[Summary]

Consider a double pendulum composed of masses m_1 and m_2 . Let us say that mass m_1 has polar coordinates (α, θ) with respect to the origin, and that mass m_2 has polar coordinates (β, ϕ) with respect to mass m_1 .

The system of differential equations governing the motion of double pendulum is as follows (note that we use the shorthand $M = m_1 + m_2$).

$$[M\alpha^2]\ddot{\theta} + [\alpha\beta\cos(\theta - \phi)]\ddot{\phi} = -m_2\alpha\beta\dot{\phi}^2\sin(\theta - \phi) - Mg\alpha\cos\theta$$
$$[m_2\alpha\beta\cos(\theta - \phi)]\ddot{\theta} + [m_2\beta^2]\ddot{\phi} = m_2\alpha\beta\dot{\theta}^2\sin(\theta - \phi) - m_2g\beta\cos\phi$$

From the above equations can derive the separate generalized forces acting on each mass. These forces are complicated to such a degree that they are beyond intuitive human understanding. Nonetheless, we give them here:

$$\ddot{\theta} = \left[\left(-m_2 \alpha \beta \dot{\phi}^2 \sin(\theta - \phi) - Mg\alpha \cos \theta \right) \left(m_2 \beta^2 \right) \right.$$

$$\left. - \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \left(m_2 \alpha \beta \dot{\theta}^2 \sin(\theta - \phi) - m_2 g \beta \cos \phi \right) \right]$$

$$\cdot 1 / \left[\left(M\alpha^2 \right) \left(m_2 \beta^2 \right) - \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \right]$$

$$\ddot{\phi} = \left[\left(M\alpha^2 \right) \left(m_2 \alpha \beta \dot{\theta}^2 \sin(\theta - \phi) - m_2 g \beta \cos \phi \right) \right.$$

$$\left. - \left(- m_2 \alpha \beta \dot{\phi}^2 \sin(\theta - \phi) - Mg\alpha \cos \theta \right) \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \right]$$

$$\cdot 1 / \left[\left(M\alpha^2 \right) \left(m_2 \beta^2 \right) - \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \right]$$

[Discovery]

Consider a double pendulum composed of masses m_1 and m_2 . Let us say that mass m_1 has polar coordinates (α, θ) with respect to the origin, and that mass m_2 has polar coordinates (β, ϕ) with respect to mass m_1 . Our goal is to find a formula for the forces on each body, such that we can numerically integrate the forces and produce a computer simulation of a swinging double-pendulum.

To begin, we say that the location of the first body is

$$x_1 = \alpha \cos(\theta)$$
$$y_1 = \alpha \sin(\theta)$$

Next we say that the location of the second body is

$$x_2 = x_1 + \beta \cos \phi$$
$$= \alpha \cos \theta + \beta \cos \phi$$
$$y_2 = y_1 + \beta \sin \phi$$

 $= \alpha \sin \theta + \beta \sin \phi$

The squared velocity of the first body is then

$$v_1^2 = \alpha^2 \dot{\theta}^2$$

and the squared velocity of the second body is

$$\dot{x}_1^2 = \left(-\alpha\dot{\theta}\sin\theta - \beta\dot{\phi}\sin\phi\right)^2$$

$$= \alpha^2\dot{\theta}^2\sin^2\theta + 2\alpha\beta\dot{\theta}\dot{\phi}\sin\theta\sin\phi + \beta^2\dot{\phi}^2\sin^2\phi$$

$$\dot{y}_1^2 = \left(+\alpha\dot{\theta}\cos\theta + \beta\dot{\phi}\cos\phi\right)^2$$

$$= \alpha^2\dot{\theta}^2\cos^2\theta + 2\alpha\beta\dot{\theta}\dot{\phi}\cos\theta\cos\phi + \beta^2\dot{\phi}^2\cos^2\phi$$

$$v_2^2 = \dot{x}_1^2 + \dot{y}_1^2$$

$$= \alpha^2\dot{\theta}^2 + \beta^2\dot{\phi}^2 + 2\alpha\beta\dot{\theta}\dot{\phi}\left[\sin\theta\sin\phi + \cos\theta\cos\phi\right]$$

$$= \alpha^2\dot{\theta}^2 + \beta^2\dot{\phi}^2 + 2\alpha\beta\dot{\theta}\dot{\phi}\cos(\theta - \phi)$$

The kinetic energy of the system can then be written as follows. Note that we use the shorthand $M = m_1 + m_2$

$$T = \frac{1}{2}m_1\alpha^2\dot{\theta}^2 + \frac{1}{2}m_2\alpha^2\dot{\theta}^2 + \frac{1}{2}m_2\beta^2\dot{\phi}^2 + m_2\alpha\beta\dot{\theta}\dot{\phi}\cos(\theta - \phi)$$
$$= \frac{1}{2}M\alpha^2\dot{\theta}^2 + \frac{1}{2}m_2\beta^2\dot{\phi}^2 + m_2\alpha\beta\dot{\theta}\dot{\phi}\cos(\theta - \phi)$$

The total potential energy of the system is just the usual gravitational potential

$$U = m_1 g \alpha \sin \theta + m_2 g \left[\alpha \sin \theta + \beta \sin \phi \right]$$
$$= M g \alpha \sin \theta + m_2 g \beta \sin \phi$$

The Lagrangian of the system is then simply

$$\mathcal{L} = T - U$$

The system's generalized momenta are then

$$p_{\theta} = \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = M\alpha^{2}\dot{\theta} + m_{2}\alpha\beta\dot{\phi}\cos(\theta - \phi)$$
$$p_{\phi} = \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = m_{2}\beta^{2}\dot{\phi} + m_{2}\alpha\beta\dot{\theta}\cos(\theta - \phi)$$

and the system's generalized forces are

$$\frac{\partial \mathcal{L}}{\partial \theta} = -m_2 \alpha \beta \dot{\theta} \dot{\phi} \sin(\theta - \phi) - Mg\alpha \cos \theta$$
$$\frac{\partial \mathcal{L}}{\partial \phi} = +m_2 \alpha \beta \dot{\theta} \dot{\phi} \sin(\theta - \phi) - m_2 g\beta \cos \phi$$

In order to apply the Euler-Lagrange equations we must differentiate the momenta with respect to time

$$\dot{p}_{\theta} = M\alpha^{2}\ddot{\theta} + m_{2}\alpha\beta\ddot{\phi}\cos(\theta - \phi) - m_{2}\alpha\beta\dot{\phi}\sin(\theta - \phi)\left[\dot{\theta} - \dot{\phi}\right]$$
$$\dot{p}_{\phi} = m_{2}\beta^{2}\ddot{\phi} + m_{2}\alpha\beta\ddot{\theta}\cos(\theta - \phi) - m_{2}\alpha\beta\dot{\theta}\sin(\theta - \phi)\left[\dot{\theta} - \dot{\phi}\right]$$

Recall that the Euler-Lagrange equations for this system are

$$\frac{\partial \mathcal{L}}{\partial \theta} = \frac{d}{dt} \left[\partial \mathcal{L} / \partial \dot{\theta} \right] \qquad \frac{\partial \mathcal{L}}{\partial \phi} = \frac{d}{dt} \left[\partial \mathcal{L} / \partial \dot{\phi} \right]$$

We have already calculated all of these terms, and so we can immediately proceed

$$-m_{2}\alpha\beta\dot{\theta}\dot{\phi}\sin(\theta-\phi) - Mg\alpha\cos\theta = M\alpha\ddot{\theta} + m_{2}\alpha\beta\ddot{\phi}\cos(\theta-\phi) - m_{2}\alpha\beta\dot{\phi}\sin(\theta-\phi) \left[\dot{\theta}-\dot{\phi}\right] m_{2}\alpha\beta\dot{\theta}\dot{\phi}\sin(\theta-\phi) - m_{2}g\beta\cos\phi = m_{2}\beta\ddot{\phi} + m_{2}\alpha\beta\ddot{\theta}\cos(\theta-\phi) - m_{2}\alpha\beta\dot{\theta}\sin(\theta-\phi) \left[\dot{\theta}-\dot{\phi}\right]$$

Thankfully, some simplification here is possible

$$-Mg\alpha\cos\theta = M\alpha^{2}\ddot{\theta} + m_{2}\alpha\beta\ddot{\phi}\cos(\theta - \phi) + m_{2}\alpha\beta\dot{\phi}^{2}\sin(\theta - \phi)$$
$$-m_{2}g\beta\cos\phi = m_{2}\beta^{2}\ddot{\phi} + m_{2}\alpha\beta\ddot{\theta}\cos(\theta - \phi) - m_{2}\alpha\beta\dot{\theta}^{2}\sin(\theta - \phi)$$

We can view the above result as a pair of linear equations in $\ddot{\theta}$ and $\ddot{\phi}$.

$$[M\alpha^2] \ddot{\theta} + [m_2\alpha\beta\cos(\theta - \phi)] \ddot{\phi} = -m_2\alpha\beta\dot{\phi}^2\sin(\theta - \phi) - Mg\alpha\cos\theta$$
$$[m_2\alpha\beta\cos(\theta - \phi)] \ddot{\theta} + [m_2\beta^2] \ddot{\phi} = m_2\alpha\beta\dot{\theta}^2\sin(\theta - \phi) - m_2g\beta\cos\phi$$

These equations can be solved for $\ddot{\theta}$ and $\ddot{\phi}$ by using Cramer's rule.

$$\ddot{\theta} = \left[\left(-m_2 \alpha \beta \dot{\phi}^2 \sin(\theta - \phi) - Mg\alpha \cos \theta \right) \left(m_2 \beta^2 \right) \right.$$

$$\left. - \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \left(m_2 \alpha \beta \dot{\theta}^2 \sin(\theta - \phi) - m_2 g \beta \cos \phi \right) \right]$$

$$\cdot 1 / \left[\left(M\alpha^2 \right) \left(m_2 \beta^2 \right) - \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \right]$$

$$\ddot{\phi} = \left[\left(M\alpha^2 \right) \left(m_2 \alpha \beta \dot{\theta}^2 \sin(\theta - \phi) - m_2 g \beta \cos \phi \right) \right.$$

$$\left. - \left(- m_2 \alpha \beta \dot{\phi}^2 \sin(\theta - \phi) - Mg\alpha \cos \theta \right) \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \right]$$

$$\cdot 1 / \left[\left(M\alpha^2 \right) \left(m_2 \beta^2 \right) - \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \left(m_2 \alpha \beta \cos(\theta - \phi) \right) \right]$$

These solutions via Cramer's rule can easily be plugged into a computer simulation and numerically integrated in order to show a visualization of the chaotic motion of a double pendulum, a chaos which arises from the tremendous complexity of the equations which we have derived here in this paper.