

Gauge ① IGAP

Invariants of mfd's from non-gauge theories - Du Pei

- why?

- some cons of gauge theories:

1) seemingly, no invariants stronger
than SW, as most can be
related to those ones

2) won't give full TQFTs,
not defined on some mfd's

- Pave's talk: 6d non-gauge th \Rightarrow 4d invs.

- this talk: 4d non-gauge \Rightarrow 3d TQFTs

- 4d $N=2$ theories have Coulomb branches

M_T : Coul. br. of theory T on $S^1 \times \mathbb{R}^3$

- hyperkähler

\Rightarrow



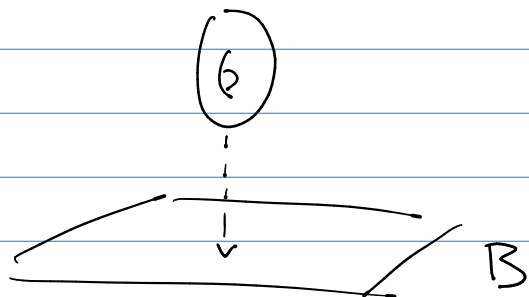
S^2 of cpx str's,
pick I

$F \hookrightarrow M_T$
 \uparrow
cpt., generically
abelian var's

B_T : Coul. br. on \mathbb{R}^4 (u-plane)
- special Kähler, w singularities
 $G = \mathbb{C}^*$

- examples: ① $U(1)$ gauge th, $N=0$
 \rightarrow labelled by an elliptic
curve Σ

$$\mathcal{M}_Y = \mathbb{T}^2 \times \mathbb{R}^2 = \text{Jac}(\Sigma) \times \mathbb{C}$$



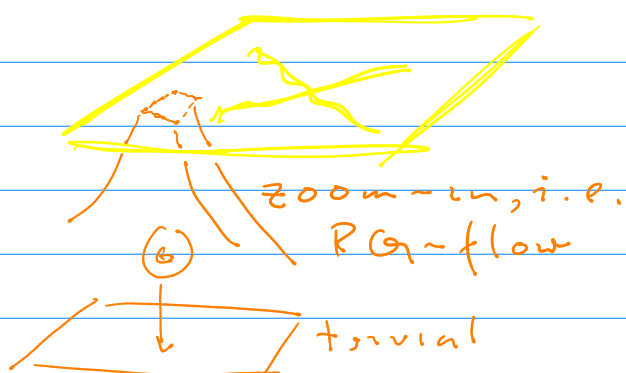
(11) $U(1)^g$ -gauge th, labelled by Σ_g

$$\mathcal{M}_Y = (\mathbb{T}^2 \times \mathbb{R}^2)^g = \text{Jac}(\Sigma_g) \times \mathbb{C}^g$$

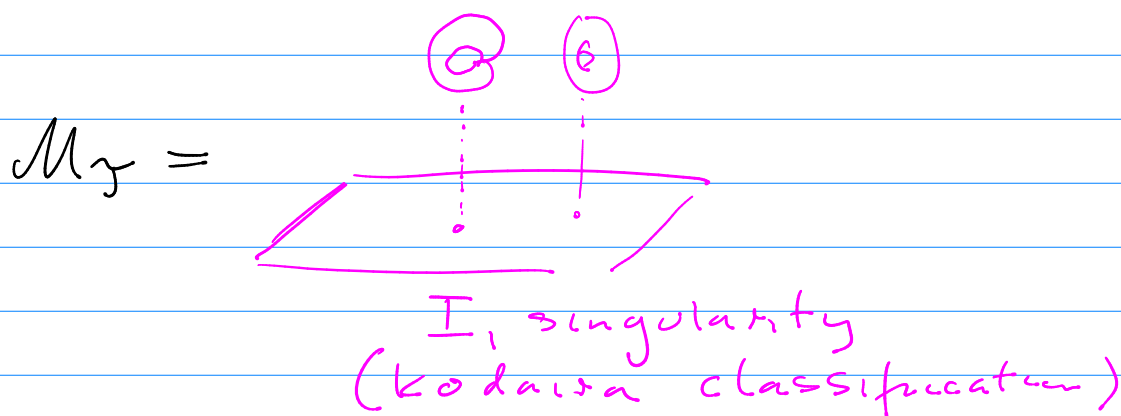
$\pi \downarrow$

$$B_Y = \mathbb{C}^g \simeq H^1(\Sigma, \mathbb{C})$$

-trivial fibration $\Leftrightarrow Y$ is a "free theory"
 \rightarrow but locally, we have triviality,
 and, physically, Y near a
 nonsingular point can be
 approximated by a free $U(1)^g$
 gauge theory (low-energy effective th.
 LBET)

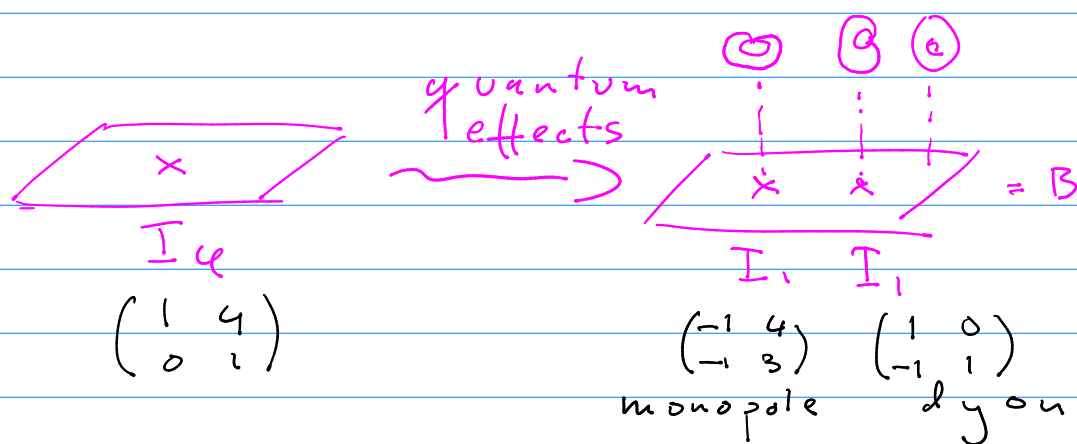


- (III) "monopole" theory
- $G = \mathbb{C}^*$, $N = \mathbb{C}$
 - SQED w 1 electron



- loop around I_1 gives $\begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \in SL_2(\mathbb{Z})$
monodromy

- (IV) $SU(2)$ pure SYM, $G = SU(2)$, $N = 0$



(III) \Rightarrow SW invs

(IV) \Rightarrow Donaldson $\xrightarrow[\text{conj.}]{\text{Witten's}}$ SW + SW'
+ u-plane

- Remark • (I), (II), (III) possess \mathcal{S}' -action (rotate $B\gamma$)
 \rightarrow look like conformal theories
- (IV) $M_\gamma \not\subset \mathcal{S}'$
 \rightarrow nonconf. (quantum anomaly of $U(1)_R$ -symmetry)
- superconformal theories:

$$\begin{array}{ccc} \mathcal{S}' & \subset & M_\gamma \\ \pi \downarrow & & \pi \text{ equivariant} \\ G & \subset & B\gamma \end{array}$$

$$-(B\gamma)^{\mathcal{S}'} = \{pt\} \quad \text{superconf. pt. } 0$$

(V) class S , labels: $g: ADE, \Sigma: RS$

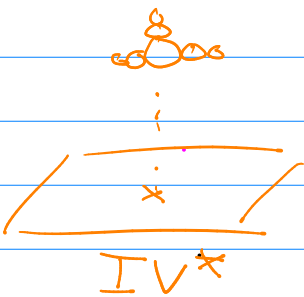
$M_\gamma = \text{m.s.p. of } G\text{-Higgs bds on } \Sigma$

π : Hitchin fibration

\mathcal{S}' : Hitchin action

- if $g = \mathbb{C} \leadsto$ (I), (II)

(VI) $g = sl(3, \mathbb{C})$, $\Sigma = \mathbb{P}^1$ w/ 3 mld pts

$$M_{\gamma} =$$


- away from sing.: $M.N(E_6)$

- can $\$'$ -action help? a little

- Hilbert sp. of γ (top. twisted)
on $\$^3 \cong \{ \text{regular functions on } B_{\gamma} \}$

- partition funct of γ on $\$' \times \$^3 \cong \dim = \infty$
 $\rightarrow \$'$ -equiv. part. funct.,
 though, $\cong \chi_{\$'}(G_{B_{\gamma}})$

- free theory $\leadsto B_{\gamma} = \mathbb{C}$
 $\mathcal{H}_{\gamma}(\$^3) \cong \langle 0, z, z^2, \dots \rangle$
 $\chi_{\$'} = 1 + t + t^2 + \dots = \frac{1}{1-t}$

- can we regularize $\$' \times M^3$ part. function?
 $\leadsto \mathbb{C}^x$ -family of 3d TQFT's

\leadsto impossible.

(Gaiotto rigidity: \nexists 1-param. family of 3d TQFTs)

\leadsto since $\$'$ is R-sym, acts on supercharges

- want: discrete flavor sym.,
won't act on superch.

Remark. $\mathbb{Z}^1 \times \mathbb{Z}^3$ worked, and it does
work for $\mathbb{Z}^1 \times \text{Seifert mfd}$

- often, it is $\mathbb{Z}_N \subset \mathbb{Z}^1$

i) $N = 1 \rightarrow$ all previous examples,
all gauge theories

ii) $N > 1 \rightarrow$ non-gauge theories

$\leadsto \mathbb{Z}_N^*$ - family of 3d TQFTs