Dimensional transition from Kitaev spin liquid to decoupled fermion chains

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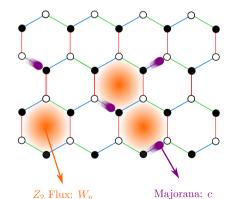


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Introduction

A. Kitaev. Ann. Phys. 321, 2–111 (2006)

$$\mathcal{H} = \sum_{\langle ij \rangle} \mathbf{K}_{x} S_{i}^{x} S_{j}^{x} + \mathbf{K}_{y} S_{i}^{y} S_{j}^{y} + \mathbf{K}_{z} S_{i}^{z} S_{j}^{z}$$



Relevant QSL materials:

- α−RuCl₃
- Y. Kasahara, T. Ohnishi, Y. Mizukami, O. Tanaka, Sixiao Ma, K. Sugii, N. Kurita, H. Tanaka, J. Nasu, Y. Motome, T. Shibauchi and Y. Matsuda. Nature 559, 227–231 (2018)

Large in-plane field to destabilize order

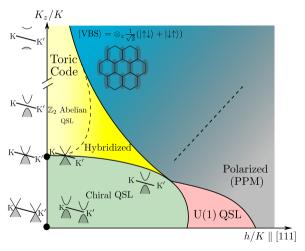
- BaCo₂(AsO₄)₂
 - X. Zhang, Y. Xu., T. Halloran, R. Zhong, C. Broholm, R. J. Cava, N. Drichko, N. P. Armitage. A magnetic continuum in the cobalt-based honeycomb magnet BaCo2(AsO4)2. Nat. Mater. 22, 58–63 (2023)

Small [111] field to destabilize order

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Phase diagram under out-of-plane [111] field

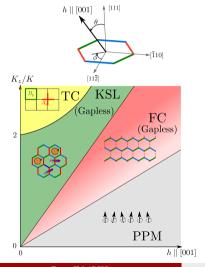


- N. Patel and N. Trivedi. Magnetic field-induced intermediate quantum spin liquid with a spinon Fermi surface. PNAS 116, 12199 (2019)
- S. Feng, A. Agarwala, S. Bhattacharjee and N. Trivedi. Anyon dynamics in field-driven phases of the anisotropic Kitaev model, arXiv:2206.12990

The simplest [001] field?

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Phase diagram under [001] field and exchange anisotropy



$$\mathcal{H} = \sum_{i} K(\sigma_{i}^{x} \sigma_{i+x}^{x} + \sigma_{i}^{y} \sigma_{i+y}^{y}) + \underline{K_{z}} \sum_{i} \sigma_{i}^{z} \sigma_{i+z}^{z} - \underline{h} \sum_{i} \sigma_{i}^{z}$$

Four phases in (K_z, h) plane:

- Gapless Kitaev spin liquid (KSL) at small h/K_z
- ② Decoupled/Weakly coupled fermion chains (FC) at intermediate h/K_z
- **3** Toric Code (TC) at large K_z
- Partially polarized magnet (PPM) at large h/K_z as Emergent decoupled boson chains

Discussion by ED, DMRG, MFT, Effective field theory.



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Phase diagram (DMRG & ED)

$$\chi = \frac{\partial^2 E_{\rm gs}}{\partial h^2}, \quad \langle W_p \rangle = \langle \sigma_1^x \sigma_2^y \sigma_3^z \sigma_4^x \sigma_5^y \sigma_6^z \rangle. \quad \langle W_p \rangle = 1 \text{ at } h = 0$$

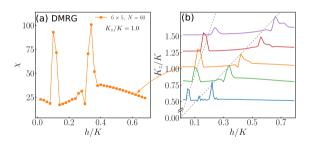


Figure: Magnetic susceptibility χ by DMRG

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Phase diagram (DMRG & ED)

$$\chi = \frac{\partial^2 E_{\text{gs}}}{\partial h^2}, \quad \langle W_p \rangle = \langle \sigma_1^x \sigma_2^y \sigma_3^z \sigma_4^x \sigma_5^y \sigma_6^z \rangle. \quad \langle W_p \rangle = 1 \text{ at } h = 0$$

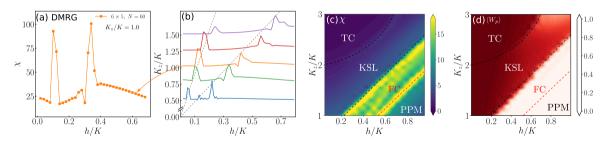
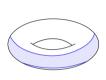


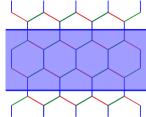
Figure: Magnetic susceptibility χ by DMRG

Figure: χ and Z_2 flux expectation by 24-site ED

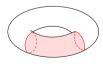
Entanglement entropy

z - bond cut





y - bond cut





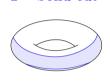
von-Neumann entropy per α bond $S_{\nu N}^{\alpha}$:

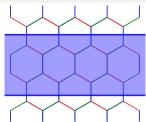
$$S_{
u N}^{lpha} = -rac{1}{|\partial A_{lpha}|} \operatorname{Tr}
ho_{A_{lpha}} \ln
ho_{A_{lpha}}, \; lpha \in \{z,y\}$$

 A_{α} : subsystem whose edges cut α bonds

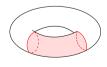
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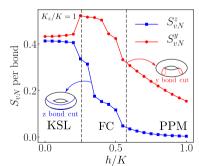




von-Neumann entropy per α bond S_{nN}^{α} :

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ho_{A_{lpha}} \ln
ho_{A_{lpha}}, \ lpha \in \{z,y\}$$

 A_{α} : subsystem whose edges cut α bonds



Decoupled bosonic chains in PPM

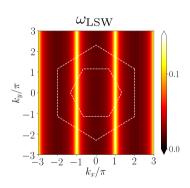


Figure: LSW at $h \sim 0.5$

$$\mathcal{H}_{LSW} = \frac{1}{2} \sum_{\mathbf{k}} \Psi_{\mathbf{k}}^{\dagger} \mathbf{H}(\mathbf{k}) \Psi_{\mathbf{k}}, \ \mathbf{H}(\mathbf{k}) \equiv \begin{pmatrix} \mathbf{M}(\mathbf{k}) & \mathbf{N}(\mathbf{k}) \\ \mathbf{N}^{\dagger}(\mathbf{k}) & \mathbf{M}(-\mathbf{k}) \end{pmatrix}$$

where we defined $\Psi_{\mathbf{k}} \equiv (a_{\mathbf{k}}, b_{\mathbf{k}}, a_{-\mathbf{k}}^{\dagger}, b_{-\mathbf{k}}^{\dagger})^{\mathrm{T}}$ with a, b boson operators of A and B sublattices, and $\mathbf{N}(\mathbf{k})$ and $\mathbf{M}(\mathbf{k})$ are:

$$\mathbf{N} = \frac{1}{4} \left(K_x e^{i\mathbf{k} \cdot \mathbf{n}_1} - K_y e^{i\mathbf{k} \cdot \mathbf{n}_2} \right) \sigma^x$$

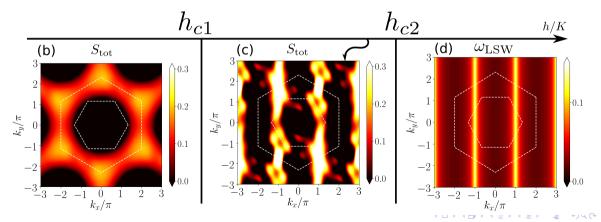
$$\mathbf{M} = \left(\frac{\mathbf{n}}{2} - \frac{1}{2} K_z \right) \sigma^z + \frac{1}{4} \left(K_x e^{i\mathbf{k} \cdot \mathbf{n}_1} + K_y e^{i\mathbf{k} \cdot \mathbf{n}_2} \right) \sigma^x$$

Near $h_{c2} \sim 0.5$, K_z exchange on z bonds vanishes in PPM! \rightarrow Decoupled boson chains.

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Structure Factor

 $S_{\mathrm{tot}}(\mathbf{k}) = \langle \mathbf{S}_i(\mathbf{k}) \cdot \mathbf{S}_j(-\mathbf{k}) \rangle - \langle \mathbf{S}_i(\mathbf{k}) \rangle \cdot \langle \mathbf{S}_j(-\mathbf{k}) \rangle, \quad \omega_{\mathrm{LSW}}(\mathbf{k}) : \text{ Linear spin wave of PPM}$



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Weakly Coupled fermionic chains in FC

Kitaev's four majorana decomposition: $\sigma^{\alpha} = ib^{\alpha}c$. Convert into canonical fermions and bond fermions:

Hence

$$K_{z}\left(b_{i,A}^{z}b_{i,A+\hat{z}}^{z}c_{i,A}c_{i,A+\hat{z}}\right) = K_{z}(2n_{i}^{f}-1)(1-2n_{i}^{z})$$

$$K_{x}\left(b_{i,A}^{x}b_{i,A+\hat{x}}^{x}c_{i,A}c_{i,A+\hat{x}}\right) = K_{x}(1-2n_{i}^{x})$$

$$\times (f_{i}f_{i-\delta_{1}} - f_{i}f_{i-\delta_{1}}^{\dagger} + f_{i}^{\dagger}f_{i-\delta_{1}} - f_{i}^{\dagger}f_{i-\delta_{1}}^{\dagger})$$

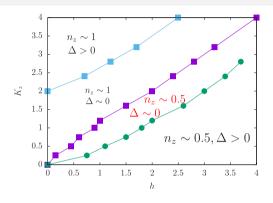


Figure: MFT phase diagram

 K_z exchange on z bonds vanishes in the MFT of FC near h_{c2}^- !



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Possible effective Theory in FC

Ignoring interchain coupling K_z , \mathcal{H} becomes that of a compass chain, which can be mapped to the critical point of *TFIM*:

$$\mathcal{H} = \sum_{i} \sigma_{i}^{x} \sigma_{i+1}^{x} + \sigma_{i+1}^{y} \sigma_{i+2}^{y} \quad \rightarrow \quad \mathcal{H}_{\mathrm{TFIM}}^{c} = -\sum_{i} \tau_{i}^{z} \tau_{i+1}^{z} - \sum_{i} \tau_{i}^{x}$$

 $\mathcal{H}_{\mathrm{TFIM}}^{c}$: 1+1D CFT with central charge $c=\frac{1}{2}$, with left and right chiral majoranas $\gamma_{L(R)}$.

$${\cal H}_L = \int dx (i\gamma_L \partial_x \gamma_L), \; {\cal H}_R = \int dx (-i\gamma_R \partial_x \gamma_R)$$

Consider weak coupling between chains via coupling between γ_L and γ_R :

$$\mathcal{H}(FC) \approx \mathcal{H}_L + \mathcal{H}_R - g \int dx (\gamma_L \partial_x \gamma_L) (\gamma_R \partial_x \gamma_R) \rightarrow \text{stable until finite } g_c$$

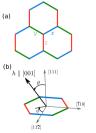
A. Rahmani, X. Zhu, M. Franz, and I. Affleck, Phys. Rev. Lett. 115, 166401 (2015)

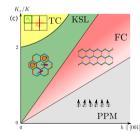
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Summary

- lacktriangle Weakly coupled fermionic chains from Kitaev model under [001] field $h\sim h_{c2}^-$
- ② Decoupled bosonic chains from Kitaev model under large [001] field $\sim h_{c2}^+$
- **3** Liftshitz transitions by depleting n_i^z bond fermions from 1 to 0.5
- **①** Proposed effective theory in FC: $\mathcal{H}(FC) \approx \mathcal{H}_L + \mathcal{H}_R g \int dx (\gamma_L \partial_x \gamma_L) (\gamma_R \partial_x \gamma_R)$





Outlook:

- How is (or is not) the FC phase connected to the proposed spinon Fermi surface under [111] field?
- ② Dynamics?

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