

Rotation About a Vector

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[Summary]

The rotation \mathbf{R} of the vector \mathbf{r} through an angle ψ about some unit vector $\hat{\omega}$ can be computed with

$$\mathbf{R}(\mathbf{r}, \hat{\omega}, \psi) = \sin \psi \cdot [\hat{\omega} \times \mathbf{r}] + \cos \psi \cdot \mathbf{r} + \hat{\omega} (1 - \cos \psi) (\hat{\omega} \cdot \mathbf{r})$$

The matrix form of this result is

$$\mathbf{R}(\mathbf{r}, \hat{\omega}, \psi) = \Psi \mathbf{r}$$

where Ψ is the rotation matrix

$$\Psi = \begin{bmatrix} c + \omega_x \omega_x (1 - c) & -\omega_z s + \omega_x \omega_y (1 - c) & +\omega_y s + \omega_x \omega_z (1 - c) \\ +\omega_z s + \omega_y \omega_x (1 - c) & c + \omega_y \omega_y (1 - c) & -\omega_x s + \omega_y \omega_z (1 - c) \\ -\omega_y s + \omega_z \omega_x (1 - c) & +\omega_x s + \omega_z \omega_y (1 - c) & c + \omega_z \omega_z (1 - c) \end{bmatrix}$$

Note that we have used the shorthand $c = \cos \psi$ and $s = \sin \psi$.

[Discovery]

We will consider the problem of taking any point in space \mathbf{r} and rotating it through some angle ψ about the unit vector $\hat{\omega}$. To begin, let us use label θ to denote the angle between \mathbf{r} and $\hat{\omega}$. We will split \mathbf{r} into \mathbf{r}_{\parallel} and \mathbf{r}_{\perp} . We will form \mathbf{r}_{\parallel} by projecting \mathbf{r} onto $\hat{\omega}$ and we will form \mathbf{r}_{\perp} by taking whatever is left over after the projection

$$\begin{aligned} \mathbf{r}_{\parallel} &= r \cos \theta \cdot \hat{\omega} \\ &= \hat{\omega} (\hat{\omega} \cdot \mathbf{r}) \end{aligned}$$

$$\begin{aligned} \mathbf{r}_{\perp} &= \mathbf{r} - \mathbf{r}_{\parallel} \\ &= \mathbf{r} - \hat{\omega} (\hat{\omega} \cdot \mathbf{r}) \end{aligned}$$

We now consider the plane which is perpendicular \mathbf{r}_{\perp} and drawn at the tip of \mathbf{r}_{\parallel} . We wish to find basis vectors in that plane. The vector \mathbf{r}_{\perp} is already one such basis vector. We know that our second basis vector $\boldsymbol{\rho}$ will be perpendicular to both \mathbf{r}_{\perp} and $\hat{\omega}$. We can therefore find $\boldsymbol{\rho}$ as follows

$$\boldsymbol{\rho} = \hat{\omega} \times \mathbf{r}_{\perp} = \hat{\omega} \times (\mathbf{r} - \mathbf{r}_{\parallel}) = \hat{\omega} \times \mathbf{r}$$

Notice that $\boldsymbol{\rho}$ is conveniently of the same magnitude as \mathbf{r}_{\perp} . To find where \mathbf{r} goes when rotated about $\hat{\omega}$ we make a circular combination of \mathbf{r}_{\perp} and $\boldsymbol{\rho}$, which we then add back to \mathbf{r}_{\parallel} , as follows:

$$\begin{aligned} \mathbf{R}(\mathbf{r}, \hat{\omega}, \psi) &= \mathbf{r}_{\parallel} + \mathbf{r}_{\perp} \cos \psi + \boldsymbol{\rho} \sin \psi \\ &= \hat{\omega} (\hat{\omega} \cdot \mathbf{r}) + \cos \psi \cdot [\mathbf{r} - \hat{\omega} (\hat{\omega} \cdot \mathbf{r})] + \sin \psi \cdot [\hat{\omega} \times \mathbf{r}] \\ &= \sin \psi \cdot [\hat{\omega} \times \mathbf{r}] + \cos \psi \cdot \mathbf{r} + \hat{\omega} (1 - \cos \psi) (\hat{\omega} \cdot \mathbf{r}) \end{aligned}$$

We give the component expansion of this result below. For the sake of readability we will use c and s for $\cos \psi$ and $\sin \psi$.

$$\begin{aligned} R_x &= s(\omega_y r_z - \omega_z r_y) + cr_x + \omega_x(1-c)(\omega_x r_x + \omega_y r_y + \omega_z r_z) \\ R_y &= s(\omega_z r_x - \omega_x r_z) + cr_y + \omega_y(1-c)(\omega_x r_x + \omega_y r_y + \omega_z r_z) \\ R_z &= s(\omega_x r_y - \omega_y r_x) + cr_z + \omega_z(1-c)(\omega_x r_x + \omega_y r_y + \omega_z r_z) \end{aligned}$$

It is clear that each of these components can be written as a linear combination of r_x , r_y , and r_z . In order to see this explicitly we first expand out all multiplications

$$\begin{aligned} R_x &= s\omega_y r_z - s\omega_z r_y + cr_x + \omega_x\omega_x r_x + \omega_x\omega_y r_y + \omega_x\omega_z r_z - c\omega_x\omega_x r_x - c\omega_x\omega_y r_y - c\omega_x\omega_z r_z \\ R_y &= s\omega_z r_x - s\omega_x r_z + cr_y + \omega_y\omega_x r_x + \omega_y\omega_y r_y + \omega_y\omega_z r_z - c\omega_y\omega_x r_x - c\omega_y\omega_y r_y - c\omega_y\omega_z r_z \\ R_z &= s\omega_x r_y - s\omega_y r_x + cr_z + \omega_z\omega_x r_x + \omega_z\omega_y r_y + \omega_z\omega_z r_z - c\omega_z\omega_x r_x - c\omega_z\omega_y r_y - c\omega_z\omega_z r_z \end{aligned}$$

Then we collect the terms in order to make the linear combinations stand out

$$\begin{aligned} R_x &= r_x(\omega_x\omega_x + c - c\omega_x\omega_x) + r_y(\omega_x\omega_y - c\omega_x\omega_y - s\omega_z) + r_z(\omega_x\omega_z - c\omega_x\omega_z + s\omega_y) \\ R_y &= r_x(\omega_x\omega_y - c\omega_x\omega_y + s\omega_z) + r_y(\omega_y\omega_y + c - c\omega_y\omega_y) + r_z(\omega_y\omega_z - c\omega_y\omega_z - s\omega_x) \\ R_z &= r_x(\omega_x\omega_z - c\omega_x\omega_z - s\omega_y) + r_y(\omega_y\omega_z - c\omega_y\omega_z + s\omega_x) + r_z(\omega_z\omega_z + c - c\omega_z\omega_z) \end{aligned}$$

This result is best written in matrix form as

$$\mathbf{R}(\mathbf{r}, \hat{\boldsymbol{\omega}}, \psi) = \boldsymbol{\Psi} \mathbf{r}$$

where $\boldsymbol{\Psi}$ is the rotation matrix

$$\boldsymbol{\Psi} = \begin{bmatrix} c + \omega_x\omega_x(1-c) & -\omega_zs + \omega_x\omega_y(1-c) & +\omega_y s + \omega_x\omega_z(1-c) \\ +\omega_zs + \omega_y\omega_x(1-c) & c + \omega_y\omega_y(1-c) & -\omega_xs + \omega_y\omega_z(1-c) \\ -\omega_y s + \omega_z\omega_x(1-c) & +\omega_xs + \omega_z\omega_y(1-c) & c + \omega_z\omega_z(1-c) \end{bmatrix}$$