

# Bestola

## KdV & Malgrange

- solitonless sol'n's

RHP on  $\mathbb{R}$

$$\Gamma_+ = \Gamma_- \begin{bmatrix} 1 - \tau^2 & -\bar{\tau}(z) e^{-2i(-)} \\ \tau(z) e^{2i(4tz^3 + xz)} & 1 \end{bmatrix} =: \Gamma_- J(z)$$

$$\Gamma(z) = \begin{pmatrix} 1 & 1 \\ 2iz & -2iz \end{pmatrix} \left( 1 + \frac{i}{z} a(x,t) \partial_z \tau G(z^{-1}) \right)$$

as  $z \rightarrow \infty$

- also a symmetry  $\Gamma(-z) = z, \Gamma(z) z,$

-  $u(x,t) = -2 \partial_x a$  solves KdV, i.e.

$$u_t - 6u u_x + u_{xxx} = 0$$

- (Dyson) if  $\tau$  is trans'n coeff. of fast decaying pot. then

$$u(x,t) = -2 \partial_x^2 \ln \det [Id - \mathcal{K}]_{(x,\infty)}$$

where  $\mathcal{K}$  has int. kernel

$$K(s,u) = F(s+u), \quad \bar{F}(s) = \frac{1}{2\pi} \int_{\mathbb{R}} \Gamma(z) e^{8itz^3 + izs} dz$$

- recall Malg. form  $\Theta = \int_{\mathbb{R}} \tau \Gamma_-^{-1} \Gamma_-^{-1} \delta J J^{-1} d\hat{z}$

$$\delta \Theta = \frac{1}{2} \int \tau \delta J J^{-1} \wedge \frac{d}{dz} \delta J J^{-1} d\hat{z}$$

$$= \frac{1}{2} \int_{\mathbb{R}} \left[ \delta(\tau e^{2\vartheta}) \wedge \frac{d}{dz} (\bar{\Gamma}(z) e^{-2\vartheta}) - \frac{d}{dz} \delta(\tau e^{2\vartheta}) \wedge \delta(\bar{\Gamma} e^{-2\vartheta}) \right] d\hat{z}$$

$$= - \int \delta \left( \frac{d}{dz} (\tau e^{2\vartheta}) \wedge \delta(\bar{\Gamma} e^{-2\vartheta}) \right) d\hat{z}$$

$$= - \delta \left( \int (\tau e^{2\vartheta})' \wedge \delta(\bar{\Gamma} e^{-2\vartheta}) \right)$$

$$\Rightarrow 2d \ln \tau := \Theta + \int \tau e^{2\vartheta} \cdot \delta(\bar{\Gamma} e^{-2\vartheta})$$

- let's compute  $\partial_x \ln \tau$

$$- \textcircled{2} (\partial_x) = \int \text{tr} \left( \Gamma^{-1} \Gamma' \partial_x \Gamma \Gamma^{-1} \right) = \dots$$

$$= \int \Delta \text{tr} (\Gamma^{-1} \Gamma' \partial_x \Gamma) dz + \partial_x (\text{tr} e^{2\varphi}) \frac{d}{dz} (\text{tr} e^{2\varphi})$$

$\hookrightarrow \text{jump } \Delta f = f_+ - f_-$

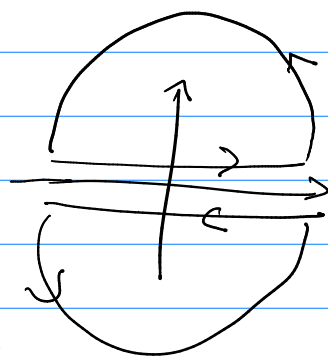
$$- \textcircled{2} (\partial_x) + \int (\dots) = \int \Delta (\dots)$$

$$= \lim_{R \rightarrow \infty} \left[ \int_{C_1} \text{tr} \dots + \int_{C_2} \text{tr} \dots \right]$$

"formal"  
residue

$$= - \text{"res"}_{\infty} \text{tr} (\Gamma^{-1} \Gamma' \partial_x \Gamma) \cdot i z$$

$$= -a(x, t)$$



$\hookrightarrow$  not a real residue

Fuchsian system of ODEs in  $\mathbb{CP}^1$

- ODE for  $n \times n$  matrix  $\psi(z)$ ,

$$\begin{cases} \psi'(z) = A(z) \psi(z), & A(z) = \sum_{j=1}^k \frac{A_j}{z - t_j}, \quad A_j \text{ fixed} \\ \psi(\infty) = \underline{1} \end{cases}$$

$\sum A_j = 0$ , i.e.  $z = \infty$  is a regular pt

- technical requirements:

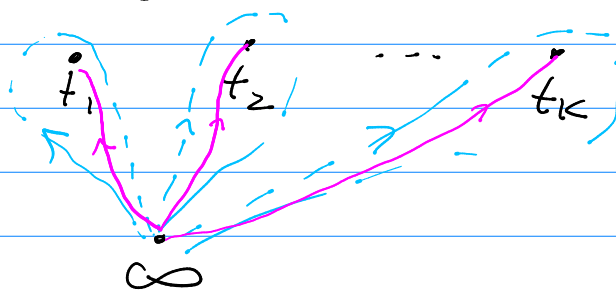
$\forall i$ , evals of  $A_i$  do not differ by nonzero integers  
(nonresonance of the Fuchsian sings.)

$$- (\det \psi)' = \text{tr } A(z) \det \psi \Rightarrow \det \psi = e^{\int \text{tr } A(z)}$$

- if  $[ \gamma ] \in \pi_1(\mathbb{CP}^1 \setminus \{t_1, \dots, t_k\}, \infty)$  and  $\psi(z\gamma)$  denotes analytic cont. of  $\psi(z)$  around  $[ \gamma ]$  then  $\psi(z\gamma) = \tilde{\psi}(z)$  solves same ODE  $\Rightarrow \tilde{\psi}(z) = \psi(z) \cdot M_{[\gamma]}^{-1}$  where  $M: \pi_1 \rightarrow GL_n$  is representation,

$$\psi(z\gamma \cdot \tilde{\gamma}) = \psi(z) M_{\tilde{\gamma}}^{-1} M_{\gamma}^{-1} = \psi(z) M_{\gamma \cdot \tilde{\gamma}}^{-1}$$

- to make it sing-val, introduce branch cuts



- $\psi$  is sing-val on  $\mathbb{CP}^1 \setminus \text{cuts}$
- on the cut  $\ell_j$  oriented towards  $t_j$ ,  
 $\psi_+(z) = \psi_-(z) M_j$
- since  $[\gamma_1 \circ \gamma_2 \circ \dots \circ \gamma_n] = [pt] \Rightarrow M_1 \dots M_n = \mathbb{1}$

Local behaviour

$$- \psi(z) = G_j \underbrace{(1 + O(z-t_j))}_{\text{analytic}} (z-t_j)^{-L_j} C_j$$

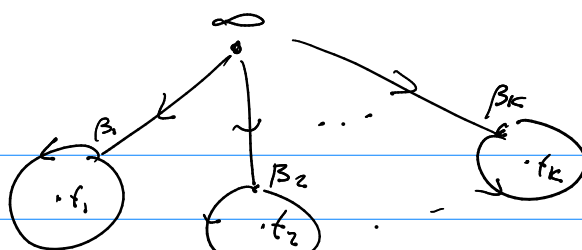
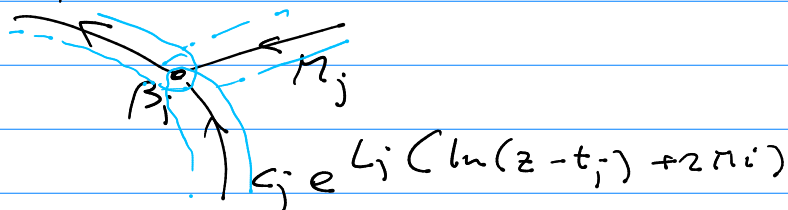
$$\text{where } A_j = G_j L_j G_j^{-1}, M_j = C_j e^{2\pi i L_j} C_j^{-1}$$

- suggests RHP

$$\Gamma(z) = \begin{cases} \varphi(z) & \text{outside } \bigcup_j D_j \\ \varphi(z) c_j (z - t_j)^{-L_j} & , z \in D_j \end{cases}$$

$c_j e^{L_j \ln(z - t_j)}$

- near  $\beta_j$ ,



## Isomonodromy deformations

- find how  $A_j$ 's should depend on  $t_1, \dots, t_k$  so that monodromy is  $\underline{t}$ -indep  
 - [Schlesinger]:

$$\begin{cases} \frac{\partial A_j}{\partial t_k} = \frac{[A_j, A_k]}{t_j - t_k} \\ \frac{\partial A_j}{\partial t_j} = - \sum_{k \neq j} \frac{[A_j, A_k]}{t_j - t_k} \end{cases} \quad \begin{matrix} (\text{nonlin. nonauton.}) \\ \text{PDEs} \end{matrix}$$

- given  $\{A_j(\underline{t}^0) = A_j^0\}_{j=1, \dots, k}$ , sol'n exists

- [80's Jimbo Miwa Ueno] define "tau fn"

$$\begin{aligned} d \ln \tau_{\text{JMU}} &:= \frac{1}{2} \sum_{j=1}^k \left( \text{res}_{z=t_j} \text{tr } A^2(z) \right) dt_j \\ &= \sum_{j=1}^k \left( \sum_{\ell \neq j} \frac{\text{tr } A_j(t) A_\ell(t)}{t_j - t_\ell} \right) dt_j \end{aligned}$$

$\rightarrow d \ln \tau_{\text{nu}}$  is closed if Schlesinger eqns hold for  $\{A_j\}$

Runk consider  $t_j \mapsto t_j' = \frac{at_j + b}{ct_j + d}$ ,  $\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL_2(\mathbb{C})$

$$\Rightarrow \tau(\underline{t}') = \tau(\underline{t}) \cdot \prod_{j=1}^k (ct_j + d)^{2\Delta_j}$$

where the conf. weights are  $\Delta_j = \frac{1}{2} \text{tr } A_j^2$   
(note the Schl. eqn's preserve  $\text{spec } A_j$ )

$\rightarrow$  CFT link:  $\tau(\underline{t}) = \langle G^{\Delta_1}(t_1) \dots G^{\Delta_k}(t_k) \rangle$

- note that RMP depends on

$$\underline{t}, \vec{L}, \vec{C} \text{ subject to } \prod_{j=1}^k C_j e^{2\pi i L_j} C_j^{-1} = \mathbb{1}$$

$\rightarrow$  "deformation space"

$$\mathcal{M} = \left\{ \underline{t}, \underbrace{\vec{L}, \vec{C}}_{\substack{\text{keep fixed} \\ \text{to obtain isomon.}}} \mid \prod C e^{2\pi i L} C^{-1} = \mathbb{1} \right\}$$

Prop  $\ominus \left( \frac{\partial}{\partial t_v} \right) = \frac{\partial}{\partial t_v} \ln \tau_{\text{nu}}$

Pf.  $J_j(z) = C_j (z - t_j)^{-L_j} \Rightarrow \partial_{t_j} J \cdot J^{-1} = C_j \frac{(-L_j)}{z - t_j} C_j^{-1}$   
 $\partial_{t_{k \neq j}} J_j^{-1} J_j = 0$

$$\Rightarrow \textcircled{2}(\partial t_j) = - \oint_{t_j} \text{tr} \underbrace{\psi^{-1} \psi'}_{A(z)} \frac{C_j L_j C_j^{-1}}{z - t_j}$$

$$= - \oint \text{tr} A(z) \cdot \frac{\psi(z) C_j L_j C_j^{-1} \psi'(z)}{z - t_j}$$

$$\frac{G_j (1 + \theta_j) L_j (1 - \theta_j) G_j^{-1}}{z - t_j}$$

$$\frac{A_j}{z - t_j} + G_j [G_j, L_j] G_j^{-1} + G(z, t_j)$$

$$= - \text{res}_{z=t_j} \text{tr} \frac{A(z) A_j}{z - t_j} = - \sum_{l \neq j} \frac{\text{tr} A_j A_l}{t_j - t_l} \stackrel{\text{claim}}{=}$$

- differential:

$$-4\pi i \delta \textcircled{2} = \sum_{l=2}^k \text{tr} (S M_{e_l} M_{e_l}^{-1} \wedge S M_{[e_l, k]} M_{[e_l, k]}^{-1})$$

$$+ \sum_{j=1}^k \text{tr} (e^{2\pi i L_j} S C_j C_j^{-1} \wedge e^{-2\pi i L_j} S C_j C_j^{-1} + 4\pi i S L_j \wedge S C_j C_j^{-1})$$

- remarks: no  $\underline{t}$ -dependence

→ closed 2-form on:

$$\mathcal{M}_0 = \left\{ \vec{C}, \vec{Q} = e^{2\pi i \vec{L}} \mid \prod C_j Q_j C_j^{-1} = \mathbb{1} \right\}$$

- $(-4\pi i) \delta \textcircled{2} =: \omega_M$  is inv under  $C_j \mapsto S C_j$ ,  
i.e.  $M_j \mapsto S M_j S^{-1}$ ,  $S \in GL_n$  arbitrary
- $\omega_M$  descends on  $\mathcal{M}_0 / GL_n$   
 $\xrightarrow{\text{Thm}}$   $\omega_M$  is symplectic on quotient
- to define  $\tau$ , look at  $\textcircled{2} - \int \omega_M \Rightarrow d \ln z$