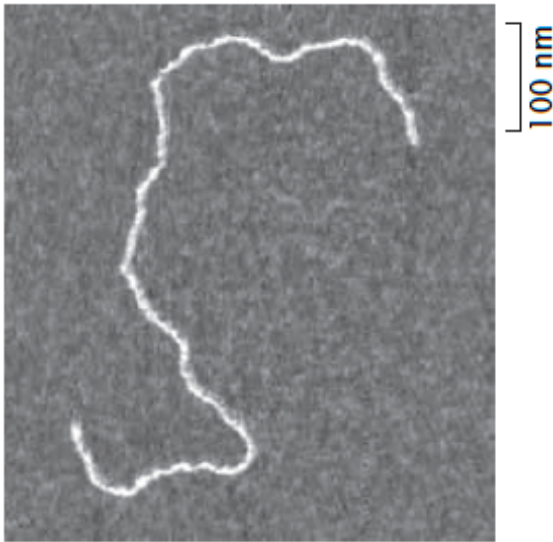


Today's class:

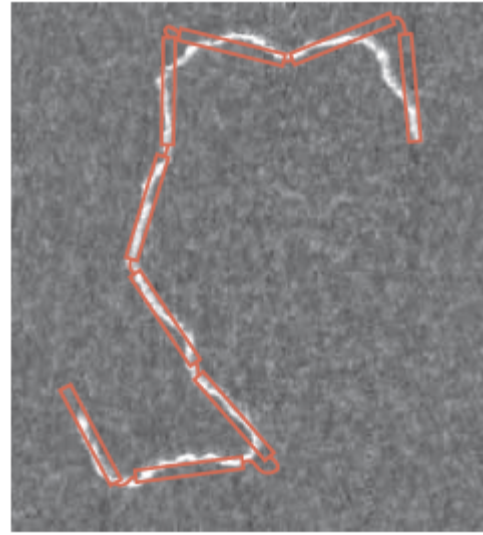
Deformable DNA described as Random Walks

*This lecture follows the parts of chapter 8 from the book
'Physical Biology of the Cell' by Philips, Kondev, Theriot and Garcia, 2nd Ed*

DNA structures can be represented as Random Walks



DNA on a surface captured by AFM
Wiggins et al Nature Nanotech 2006



DNA approximated by an array of
randomly oriented segments



Each macromolecular configuration is a random walk along contour

Each step of random walk = one 'rigid' segment

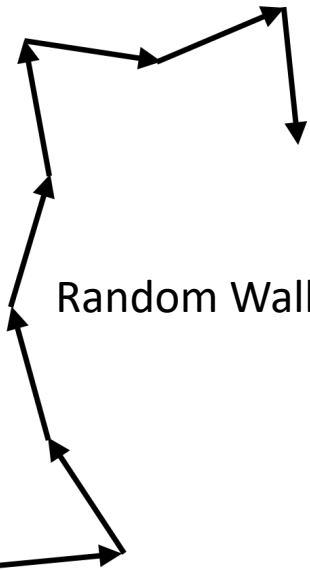
Independent steps, no dependence on history

For a 1D random walker, $p_R = p_L = 1/2$

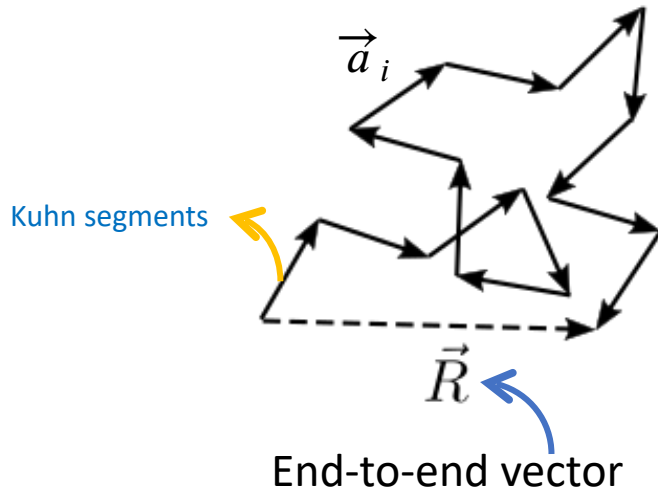
For a chain of N segments, 2^N macromolecular configurations



Random Walk



The Freely Jointed Chain model



Non-interacting polymer of N Kuhn segments of length a

Fully unfolded length, $L = Na$

Some important results:

$$\langle \vec{R} \rangle = \left\langle \sum_{i=1}^N \vec{a}_i \right\rangle = \sum_{i=1}^N \langle \vec{a}_i \rangle = 0$$

$$\begin{aligned} \langle R^2 \rangle &= \langle \vec{R} \cdot \vec{R} \rangle \\ &= \left\langle \sum_{i,j} \vec{a}_i \cdot \vec{a}_j \right\rangle = \left\langle \sum_i a_i^2 \right\rangle + \left\langle \sum_{i \neq j} \vec{a}_i \cdot \vec{a}_j \right\rangle \end{aligned}$$

$= 0$

RMS end-to-end distance

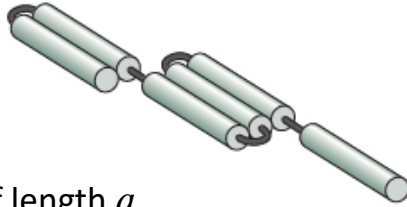
$$\sqrt{\langle R^2 \rangle} = a\sqrt{N}$$

$$\langle R^2 \rangle = Na^2$$

Volume occupied by the polymer

$$V_{poly} \propto \left(\sqrt{\langle R^2 \rangle} \right)^3 = a^3 N^{\frac{3}{2}} > \text{total volume of the unjointed monomers}$$

The probabilistic interpretation of the FJC

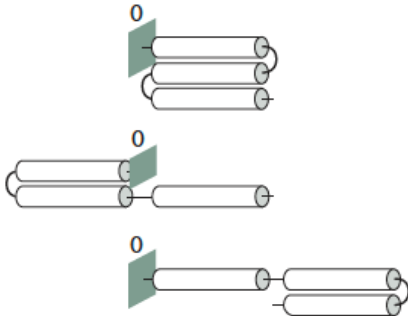


For a 1D chain

N segments of length a

n_r right directed segments

$n_l = N - n_r$ left directed segments



$N = 3, n_r = 2$
Degeneracy = 3

At thermodynamic equilibrium, all states are equally probable for the FJC

We can use this to compute the probability distribution function of the end-to-end distance

Each step occurs with probability $1/2$

Probability of any sequence of N steps $= \left(\frac{1}{2}\right)^N$

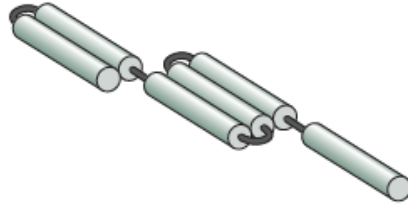
Number of ways to realize $n_r \rightarrow \frac{N!}{n_r!(N - n_r)!}$ From degeneracy

$$\text{Total probability, } p(n_r) = \frac{N!}{n_r!(N - n_r)!} \left(\frac{1}{2}\right)^N$$

Problem: Show $p(n_r)$ is normalized.

Hint: use binomial theorem

How to get $p(R)$ from $p(n_r)$?



For a 1D chain

N segments of length a

n_r right directed segments

$n_l = N - n_r$ left directed segments

$$R = (n_r - n_l)a$$

$$n_r = \left(\frac{N}{2} + \frac{R}{2a} \right)$$

Making replacements in

$$p(n_r) = \frac{N!}{n_r!(N - n_r)!} \left(\frac{1}{2} \right)^N$$

$$p(R) = \frac{N!}{\left(\frac{N}{2} + \frac{R}{2a} \right)! \left(\frac{N}{2} - \frac{R}{2a} \right)!} \left(\frac{1}{2} \right)^N$$

When unfolded length $L = Na \gg R$



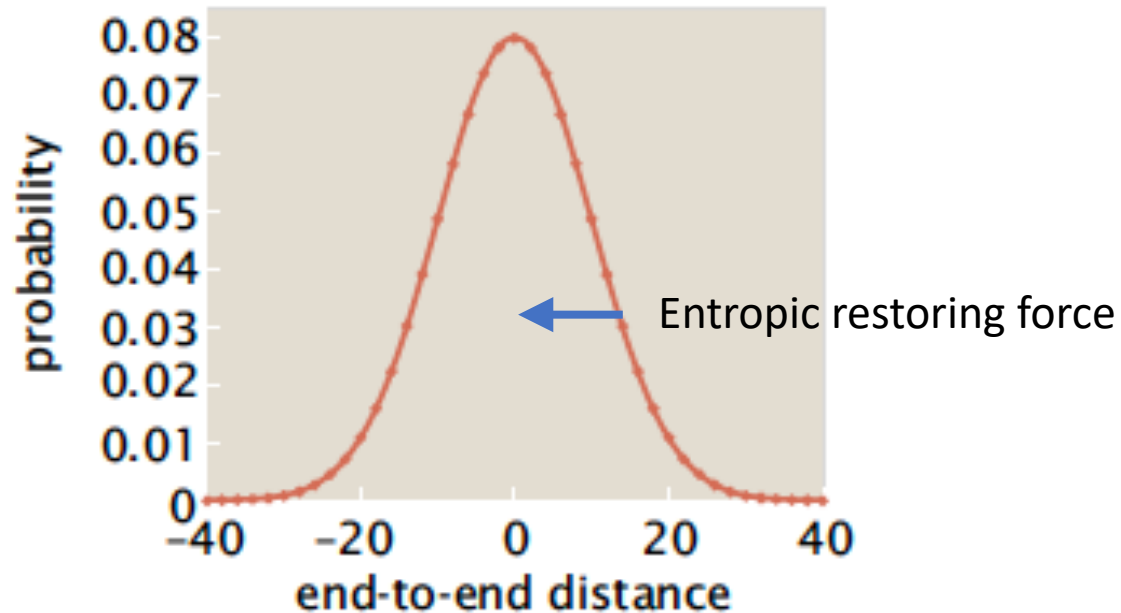
$$p(R) = \frac{2}{\sqrt{2\pi N}} e^{-R^2/2Na^2}$$

Gaussian distribution!

Here $p(R)$ has zero mean and variance Na^2

Analysis of the end-to-end probability

End-to-end probability for a 1D polymer with 100 segments



Most probable state, $R = 0$

If the chain is stretched ($R > 0$) it experiences an 'entropic restoring force'

Entropic Elasticity

End-to-end probability for E.coli genome

For E.coli genome $\rightarrow 5 \times 10^6$ bp

If this is modeled as an open DNA chain it will have $N \approx 15000$ segments, since estimated Kuhn segment ≈ 300 bp long

Deduce the probabilities of $R = 0, 100a, 1000a$

Hint: probability is given by: $p(R) = \frac{2}{\sqrt{2\pi N}} e^{-R^2/2Na^2}$

Probability of $R = 0 \approx 0.007$

Probability of $R = 100a \approx 0.005$

Probability of $R = 1000a \approx 2 \times 10^{-17}$