

# Porta.

- what we did:

\* )  $\langle \rangle_{\beta, L}$  analytic for  $|\lambda| < \lambda_{\beta, L}$

\* ) Improve? using info on  $h$  (like gap)  
→ part. theory. Problem? combinatorics causes  
 $G(n!)$  growth

BBF formula - good if  $g(z, z') = \langle A_z, B_{z'} \rangle$

- not really the case due to disc. in  $T$

\* ) solution? regularized theory  $\lim_{N \rightarrow \infty} g^{(\leq N)} = g$

- however, now  $\|A_z^{(\leq N)}\| \cdot \|B_z^{(\leq N)}\| \leq C 2^{3N}$

- maybe not an optimal bound. try:

Task: multiscale analysis  $g^{(\leq N)} = g^{(\leq 0)} + \sum_{h=0}^N g^{(h)}$   
instead of all scales  $(\leq N)$  at once

- free energy  $f_{\beta, L, N} = -\frac{1}{\beta \Lambda_L} \log \int P_{\leq N}(d\varphi^{(\leq N)}) e^{-V(\varphi^{(\leq N)})}$

- now, by fermion magic,

$$\int P_{\leq N}(d\varphi^{(\leq N)}) e^{-V(\varphi^{(\leq N)})} = \int P_N(d\varphi^N) e^{-V(\varphi^{(N)})} \cdot \int P_{\leq N-1}(d\varphi^{(\leq N-1)}) e^{-V^{(N-1)}(\varphi^{(\leq N-1)})}$$

$$-V^{(N-1)}(\varphi^{(\leq N-1)}) = \sum_{\substack{p \in \mathbb{Z}^2 \\ \text{even}}} \int dz \mathcal{W}_p^{(N-1)}(z) \varphi_z^{(\leq N-1)}(p)$$

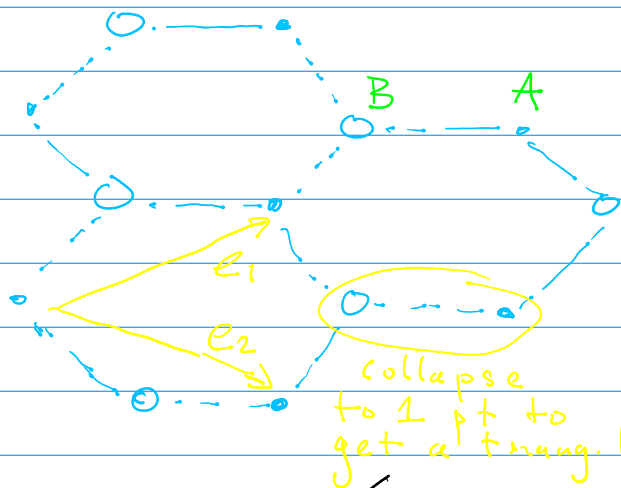
$\uparrow$  "  $\varphi_z^{(\leq N-1)}(p)$   
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- inductively,  $E_{\leq N}(e^{-V(\varphi^{(\leq N)})}) = e^{\sum_{h=1}^N E_h} E_{\leq h}(e^{-V^{(h)}(\varphi^{(\leq h)})})$

$$\text{where } V^{(h)}(\varphi^{(\leq h)}) = \sum_p \int dz \mathcal{W}_p^{(h)}(z) \varphi_z^{(\leq h)}(p)$$

- some estimates were being made, did not write
- in conclusion,  $f_{\beta, L, N} = f_{\beta, L, 0} + \sum_{T=0}^{N-1} \frac{E_T}{\beta(L)}$   
uniform in  $\beta, L, N$ , analytic for  $|\lambda| < \lambda'$ ,  
the limit  $\beta, L, N \rightarrow \infty$  exists
- however, the existence of a gap was crucial  
- serious infrared problems
- let's look at a specific model

## Interacting graphene



$d=2$ ,  $H = \text{Laplacian}$   
on honeycomb lattice

$$\Lambda \cong \mathbb{Z}^2 \times \{A, B\}, \quad \psi(x) = \begin{pmatrix} \psi_A(x) \\ \psi_B(x) \end{pmatrix}$$

$$(H\psi)(x) = \begin{pmatrix} \psi_B(x) + \psi_B(x+e_1) + \psi_B(x+e_2) \\ \psi_A(x) + \psi_A(x-e_1) + \psi_A(x-e_2) \end{pmatrix}$$

A  
B

$$\hat{(H\psi)}(k) = \hat{H}(k) \hat{\psi}(k), \quad \hat{H}(k) = \begin{pmatrix} 0 & -t\Omega(k) \\ -t\overline{\Omega(k)} & 0 \end{pmatrix}$$

$$\Omega(k) = 1 + e^{-ik \cdot e_1} + e^{-ik \cdot e_2}$$

- eigenvalues  $\epsilon_{\pm}(k) = \pm t |\Omega(k)|$

$$\Omega(\underline{k}) = 0 \iff \underline{k} = \underline{k}_F^\omega, \omega = \pm$$

$$\Omega(\underline{k}' + \underline{k}_F^\omega) = \frac{3}{2} (i k'_1 + \omega k'_2) + O(|\underline{k}'|^2)$$

$$\bigcup_{k_2} \mathbb{Z}(\hat{H}(k_1, k_2))$$

