

Nonlinear diffusion equation with finite elements

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About the problem

In this project we will solve a nonlinear diffusion equation using the finite element method and the FEniCS software package. More precicely we will be looking at the equation

$$\rho u_t = \nabla \cdot (\alpha(u) \nabla u) + f(\mathbf{x}, t) \quad (1)$$

where ρ is a constant, $\alpha(u)$ is a known function of u , and where we have the initial condition $u(\mathbf{x}, 0) = I(\mathbf{x})$, and Neuman boundary conditions $\frac{\partial u}{\partial n} = 0$.

Discretization

In order to solve equation (1) we will need to discretize it. The standard approach is then to pick a nummerical approximation to the derivatives. We will use a finite difference discertization in time, more specifically the Backward Euler (BE) discretization, and a finite element approximation in space. The BE discertization gives us

$$\begin{aligned} \rho u^n &= \Delta t \nabla \cdot (\alpha(u) \nabla u) + \Delta t f(\mathbf{x}, t) + \rho u^{n-1} \\ \rho u^n - \Delta t \nabla \cdot (\alpha(u) \nabla u) - \Delta t f(\mathbf{x}, t) - \rho u^{n-1} &= R \end{aligned}$$

we now approximate R on a functionspace $V = \text{span}\{\phi_i\}$, where ϕ_i denotes the P1 elements, and minimize the error by demanding

$$(R, v) = 0 \quad \forall v \in V$$

which gives us

$$\rho(u^n, v) - \Delta t(f(\mathbf{x}, t_n), v) - \rho(u^{n-1}, v) - \Delta t(\nabla \cdot (\alpha(u) \nabla u^n), v) = 0$$

The first three terms are ok for now, but the last term has a double derivative in u , which is something we dont want seeing as we are trying to approximate u by P1 elements as $u \simeq \sum_k u_k \phi_k$, and the double derivative of a linear function is simply zero. To get around this problem we try to integrate by parts

$$\begin{aligned} \Delta t(\nabla \cdot (\alpha(u) \nabla u^n), v) &= \Delta t \left(\int_{\Omega} \nabla \cdot (\alpha(u) \nabla u^n) v dx \right) \\ &= \Delta t \left([\alpha(u) \nabla u v]_{\partial \Omega} - \int_{\Omega} \alpha(u) \nabla u^n \nabla v dx \right) = -\Delta t \int_{\Omega} \alpha(u) \nabla u^n \nabla v dx \end{aligned}$$

where we have inserted for the boundary conditions $\frac{\partial u}{\partial n} = \mathbf{n} \cdot \nabla u = 0$ on $\partial \Omega$. We have now arrived on a linear variational problem to be solved at each timestep

$$\begin{aligned} \rho(u^n, v) - \Delta t(f(\mathbf{x}, t_n), v) - \rho(u^{n-1}, v) + \Delta t(\alpha(u) \nabla u^n, \nabla v) &= 0 \\ \underbrace{\rho(u^n, v) + \Delta t(\alpha(u) \nabla u^n, \nabla v)}_{a(u, v)} &= \underbrace{\Delta t(f(\mathbf{x}, t_n), v) + \rho(u^{n-1}, v)}_{L(u, v)} \end{aligned}$$

Notice that we have a nonlinearity in $\alpha(u)$. There are several way to get around this nonlinearity, we will use Picard iteration meaning that we start out with $\alpha(u^{n-1})$ and solve the linear system to get a solution \tilde{u}^n . Then we update $\alpha(\tilde{u}^n)$ and solve the system again. We keep dionig this untill $\|\tilde{u}^n - \tilde{u}^{n-1}\| \leq \epsilon$ where ϵ is a predefined

tolerance. Hopefully this converges to the correct solution, and we can set our $\tilde{u}^n = u^n$. Picard iteration leads us to the following linear system for each iteration

$$\rho \int_{\Omega} u^n v dx + \Delta t \int_{\Omega} \alpha(u^{n-1}) \nabla u^n \nabla v dx = \rho \int_{\Omega} u^{n-1} v dx + \Delta t \int_{\Omega} f(x, t_n) v dx$$

we will now insert for $u \simeq \sum_k u_k \phi_k$ and because we are working with an undefined number of cells of potentially variable size we will transform all the integrals to a reference cell where $-1 \leq X \leq 1$. We will also do the calculations in 1 dimension, but they easily generalize to multiple dimensions by simply setting each element in say the vector \mathbf{u}^n equal to a vector. The same approach gives us a block matrix for M and K.

$$\rho \sum_{i=1}^{N_x} \int_{-1}^1 (\phi_i \phi_j) u_i^n dx + \Delta t \sum_{i=1}^{N_x} \int_{-1}^1 \alpha(u_i^{n-1}) (\phi_i' \phi_j') u^n dx = \rho \sum_{i=1}^{N_x} \int_{-1}^1 (\phi_i \phi_j) u_i^{n-1} dx + \Delta t \sum_{i=1}^{N_x} \int_{-1}^1 f_i(x, t_n) \phi_j dx$$

If we now set $M_{ij} = \rho \int_0^L (\phi_j \phi_i) dx$ and $K_{ij} = \Delta t \int_0^L \alpha(u^{n-1}) (\phi_j' \phi_i') dx$ we get

$$M_{ij} = \int_0^L (\phi_j \phi_i) dx \implies M_{00} = \int_{-1}^1 (\phi_0 \phi_0) dx$$

We are using P1 elements, and so the only nonzero integrals we will get are when $i = j$ and $i \pm 1 = j$, meaning M and K will become tridiagonal.
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we are left with

$$M \mathbf{u}^n + K(u^{n-1}) \mathbf{u}^n = M \mathbf{u}^{n-1} + \mathbf{b} \\ \implies \mathbf{u}^n = (M + K(u^{n-1}))^{-1} (M \mathbf{u}^{n-1} - \mathbf{b})$$

to be solved for each Picard iteration. The crudest Picard iteration will be just one iteration, which is the same as solving the system using $\alpha(u^{n-1}) \simeq \alpha(u^n)$ and this is our chosen approach.

Group Finite Element method

We can also approximate the nonlinear term by the group finite element method. Looking at the variational form of our equation and inserting for $u \simeq \sum_{i=1}^N \phi_i u_i$ we get the following

$$\rho \int_{\Omega} \sum_i (\phi_i \phi_j) u_i^n d\Omega + \Delta t \int_{\Omega} \alpha \left(\sum_i \phi_i u_i^n \right) \sum_j (\nabla \phi_j \nabla \phi_k) u_j^n d\Omega = \rho \int_{\Omega} \sum_i (\phi_i \phi_j) u_i^{n-1} d\Omega + \Delta t \int_{\Omega} f(x, t_n) v d\Omega$$

If we now set $\alpha(\sum_i \phi_i u_i^n) \approx \sum_i \phi_i \alpha(u_i^n)$ and focus on the integral with α , we are left with (in one dimension)

$$\int_{\Omega} \sum_i \phi_i \alpha(u_i^n) \sum_j (\phi_j' \phi_k') u_j^n dx$$

Which we transfer to a reference element $X \in [-1, 1]$ where $\frac{d\phi}{dX} = \frac{d\phi}{dx} \frac{dx}{dX} = \frac{d\phi}{dX} \frac{2}{h}$ and $\frac{dx}{dX} = \frac{h}{2}$. These calculations will result in an element matrix which we can assemble so that we get a linear system

$$\frac{(-1)^r (-1)^s}{h} \int_{-1}^1 \alpha(u_i) \phi_i \phi_j' \phi_k' dX \\ \approx \frac{(-1)^r (-1)^s}{2h} \left(\underbrace{\sum_r \alpha(u_{i+r}) \cdot \phi_{i+r}}_{\neq 0 \text{ when } r=1} + \underbrace{\sum_s \alpha(u_{i+s}) \cdot \phi_{i+s}}_{\neq 0 \text{ when } s=0} \right) \\ \implies \frac{1}{2h} (\alpha(u_{i+1}) + \alpha(u_i)) \cdot \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} = K^{(i)}$$

The entire linear system is then

$$M \mathbf{u}^n + \Delta t K \mathbf{u}^n = M \mathbf{u}^{n-1} + \mathbf{b}$$

where M is the same matrix as in (SOMEREFERENCE). This means that the i'th equation in the system is

Newton's Method

As a different approach we can use Newton's method to estimate the nonlinear term in our equation (1). Newton's method is quite simply an iterative method where we linearize our problem through a first order Taylor expansion, and can be expressed as

$$\mathbf{F}(\mathbf{u}) = 0 \simeq M(u; u^q) = F(u^q) + J(u - u^q)$$

where $J = \nabla F$ which means that $J_{i,j} = \frac{\partial F_i}{\partial u_j}$ and $\mathbf{F}(\mathbf{u})$ is the equation we want to solve only written on a general form. If we reformulate a bit we find

$$\begin{aligned} J(u^{q+1} - u^q) &= -F(u^q) \\ u^{q+1} &= u^q - J^{-1}F(u^q) \end{aligned}$$

which is the very same formulation as in 1D $x^{k+1} = x^k - \frac{f(x^k)}{f'(x^k)}$. We wish to apply this on the variational form of the equation which ends up as our F : (notice that we set $\rho = 1$)

$$\begin{aligned} F_i(u) &= \int_{\Omega} (u_i^n v + \Delta t (\alpha(u_i^n) \nabla u_i^n \nabla v) - u_i^{n-1} v - \Delta t f(\mathbf{x}, t) v) d\Omega \\ J_{i,j} &= \frac{\partial}{\partial u_j^n} F_i \\ J_{i,j} &= \int_{\Omega} (\phi_i \phi_j + \Delta t (\alpha'(u_i^n) \phi_i \nabla u_i^n \nabla \phi_j + \alpha(u_i^n) \nabla \phi_i \nabla \phi_j)) d\Omega \end{aligned} \quad (2)$$

where we have used $u_i^n = \phi_i u_i^n \implies \frac{\partial}{\partial u_j^n} u_i^n = \phi_i$, $v = \phi_j$ and $\frac{\partial}{\partial u_j^n} \alpha(u_i^n) = \alpha'(u_i^n) \phi_i$. And there we have the expressions for the entries in the jacobian matrix we can use for a Newton's method on our equation. The remaining steps are to assemble the linear problem, and iterate in the same way as we did with Picard iteration. These are both tasks well suited for the FEniCS software.

Simplifying (2) to one spatial dimension lets us compute the actual entries in the jacobian in a simple manner so that we can compare with something else.

Implementation

To make the program as general as possible, we will make it dimension independent by providing the number of spatial points on the commandline as n_x, n_y, n_z and choose a unit interval, square or cube accordingly. Also, even though we will only use one Picard iteration, we will implement this as a function where the maximum number of iterations is given as an argument. The resulting code can be found in the appendix.

Verification

For the first verification of the program we will assume $\rho = \alpha(u) = 1$, $f(\mathbf{x}, t) = 0$ and $I(\mathbf{x}) = \cos(\pi x)$. We will work on the domain $\Omega = [0, 1] \times [0, 1]$. This should give an analytic solution of $u(x, y, t) = e^{-\pi^2 t} \cos(\pi x)$

Errors

As in any (numerical) approximation there will be an error in our solution. We can distinguish two types of errors, there is the error we do in the approximation of the derivatives, and there is the numerical error.

Our approximation using the Backward Euler scheme has an error which goes like $\mathcal{O}(\Delta t)$ for other reasons the spatial error goes like $\mathcal{O}(\Delta x^2)$. Other potential errors which will arise in the FEniCS program are errors due to loss of precision in the numerical integration we do upon assembly of the matrices M and K . Depending on what kind of numerical integration we use, the error from integration could be $\mathcal{O}(\Delta x)$ for the rectangle (midpoint) rule or it could even be zero if the solution is a polynomial of degree $2n - 1$ and we integrate by some form of Gaussian quadrature. n is here the order of elements used. There is always a possibility of losing numerical precision because of adding a small and a large number since we only have 16 valid decimals. In our case we also have an obvious error in the approximation of the nonlinear term.

Analytic solution

Results

Final comments