = 1 + \frac{\varepsilon \cos\theta}{2}$ ,  $g^{\theta\theta} = \frac{1}{r^2} (1 - \frac{5}{2}\varepsilon \cos\theta)$ ,  $g^{r\theta} = -\frac{3}{2r}\varepsilon \sin\theta$ .

$$\Theta(\theta, r, \Lambda) = \int_0^\theta d\theta' \frac{1}{b\sqrt{(1 - \Lambda_b)}} = \int_0^\theta B_\Lambda \quad (12)$$

## 2.2 The mode structure

$$\xi_r(r, \theta) = \xi_0(r) \exp(-i\theta), \quad \xi_\theta(r, \theta) = -i\xi_0(r) r \exp(-i\theta)$$

$$\xi_0(r) = \begin{cases} 1 & r \leq r_s - \Delta r/2 \\ \frac{\Delta r - 2r + r_s - \Delta r/2}{\Delta r} & r_s - \frac{\Delta r}{2} < r < r_s + \frac{\Delta r}{2} \\ 0 & r \geq r_s + \frac{\Delta r}{2} \end{cases}$$

## 2.3 The normalized quantities used for $\delta W_k$ :

$v_0 = \sqrt{2T_h/M}$ ,  $T_h$  is the fast ions injection energy,  $M$  is the fast ion's mass.  $\omega_0 = \frac{v_0}{R_0}$ .  $F_0 = \frac{n_0}{v_0^3}$ ,  $n_0$  is the fast ion density at axis.  $r_0 = a$  is the minor radius,  $\varepsilon = a/R_0$ ,  $E_0 = T_h/M$ ,  $B_0$  is the torodial magnetic field at magnetic axis.

## 3 Fishbone dispersion relation

The dwk++ code calculate the fishbone dispersion relation:

$$\delta W_{mhd} + \delta W_k + \delta I = 0 \quad (13)$$

Assume  $\delta W_{mhd} = 0$ , the equation reduced to:

$$\frac{4}{\pi} \left( \frac{r_s}{R_0} \right)^2 \left| \frac{\xi_s}{\xi_0} \right|^2 \left( -i \frac{\omega}{\omega_A} \right) + \beta_h \delta \bar{W}_k = 0 \quad (14)$$

$$\omega = -iC\beta_h \delta W_k \quad (15)$$

$$i\omega = C\beta_h \delta W_k \quad (16)$$

$$C = \frac{\omega_A}{\omega_0} \frac{1}{\frac{4}{\pi} \left( \frac{r_s}{R_0} \right)^2 \left| \frac{\xi_s}{\xi_0} \right|^2} \quad (17)$$

$$\omega_A = \frac{2}{\tau_{A\theta s}} \quad (18)$$

$$\tau_{A\theta} = \frac{3^{1/2} r_s}{(B_{\theta s}^2 / \mu_0 \rho_m)^{1/2}} \quad (19)$$

$$s = r_s \frac{dq}{dr(r = r_s)} \quad (20)$$

$$\beta_h = 2\mu_0 n_0 T_h / B_t^2 \quad (21)$$

## 4 How to run dwk++

### 4.1 Compile the code

dwk++ code is using c++ language as the main language, it need a c++ compiler to compile the code, gcc/g++ on max os/linux and intel compiler on linux was tested. To compile the code, libconfig with version >1.5 is needed, and set environment variable LIBCONFIG\_DIR to the path where libconfig located.

### 4.2 Input file

dwk.cfg:

```
/*the input file for dwk++*/
//tokamak parameters
tokamak=
{
    a=0.4;           //minor radius (m).
    R0=1.65;         //major radius (m).
    Bt =1.34;        //Toroidal magnetic field at axis without plasma, (Tesla).
    n0 =1.31e19;     //thermal plasma density, (m^-3).
    mi =2.0;         //ion mass, (protom mass,m_p).
    E_i0 =41.0;      //Fost ion injection energy, (KeV).
    m_ep =2.0;       //fat ion mass
    //q profile, polynomial coefficient, qc[5];
    //q=qc[0] +qc[1]*r +qc[2]*r^2 +qc[3]*r^3 +qc[4]*r^4;
    qc=[0.5, 2.0, 0.0, 0.0, 0.0];
}
//grid parameters
grid=
{
    nx=400; //grid size should be 3n + 1, n is a positive integer.
    nL=100;
    nE=1000;
    ntheta=400;
}
//fast ion distribution
slowing=
{
    r0=0.0;          // (a)
    rd=0.2;          // (a)
    L0=0.28;
    Ld=0.1;
    E0=1.0;           // (E_i0)
    Ed=0.3; // (E_i0)
    Ec=0.25;          // (E_i0)
    rho_d=0.00;       // drift orbit width, (a)
    sigma=1;          // co(1) ? count (-1)
}
//perturbation and omega range for the souldtion
mode=
{
    n=1;              //toroidal mode number
    m=1;              //poloidal mode number
    pa=0;             //resource: sum Yps/(n*omega_phi +p*omega_theta -\omega), sum from pa to pb
    pb=0;
}
```

```

r_s=0.5; //q==1 surface.
q_s=1.0; //resonance surface
delta_r=0.001; // step function width for kink, (a).
omega_0=0.1; //find omega_0(the solution) and scan dwk between omega_0 and omega_1,
// (v_0/R_0), v_0 is the injection speed.
omega_1=1.1;
omega_i=0.002; //image part of omega
omega_n=120; //scan steps
omega_err=1.0e-5; //residual of omega_0
max_iter =100; // maximum iteration number to find the omega_0
max_iterg =2;
dw_f=-0.00; //dw_mhd
}

```

### 4.3 Output file

### 4.4 Utilities to plot results

## 5 Examples

$q$  profile:

$$\begin{aligned}
q &= 0.5 + 2r^2 \\
r_s &= 0.5a = 0.2m \\
s &= 1 \\
R_0 &= 1.65m \\
B_{t0} &= 4T \\
B_{\theta s} &= \frac{r_s B_t}{R q_s} = 0.6T \\
n_0 &= 1 \times 10^{20} m^{-3} \\
\rho_m &= 3.3452 \times 10^{-7} kg m^{-3} \\
\tau_{A\theta} &= 4.6328 \times 10^{-7} s \\
\omega_A &= 4.317 \times 10^6 \\
T_h &= 40KeV \\
v_h &= 1.9575 \times 10^6 m/s \\
\omega_0 &= 1.1863 \times 10^6 \\
C &= 28.2486
\end{aligned}$$

$$\frac{\pi B_{\theta s}^2}{\mu_0} R_0 |\xi_s|^2 \left( \frac{-i\omega}{\omega_A} \right) + \frac{1}{2} \pi^2 |\xi_0|^2 R_0 n_0 T_h \delta W_k = 0 \quad (22)$$

$$\frac{\pi B_{\theta s}^2}{\mu_0} |\xi_s|^2 \left( \frac{-i\omega}{\omega_A} \right) + \frac{1}{2} \pi^2 |\xi_0|^2 n_0 T_h \delta W_k = 0 \quad (23)$$

$$\frac{B_{\theta s}^2}{\mu_0} |\xi_s|^2 \left( \frac{-i\omega}{\omega_A} \right) + \frac{1}{2} \pi |\xi_0|^2 n_0 T_h \delta W_k = 0 \quad (24)$$

$$\frac{B_{\theta s}^2}{\mu_0} |\xi_s|^2 \left( \frac{i\omega}{\omega_A} \right) = \frac{1}{2} \pi |\xi_0|^2 n_0 T_h \delta W_k \quad (25)$$

$$\frac{B_{\theta s}^2}{\mu_0} |\xi_s|^2 \left( \frac{\omega}{\omega_A} \right) = -\frac{1}{2} i \pi |\xi_0|^2 n_0 T_h \delta W_k \quad (26)$$

$$\frac{\omega}{\omega_A} = \frac{-\frac{1}{2} i \pi |\xi_0|^2 n_0 T_h \delta W_k}{\frac{B_{\theta s}^2}{\mu_0} |\xi_s|^2} \quad (27)$$

$$\frac{\omega}{\omega_A} \frac{\omega_0}{\omega_0} = \frac{-\frac{1}{2} i \pi |\xi_0|^2 n_0 T_h \delta W_k}{\frac{B_{\theta s}^2}{\mu_0} |\xi_s|^2} \quad (28)$$

$$\frac{\omega i}{\omega_0} = \frac{\frac{1}{2} \pi |\xi_0|^2 n_0 T_h}{\frac{B_{\theta s}^2}{\mu_0} |\xi_s|^2} \frac{\omega_A}{\omega_0} \delta W_k \quad (29)$$

$$\frac{\omega i}{\omega_0} = \frac{1}{2} \frac{\pi \frac{|\xi_s|^2}{|\xi_0|^2}}{B_{\theta s}^2} \mu_0 n_0 T_h \frac{\omega_A}{\omega_0} \delta W_k \quad (30)$$

$$\frac{\omega i}{\omega_0} = \frac{\pi}{2} \frac{|\xi_s|^2}{|\xi_0|^2} \frac{\mu_0 n_0 T_h}{B_{\theta s}^2} \frac{\omega_A}{\omega_0} \delta W_k \quad (31)$$

$$\frac{\omega i}{\omega_0} = \frac{\pi}{2} \frac{|\xi_s|^2}{|\xi_0|^2} \frac{\mu_0 n_0 T_h R_0 q_s}{B_t^2 r_s^2} \frac{\omega_A}{\omega_0} \delta W_k \quad (32)$$

$$\frac{\omega i}{\omega_0} = \frac{\pi}{4} \frac{|\xi_s|^2}{|\xi_0|^2} \frac{2 \mu_0 n_0 T_h R_0^2 q_s^2}{B_t^2 r_s^2} \frac{\omega_A}{\omega_0} \delta W_k \quad (33)$$

$$\frac{\omega i}{\omega_0} = \frac{\pi}{4} \frac{|\xi_s|^2}{|\xi_0|^2} \frac{R_0^2 q_s^2}{r_s^2} \frac{\omega_A}{\omega_0} \beta_h \delta W_k \quad (34)$$

$$\frac{\omega i}{\omega_0} = C \beta_h \delta W_k \quad (35)$$

$$C = \frac{\pi}{4} \frac{|\xi_s|^2}{|\xi_0|^2} \frac{R_0^2 q_s^2}{r_s^2} \frac{\omega_A}{\omega_0}$$