

Brownian motion

Gypsum cristals in a *closterium moniliferum*

Movie

The ring of Brownian motion: Past - Presence and Future Trends

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Why you should **not** do Brownian motion

- You know nothing about the subject
- Many very good people worked on it
(Einstein, Langevin, Smoluchowski, Ornstein, Uhlenbeck, Wiener, Onsager, Stratonovich, ...)
- You don't have your own pet theory yet

Why you should do Brownian motion

- You know nothing about the subject
- Many very good people worked on it
- You still can do your own pet theory

Robert Brown (1773-1858)



Source: www.anbg.gov.au



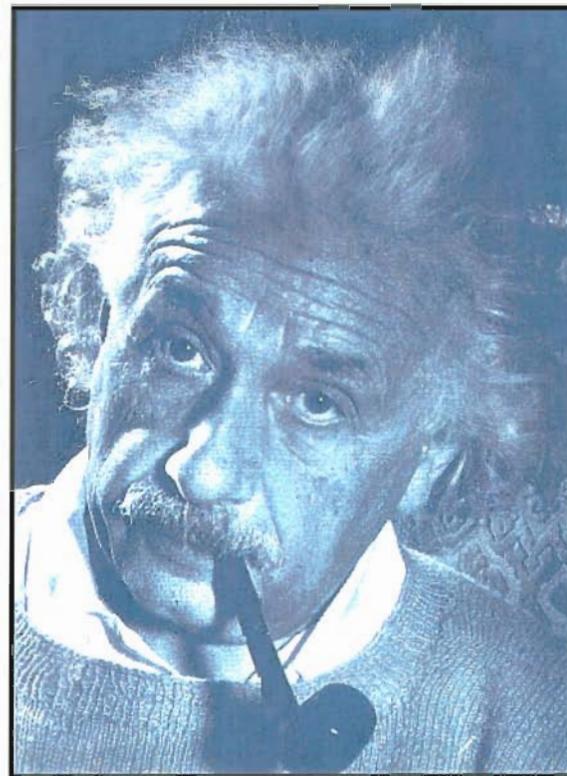
Source: permission kindly granted by Prof. Brian J. Ford
<http://www.brianjford.com/wbbrown.htm>

1827 – irregular motion of granules of pollen in liquids

- Brown, Phil. Mag. **4**, 161 (1928)
- Deutsch: *Did Robert Brown observe Brownian Motion: probably not*, Sci. Am. **256**, 20 (1991)
- Ford: “*Brownian movement in clarkia pollen: a reprise of the first observations*”,
The Microscope **39**, 161 (1991)



Robert Brown



Albert Einstein



Marian Smoluchowski

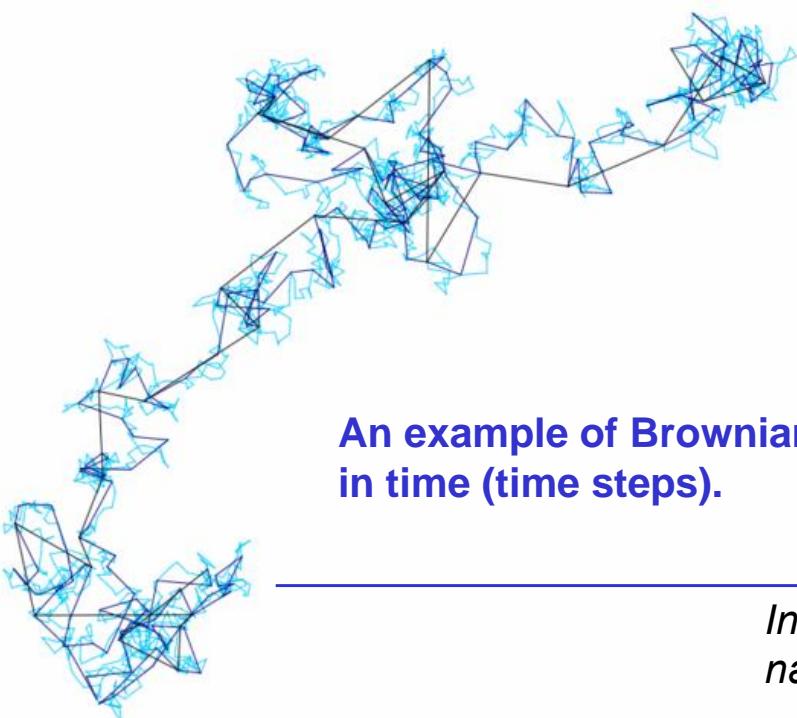
Brownian motion



Robert Brown (1773-1858)

In 1827, the botanist Robert Brown published a study "*A brief account of microscopical observations on the particles contained in the pollen of plants...*", where he reported his observations of irregular, jittery motion of small (clay) particles in pollen grains.

He repeated the same experiment with particles of dust, showing that the motion could not be due to the pollen particles being alive.



An example of Brownian motion of a particle, recorded for three different resolutions in time (time steps).

Incidentally, Robert Brown was also the first to note the ubiquitous nature of a part of eukaryotic cells which he named the "cell nucleus".





Jan Ingen-Housz (1730-1799)



William Sutherland (1859-1911)

Jan Ingen-Housz (1730-1799)



Source: www.americanchemistry.com



To see clearly how one can deceive one's mind on this point if one is not careful, one has only to place a drop of alcohol in the focal point of a microscope and introduce a little finely ground charcoal therein, and one will see these corpuscles in a confused, continuous and violent motion, as if they were animalcules which move rapidly around.

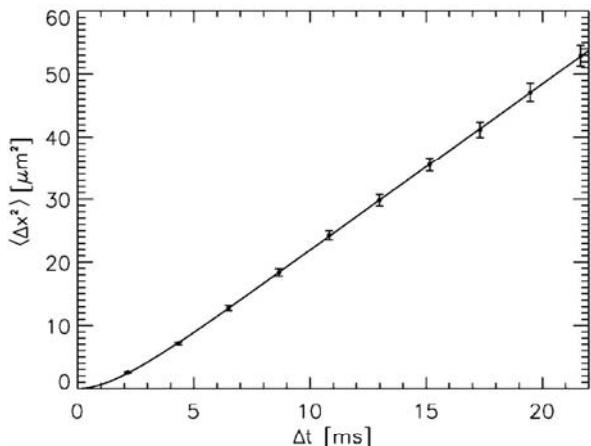
Measurement of the Translational and Rotational Brownian Motion of Individual Particles in a Rarefied Gas

Jürgen Blum,* Stefan Bruns, Daniel Rademacher, Annika Voss, Björn Willenberg, and Maya Krause

*Institut für Geophysik und Extraterrestrische Physik, Technische Universität Braunschweig,
Mendelssohnstraße 3, 38114 Braunschweig, Germany*
(Received 2 August 2006; published 4 December 2006)

We measured the free Brownian motion of individual spherical and the Brownian rotation of individual nonspherical micrometer-sized particles in rarefied gas. Measurements were done with high spatial and temporal resolution under microgravity conditions in the Bremen drop tower so that the transition from diffusive to ballistic motion could be resolved. We find that the translational and rotational diffusion can be described by the relation given by Uhlenbeck and Ornstein [Phys. Rev. **36**, 823 (1930)]. Measurements of rotational Brownian motion can be used for the determination of the moments of inertia of small particles.

$$\langle \Delta x^2 \rangle = 2D\Delta t \left(1 - \frac{\tau_f}{\Delta t} + \frac{\tau_f}{\Delta t} \exp \left[-\frac{\Delta t}{\tau_f} \right] \right)$$



Theory of Brownian motion

W. Sutherland (1858-1911)

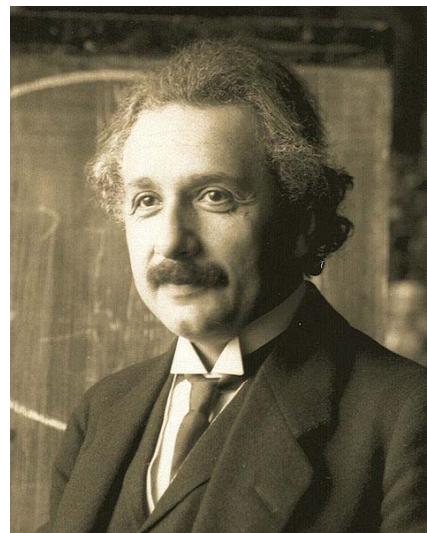


Source: www.theage.com.au

$$D = \frac{RT}{6\pi\eta aC}$$

Phil. Mag. **9**, 781 (1905)

A. Einstein (1879-1955)



Source: [wikipedia.org](https://en.wikipedia.org)

$$\langle x^2(t) \rangle = 2Dt$$

$$D = \frac{RT}{N} \frac{1}{6\pi k P}$$

Ann. Phys. **17**, 549 (1905)

M. Smoluchowski (1872-1917)



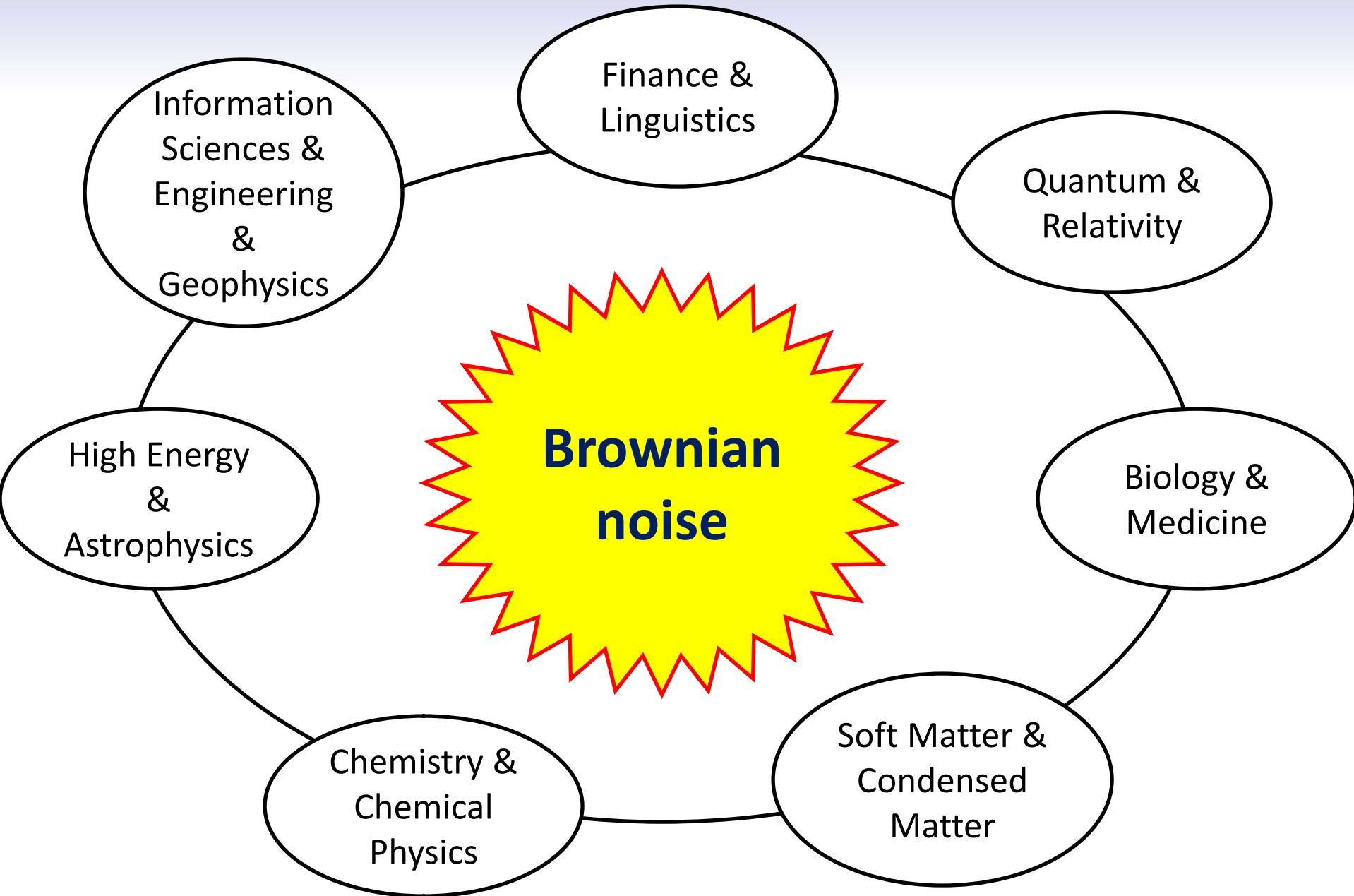
Source: [wikipedia.org](https://en.wikipedia.org)

$$D = \frac{32}{243} \frac{mc^2}{\pi\mu R}$$

Ann. Phys. **21**, 756 (1906)

<http://www.physik.uni-augsburg.de/theo1/hanggi/>

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- [Das Universum bremst nicht](#)
- [100 years of Brownian motion: historical items and surveys](#)
- [Thermodynamics and statistical physics of small systems](#)
- [Palermo Lecture Series, November 2016 "Classical and Quantum Thermodynamics and Statistical Physics"](#)



Quantum-Mechanics

= Brownian Motion ?

= Stochastic Mechanics ?

$$p(x, t) = |\Psi(x, t)|^2$$

(E. Nelson; 1966, 1986)

Schrödinger-
gleichung

$$\Psi(x, t) = |\Psi(x, t)| e^{iS(x, t)}$$

$$\dot{p}(x, t) = -\frac{\hbar}{m} \nabla [(\nabla \ln |\Psi(x, t)| + \nabla S(x, t)) p(x, t)] + \frac{\hbar}{2m} \nabla^2 p(x, t)$$

$$\mathbf{f}_1 \geq 0, \mathbf{f}_2 \geq 0 : \quad \frac{1}{2} \langle \mathbf{f}_1(t_1) \mathbf{f}_2(t_2) + \mathbf{f}_2(t_2) \mathbf{f}_1(t_1) \rangle \geq 0$$

QM: NO !

Diffusion: space-time only

classical diffusion (Markovian)

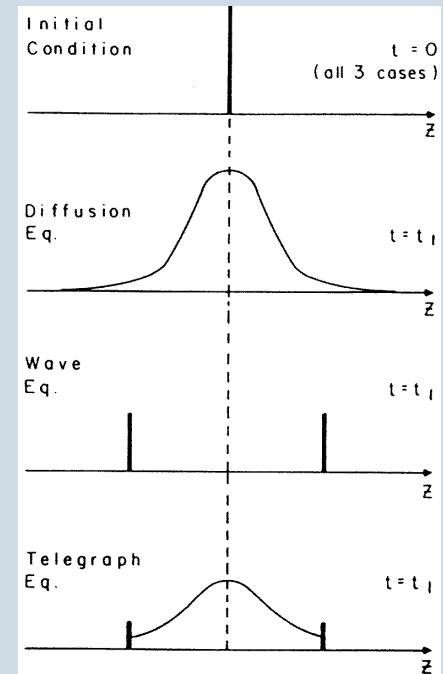
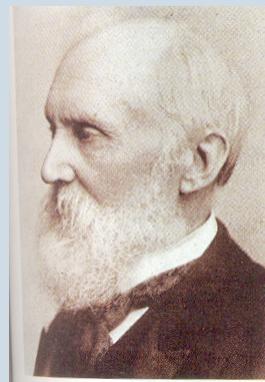
$$p(t, x|t_0, x_0) = \left[\frac{1}{4\pi \mathcal{D}(t - t_0)} \right]^{1/2} \exp \left[-\frac{(x - x_0)^2}{4\mathcal{D}(t - t_0)} \right].$$

telegraph equation (non-Markovian)

$$\tau_v \frac{\partial^2}{\partial t^2} \varrho + \frac{\partial}{\partial t} \varrho = \mathcal{D} \nabla^2 \varrho,$$

alternative approach

$$p(\bar{x}|\bar{x}_0) \propto \int_{a_-(\bar{x}|\bar{x}_0)}^{a_+(\bar{x}|\bar{x}_0)} da \exp\left(-\frac{a}{2\mathcal{D}}\right)$$



J Masoliver & G H Weiss
Eur J Phys 17:190 (1996)

PRD 75:043001 (2007)

Relativistic thermodynamics & BM

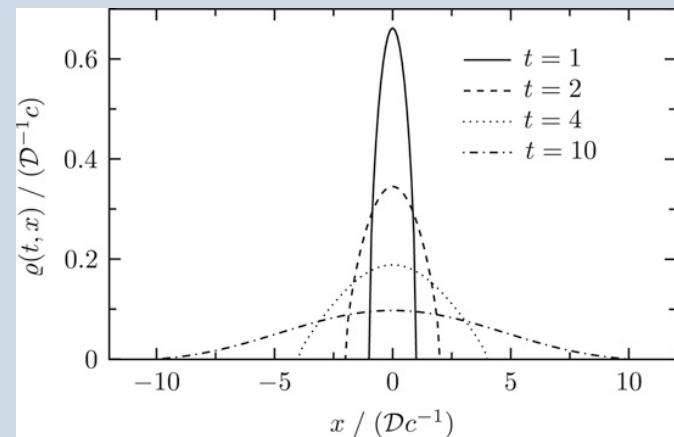
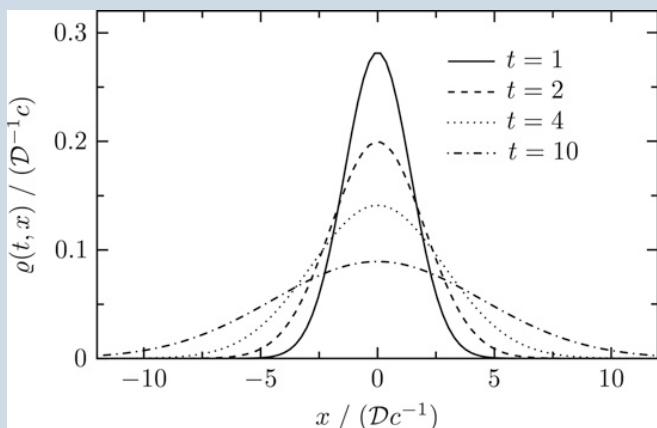
Diffusion propagator

JD et al, PRD 75:043001 (2007)

$$a(\bar{x}|\bar{x}_0) = \frac{1}{2} \int_{t_0}^t dt' v(t')^2,$$

$$a(\bar{x}|\bar{x}_0) = - \int_{t_0}^t dt' [1 - v(t')^2]^{1/2}$$

$$p(\bar{x}|\bar{x}_0) \propto \int_{a_-(\bar{x}|\bar{x}_0)}^{a_+(\bar{x}|\bar{x}_0)} da \exp\left(-\frac{a}{2\mathcal{D}}\right)$$



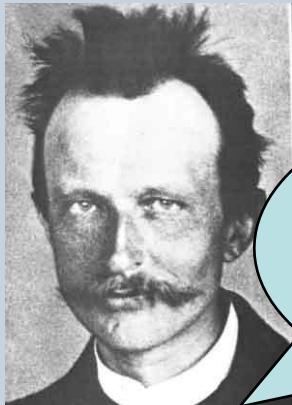
$$a_+ = +\infty$$

$$a_-(\bar{x}|\bar{x}_0) = \frac{(x - x_0)^2}{2(t - t_0)}$$

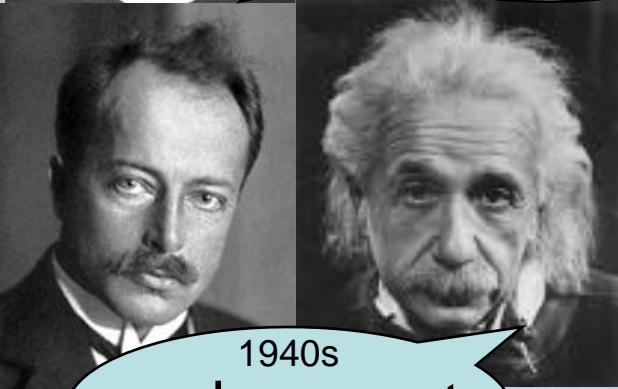
$$a_+ = 0$$

$$a_-(\bar{x}|\bar{x}_0) = -[(t - t_0)^2 - (x - x_0)^2]^{1/2}$$

“Temperature” problem in RTD?



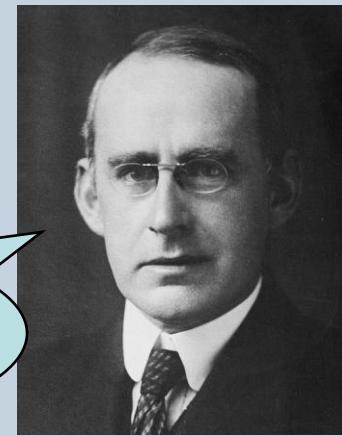
1907/08
moving bodies
appear cooler



1940s
maybe ... not

1923/1963

.. hotter!



$$T'(w) = T (1 - w^2)^{\alpha/2}$$

$$\alpha = \begin{cases} +1 & \text{Planck, Einstein} \\ 0 & \text{Landsberg, van Kampen} \\ -1 & \text{Ott} \end{cases}$$

$T=T'$ 1966-69

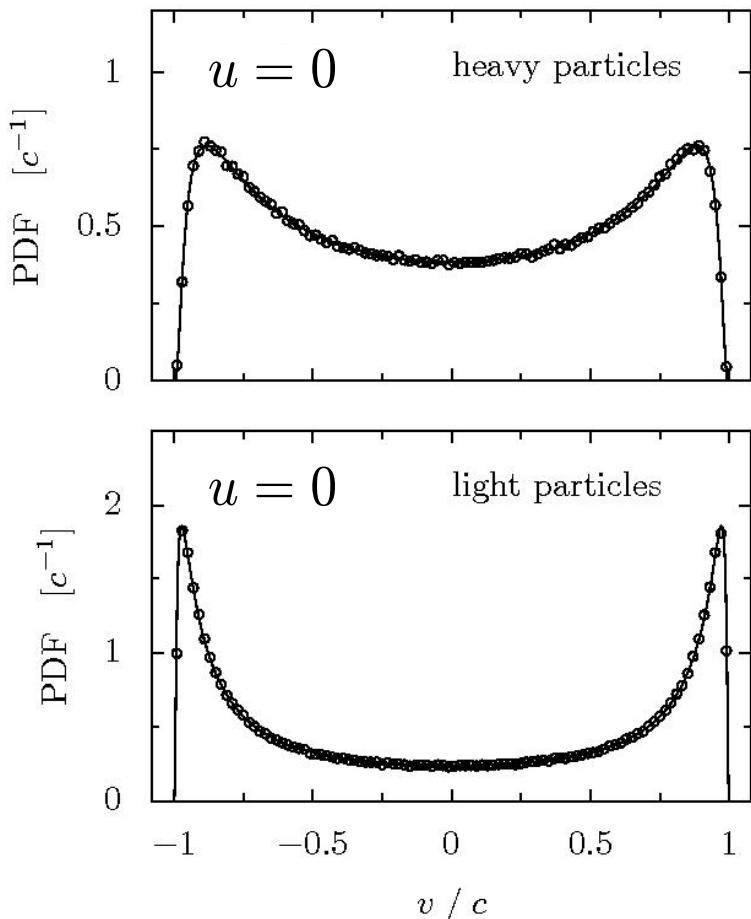
CK Yuen, Amer. J. Phys. 38:246 (1970)



Jüttner Gas

$$f_{\text{Maxwell}}(\vec{p}) = [\beta/(2\pi m)]^{d/2} \exp(-\beta p^2/2m)$$

$$f_{\text{Jüttner}}(\vec{p}) = Z_d^{-1} \exp\left[-\beta_J(m^2 c^4 + p^2 c^2)^{1/2}\right]$$



$\langle \vec{p} \cdot \vec{v} \rangle = dk_B \mathcal{T} = d/\beta_J$

statistical relativistic temperature

$$T = \mathcal{T} = (k_B \beta_J)^{-1}$$

Two prominent examples

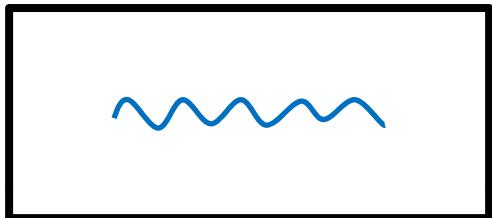
Stochastic
Resonance

Brownian
Motors

Stochastic Resonance

(in a nutshell)

Weak signal

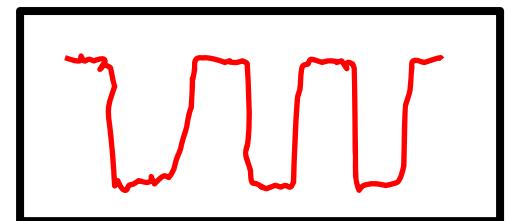


Noise source



System

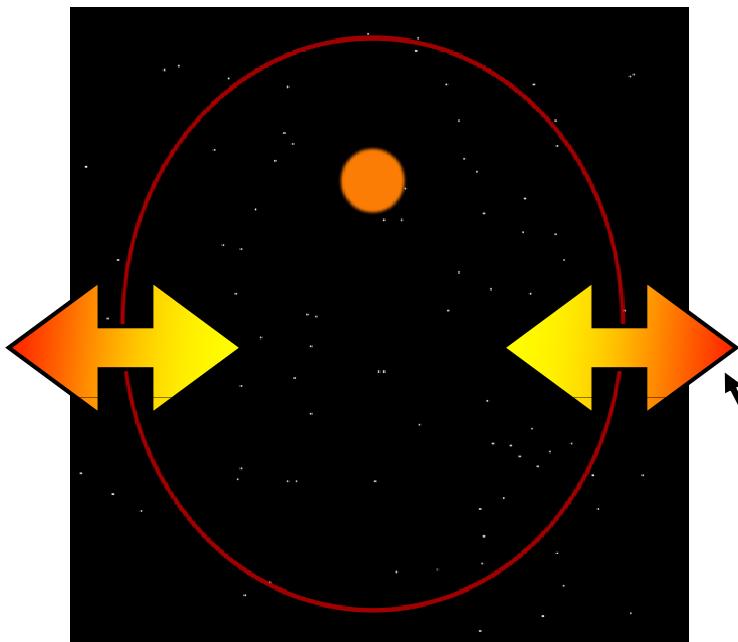
Output signal



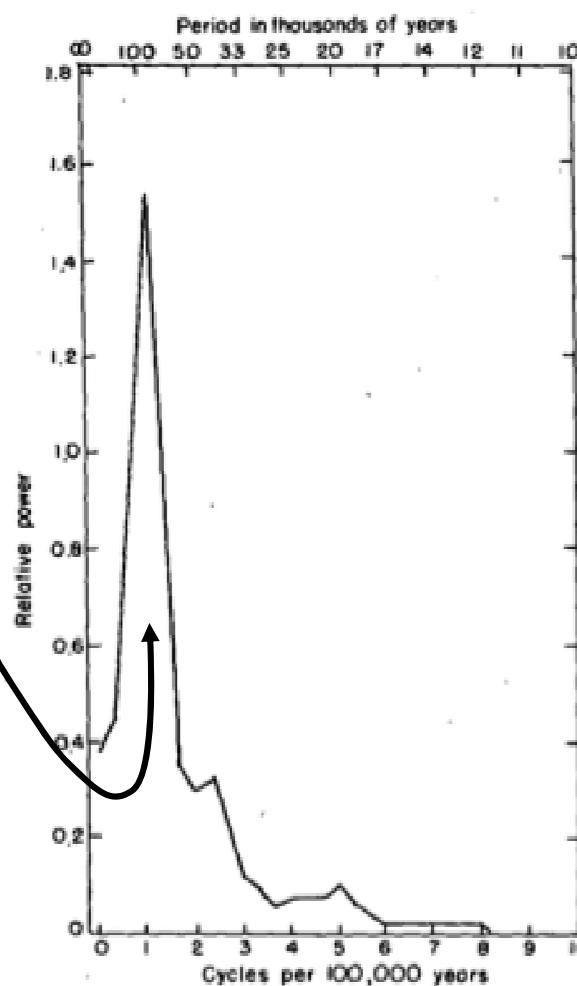
Why are the ice-ages so periodic ?

Milankowitch cycles:

Small changes in earth orbit eccentricity with 100k year periodicity



M. Milankowitch, Handbuch der Klimatologie I
(1930)

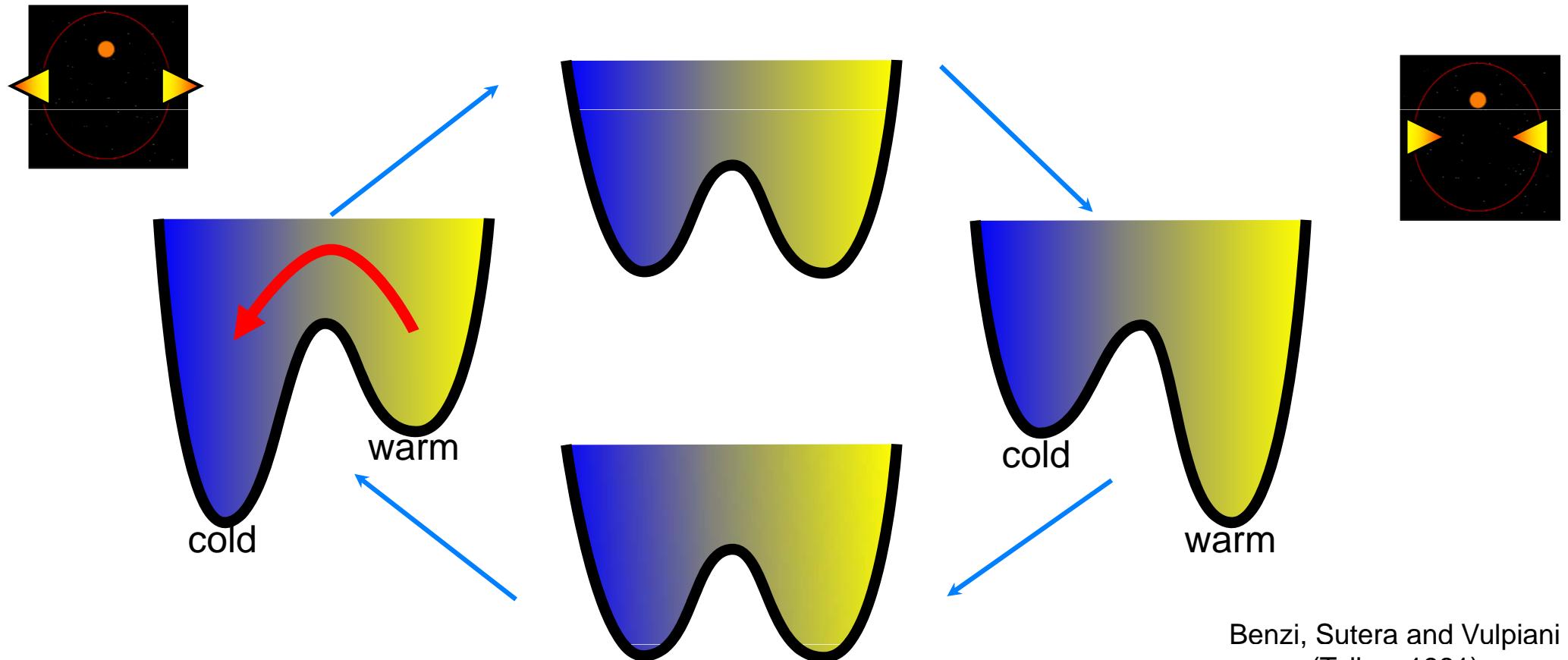


Changes are small!
(<0.1% of solar constant)

What can amplify those small changes ?

Milankowitch Cycles and Bistability

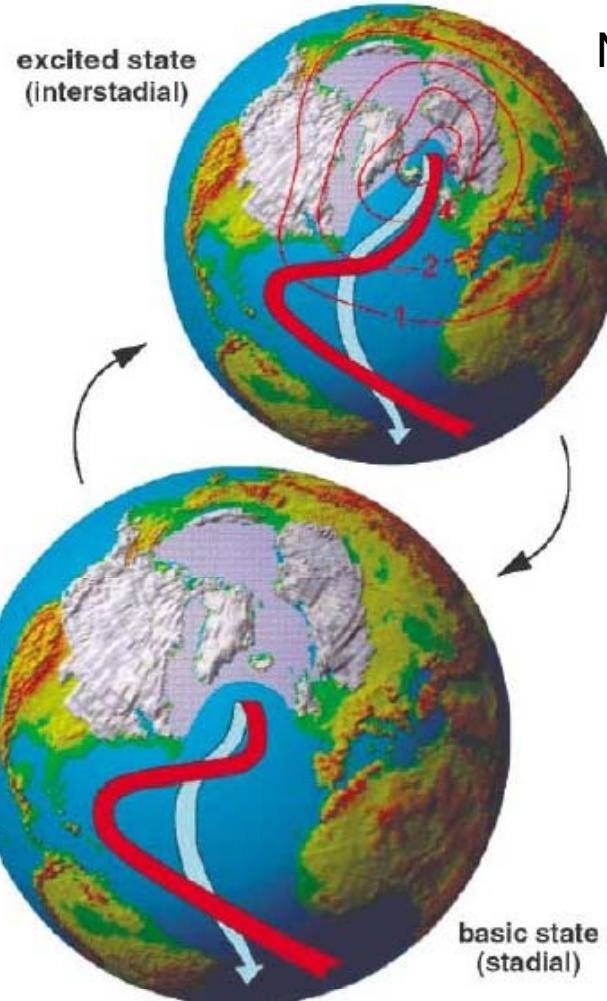
Climate “landscape”



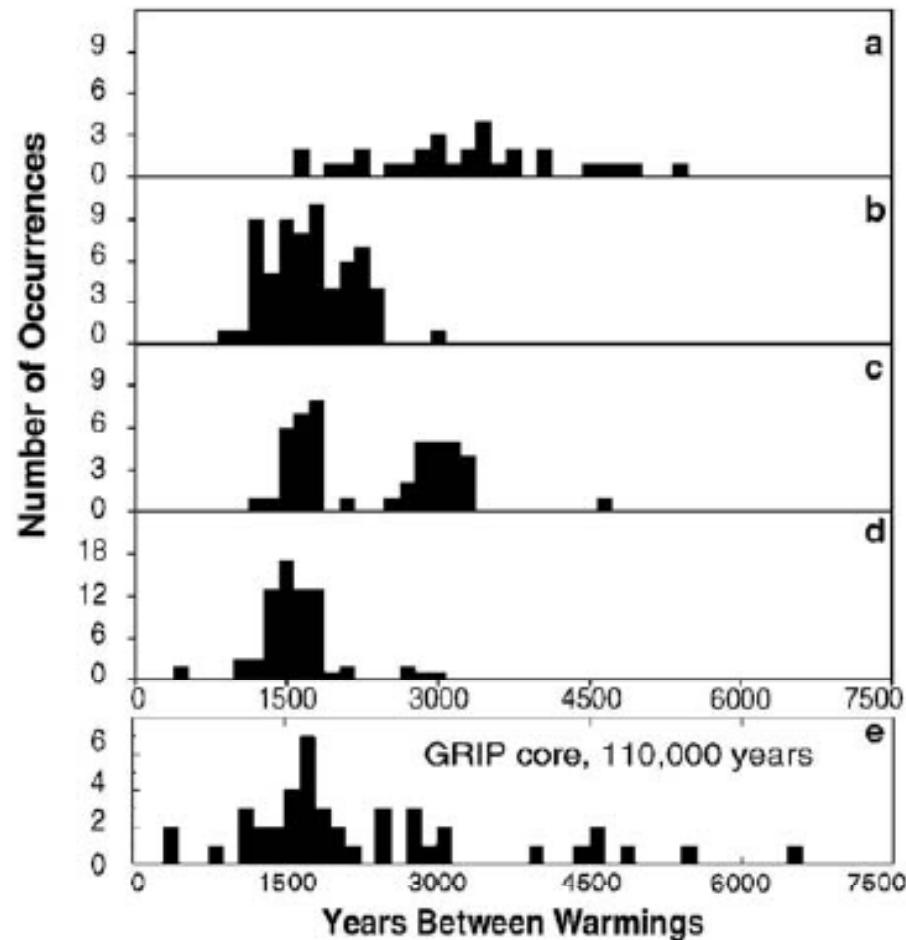
Benzi, Sutera and Vulpiani
(Tellus, 1981)
C. Nicolis and G. Nicolis
(Tellus, 1981)

- The 100ky cycles only bias the climate
- Fluctuations make climate switch
- small changes of conditions can have huge impact

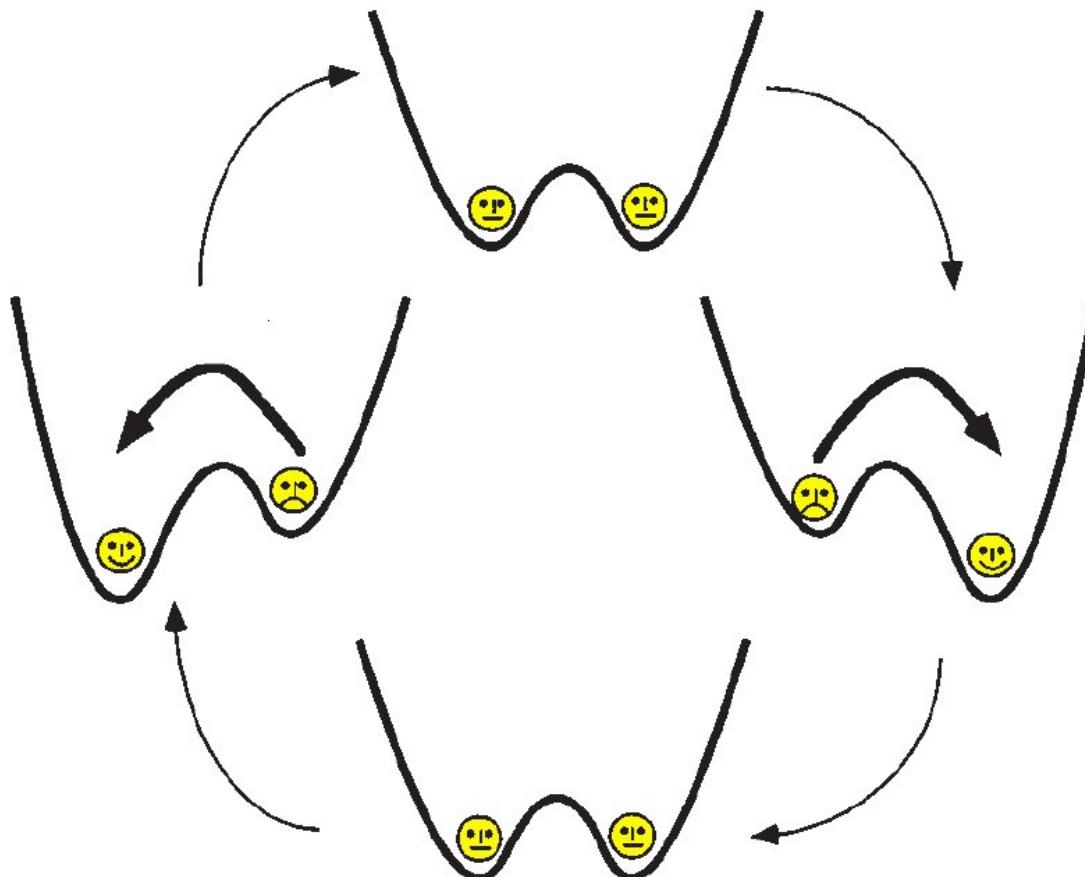
Occurrence of ice ages



Noise and periodic of solar origin



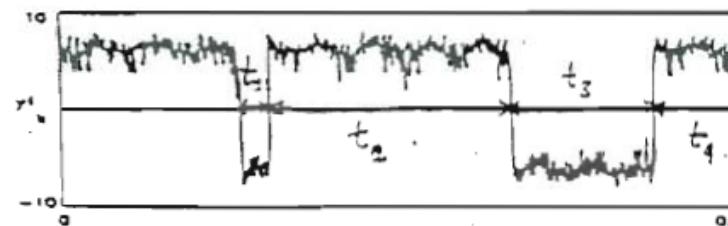
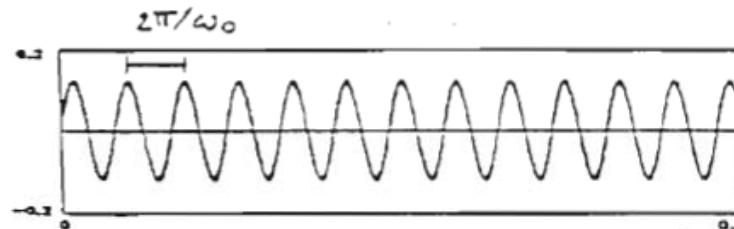
Noise-assisted synchronized hopping



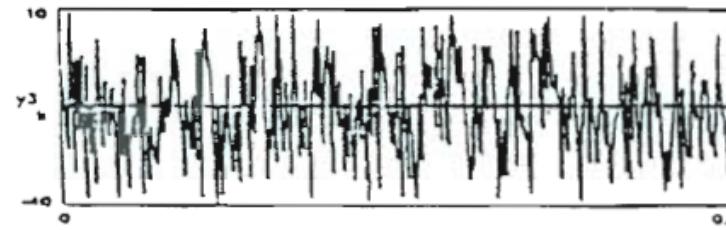
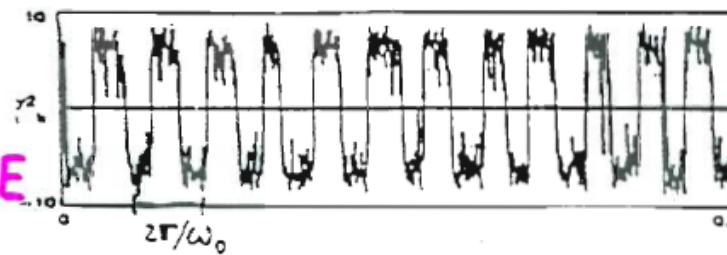
$$T_{\text{period}} \simeq 2T_{\text{escape}}$$

Synchronization

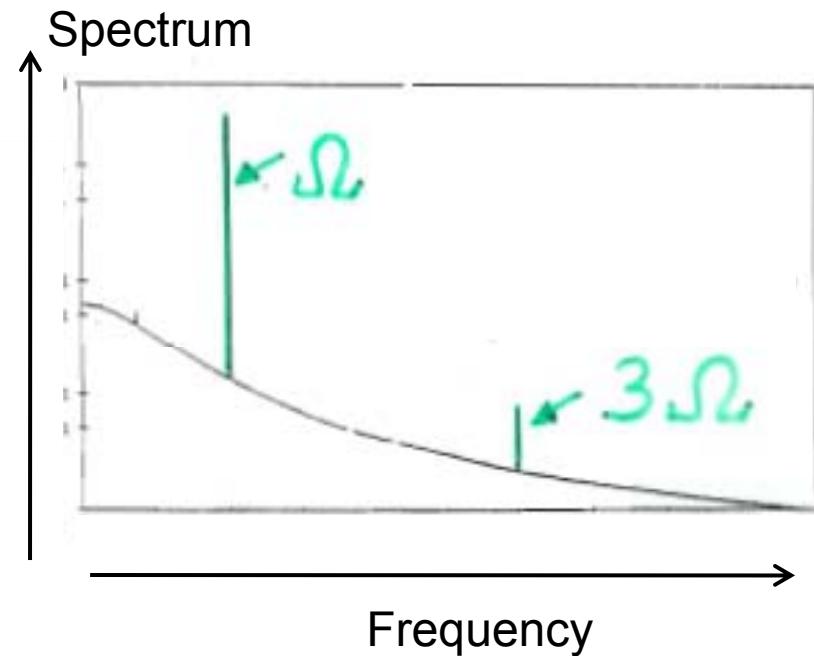
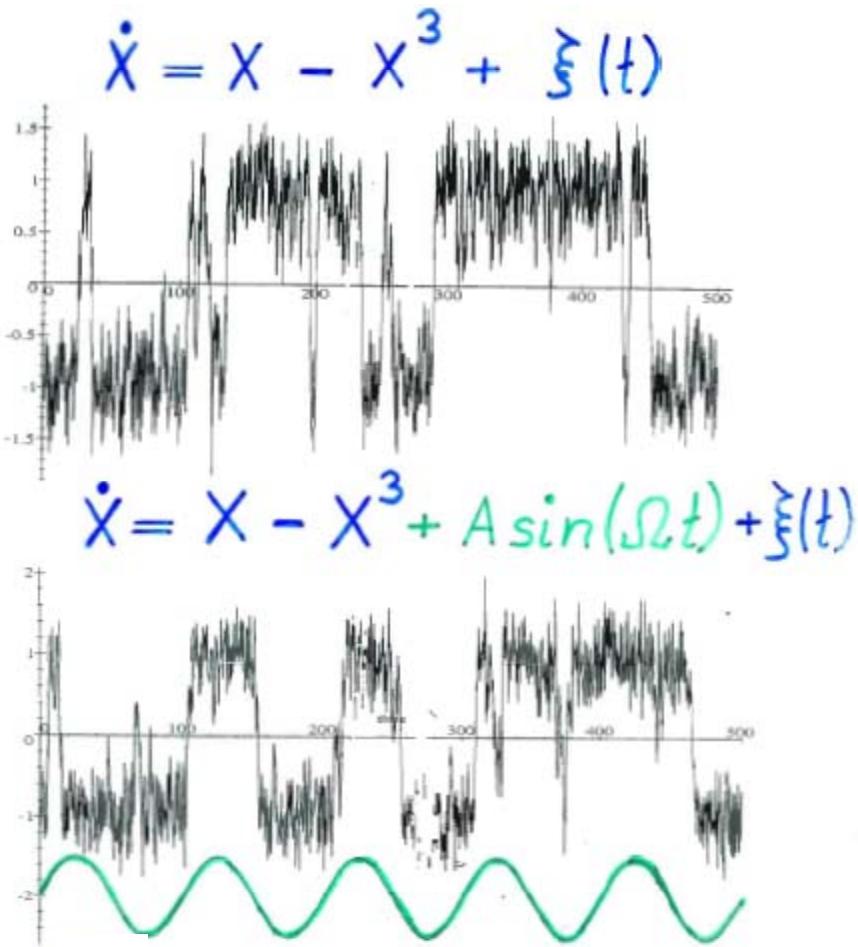
SIGNAL



$T_e \sim 2\Gamma^{-1}$
ESCAPE



Power spectral density



Measuring SR

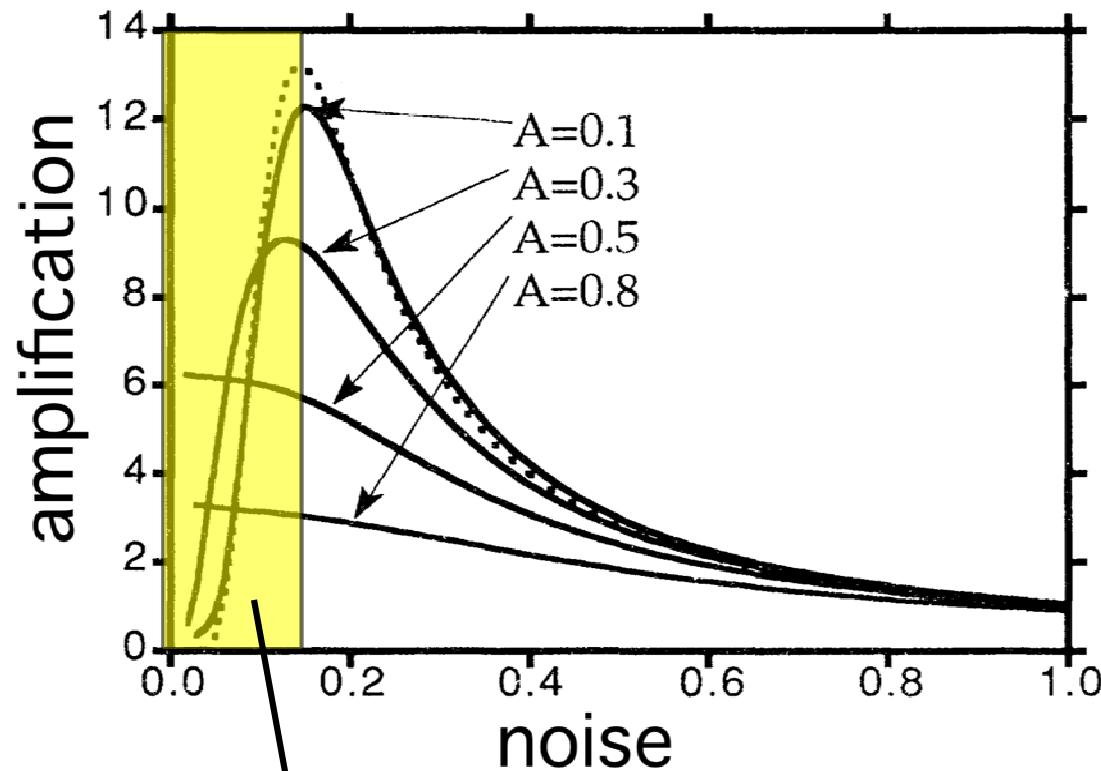
- Signal to noise ratio
- Spectral amplification
- mutual information
- cross-correlation: input \leftrightarrow output
- peak area, (phase-) synchronization, ...

SR-reviews:

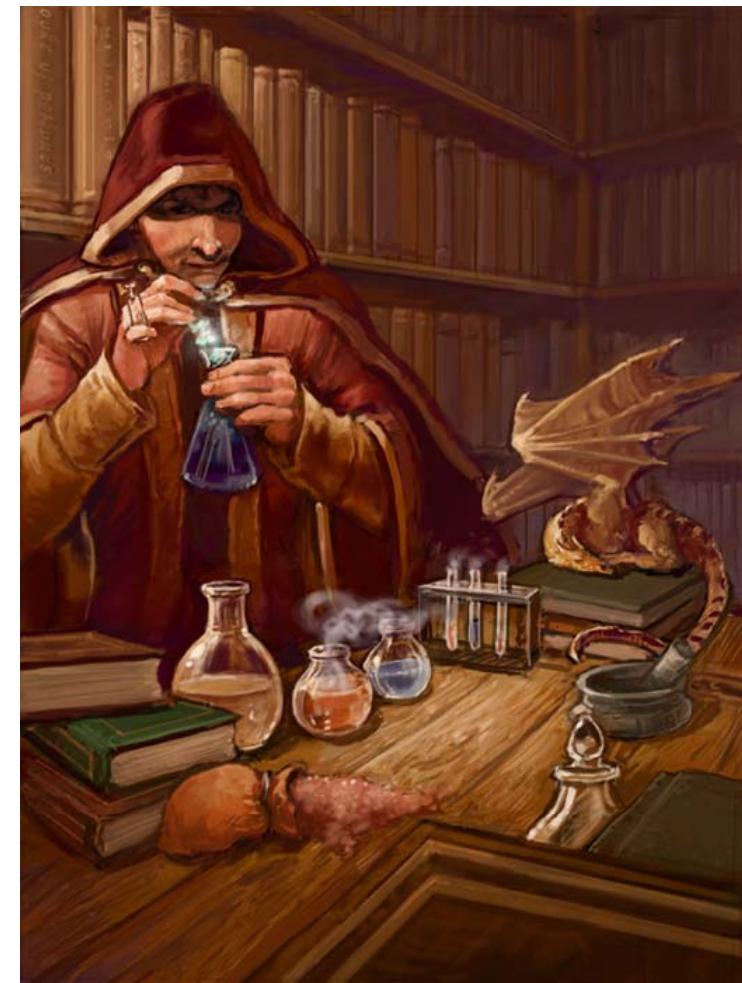
L. Gamaitoni, P. Hänggi, P. Jung, F. Marchesoni, Rev. Mod. Phys. **70**, 223 (1998)
P. Hänggi, ChemPhysChem **3**, 285 (2002)

Amplification of small signals by noise

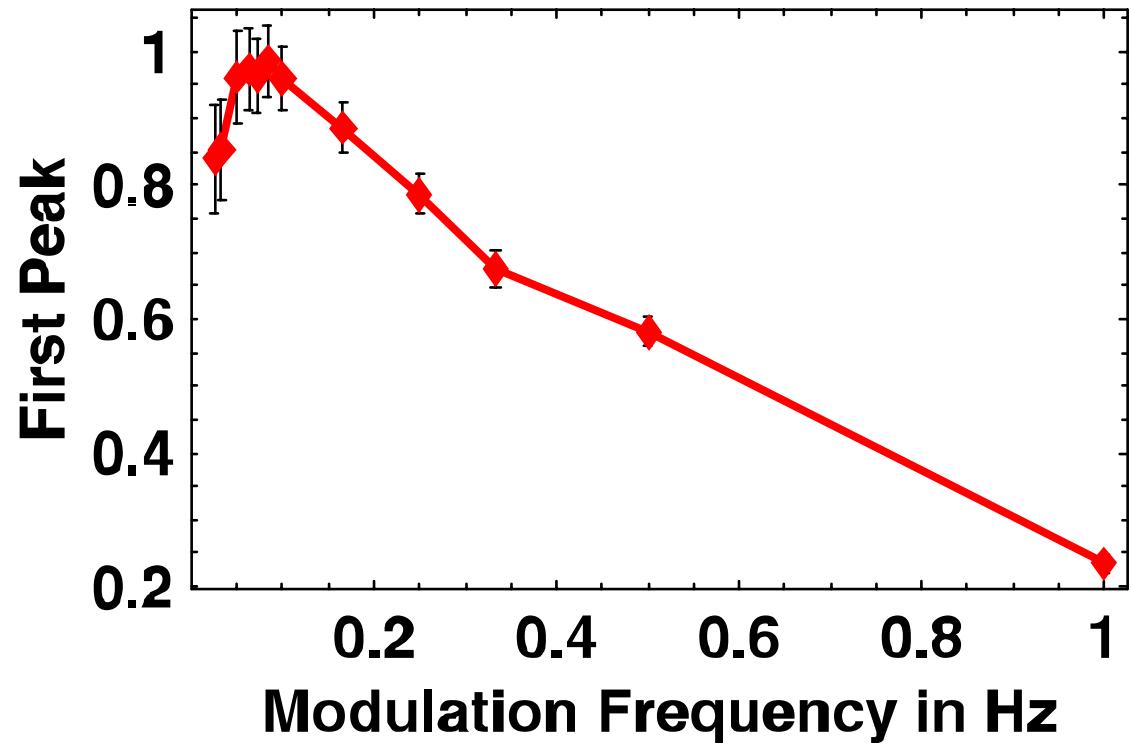
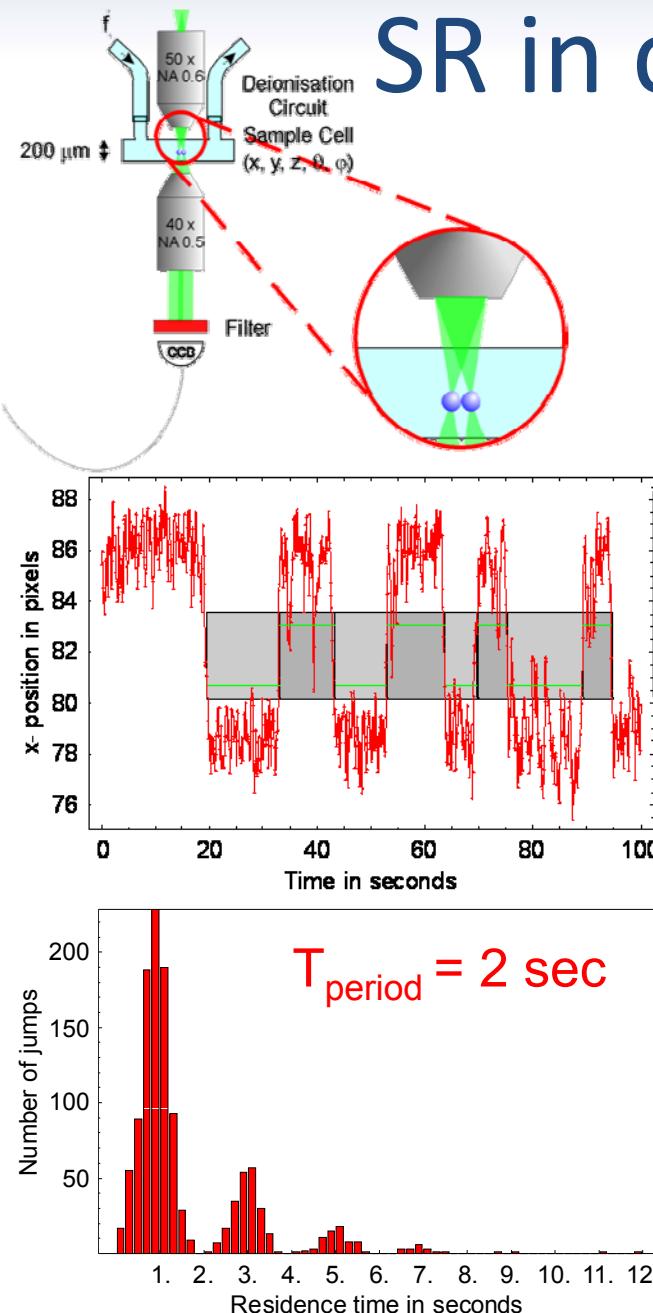
(P. Jung, P. Hänggi, Phys. Rev. A **44**, 8032 (1991))



More noise , more signal !!



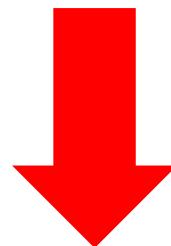
SR in colloidal systems



D. Babic, C. Schmitt, I. Poberaj, C. Bechinger,
Europhys. Lett. **67**, 158 (2004)

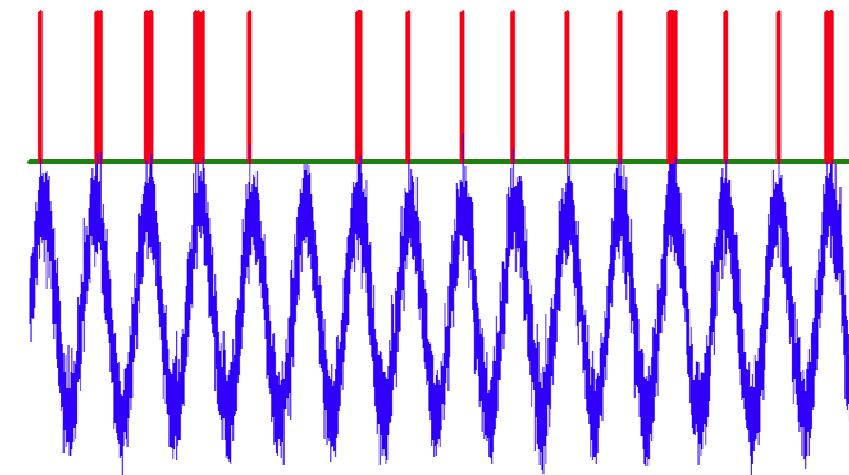
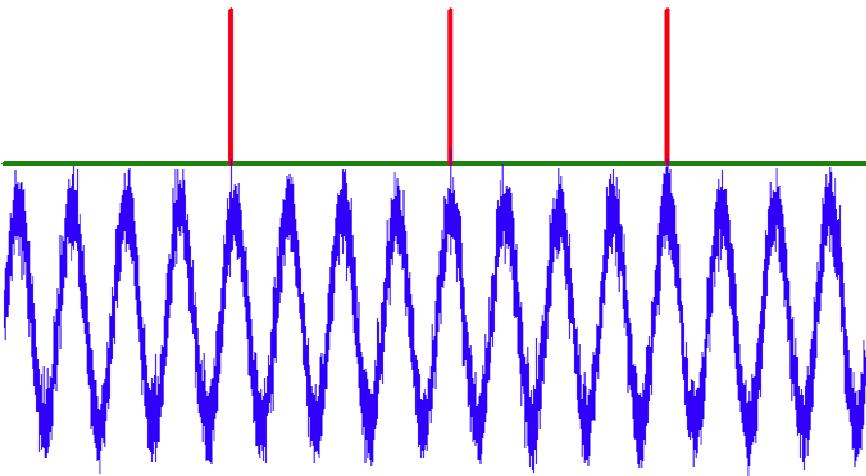
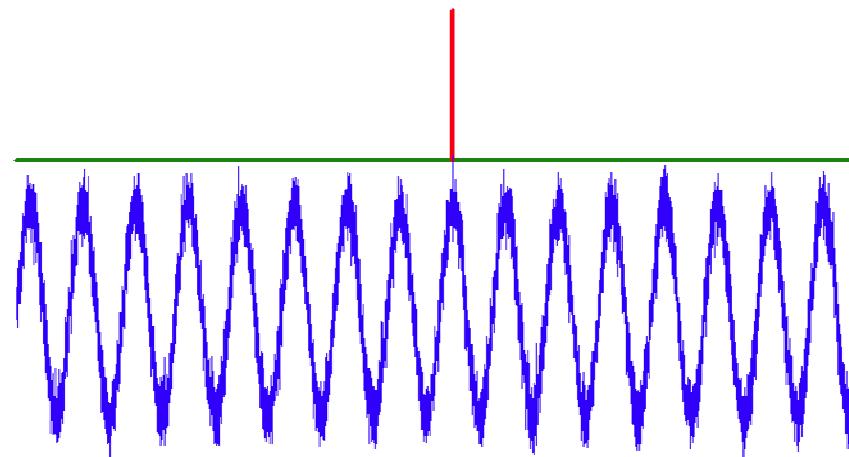
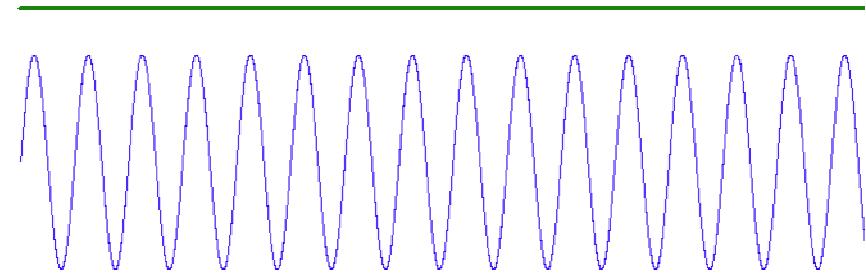
SR - Ingredients

- ✓ Threshold system
- ✓ Weak (subthreshold) signal
- ✓ Noise



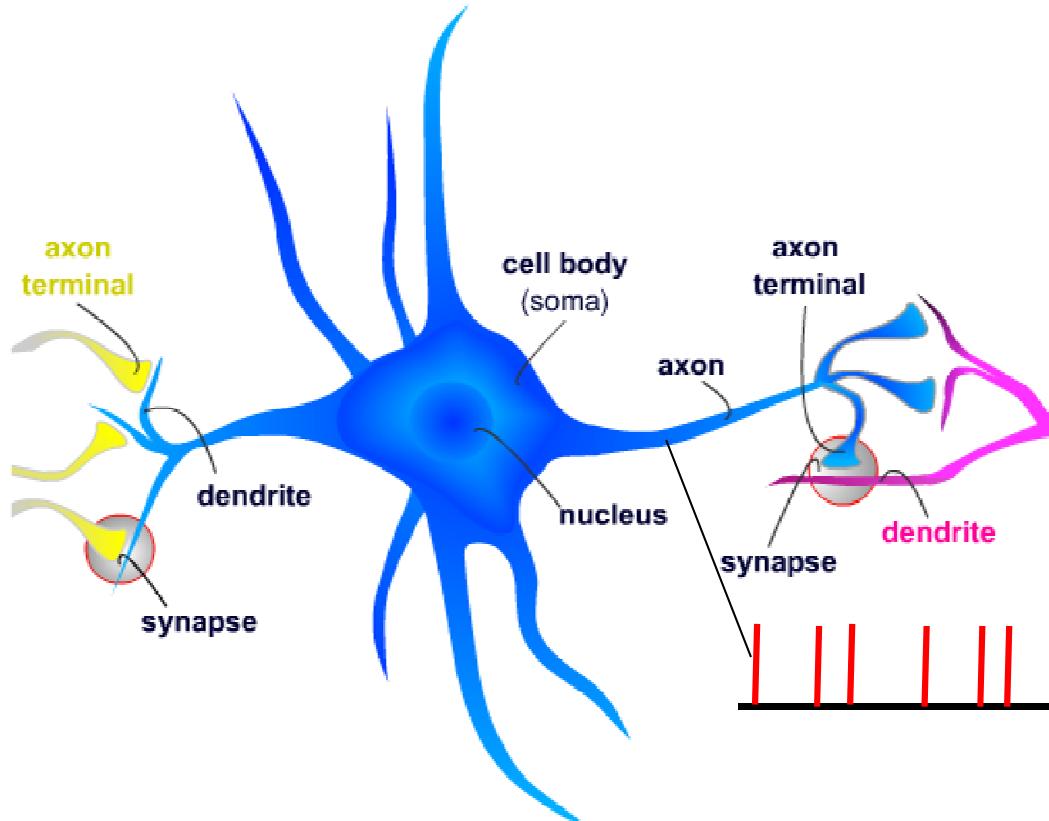
Anomalous amplification properties

Thresholds and Stochastic Resonance



P. Jung, Phys. Rev. E50, 2513 (1994), F. Moss and L. Kiss, EPL, 29 (1995)

Stochastic Resonance in Neurobiology



Input: currents at synapses

Processing: action potential if the sum of currents exceeds threshold

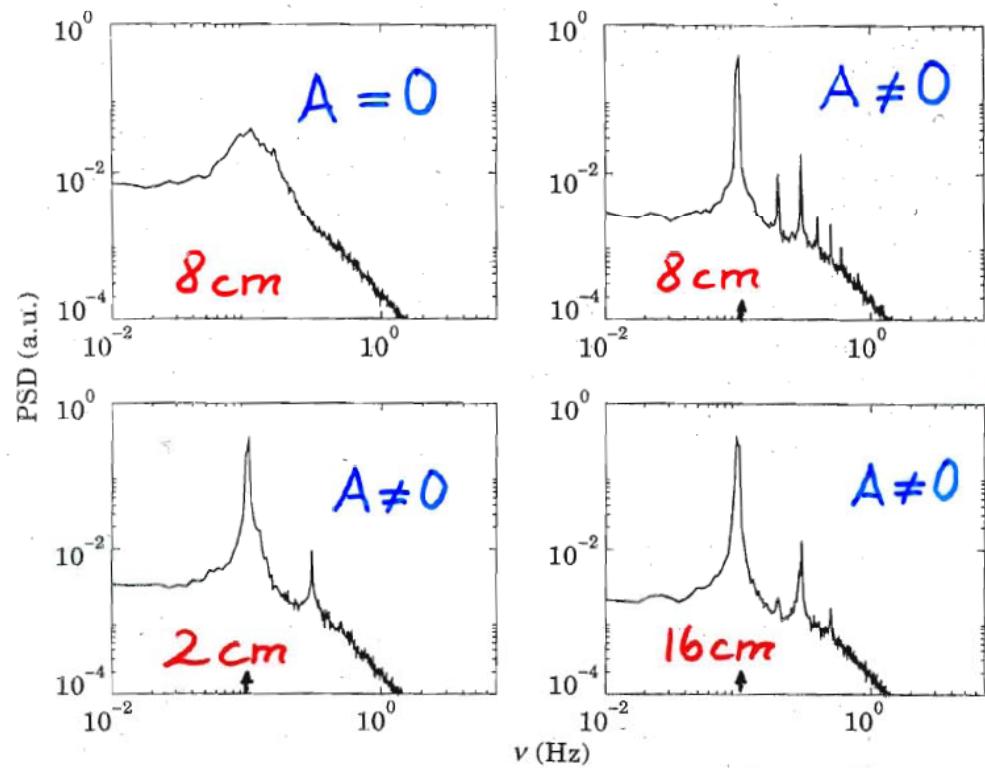
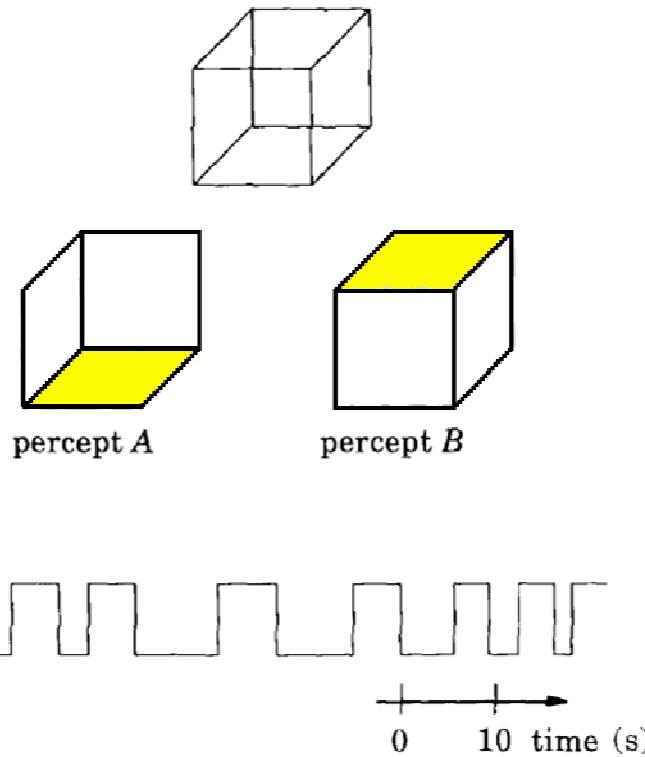
Output: electric pulses traveling down the axon

source: Consortium on Cognitive Science Instruction (CCSI)

Basic idea: Signals below threshold can be detected in the presence of additional noise

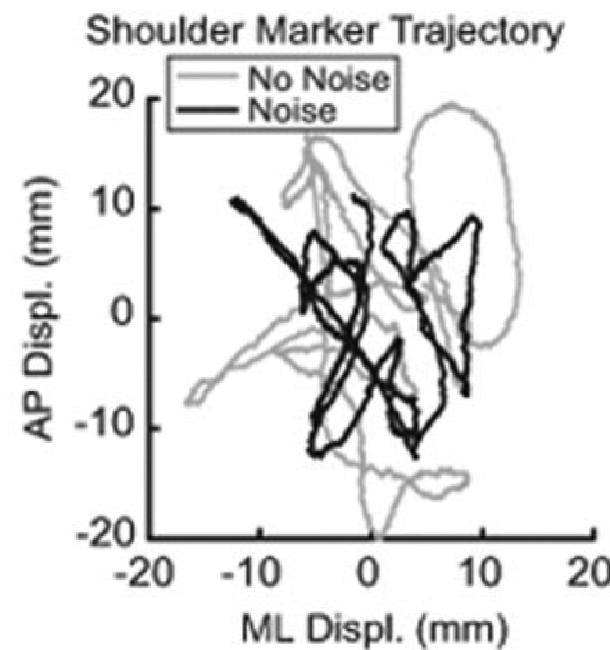
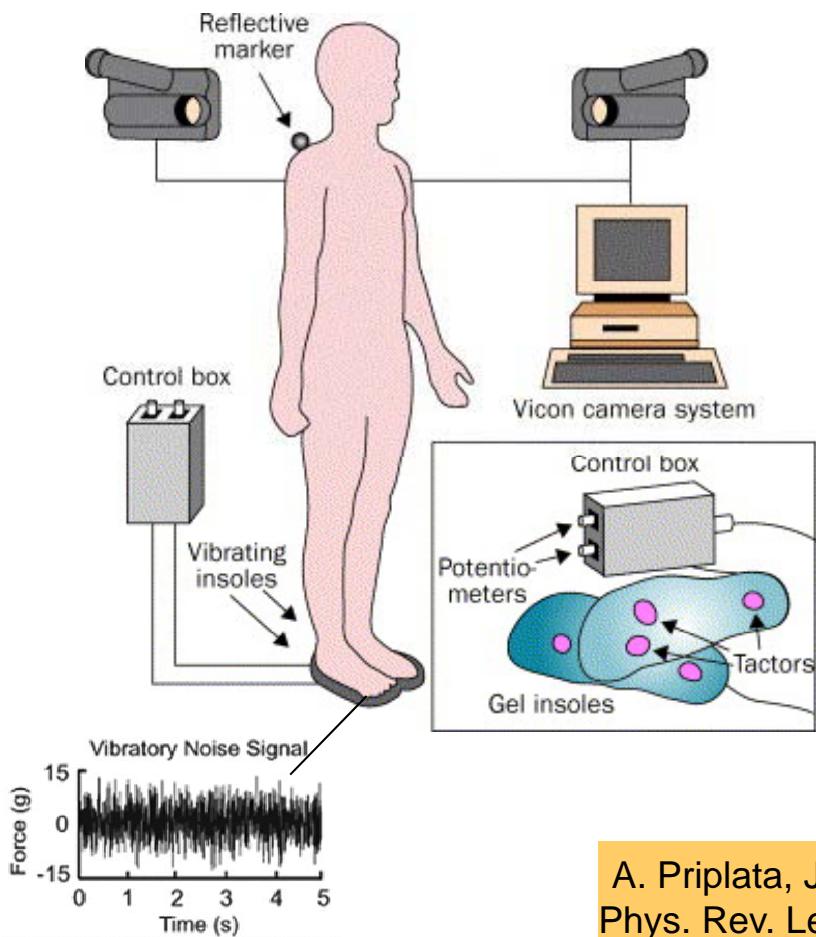
SR in Visual Perception

Experiment:



SR and human posture control

Somatosensory function declines with age and in diabetic patients. Can additional noise help restore function?



Reduction in sway of person

A. Priplata, J. Niemi, M. Salen, J. Harry, L.A. Lipsitz and J.J. Collins
Phys. Rev. Lett. 89 (2002)



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- Multiple sclerosis (MS)
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- Cross-section paralysis
- Depression
- Pain
- ...

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- Parkinson
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- Depression
- Pain
- ...

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SR trends

- Spatio – temporal SR
- Aperiodic SR
- Quantum SR

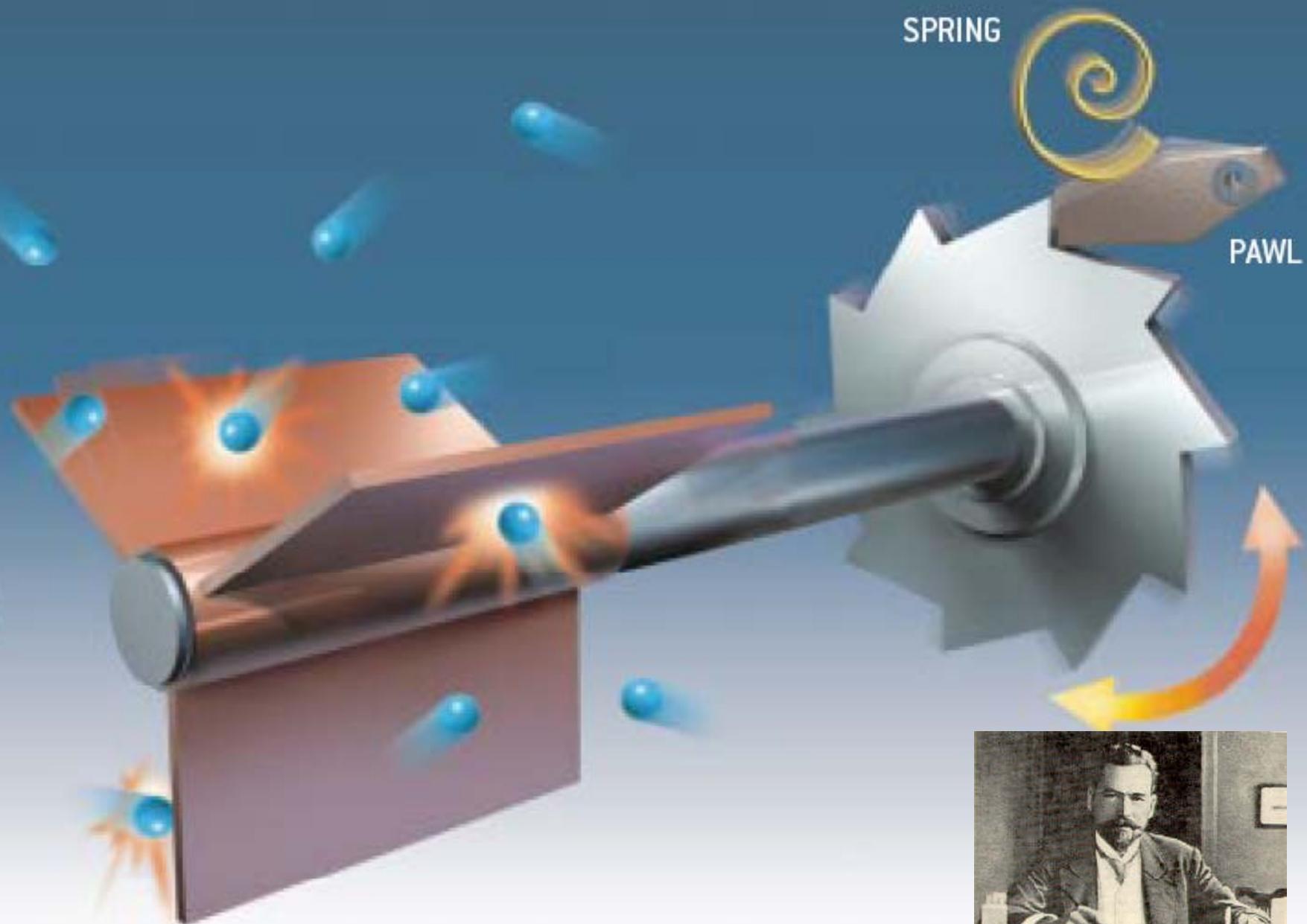
Motors \implies Brownian motors

Two heat reservoirs

One heat reservoir

Perpetuum mobile of the second kind?

NO !



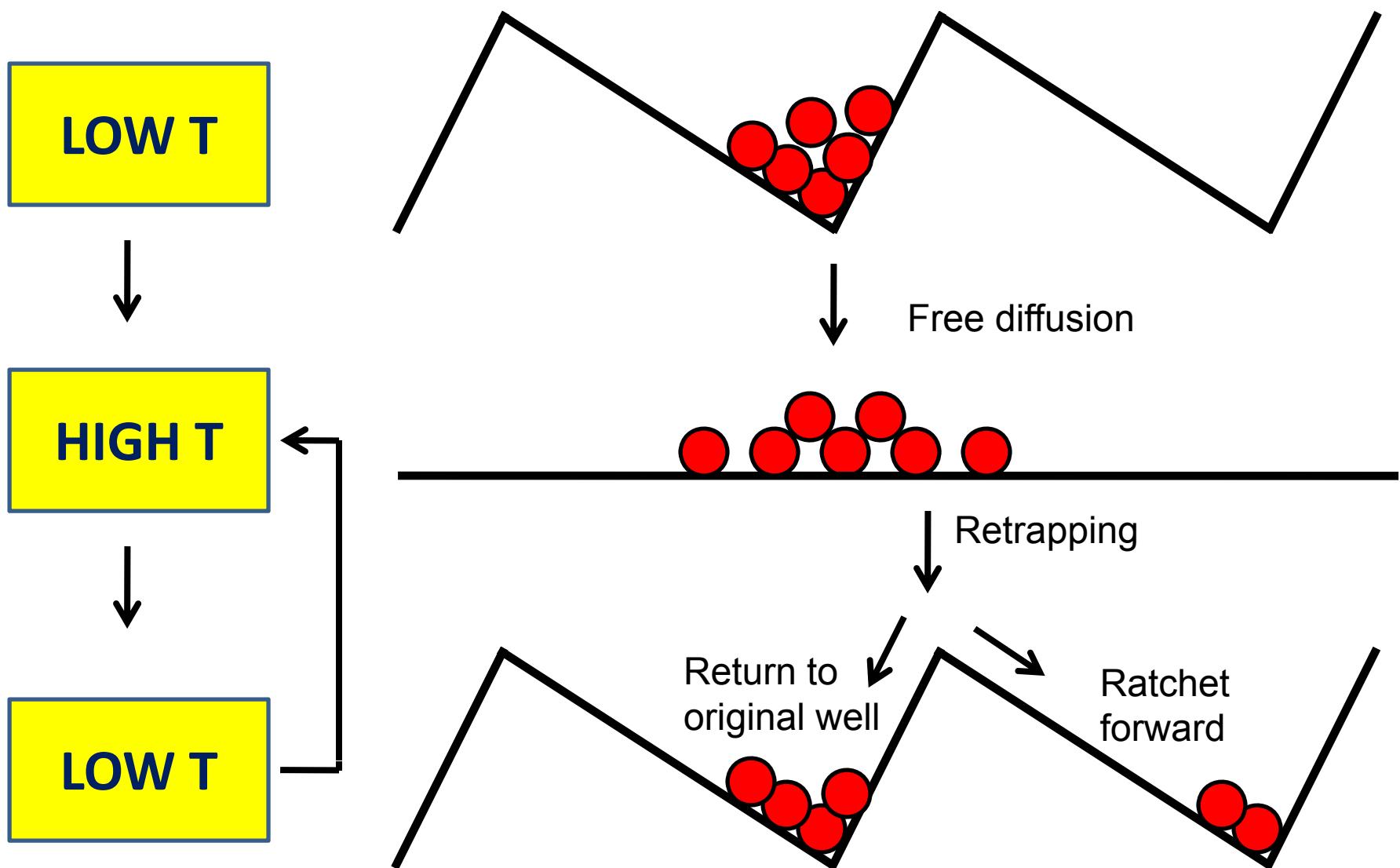
Source: Scientific American (2001)



Brownian motor

Movie

Temperature / Flashing Ratchet



Brownian motors - Characteristics

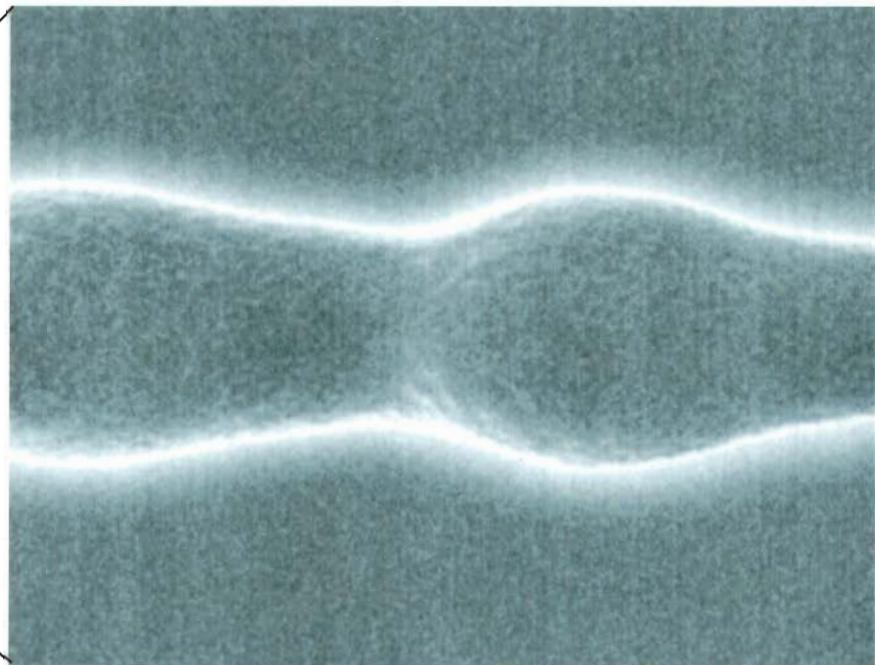
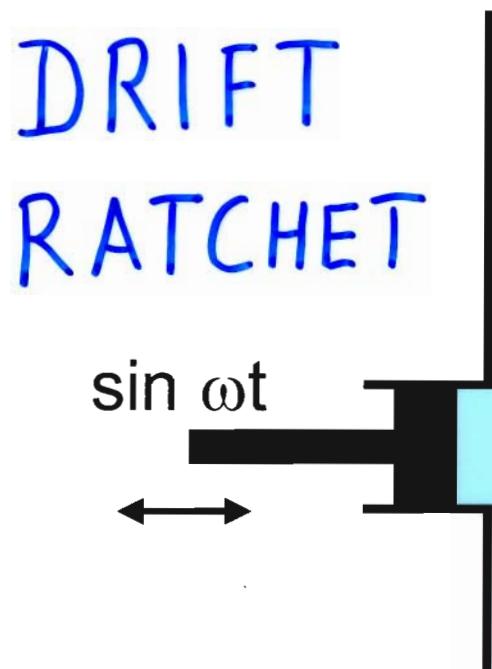
- Noise & AC-Input → **DC-Ouput**
- Non-equilibrium Noise → **Directed Transport**
- **Current reversals**
- **Applications:**
 - Novel pumps and traps for charged or neutral particles
 - Brownian diodes & transistors

Ask not what physics can do
for biology, ask what biology
can do for physics

REVIEWS OF MODERN PHYSICS, VOLUME 81, JANUARY–MARCH 2009

Artificial Brownian motors: Controlling transport on the nanoscale
P.H. and F. Marchesoni

Micro Pump based on Macroporous Silicon



C.H. Kettner, P. Reimann, P. H., F. Müller, PHYS. REV. E 61: 312 (00)



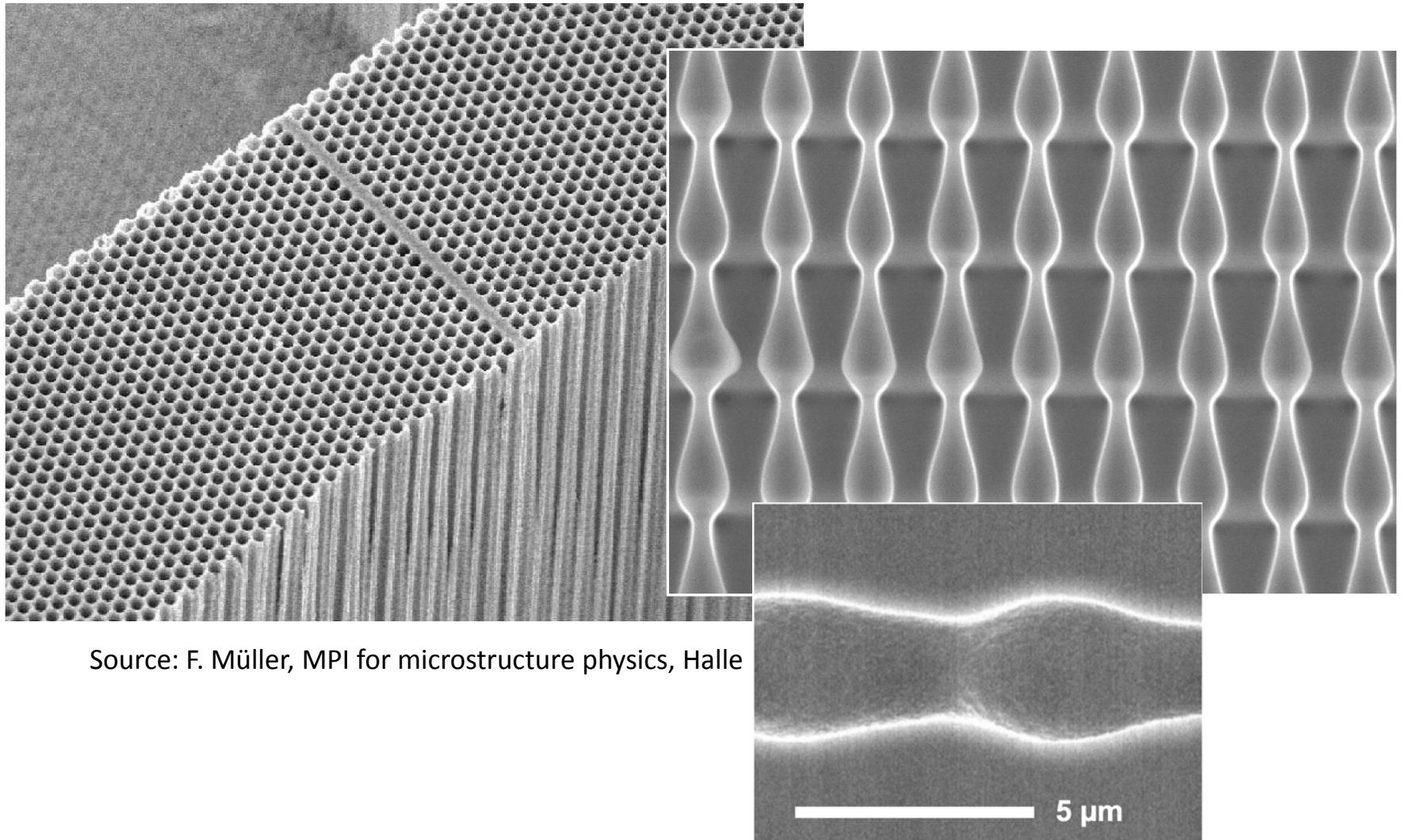
F. Müller, A. Birner, U. Gösele
MPI of Microstructure Physics, Halle/Saale, Germany

LANGEVIN EQ. FOR BROWNIAN PARTICLES

$$\vec{\dot{x}}(t) = \vec{v}(\vec{x}(t), t) + \sqrt{2D_{th}} \vec{\xi}(t)$$

$\frac{k_B T}{\eta}$

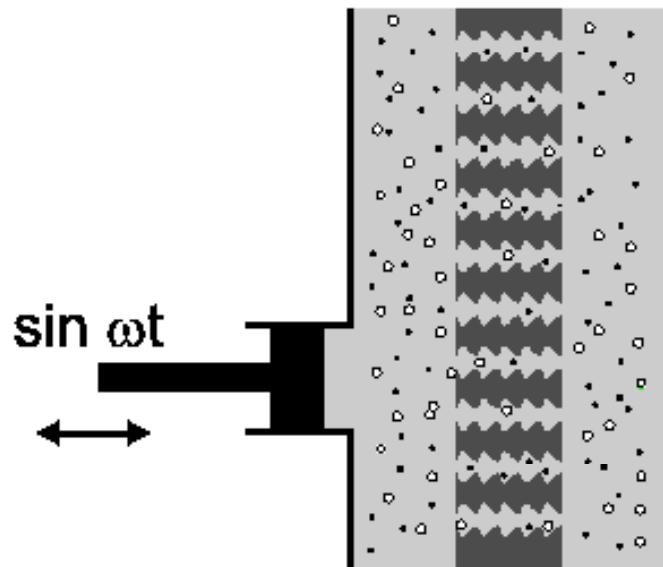
Drift Ratchet - Device



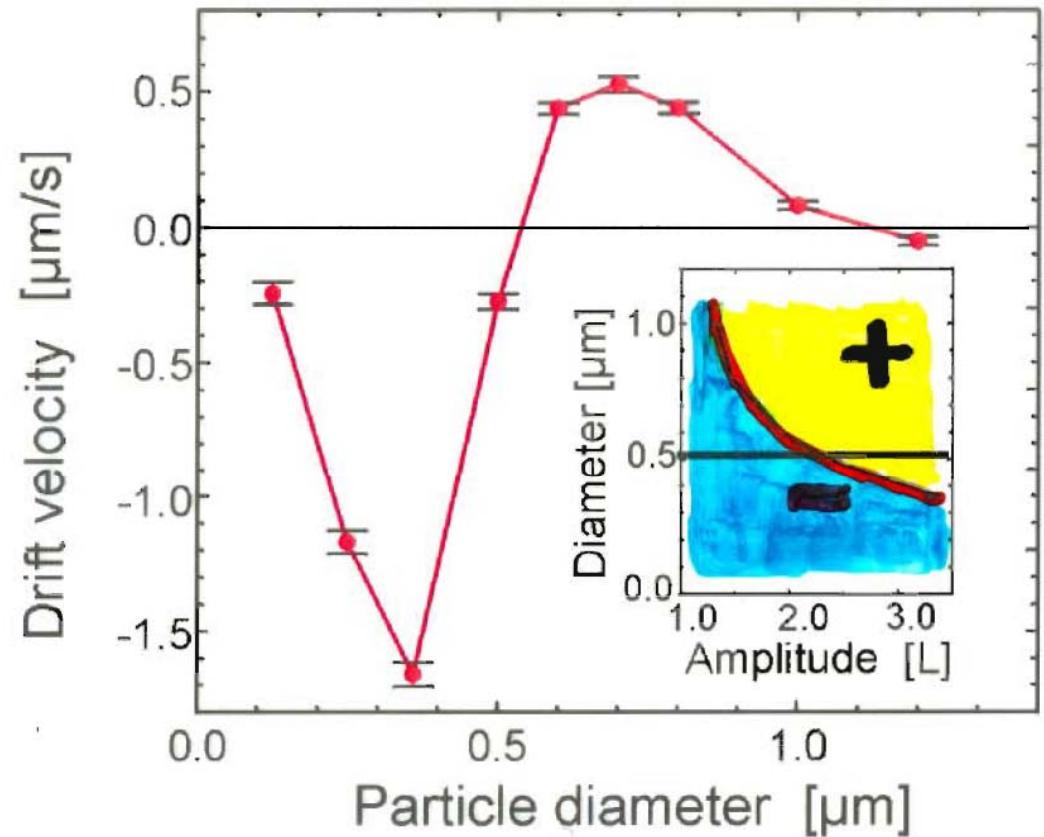
Drift Ratchet - Theory

C. Kettner, P. Reimann, P. H., F. Müller, Phys. Rev. E **61**, 312 (2000)

Setup

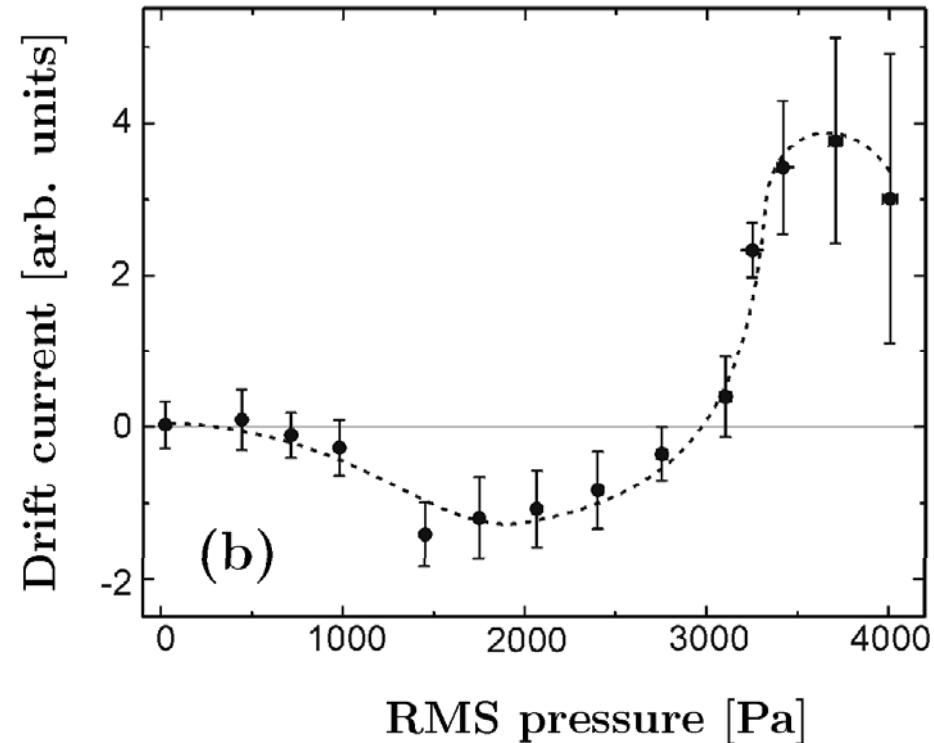
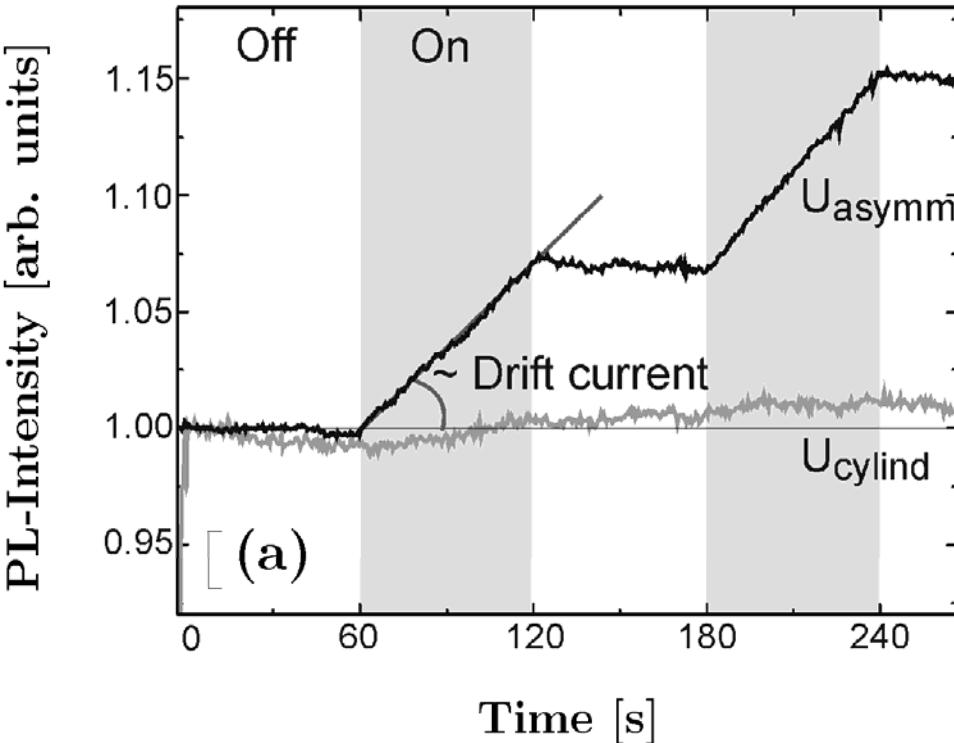


Particle Separation



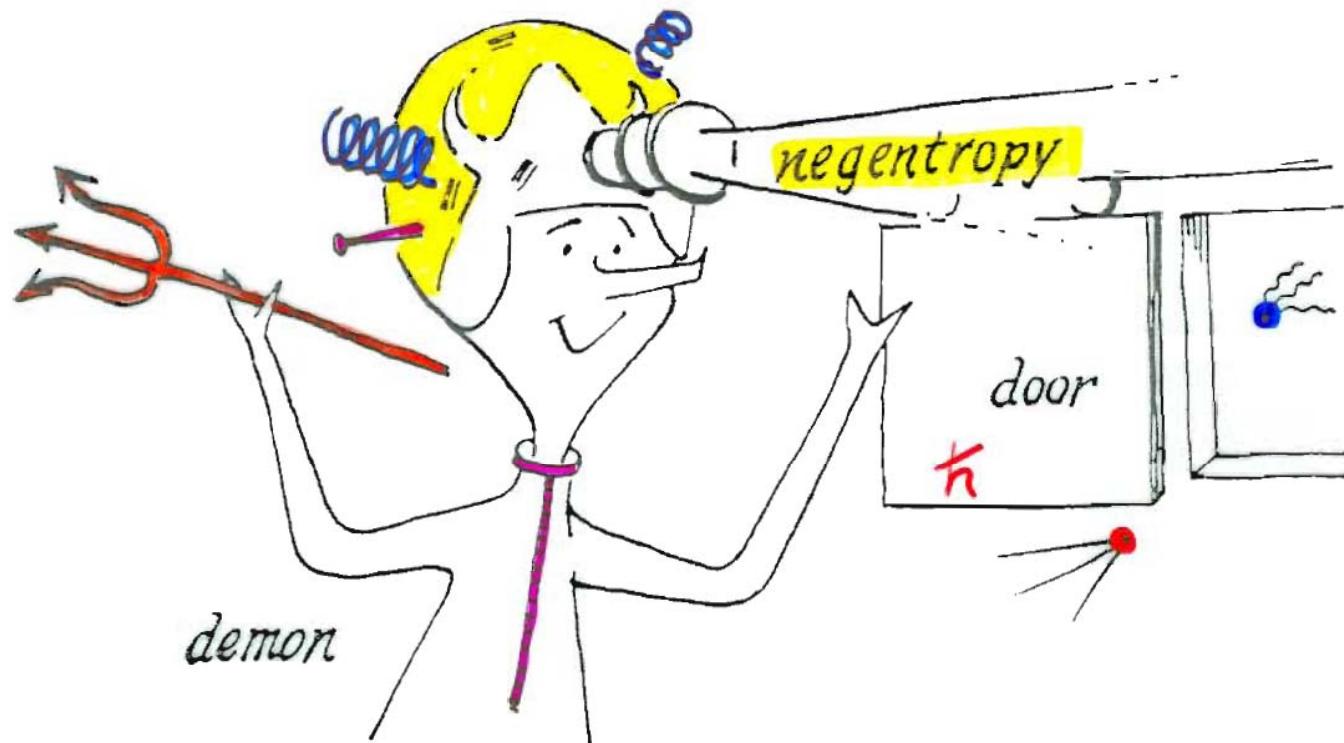
Drift Ratchet – Experiment

S. Matthias, F. Müller, Nature **424**, 53 (2003)



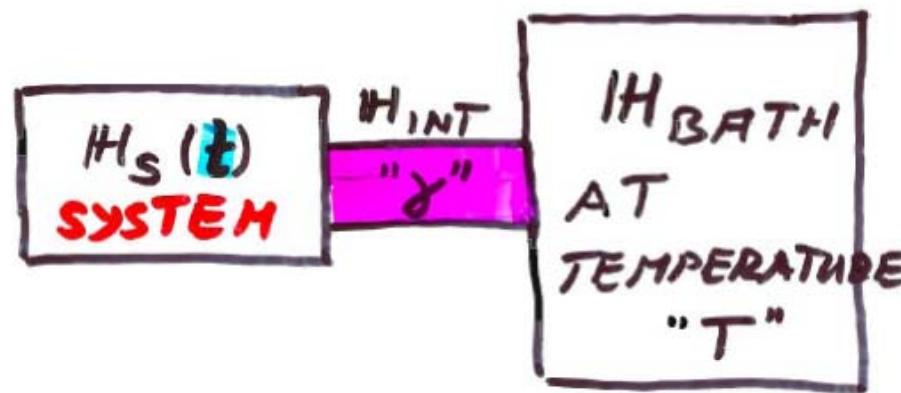
Quantum Demon ?

A measurement → Increase information → Reduction of entropy



Source: H.S. Leff, *Maxwell's Demon* (Adam Hilger, Bristol, 1990)

Quantum Brownian Motors



$$i\hbar \dot{\varrho} = [H_S(t) + H_{INT} + H_{BATH}, \varrho]$$

Hilbert space: $SYSTEM \otimes BATH$

SUPER-
BATH

Quantum-Langevin-equation

$$m\ddot{\mathbf{x}}(t) + \int_{-\infty}^t \gamma(t-t') \dot{\mathbf{x}}(t') dt' + V'(\mathbf{x}; t) = \boldsymbol{\xi}(t)$$

$$\frac{1}{2} \langle \boldsymbol{\xi}(t) \boldsymbol{\xi}(s) + \boldsymbol{\xi}(s) \boldsymbol{\xi}(t) \rangle_{\text{bath}} = \\ \cdot \frac{m}{\pi} \int_0^\infty \text{Re} \hat{\gamma}(-i\omega + 0^+) \hbar\omega \coth \left(\frac{\hbar\omega}{2k_B T} \right) \cos [\omega(t-s)] d\omega$$

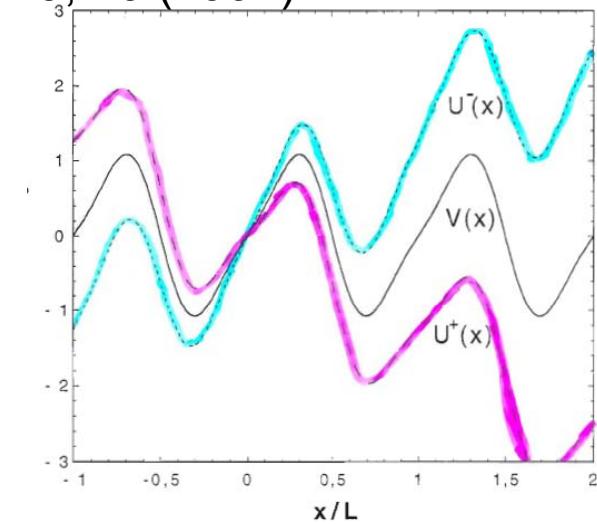
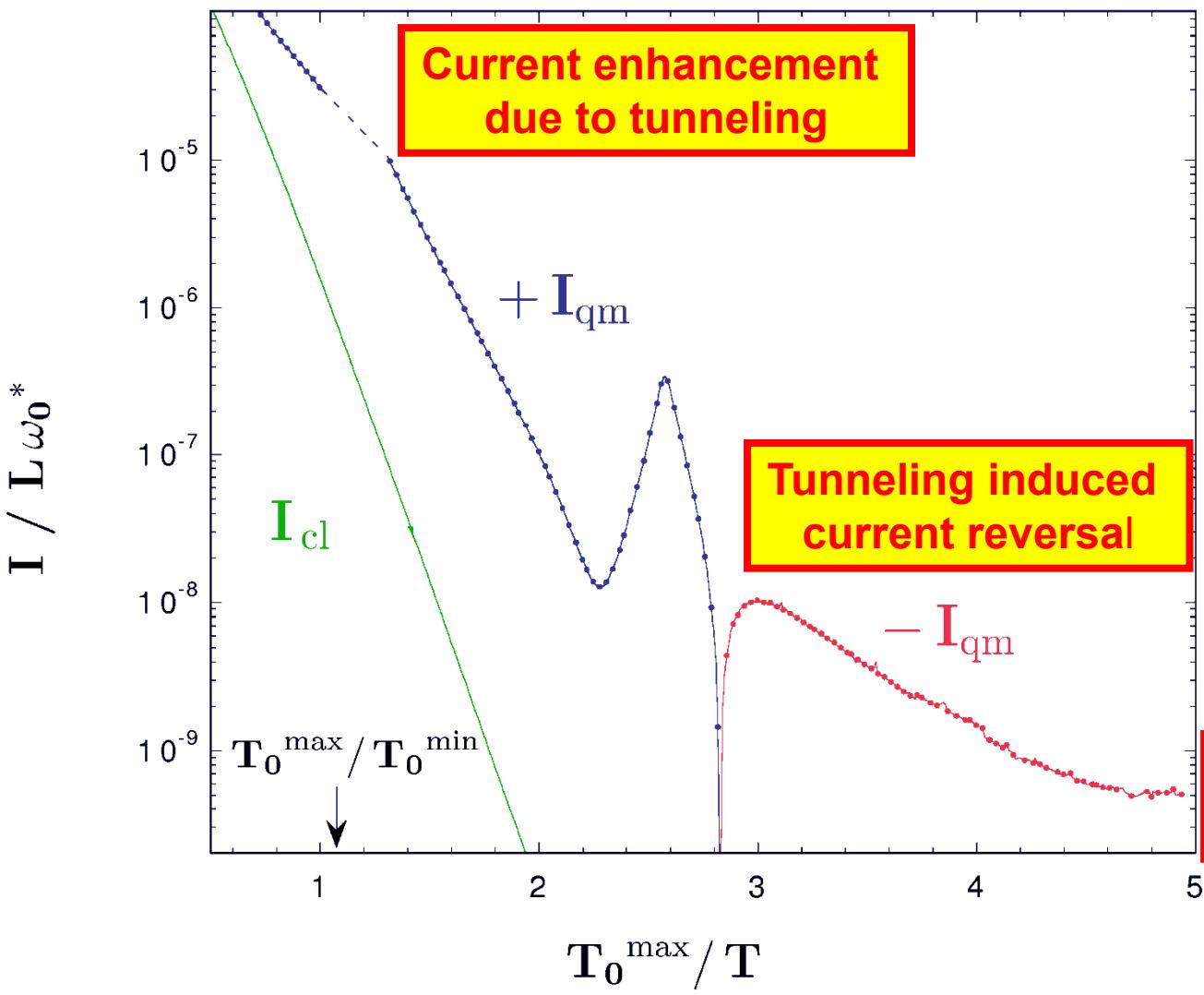
And:

$$[\boldsymbol{\xi}(t), \boldsymbol{\xi}(s)] = -i\hbar \dots \neq 0$$



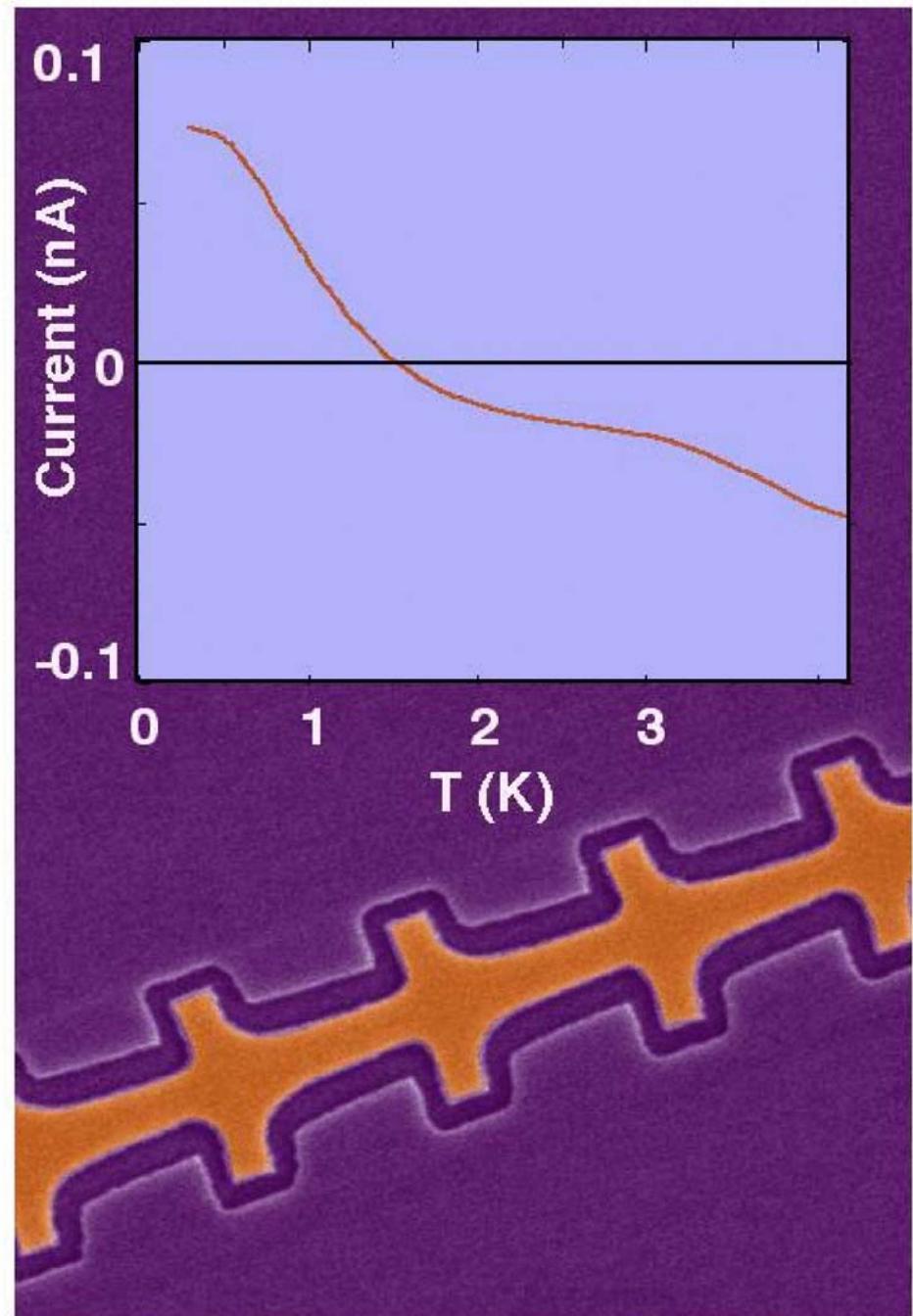
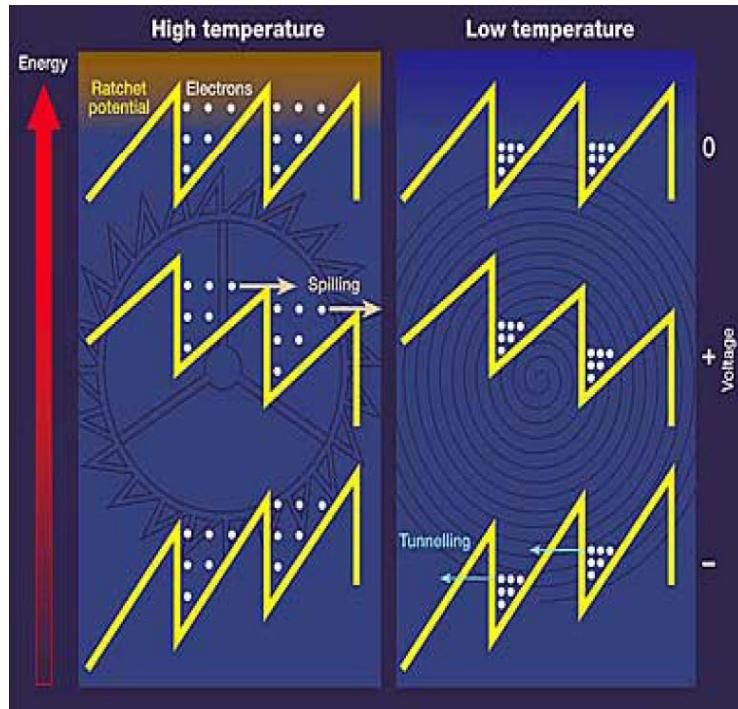
Rocking Ratchet - Theory

P. Reimann, M. Grifoni, P. H., Phys. Rev. Lett. **79**, 10 (1997)



Rocking QM Ratchet – Experiment

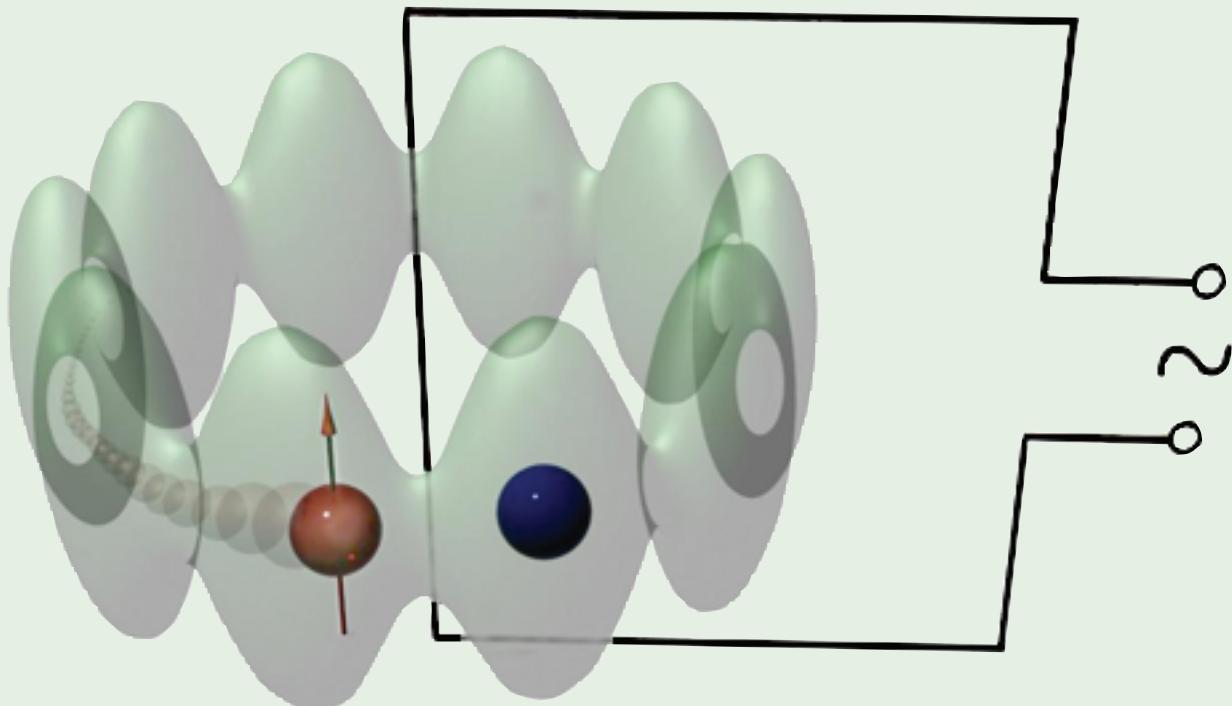
H. Linke, *et al.*,
SCIENCE **286**, 2314 (1999)



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Generalizations of Brownian Motion

Brownian motion: Generalized Langevin-equation

Hamiltonian: $H_{\text{tot}} = H_{\text{sys}} + H_{\text{env}} + H_{\text{ww}}$

→ $m\ddot{\mathbf{x}}(t) + \int_{-\infty}^t \gamma(t-t') \dot{\mathbf{x}}(t') dt' + V'(\mathbf{x}; t) = \xi(t)$

Asymptotically normal, anomalously fast, or anomalously slow
– via fractional Brownian motion –

$$\int_0^\infty \gamma(t) dt = \begin{cases} \text{const} & \Rightarrow \text{normal} \\ 0 & \Rightarrow \text{superfast} \\ \infty & \Rightarrow \text{superslow} \end{cases}$$

Connection to the fractional Fokker-Planck-equation

Confined Diffusion of Brownian particles: Entropic versus hydrodynamic interactions

PNAS

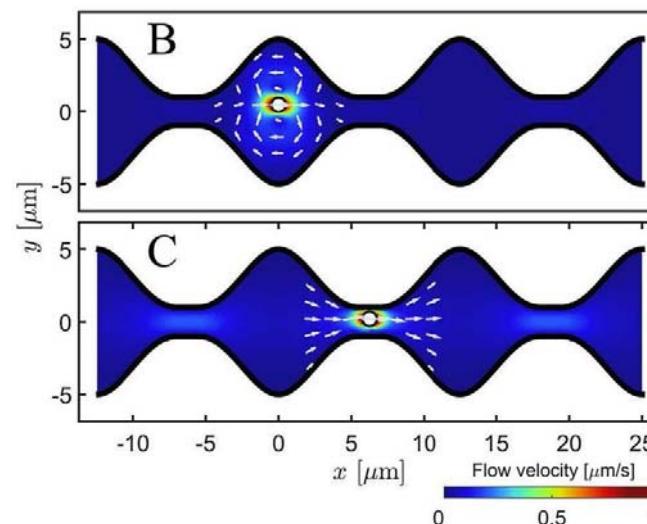
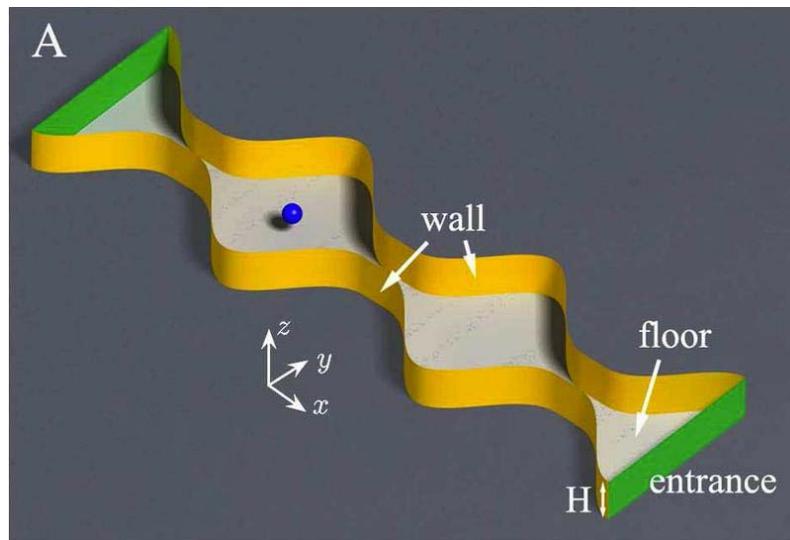
Hydrodynamic and entropic effects on colloidal diffusion in corrugated channels

Xiang Yang^{a,b}, Chang Liu^{a,b}, Yunyun Li^{c,d}, Fabio Marchesoni^{c,e}, Peter Hänggi^{f,g}, and H. P. Zhang^{a,b,h,1}

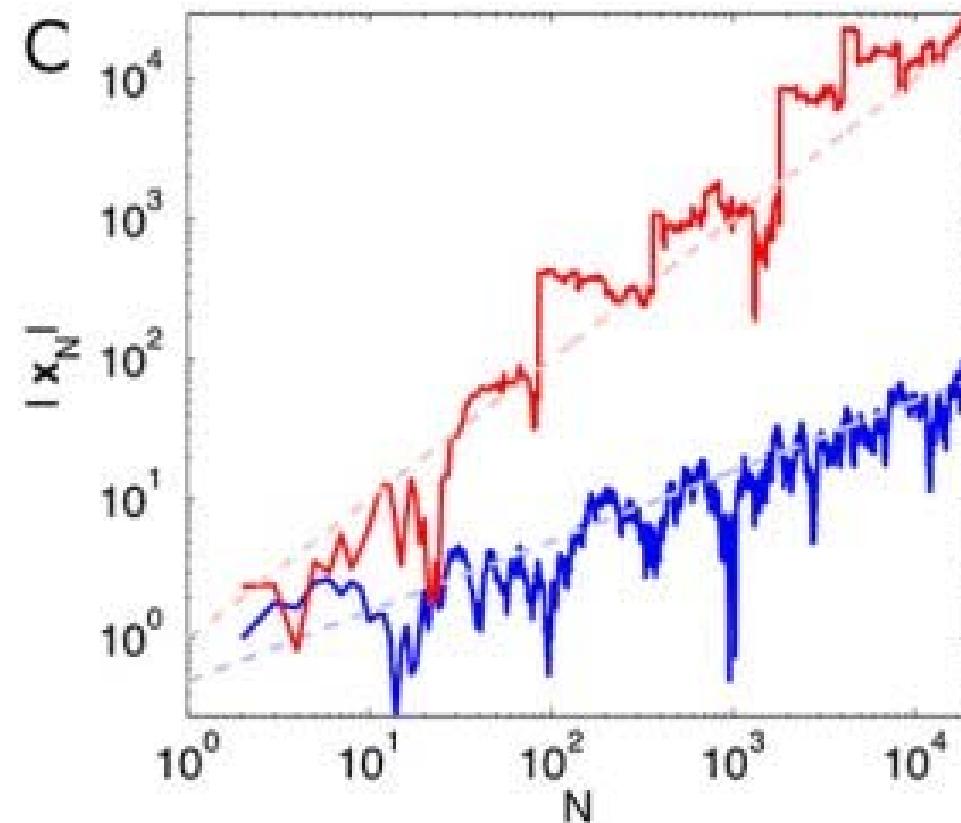
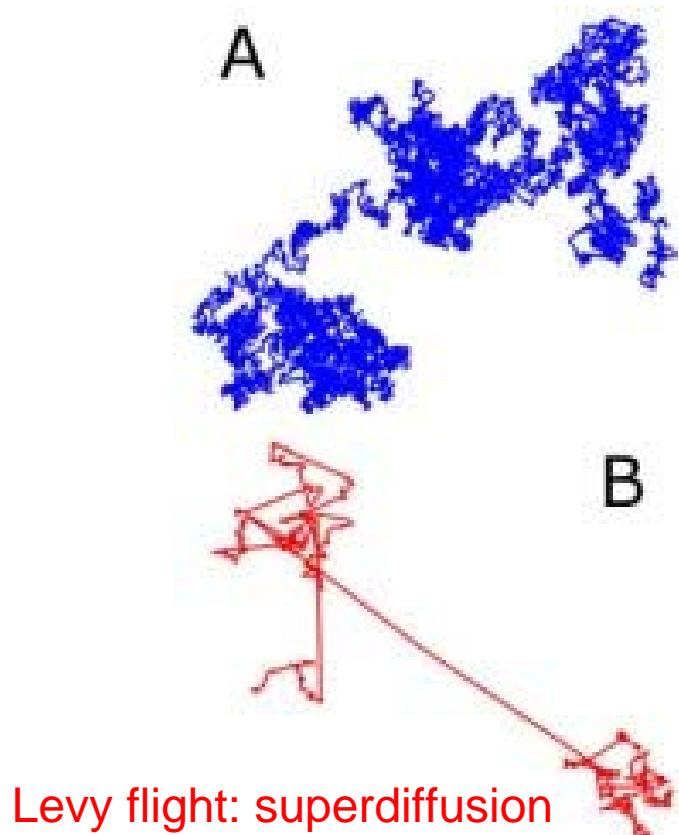
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Edited by David A. Weitz, Harvard University, Cambridge, MA, and approved July 25, 2017 (received for review May 10, 2017)

Proc. Natl. Acad. Sci. 114, 9564–9569 (2017)



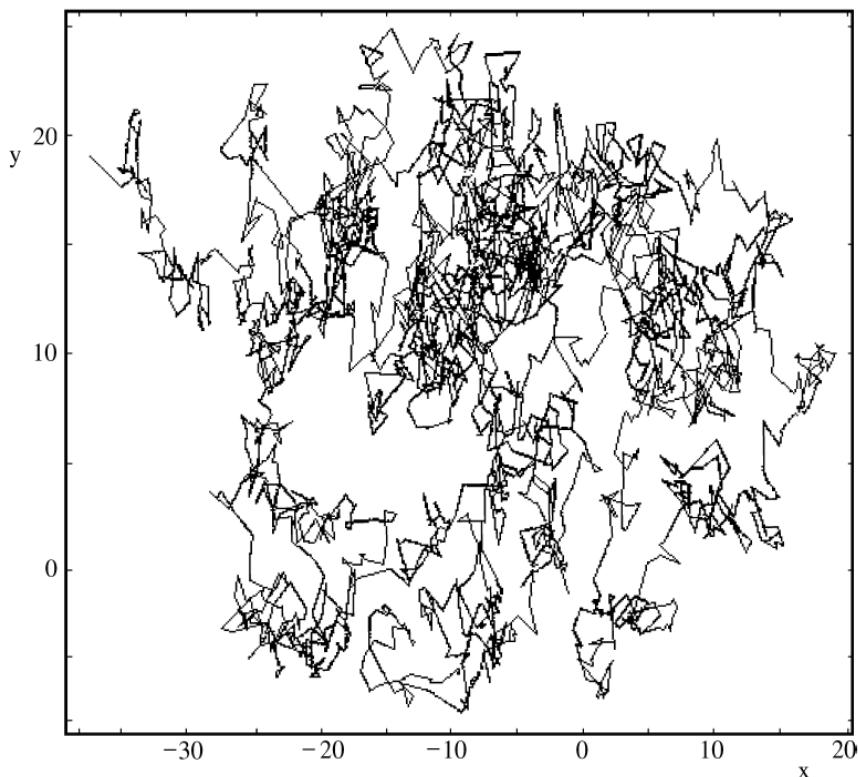
normal Brownian motion



Mean squared displacement

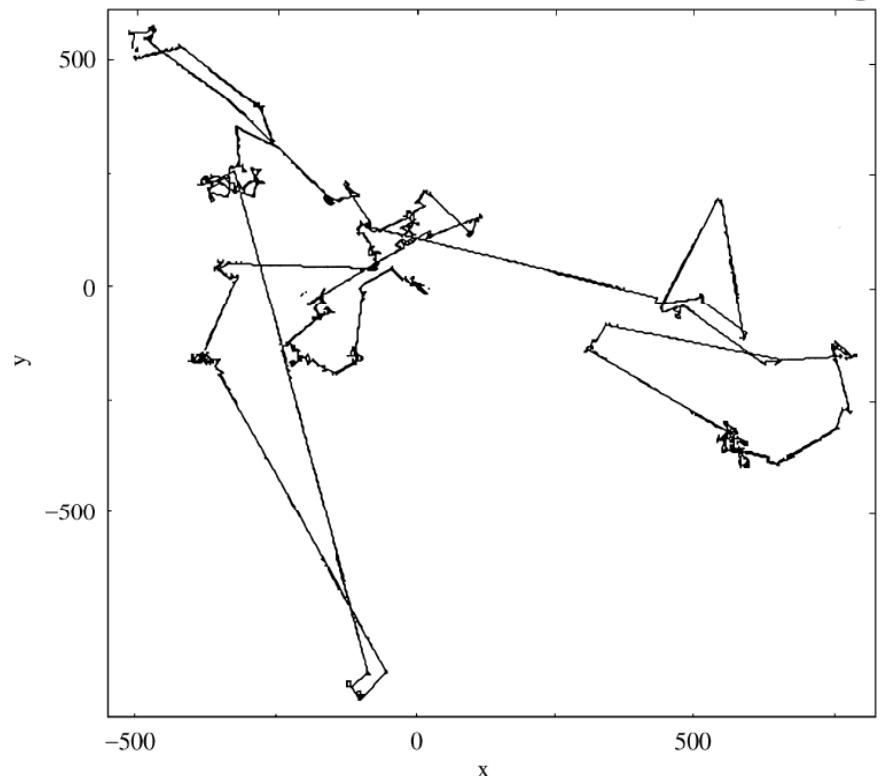
$$\langle x^2(t) \rangle \propto t^\alpha$$

Brownian movement $\alpha = 1$



Source: Physica A **282**, 13 (2000)

Lévy-Brownian movement $\alpha = \frac{4}{3}$



Source: Physica A **282**, 13 (2000)

Fractional Fokker-Planck equation

Subdiffusion ($\alpha < 1$):

$$\frac{\partial P(x,t)}{\partial t} = \left[\frac{\partial}{\partial x} \frac{V'(x,t)}{\gamma_\alpha} + K_\alpha \frac{\partial^2}{\partial x^2} \right] {}_0D_t^{1-\alpha} P(x,t)$$

Fat tails in the distribution of the residence times

Riemann-Liouville Operator



Superdiffusion ($\alpha > 1$):

$$\frac{\partial P(x,t)}{\partial t} = \left[\frac{\partial}{\partial x} \frac{V'(x,t)}{\gamma_\alpha} + K_\alpha \frac{\partial^{2/\alpha}}{\partial |x|^{2/\alpha}} \right] P(x,t)$$

Fat tails in the distribution of the jump lengths

Riesz-derivative

Fractional Fokker-Planck equation

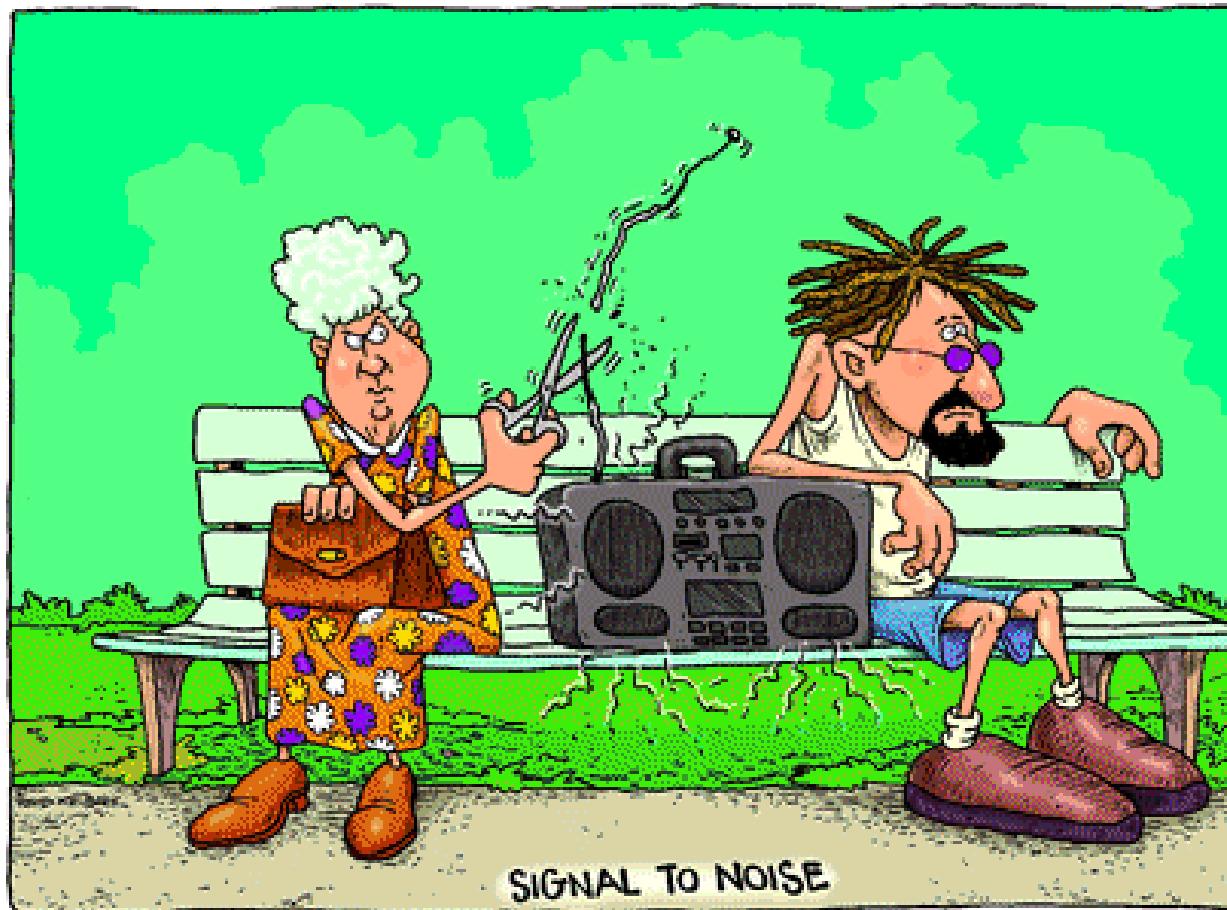
subdiffusive ($\alpha < 1$)

$$\frac{\partial P(x, t)}{\partial t} = \left[\frac{\partial}{\partial x} \frac{V'(x, t)}{\eta_\alpha} + K_\alpha \frac{\partial^2}{\partial x^2} \right] {}_0 D_t^{1-\alpha} P(x, t)$$

Riemann-Liouville Operator

$${}_0 D_t^{1-\alpha} f(t) = \frac{1}{\Gamma(\alpha)} \frac{d}{dt} \int_0^t \frac{f(t')}{(t-t')^{1-\alpha}} dt'$$

Noise – always bad ?



Source: Agilent Technologies

A QUESTION ?

