

Cecotti

Review of Swampland

1) finiteness conj.

- there is a uniform bound (in theory space)
 $\# \text{ massless spin } s \text{ fields} \leq N_d(s)$

spacetime dim.

- clearly $N_d(\geq 2) = 0$, $N_d(2) = 1$, $N_d(\frac{3}{2}) = \frac{32}{2^{1+\frac{1}{2}}}$
but also $N_d(0, \frac{1}{2}, 1)$?

II) no parameter

strong: couplings are vevs $\lambda = \langle \varphi \rangle$

weak: \mathcal{L}_{eff} is rigid, i.e. \nexists continuous parameters

- of course, if $\lambda = \langle \varphi_{\text{heavy}} \rangle$, where φ_{heavy} is in some UV-completion, then $\lambda \in \{\text{critical pts of } V(\varphi_{\text{heavy}})\}$ which is necessarily a set of isolated pts \rightarrow if not, we would have light particles along valleys

- e.g. type IIB

10d

$M_2 \times \mathbb{C}P^1$
rigid, $h^{1,1} = 0$

1 vector

graviphoton $N=2$

$$\tau = \frac{\vartheta}{2\pi} + \frac{4\pi i}{g^2}$$

$\rightarrow \tau$ is not a generic thing, but determined as $\vartheta = \int_B \Omega / \int_A \Omega$, $g = \text{something not understood, integrals of mod. forms}$

- of course, this comes from a heavy field of mass $>$ cutoff

Rmk "Opposite" of naturalness. We claim thing is not consistent, unless very specially fine-tuned

III) no global symmetry

strong; consistent th. has no global symmetries (cont., disc., ...)

weak: $|G_{gl}| < k d$

- of course \mathcal{L}_{eff} may have emergent symmetries, exact up to cutoffs

- let G gauge gp, $H = G/G^0$ disc. gauge gp

IV) G_0 is cpt. $\xRightarrow{\text{Dirac}} (e, m) \in \mathbb{Z}$

- essentially, we want integral fluxes

$$\Lambda = \bigoplus_{d=1}^d \Lambda_d, d \neq 4, 5+2$$

$$\Lambda = \Lambda_+ \oplus \Lambda_- \oplus \left(\bigoplus_{n \neq d/2} \Lambda_n \right) d = 4, 5+2$$

$$\begin{array}{ccc} & H & \\ \nearrow & f & \searrow \\ G & \longrightarrow & \Lambda_+ \oplus \Lambda \\ & f(H) & \longrightarrow \text{U-duality gp} \end{array}$$

V) Completeness

Strong: every charge allowed by Dirac q.
is realised

- black hole argument \rightarrow if charges fall
into BH, when it evaporates we get som

- now we come to the geometric conjectures
- $\mathcal{L}_{\text{eff}} = \dots - \frac{1}{2} \sqrt{-g} G_{ij}(\varphi) \partial^\mu \varphi^i \partial_\mu \varphi^j + \dots$
 $\Rightarrow \{ \varphi^i \}$ chart on \mathcal{M}

Original conj:

$\dim \mathcal{M} \geq 1$. Then \mathcal{M} complete, noncpt.,
finite-volume

- better: every subharmonic, bounded
function on \mathcal{M} is constant

Original: $\pi_1(\mathcal{M}) = 0$

- better: Let $\tilde{\mathcal{M}}$ fin. cover of \mathcal{M} .
 $\forall \varepsilon > 0 \quad \forall [\gamma] \in H_1(\tilde{\mathcal{M}}, \mathbb{Z})$, we can find
loop $\tilde{\gamma}$ s.t. $[\tilde{\gamma}] = [\gamma]$ and $\text{len } \tilde{\gamma} < \varepsilon$.