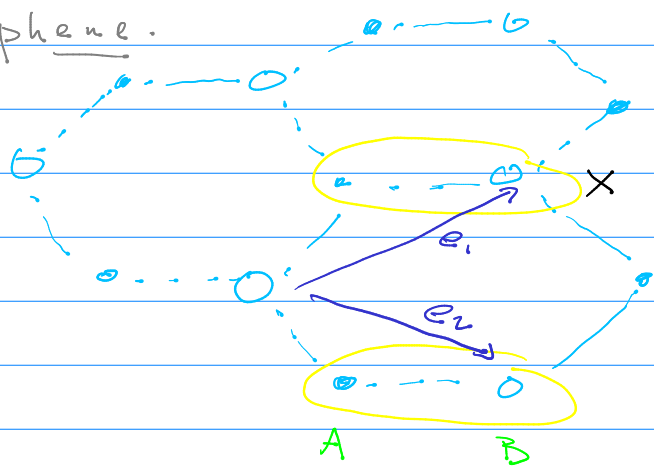


# Portals.

graphene.



$$\Lambda = \Lambda_A + \Lambda_B$$

$$\psi(x) = \begin{pmatrix} \psi_A(x) \\ \psi_B(x) \end{pmatrix}$$

$$(H\psi)(x) = -t \begin{pmatrix} \psi_B(x) + \psi_B(x - e_1) + \psi_B(x - e_2) \\ \psi_A(x) + \psi_A(x + e_1) + \psi_A(x + e_2) \end{pmatrix}$$

$\hat{H} = \bigoplus_{\underline{k} \in \pi^2} \hat{H}(\underline{k})$ ,  $\hat{H}(\underline{k})$  called Bloch Hamiltonian

$$\hat{H}(\underline{k}) = \begin{pmatrix} 0 & -t \Sigma(\underline{k}) \\ -t \overline{\Sigma(\underline{k})} & 0 \end{pmatrix}, \quad \Sigma(\underline{k}) := 1 + e^{-i\underline{k} \cdot \underline{e}_1} + e^{-i\underline{k} \cdot \underline{e}_2}$$

- chemical pot.  $\mu = 0$

- Fermi s.f.c.  $F_\mu := \{ \underline{k} \in \pi^2 \mid \hat{H}(\underline{k}) - \mu \text{ has } 0 \text{ eigenvalue} \}$   
 $= \{ k_F^+, k_F^- \}$

Transport:  $\mathcal{Z}_{12} = \mathcal{Z}_{21} = 0$ ,  $\mathcal{Z}_{11} = \mathcal{Z}_{22} = 1/4$  (universal, i.e.  $t$ -indep.)

- was found for nonint. systems

- interacting ones?

Goal, construct  $\langle \cdot \rangle_{\beta, L}$  with  $H = H_0 + \lambda V$ , analyticity?

$$g(\underline{z}, \underline{z}') = \frac{1}{\beta} \sum_{\underline{k} \in \pi^2} e^{-i\underline{k} \cdot (\underline{z} - \underline{z}')} \int \frac{d^2 \underline{k}}{(2\pi)^2} e^{-i\underline{k} \cdot (\underline{z} - \underline{z}')} \hat{g}(\underline{k})$$

where  $\hat{g}(\underline{k}) = \frac{1}{-i k_0 + \hat{H}(\underline{k})}$  (writing  $\underline{k} := (k_0, \underline{k})$ ,  $k_F^\omega := (0, \underline{k}_F^\omega)$ )

$\underline{k} = \underline{k}' + \underline{k}_F^\omega$ ,  $\underline{k}'$  small  $\Rightarrow \|\hat{g}(\underline{k})\| \simeq \frac{1}{\|\underline{k}'\|}$  for  $\|\underline{k}'\|$  small  
 $\Rightarrow g \notin L^1$

# IR multiscale analysis.

-  $\hat{g}(k) = \hat{g}_{uv}(k) + \hat{g}_{ir}(k)$

where  $\hat{g}_{uv}(k) = \hat{g}(k) \chi(|k - k_F^\omega| > \delta)$ ,

$\hat{g}_{ir}(k) = \hat{g}(k) \sum_{\omega \in \mathbb{Z}} \chi(|k - k_F^\omega| < \delta)$

- decompose  $\varphi = \varphi^{(uv)} + \varphi^{(ir)}$

where  $\varphi^{(ir)\pm}(x) = \sum_{\omega} e^{\pm i k_F^\omega x} \varphi_{\omega, x}^{(ir)\pm}$

$\|\hat{g}_\omega(k')\| \approx \frac{1}{\|k'\|}$  means  $\|g_\omega(z, z')\| \leq \frac{1}{\|z - z'\|^2}$

- pick  $\gamma > 1, h \in \mathbb{Z}_+$  and set cutoff

$f_h(k') = \chi(\gamma^{-h} \|k'\|) - \chi(\gamma^{-(h+1)} \|k'\|)$

so

$\hat{g}_\omega(k') = \sum_{h=h_\beta}^0 \underbrace{\hat{g}_\omega^{(h)}(k')}_{:= f_h(k') \hat{g}_\omega(k')}$

-  $k_0 = \frac{2\pi}{\beta}(n+1), k' = (k_0, \underline{k}'), \|k'\| \geq \frac{\pi}{\beta} =: \gamma^{h_\beta}$

$|g_\omega^{(h)}(z, z')| \leq \int d^3k |\hat{g}_\omega^{(h)}(k')| \leq C \gamma^{2h}$

$(1 + \|z - z'\|^M \gamma^{hM}) |g_\omega^{(h)}(z, z')| \leq C_M \gamma^{2h}$

$\Rightarrow \|g_\omega^{(h)}(z, z')\| \leq \frac{C_M \gamma^{2h}}{(1 + \gamma^h \|z - z'\|)^M}$

- write  $\varphi_\omega^{(\leq 0)} = \sum_{h=h_\beta}^0 \varphi_\omega^{(h)}$

$\rightsquigarrow V^{(h)}(\varphi^{(\leq h)}) = \sum_P \int dz W_\dagger^{(h)}(z) \varphi_z^{(\leq h)}(P)$

- power-counting.  $\sum_{s, s'} \int dx \psi_{x, s}^{(\leq 0)+} \varphi_{x, s}^{(\leq 0)-} \varphi_{x, s'}^{(\leq 0)+} \varphi_{x, s'}^{(\leq 0)-}$

becomes, by writing  $\varphi^{(\leq 0)} = \varphi^{(\leq -1)} + \varphi^{(0)}$ :

$\rightarrow \sum_{s, s'} \int dx \psi_{x, s}^{(\leq -1)+} \varphi_{x, s}^{(\leq -1)-} \varphi_{x, s'}^{(\leq -1)+} \varphi_{x, s'}^{(\leq -1)-} + \text{h.a.t.}$

$$-g^{(\leq -1)} = \gamma^{-2} \tilde{g}^{(\leq 0)}(\gamma^{-1} \cdot) \quad (\gamma > 1)$$

and  $\varphi^{(\leq -1)} = \gamma^{-1} \tilde{\varphi}^{(\leq 0)}_{\gamma^{-1} \cdot}$

- so our sum becomes

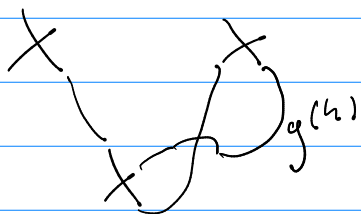
$$\sum_{s,s'} \underbrace{\gamma^{-4} \gamma^3}_{\gamma \gamma^{-1} \rightarrow \text{RG - irrelevant}} \int dx (\tilde{\varphi}^{(\leq 0)})^4$$

$$\Rightarrow \|g^{(h)}_{\omega}(z, z')\| \leq \frac{C_n \gamma^{2h}}{1 + (C_n \gamma^{2h} |z - z'|)^n}$$

-  $\varphi^2$ : RG - relevant

$$\mathbb{E}_{\leq 0}(e^{-V^{(0)}(\varphi^{(\leq 0)})}) = e^{\sum_{T \leq h}^{-1} E_T} \mathbb{E}_{\leq h}(e^{-V^{(h)}(\varphi^{(\leq h)})})$$

$$\text{where } V^{(h)}(\varphi^{(\leq h)}) = \sum_p \int dz \mathcal{W}_p^{(h)}(z) \varphi_z^{(\leq h)}(p)$$



$$\|g^{(h)}\|_r^{s-1} \|g^{(h)}\|_{\infty}^{s+1 - \frac{|P|}{2}} \leq \gamma^{-h(s-1)} \gamma^{2h(s+1 - \frac{|P|}{2})}$$

$$\gamma^{h(-|P|+3+s)} \leq \gamma^{h(-|P|+4)} \quad (*)$$

- however  $\rightarrow \dots \rightarrow \dots \rightarrow$  bud  $\delta$  isolate  $|P|=2$

$$\text{Idea: } V^{(h)}(\varphi^{(\leq h)}) = \underbrace{\sum_{|P|=2} V^{(h)}(\varphi^{(\leq h)})}_{\text{bud } \delta} + \underbrace{\sum_{|P| \geq 4} V^{(h)}(\varphi^{(\leq h)})}_{\text{isolated}}$$

$$\sum V^{(h)}(\varphi^{(\leq h)}) = (\varphi^{(\leq h)+}, \mathcal{W}_2^{(h)} \varphi^{(\leq h)-})$$

$$- \sum V^{(h)} \text{ in } P_{\leq h}(d\varphi^{(\leq h)})$$

$$\int P_{\leq h}(d\varphi^{(\leq h)}) e^{-(\varphi^{(\leq h)}, \mathcal{W}_2^{(h)} \varphi^{(\leq h)})} \quad (1)$$

$$= e^{t_h} \int \tilde{P}_{\leq h}(d\varphi^{(\leq h)}) \quad (\dots)$$

where  $\tilde{P}_{\leq h}(d\varphi^{(\leq h)}) = D\varphi^{(\leq h)} \exp\left\{-(\varphi^{(\leq h)}, (g^{(\leq h)} + \mathcal{W}_2^{(h)})\varphi^{(\leq h)})\right\}$

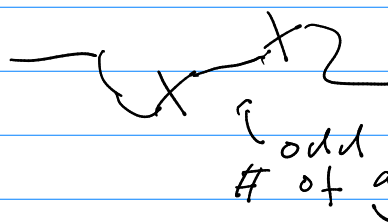
$$\frac{1}{1+V-z} = \frac{1}{1-z} - \frac{1}{1-z} \vee \frac{1}{1-z} + \dots$$

$$-\hat{g}^{(\leq h)}_{\omega}(k) = \chi_{\leq h} \begin{pmatrix} i k_0 & v_n(i k_1 + \omega k_2') \\ v_n(-i k_1 + \omega k_2') & i k_0 \end{pmatrix}^{-1} \frac{1}{z_n}$$

where  $v_n, z_n$  given iteratively (from interaction)

and  $z_0 = 1, v_0 = \frac{3}{2}t$

$$-\hat{\mathcal{W}}_2^{(h)}(k') \Big|_{k'=0} = 0, \quad k' \cdot \nabla_{k'} \hat{\mathcal{W}}_2^{(h)}(k') \Big|_{k'=0} \\ \begin{pmatrix} z_{n,0} i k_0 & z_{n,1} (i k_1 + \omega k_2') \\ z_{n,1} (-i k_1 + \omega k_2') & z_{n,0} i k_0 \end{pmatrix}$$

  $g = \frac{1}{k}$

$$z_{n+1} = z_n + z_{0,k}$$

$$z_{n+1} v_{n+1} = z_n v_n + z_{1,h}$$

$$\begin{pmatrix} z_{n+1} \\ v_{n+1} \end{pmatrix} = \begin{pmatrix} z_n \\ v_n \end{pmatrix} + \hat{\beta}_n \leftarrow \text{beta function from RG-flow}$$

$$|\hat{\beta}_h| \leq C \|1\| \gamma^h \Rightarrow \|z_{n+1}\| \leq \sum_{i=0}^n \|1\| \gamma^i = O(1) \\ \Rightarrow \|\mathcal{W}_p^{(h)}\|_1 \leq \|1\|^{(1-\delta)(\frac{1}{2}-1)} \gamma^{h(-1+\delta)} \\ \|\mathcal{W}_p^{(h)} - \mathcal{X}\| \leq \|1\|^{2(1-\delta)} \gamma^{h(-1+\delta)}$$

-  $|P| \geq 6$ ,  $\langle \cdot \rangle_{\beta, L}$  is analytic in  $\beta, L$  (result from 2010.)

Applications to transport.

$$\mathcal{Z}_{ij} = \lim_{\gamma \rightarrow 0^+} \lim_{\beta, L \rightarrow \infty} \frac{i}{L^2} \int_{-\infty}^0 dt e^{\gamma t} \langle [J_i(t), X_j] \rangle_{\beta, L}$$

in 2<sup>nd</sup> quantization, i.e.  $X = \sum_x x a_x^\dagger a_x$ ,  $J = j[H, x]$

- one can prove universality:

$\mathcal{Z}_{ij} = \mathcal{Z}_{ij}|_{\lambda=0}$   $\rightarrow$  for graphene-like models (w/ Dirac cones)  
 $\rightarrow$  for gapped models

$\rightarrow$  graphene,  $\mathcal{Z}_{ii} = 1/4$

$\rightarrow$  gapped  $H$ ,  $\mathcal{Z}_{ii} \in \mathbb{Z}/2\pi$ ,  $\mathcal{Z}_{ii} = 0$

Wick rotation

$$\int_{-\infty}^0 dt e^{\gamma t} \langle [J_i(t), X_j] \rangle \quad J_i(t) = e^{iHt} J_i e^{-iHt}$$

$$= \frac{1}{\gamma} \left( \int_{-\infty}^0 dt e^{\gamma t} \langle [J_i(t), J_j] \rangle_{\beta, L} - \langle [J_i, X_j] \rangle_{\beta, L} \right)$$

Claim.  $\lim_{\beta, L \rightarrow \infty} \frac{1}{L^2} \int_{-\infty}^0 dt e^{\gamma t} \langle [J_i(t), J_j] \rangle_{\beta, L}$

$$= \lim_{\beta, L \rightarrow \infty} \frac{-i}{L^2} \int_{-\beta/2}^{+\beta/2} dt e^{-i\gamma t} \langle T J_i(-it) J_j \rangle_{\beta, L}$$

- idea: suppose  $\gamma = \frac{2\pi}{\beta} \cdot u$ ,  $u \in \mathbb{Z}$  ("bosonic Matsubara")

- write  $\int_{-T}^0 dt e^{\eta t} \langle [J_i(t), J_j] \rangle_{\beta L}$

$$= \int_{-T}^0 dt e^{\eta t} \langle J_i(t) J_j \rangle_{\beta L} - \int_{-T}^0 dt e^{\eta t} \langle J_j J_i(t) \rangle$$

$$- \text{use } e^{-\beta H} J_j J_i(t) = T \text{r } e^{-\beta H} \left( \underbrace{e^{\beta H} J_i(t) e^{-\beta H}}_{J_i(t - i\beta)} \right) J_j$$

- so, we get

$$\int_{-T}^0 dt e^{\eta t} \langle J_i(t) J_j \rangle_{\beta L} - \int_{-T}^0 dt \underbrace{e^{\eta t}}_{= e^{\eta(t - i\beta)}} \langle J_i(t - i\beta) J_j \rangle_{\beta L}$$

$$= - \int_0^{-i\beta} dt e^{\eta t} \langle J_i(t), J_j \rangle_{\beta L} + \text{error}_T$$

$$= i \int_0^\beta dt e^{-it\eta} \langle J_i(-it), J_j \rangle_{\beta L}$$

$$= i \int_{-\beta/2}^{\beta/2} dt e^{-it\eta} \langle T J_i(-it), J_j \rangle_{\beta L}$$

$$\text{so } \mathcal{G}_{ij} = \lim_{\eta \rightarrow 0} \lim_{\beta L \rightarrow \infty} \frac{1}{\eta} \frac{i}{L^2} \int dt (e^{-it\eta} - 1) \langle T J_i(-it) J_j \rangle_{\beta L}$$

$$= \partial_\eta K_{ij}(\eta) \big|_{\eta=0}$$

where  $K_{ii}(-\eta) = K_{ii}(\eta)$  if  $K \in C^1$   $\mathcal{G}_{ii} = 0$

$$K_{ii}(\eta) = K_{ii}^{\text{rel.}}(\eta) + K_{ii}^{\text{even}}(\eta) \quad \rightarrow \partial_\eta (-) \big|_{\eta=0} = 0$$

$\hookrightarrow \eta \rightarrow 1/\eta$