

Binary Black Holes mergers from Population III stars

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Jacopo Tissino, Giuliano Iorio, Michela Mapelli, M. Celeste Artale *et al.*

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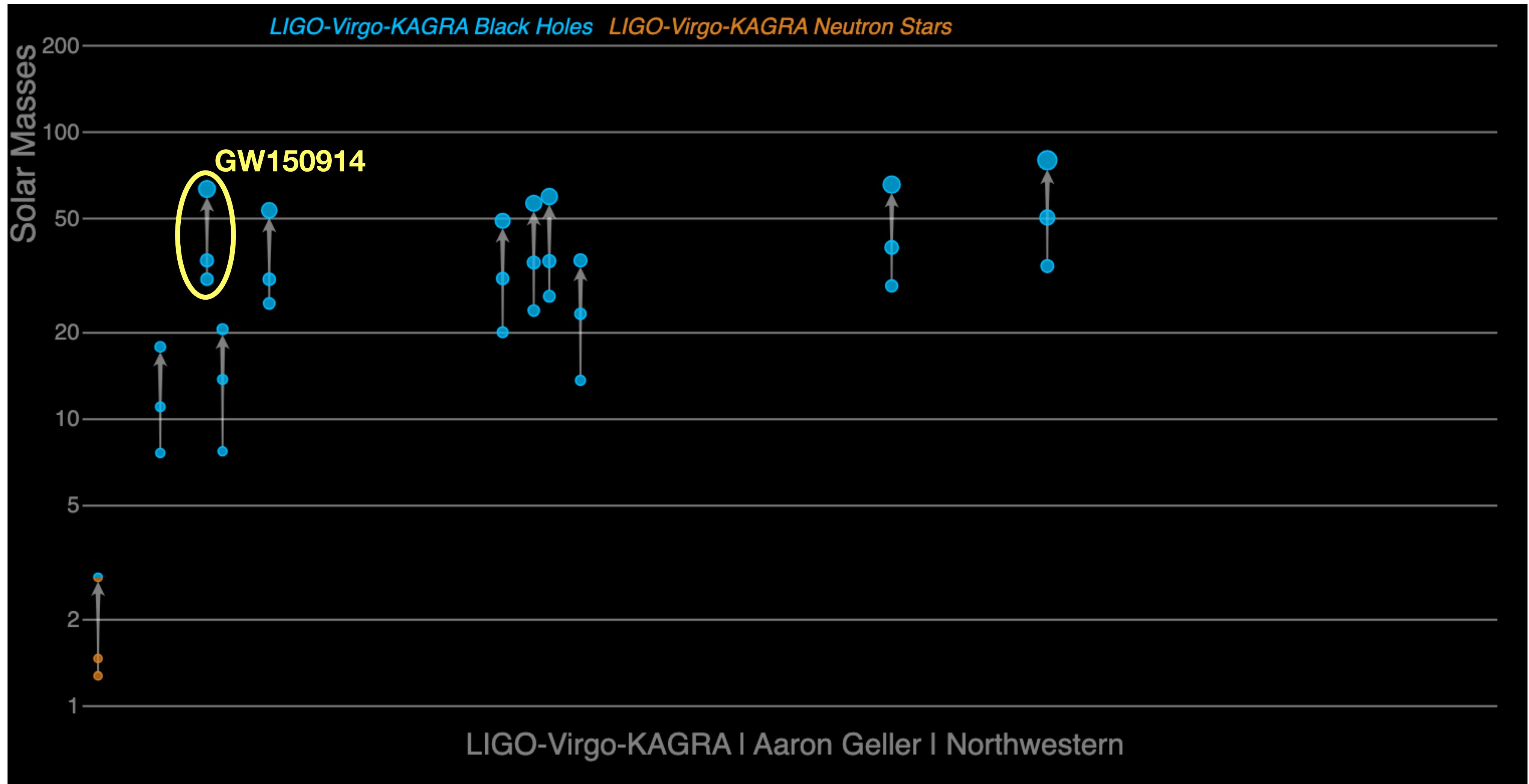
I introduce myself



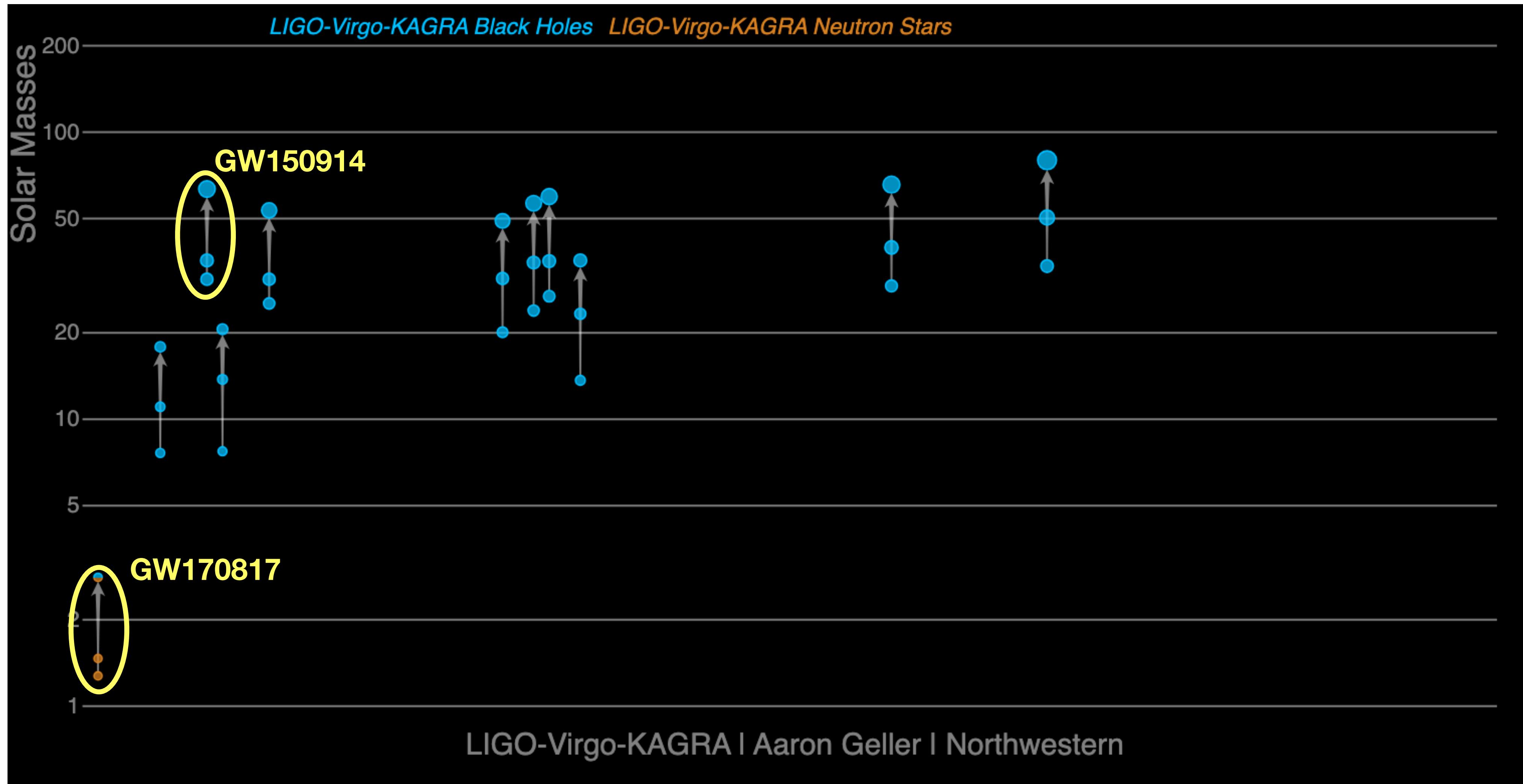
And many others...

Contact me at filippo.santoliquido@gssi.it or visit <https://filippo-santoliquido.github.io/>

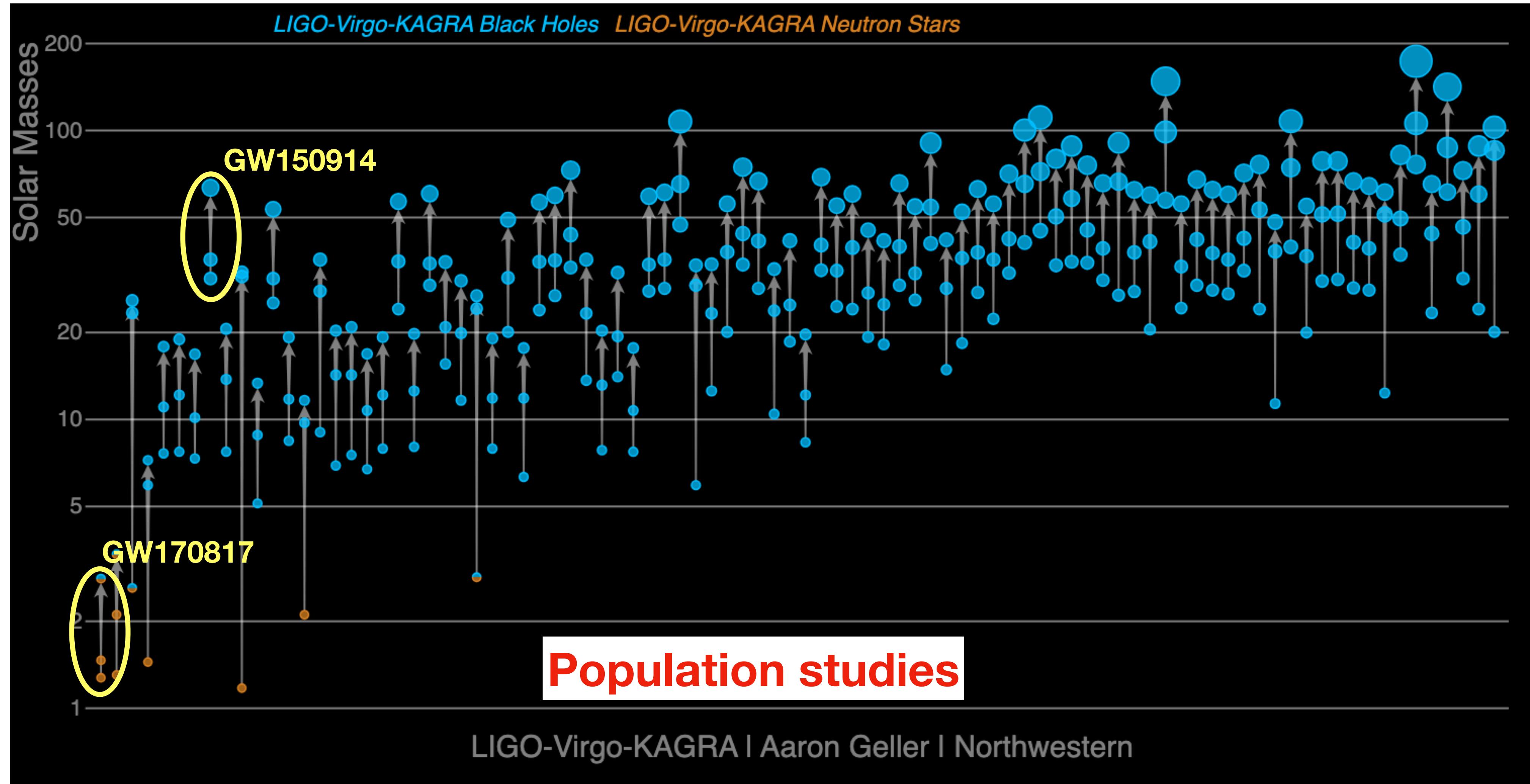
Gravitational Wave Astrophysics



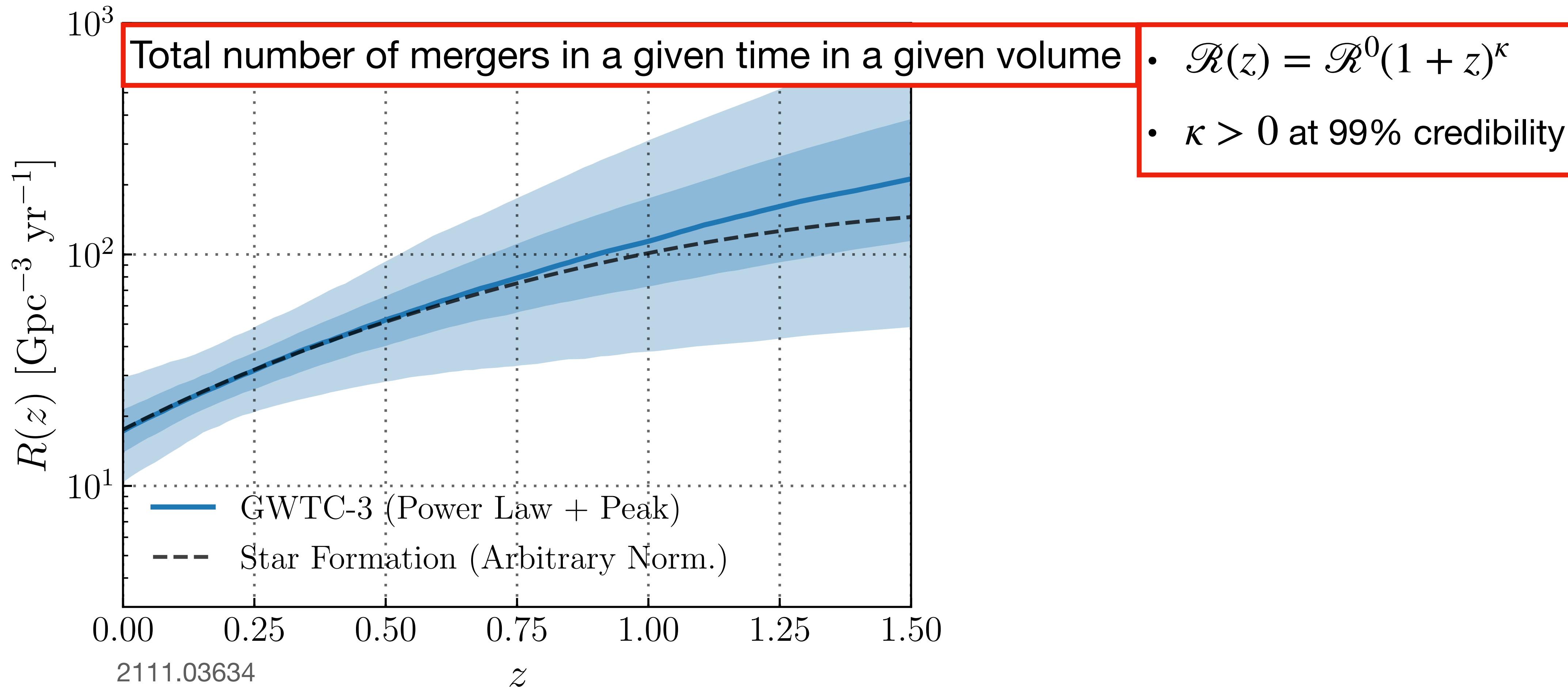
Gravitational Wave Astrophysics



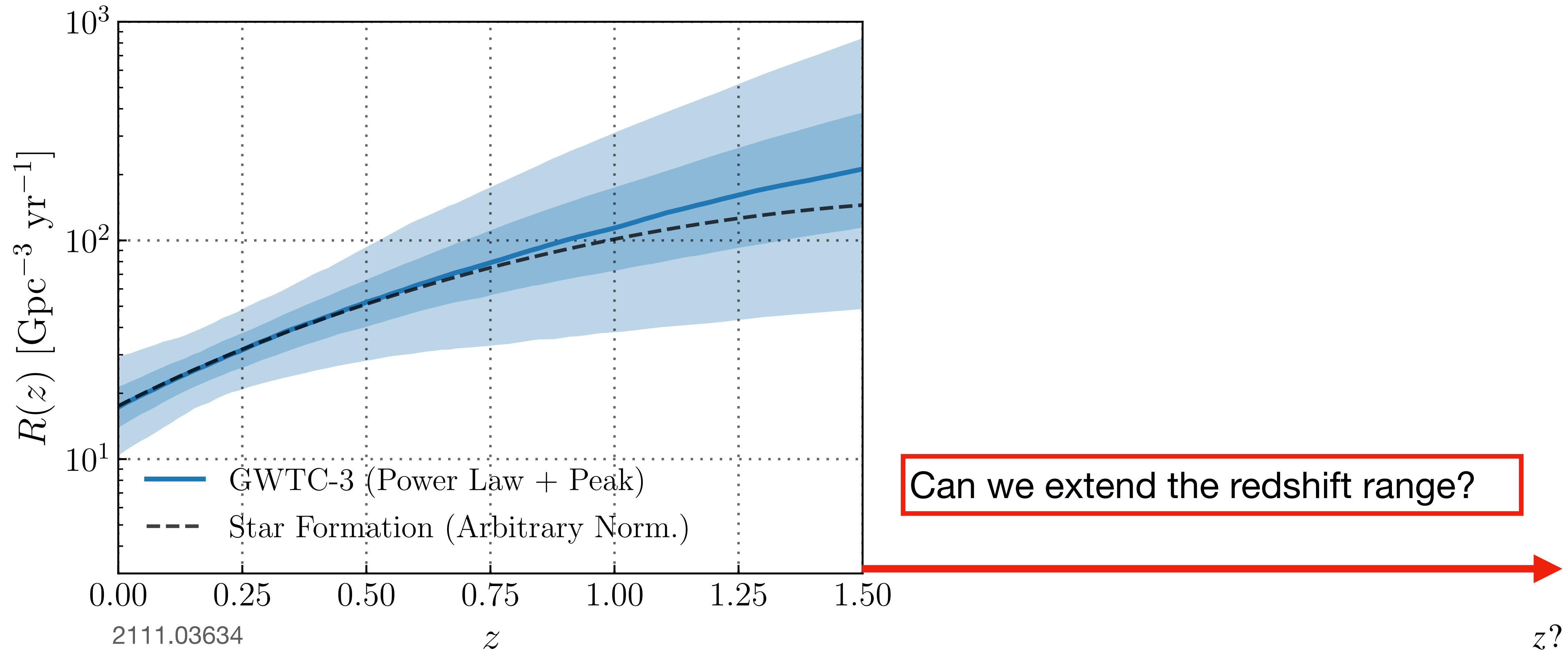
Gravitational Wave Astrophysics



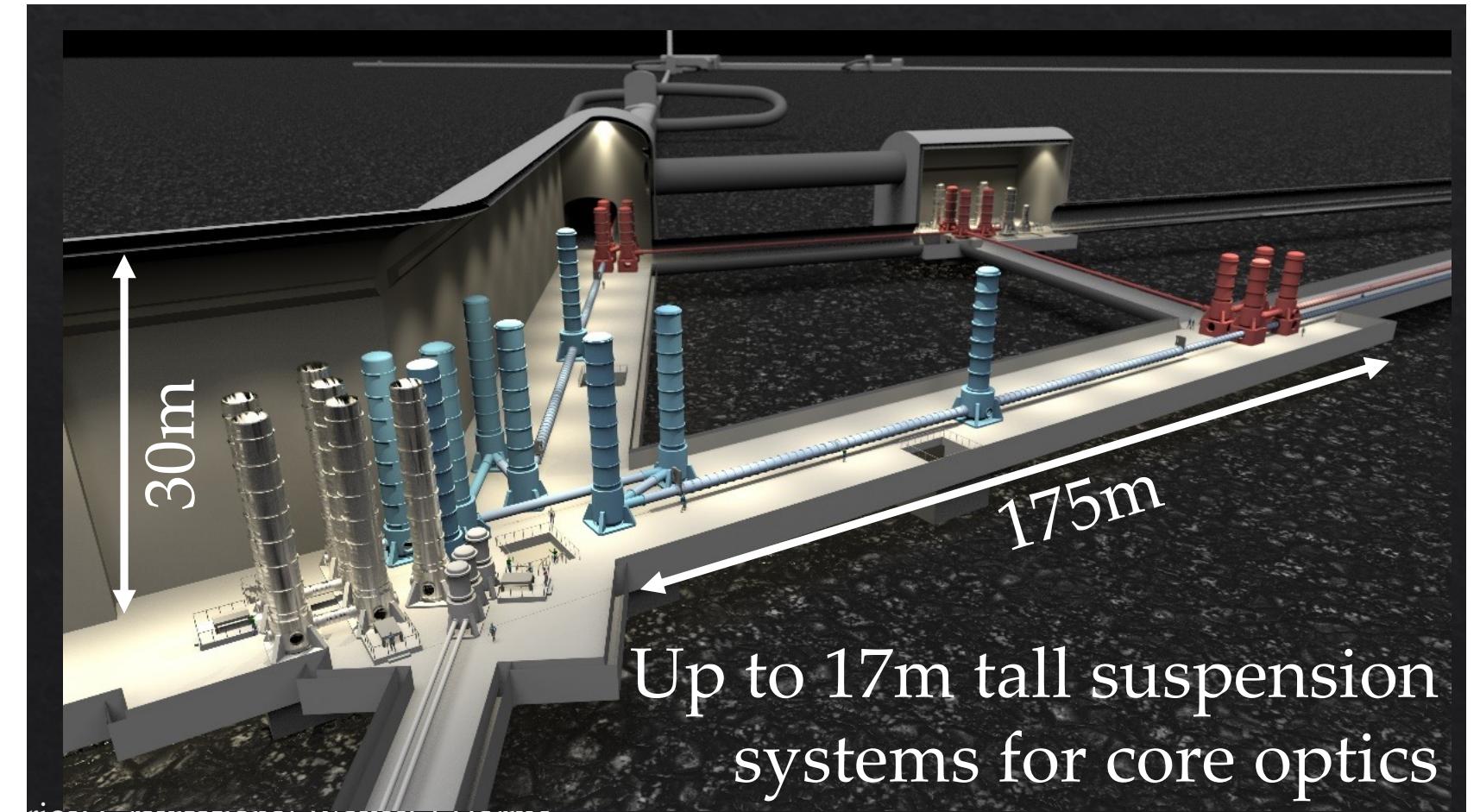
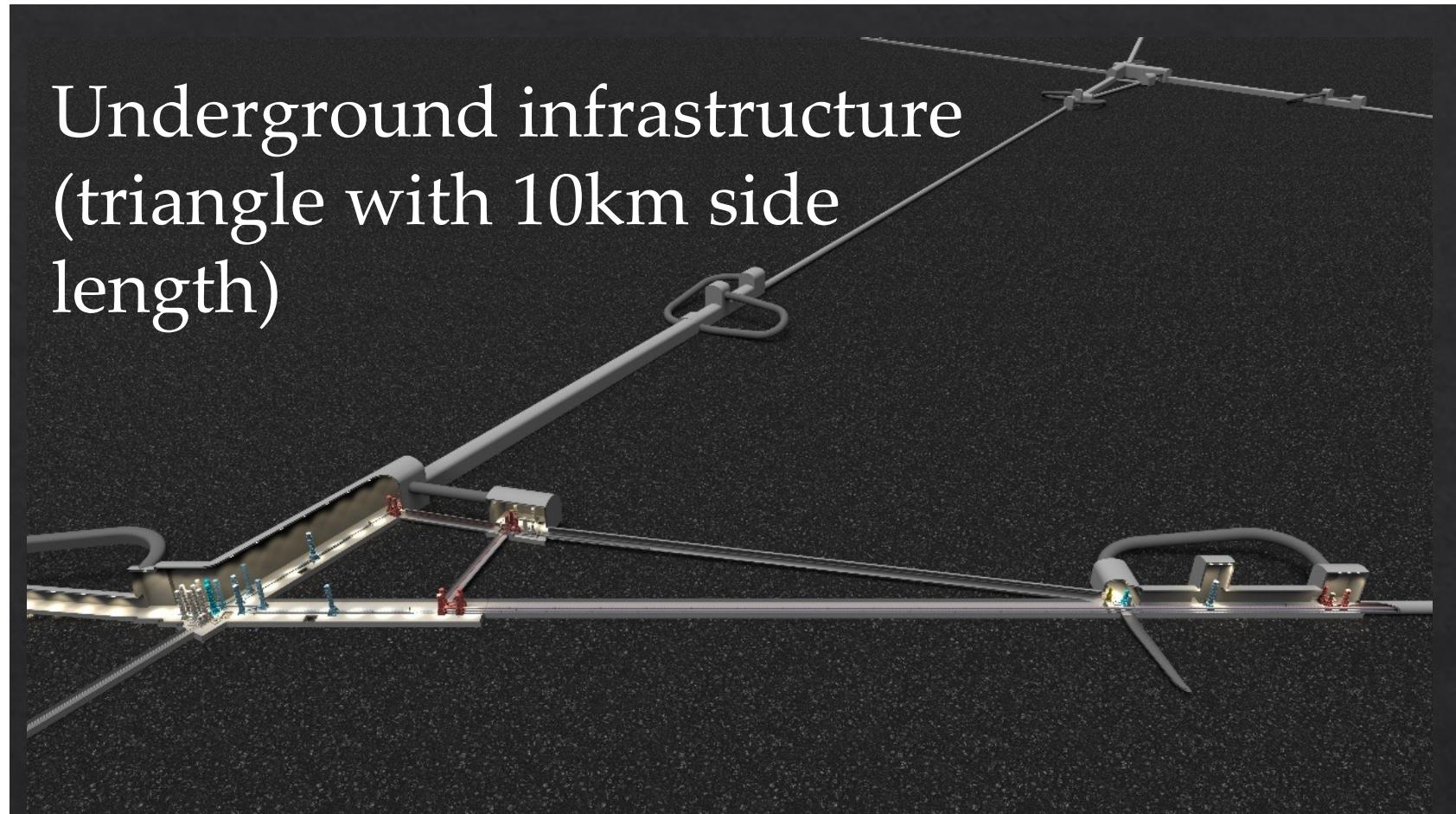
Merger rate density



Merger rate density

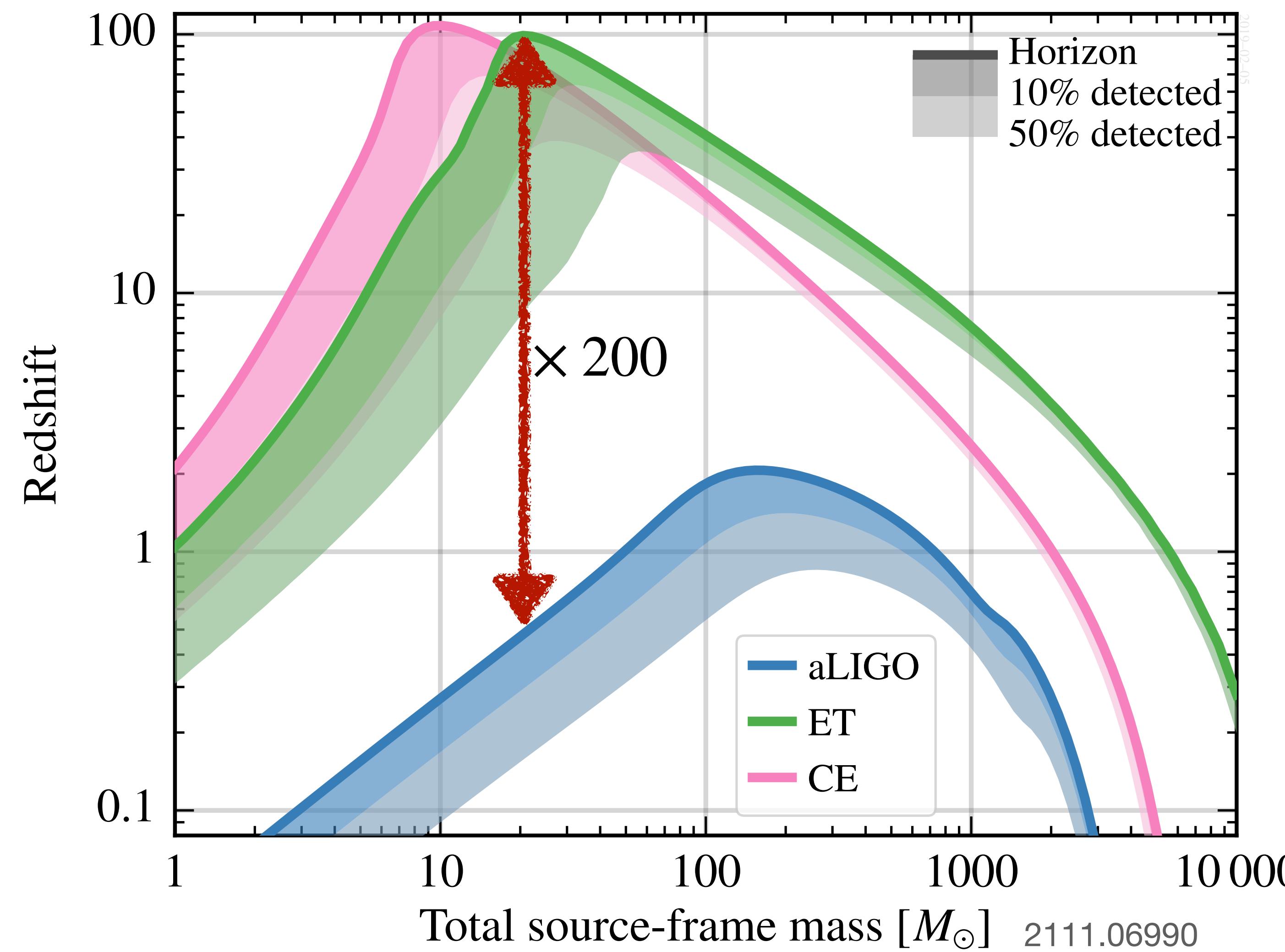


The Einstein Telescope



Credits: Jan Harms

Third-generation detectors



Population III stars

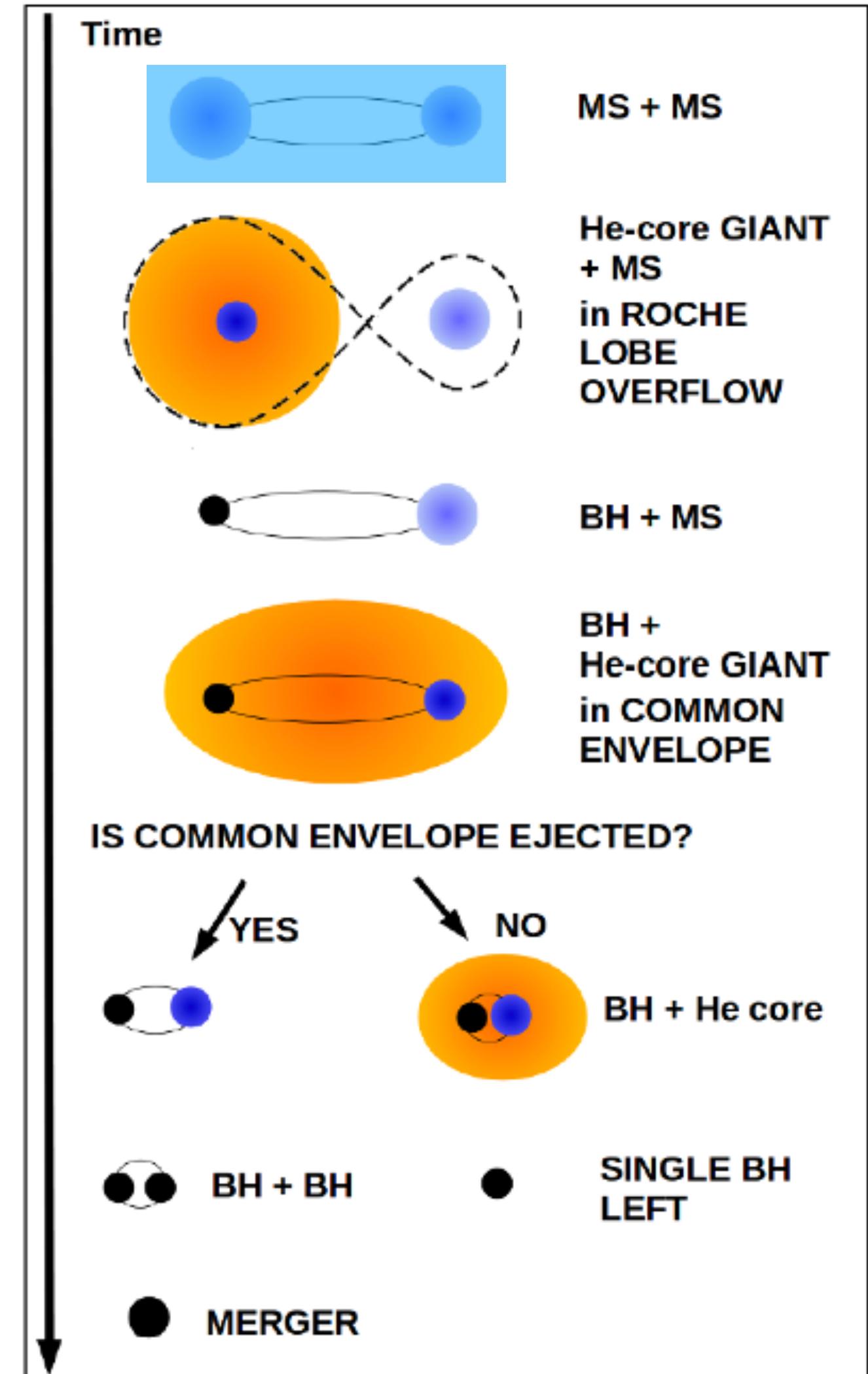
- Population III stars are believed to be the first generation of stars formed at **high redshift** ($z > 20$)
- They are massive and formed from pristine gas (i.e. **zero metallicity**)
- They are still undetected (*traces of their existence with JWST*)
- Modelled through cosmological simulations

Our goals

- Large **parameter space exploration** of Pop. III BBHs
- **merger rate density**
- **evolution of mass spectrum with redshift**
- expected **detection rates** with Einstein Telescope

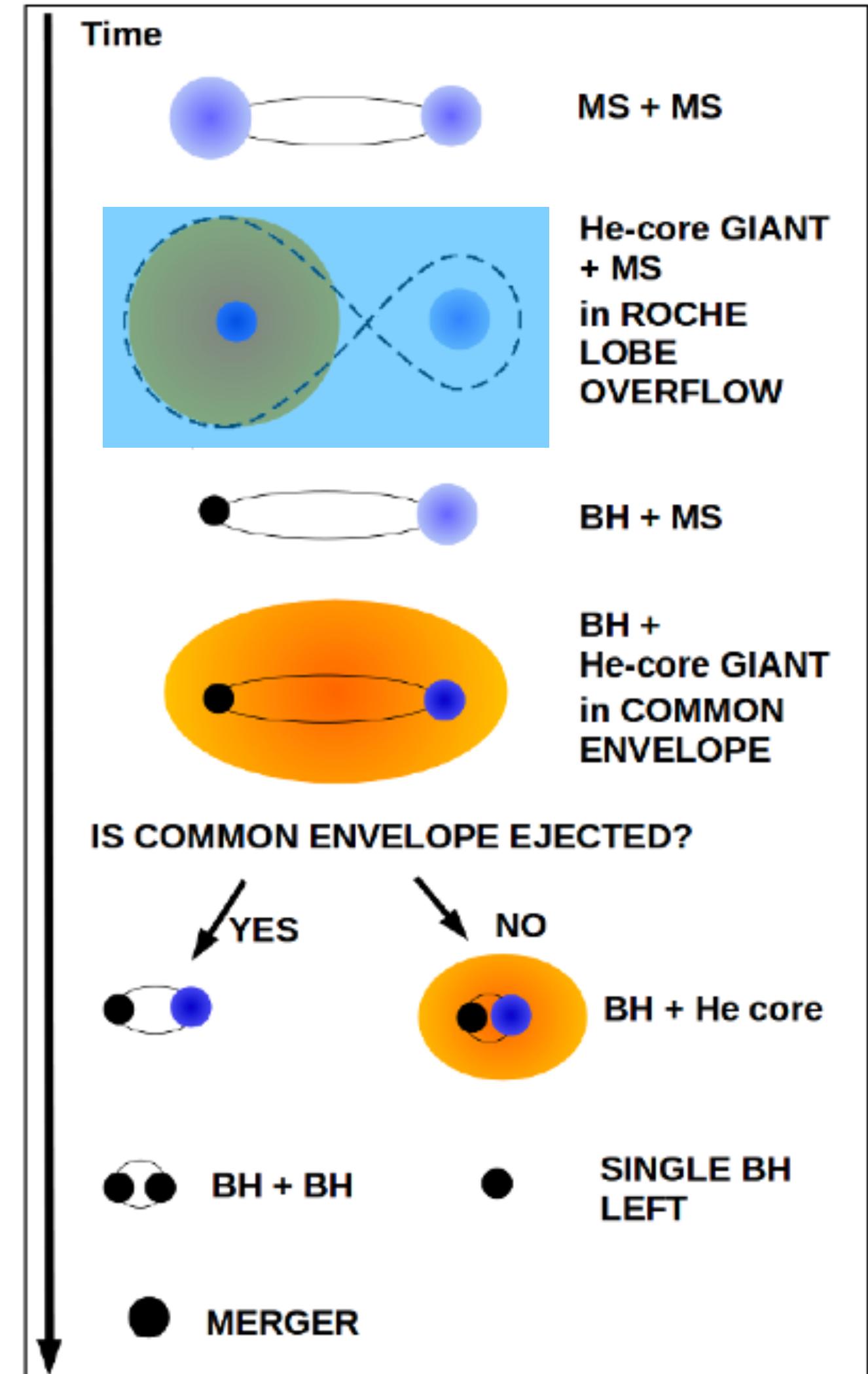
Isolated formation channel

- Two massive stars form in the **same binary system** and evolve together
- *Let's see an example*



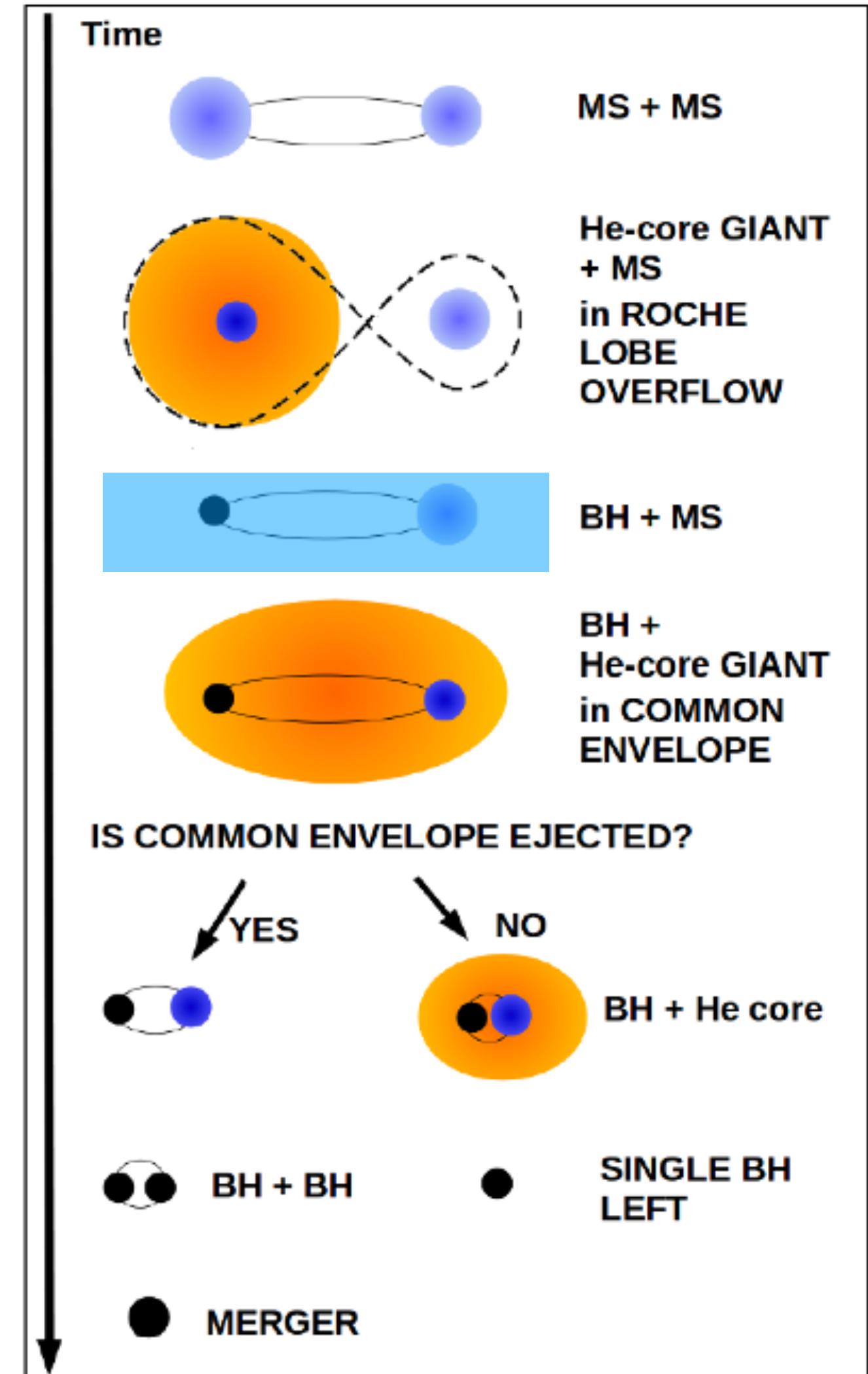
Isolated formation channel

- Following radius expansion, one star can transfer mass to the other in a **stable mass transfer**
- This shrinks the orbit of the binary system



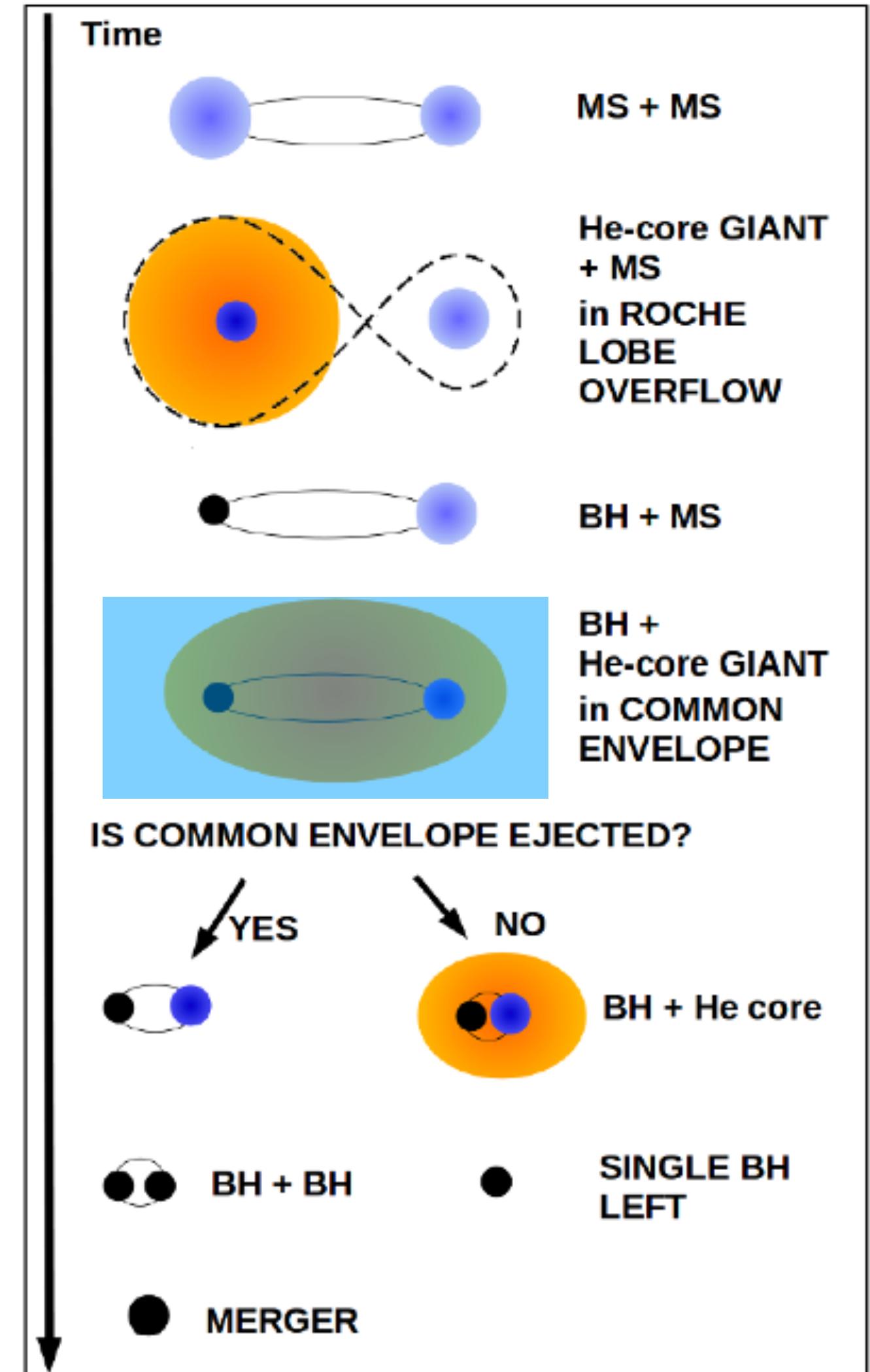
Isolated formation channel

- The first black is formed after core **collapse supernova**



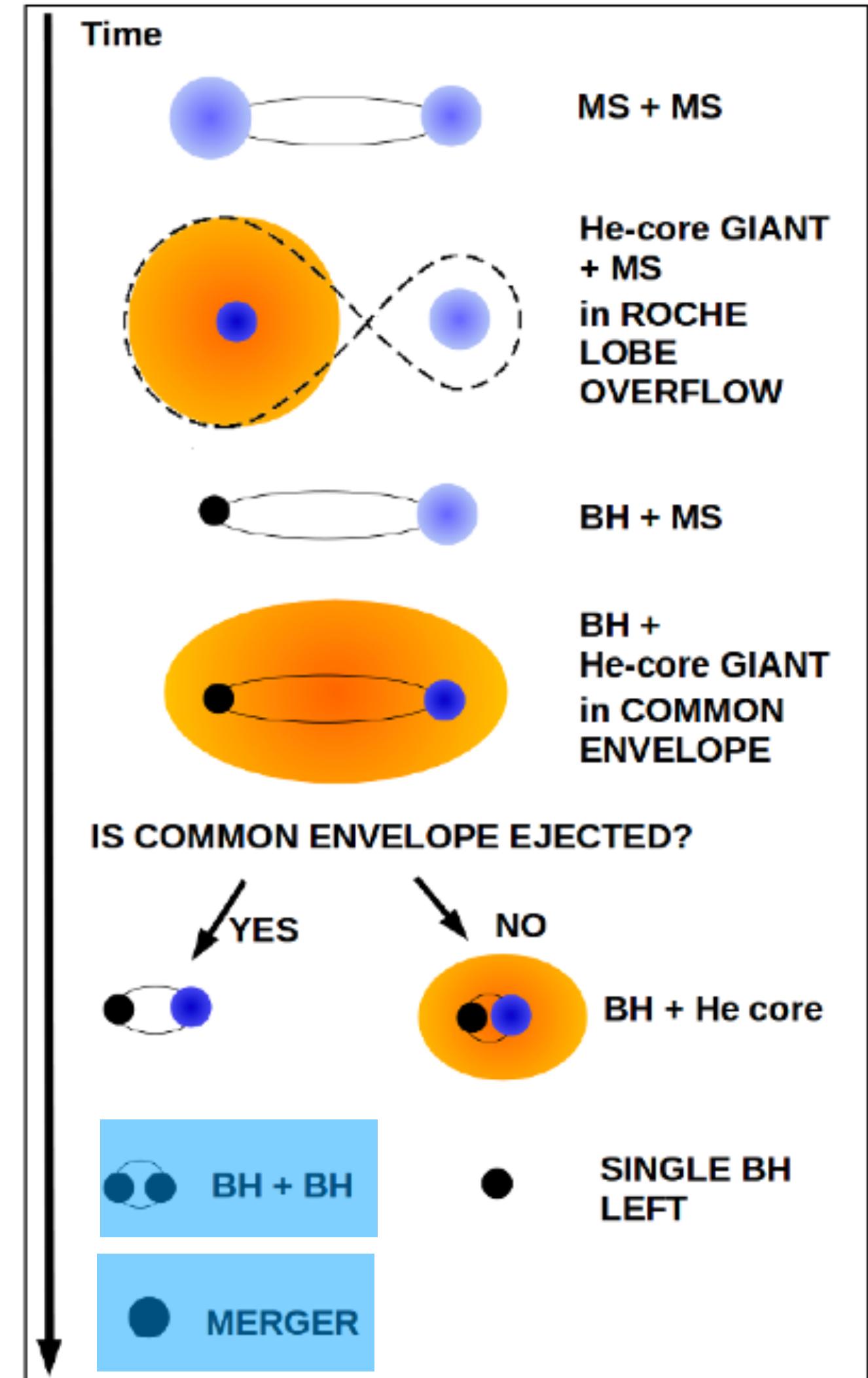
Isolated formation channel

- The other star starts transferring mass through **unstable mass transfer**
- Beginning of the **common envelope phase**
- A **drag force** is exerted between the black hole and the stellar core.
- This phase further shrinks the binary systems



Isolated formation channel

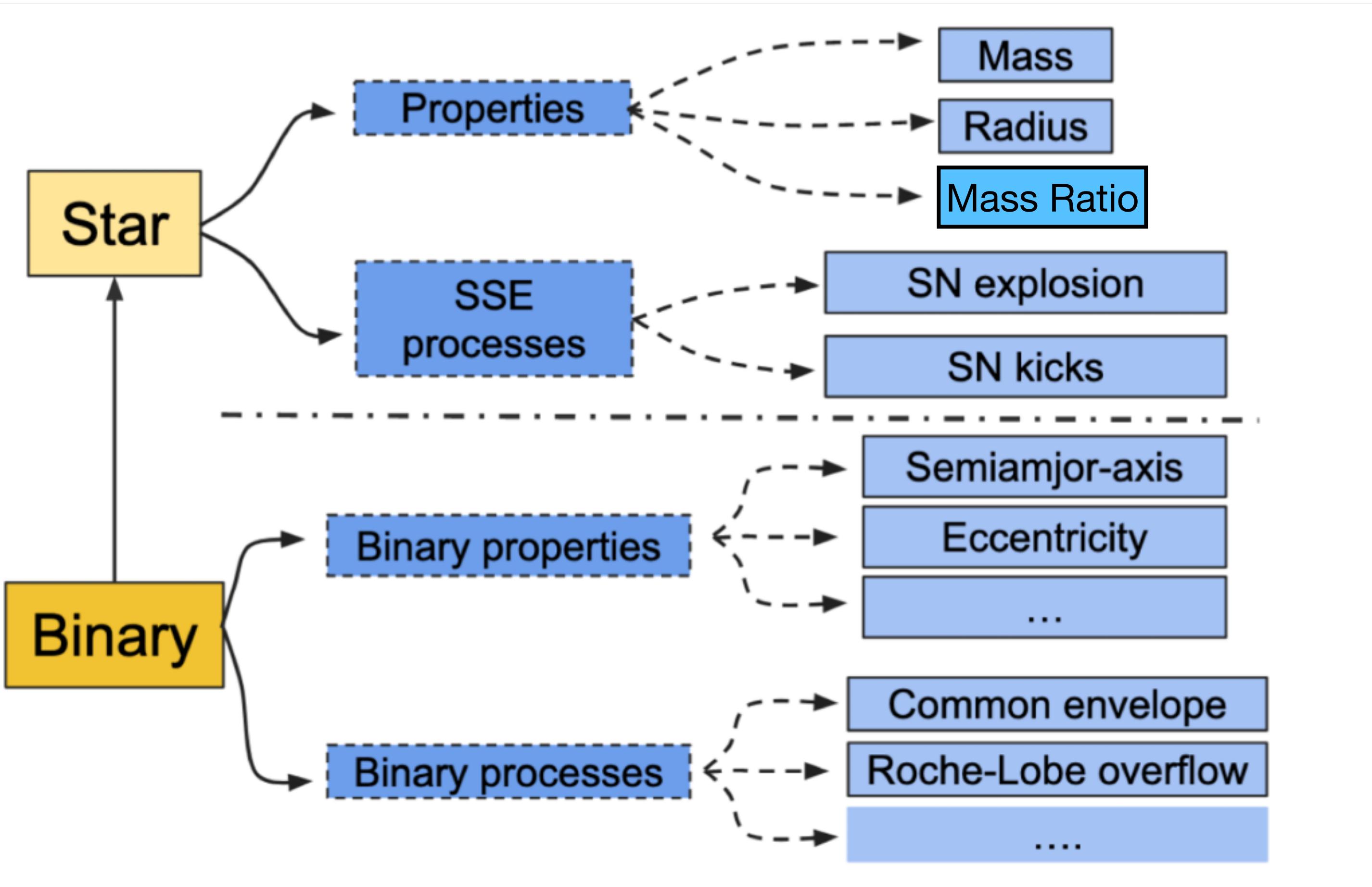
- The second black hole forms
- The binary system starts to loose energy through **gravitational wave radiation emission**
- The **two black holes** might **merge** within an Hubble time



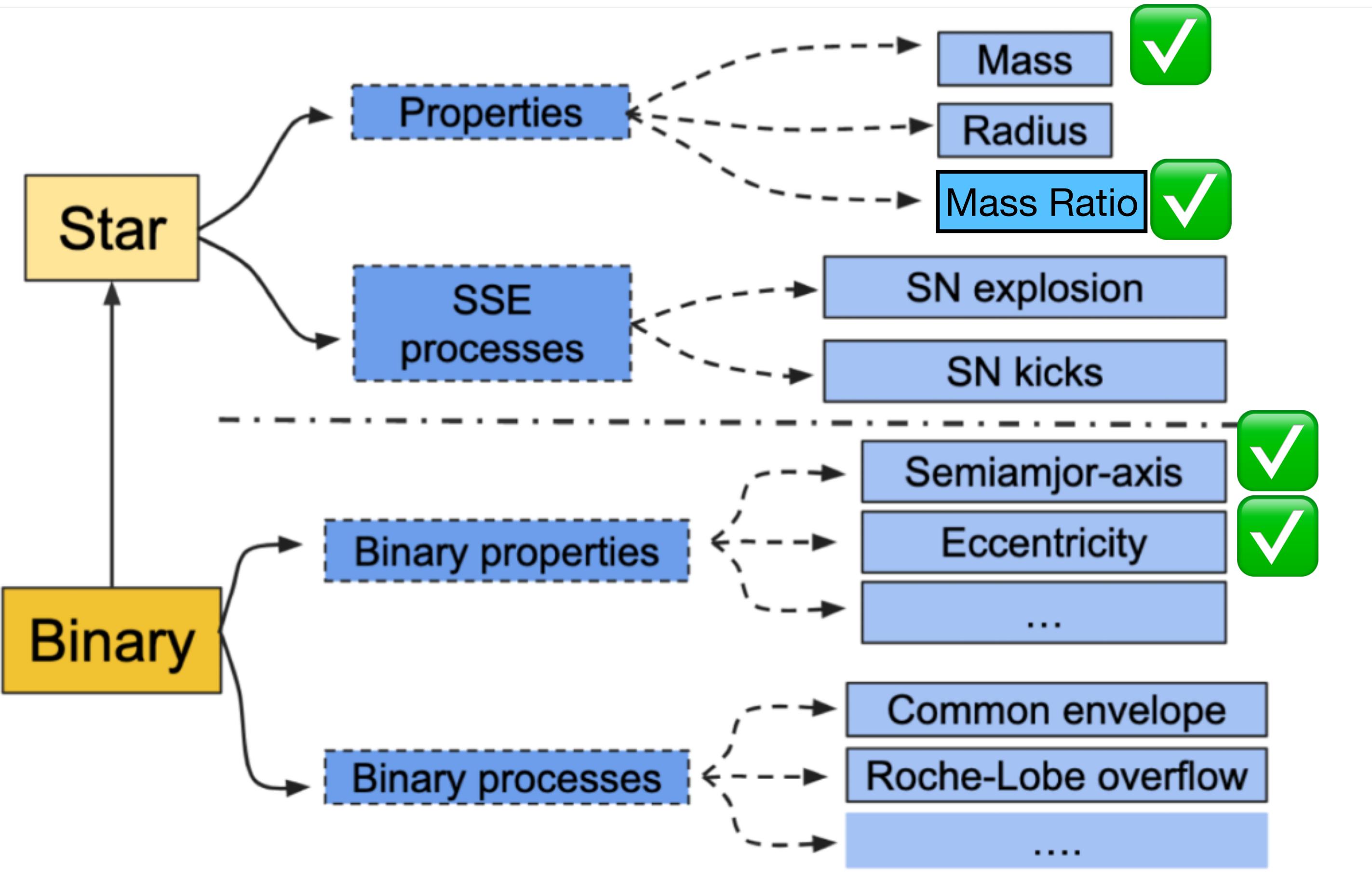
Population synthesis

- SEVN: population-synthesis code with **stellar tracks**
- [Costa et al. 2023](#) generated a new set of Pop. III stellar tracks
- We evolve a large set of Pop III. binary stars

Population synthesis

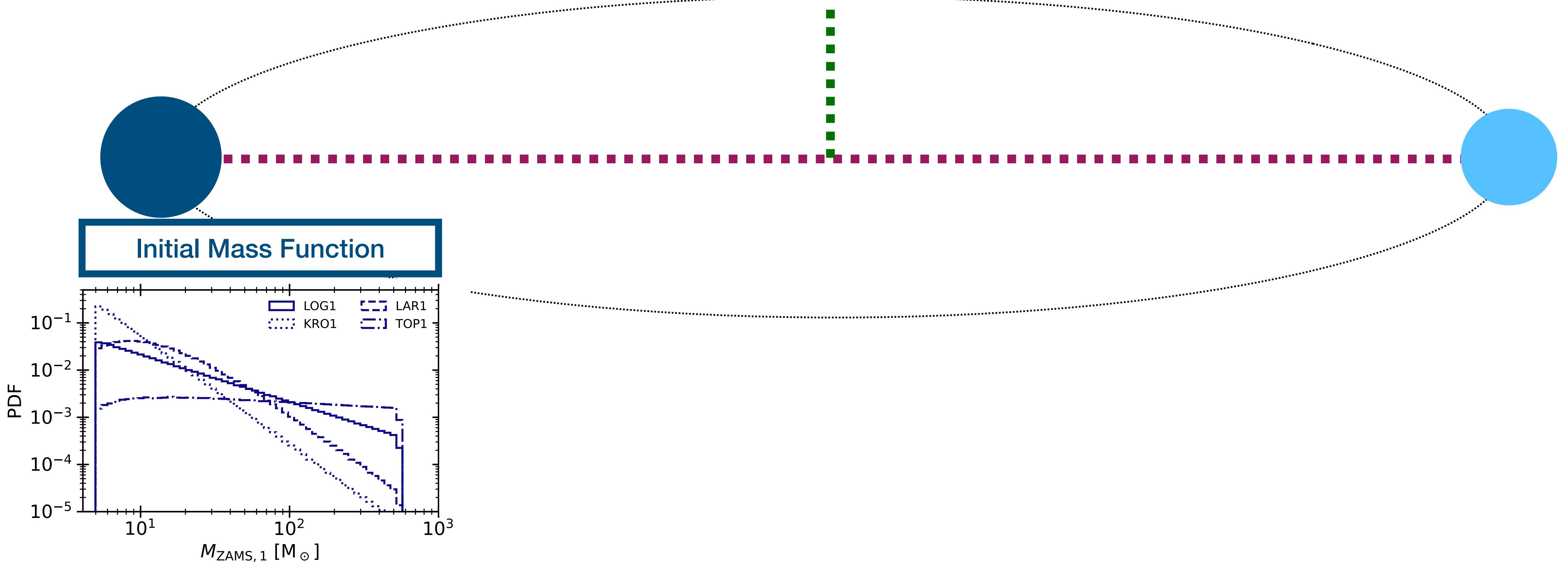


Population synthesis

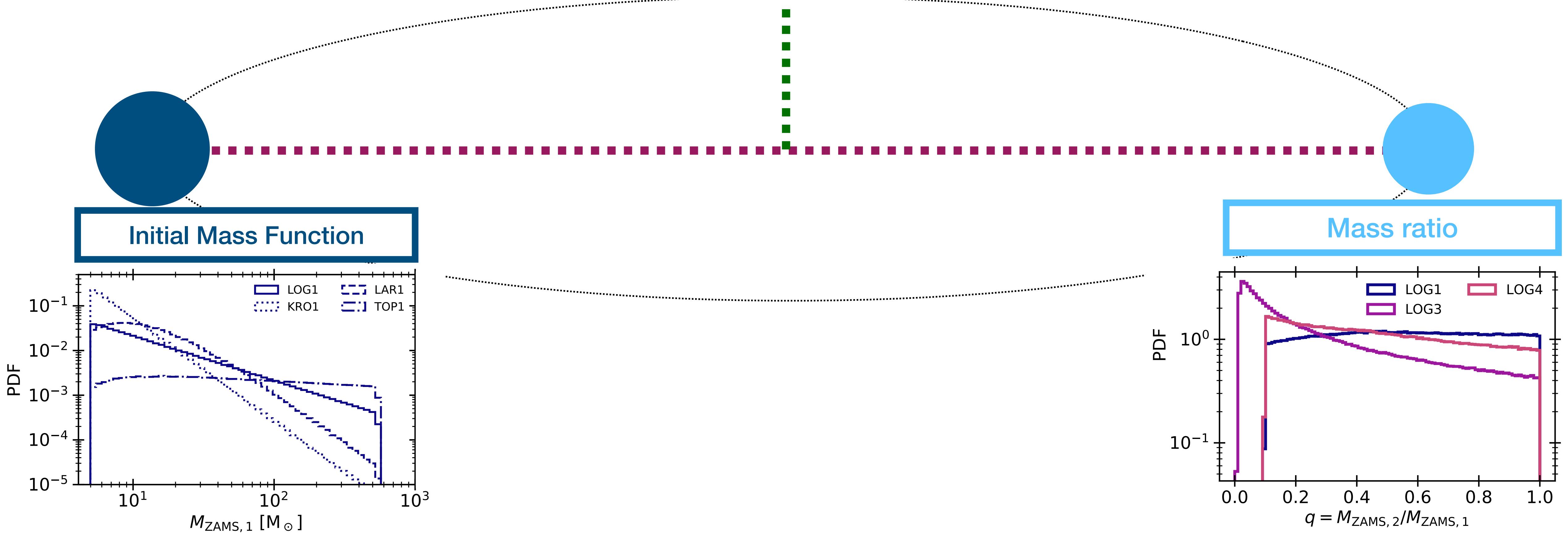


Exploring initial conditions

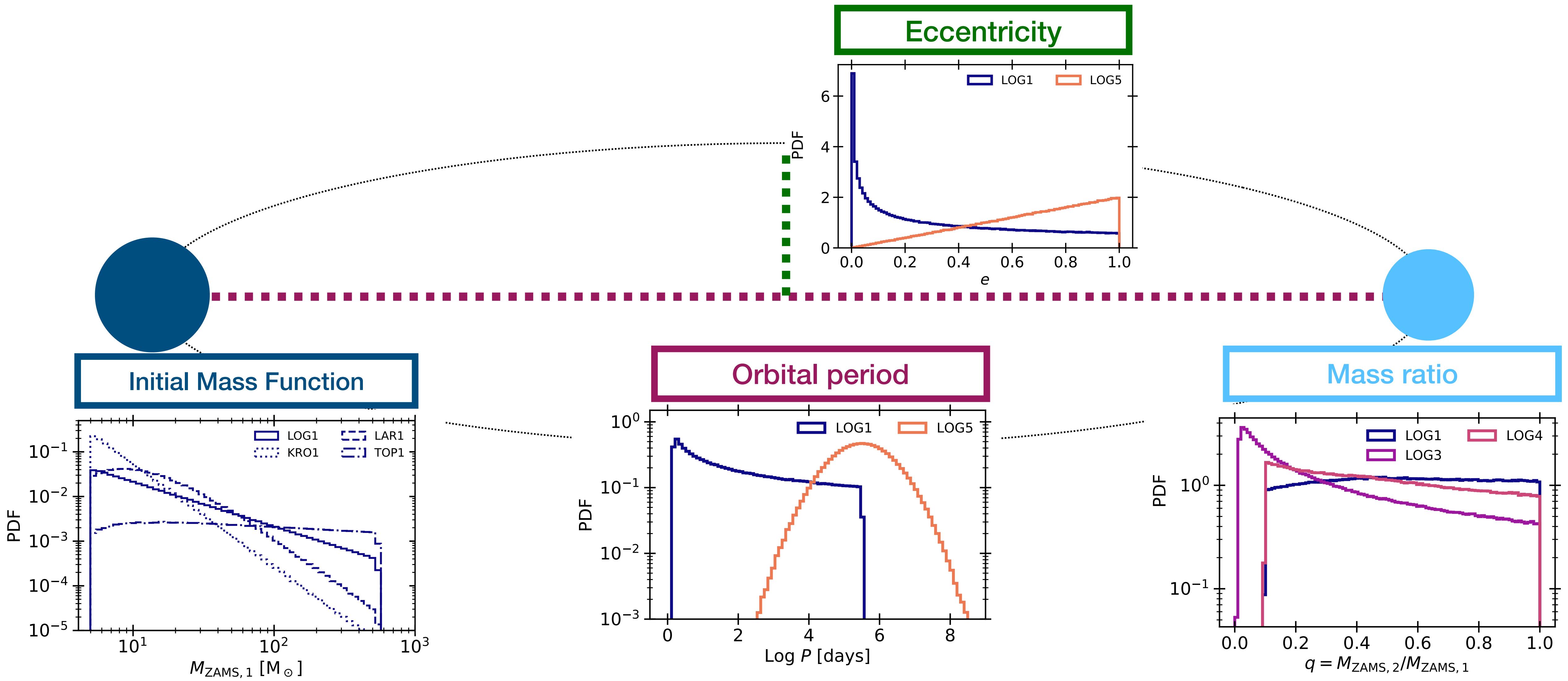
Initial conditions

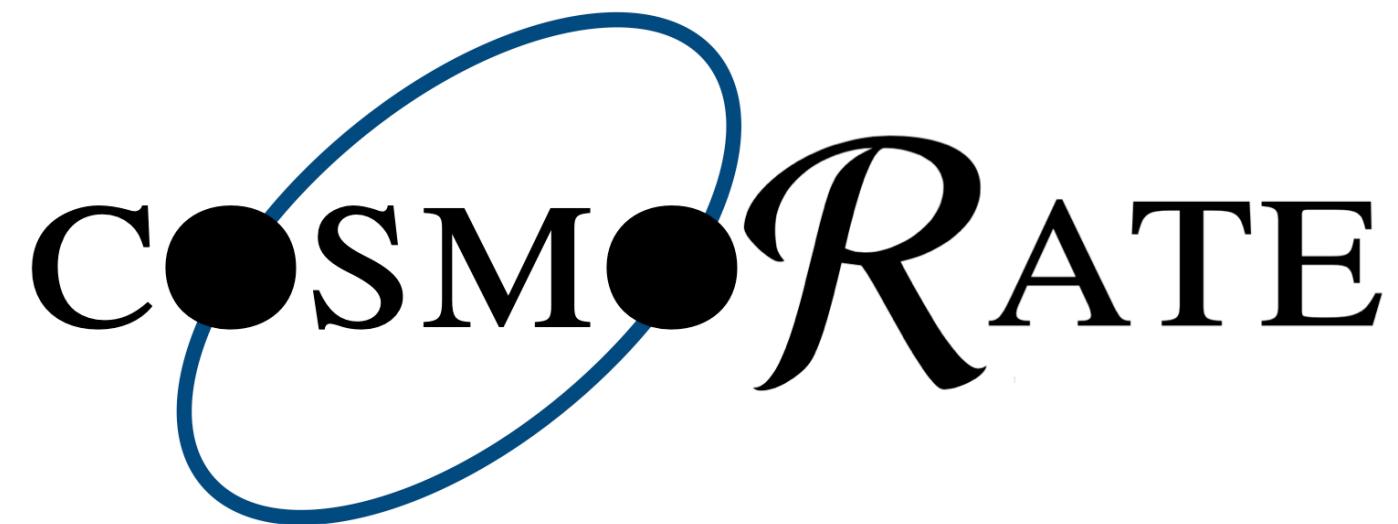


Initial conditions



Initial conditions



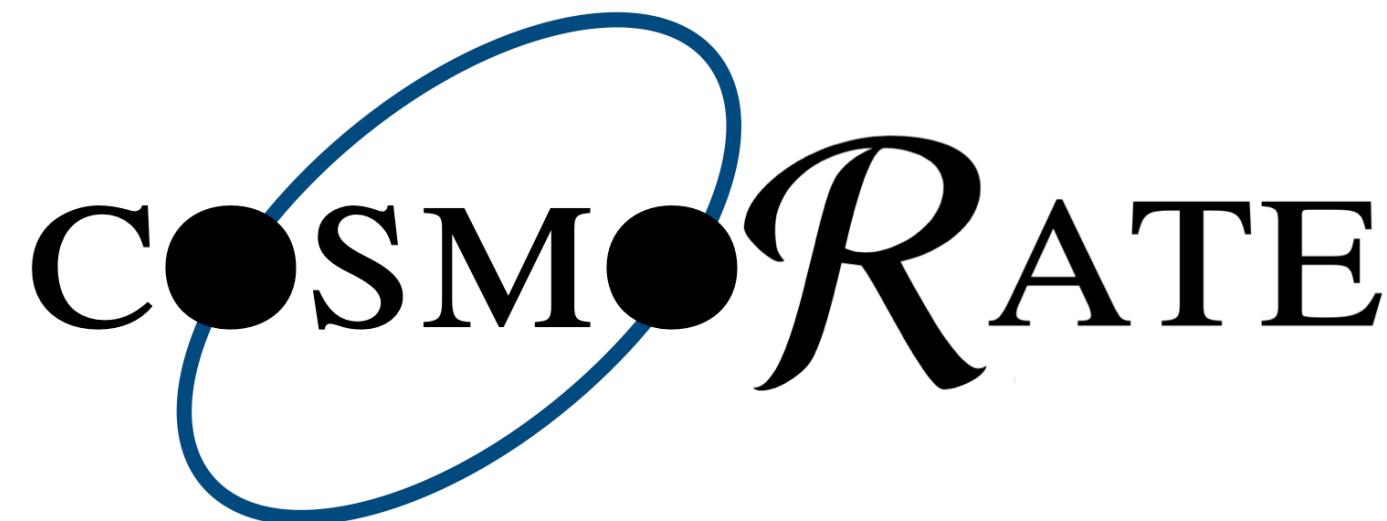


$$\mathcal{R}(z) = \int_{z_{\max}}^z \left[\int_{Z_{\min}}^{Z_{\max}} \text{SFRD}(z', Z) \mathcal{F}(z', z, Z) dZ \right] \frac{dt(z')}{dz'} dz'$$

SFRD(z', Z)
 $\mathcal{F}(z', z, Z)$

Output from SEVN

Catalogs of merging BBHs:
primary mass, secondary mass,
delay time



$$\mathcal{R}(z) = \int_{z_{\min}}^z \left[\int_{Z_{\min}}^{Z_{\max}} \text{SFRD}(z', Z) \mathcal{F}(z', z, Z) dZ \right] \frac{dt(z')}{dz'} dz'$$

The equation shows the calculation of the rate $\mathcal{R}(z)$ as a function of redshift z . It is a double integral where the inner integral is over the mass Z from Z_{\min} to Z_{\max} , and the outer integral is over time t from z_{\min} to z . The integrand is the product of the Star Formation Rate Density (SFRD) $\text{SFRD}(z', Z)$ and the selection function $\mathcal{F}(z', z, Z)$.

Two components of the integrand are highlighted with colored boxes:

- $\text{SFRD}(z', Z)$ is highlighted with a teal box and has a teal arrow pointing down to the definition below.
- $\mathcal{F}(z', z, Z)$ is highlighted with a magenta box and has a magenta arrow pointing down to the source information below.

Below the equation, the SFRD is defined as:

$$\text{SFRD}(z, Z) = \psi(z) p(Z|z)$$

And the source information is:

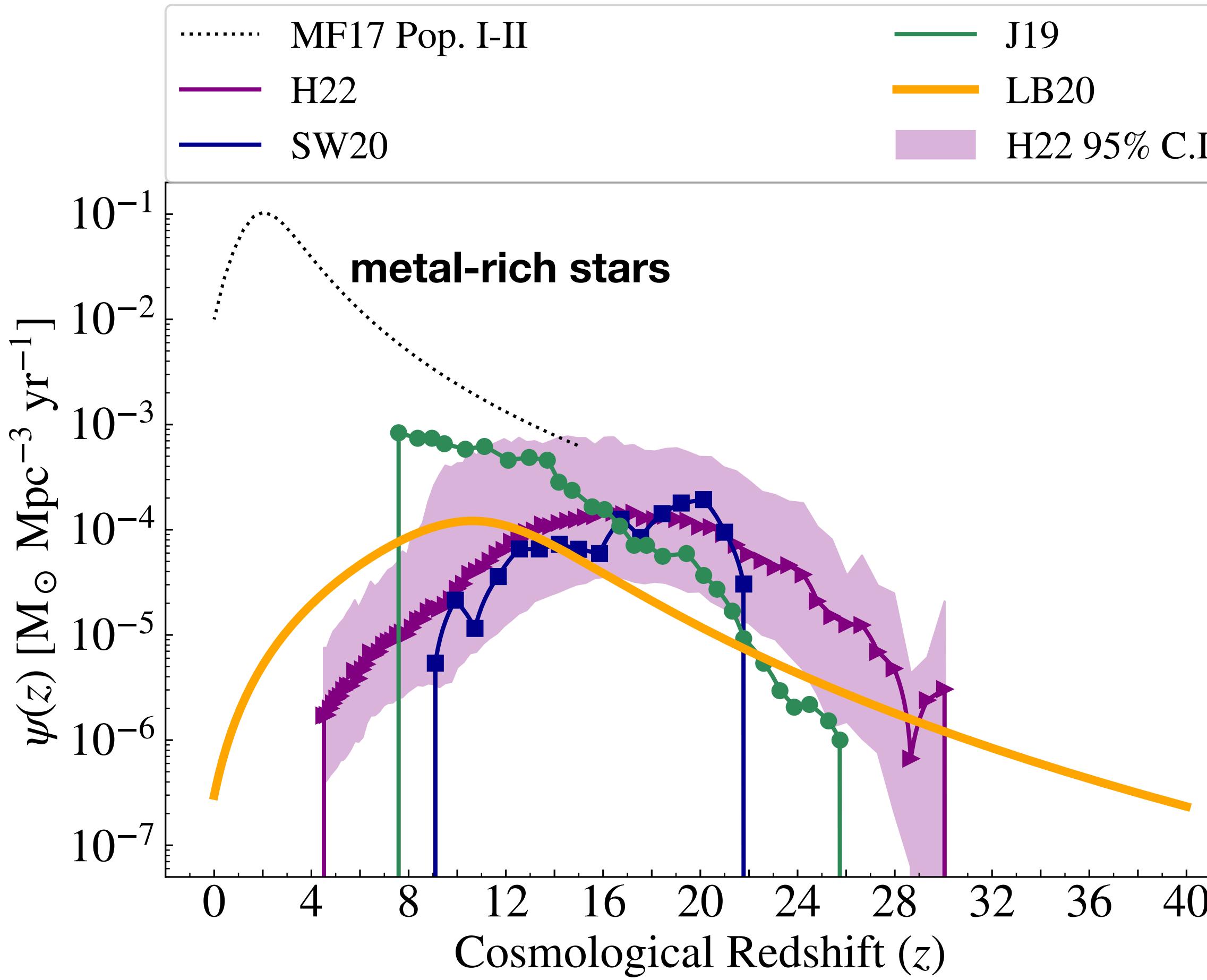
Evaluated from SEVN catalogs

COSMORATE

$$\mathcal{R}(z) = \int_{z_{\max}}^z \left[\int_{Z_{\min}}^{Z_{\max}} \text{SFRD}(z', Z) \mathcal{F}(z', z, Z) dZ \right] \frac{dt(z')}{dz'} dz'$$

SFRD(z, Z) = $\psi(z) p(Z|z)$

Evaluated from SEVN catalogs

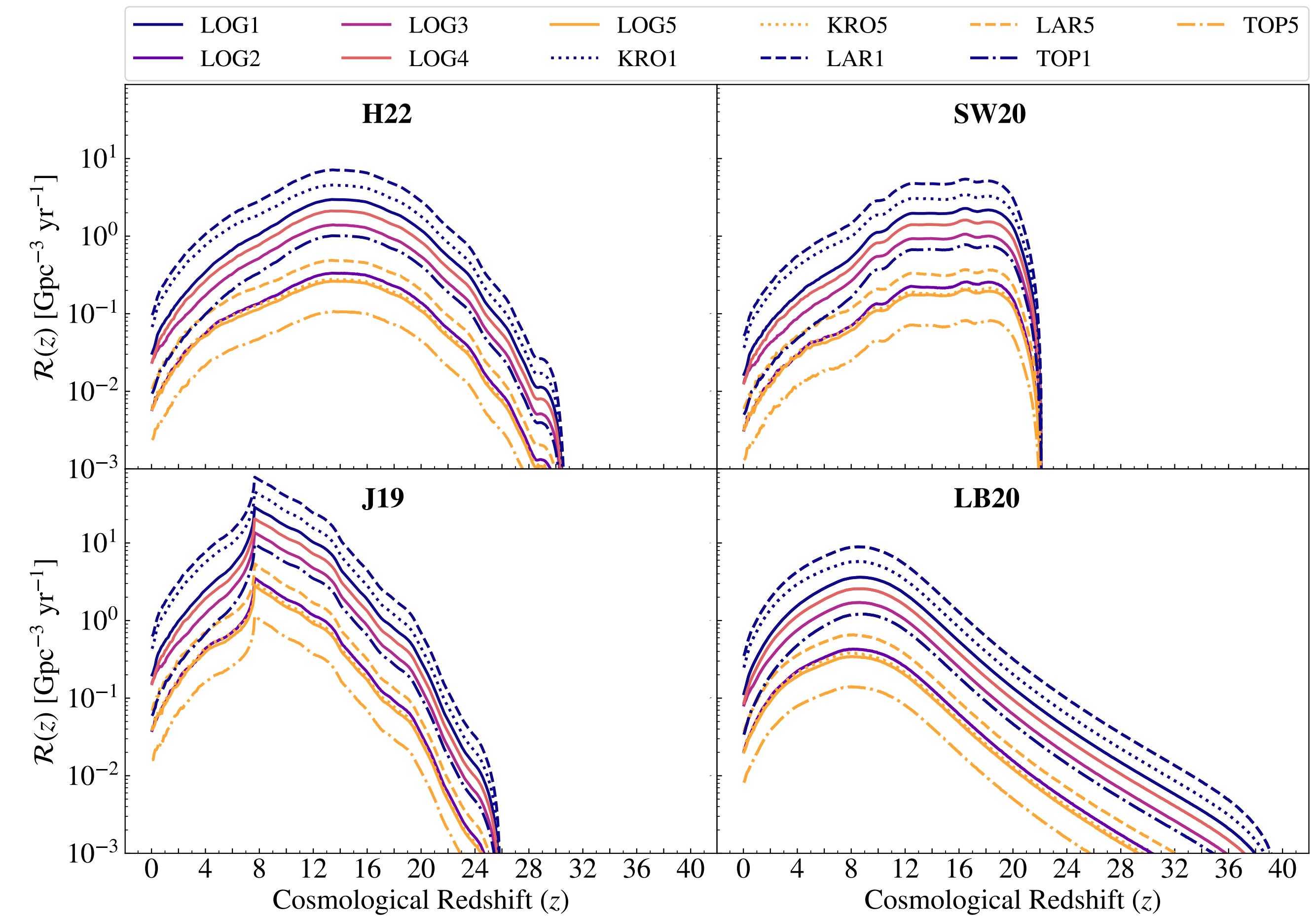


4 different Pop. III SFRD:

- H22** - [Hartwig et al. 2022](#)
- J19** - [Jaacks et al. 2019](#)
- LB20** - [Liu & Bromm 2020](#)
- SW20** - [Skinner & Wise 2020](#)

different assumptions on baryonic physics
+ cosmic variance

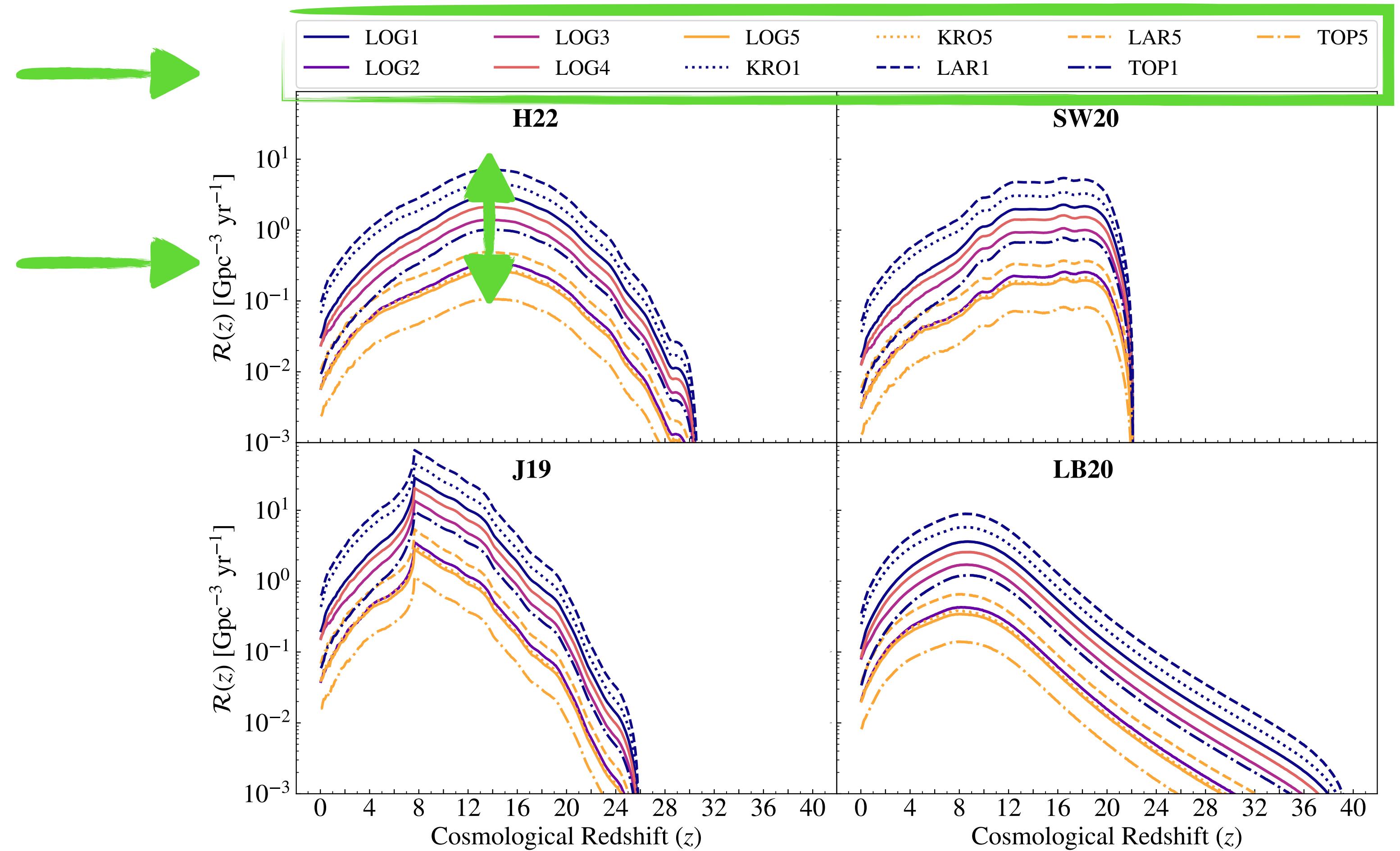
Merger rate density



Merger rate density

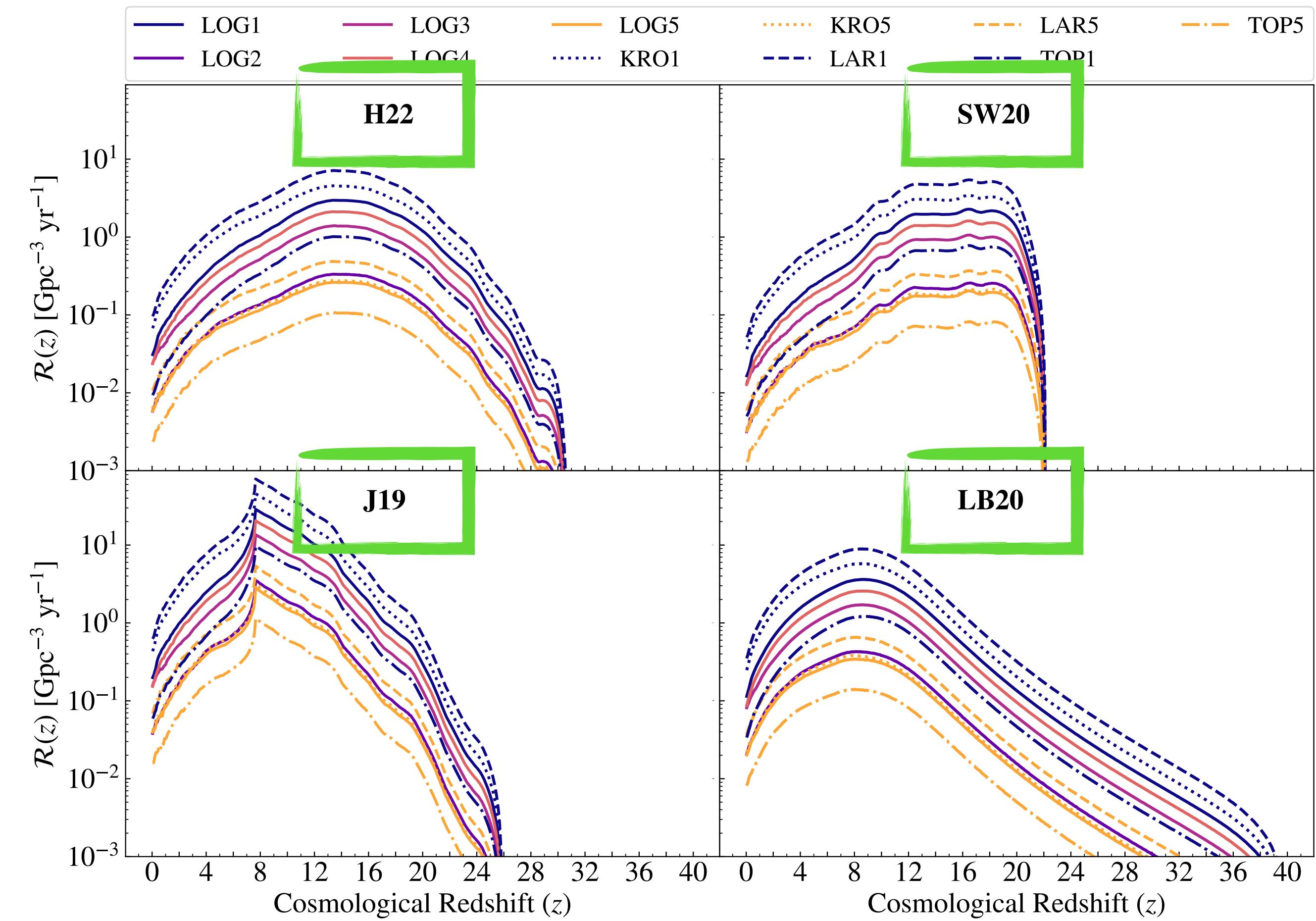
initial conditions

impact ~2 orders of magnitude

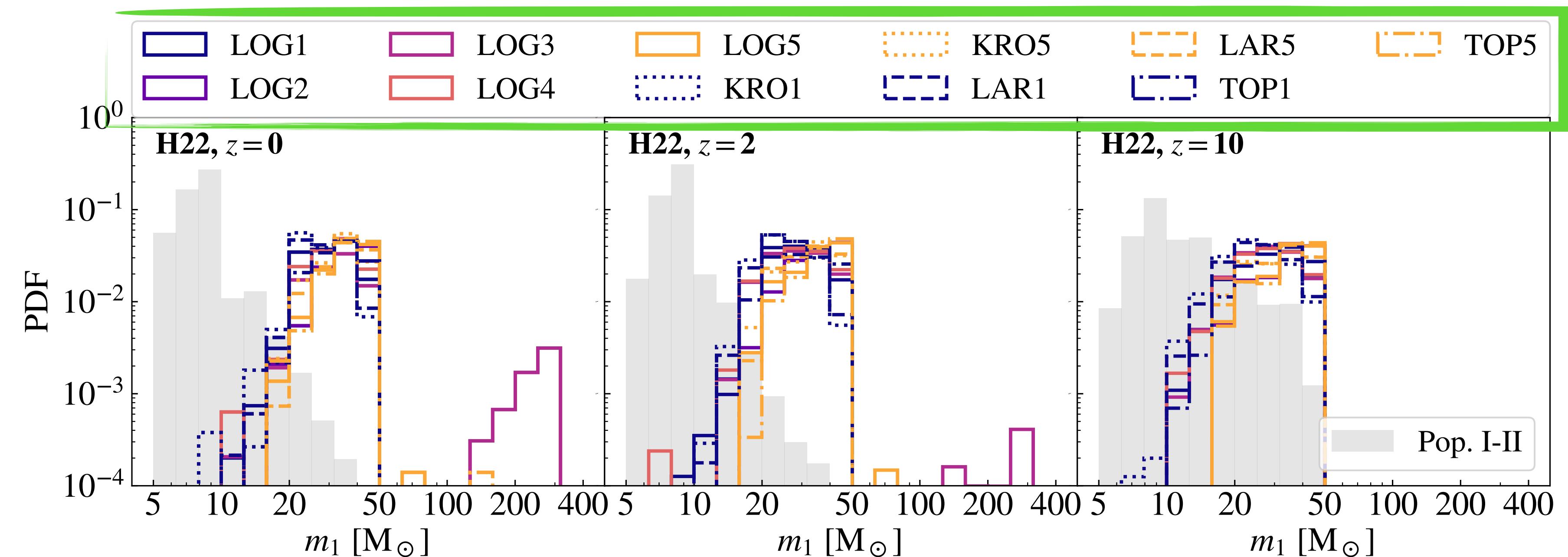


Merger rate density

star formation history impacts
shape and normalisation

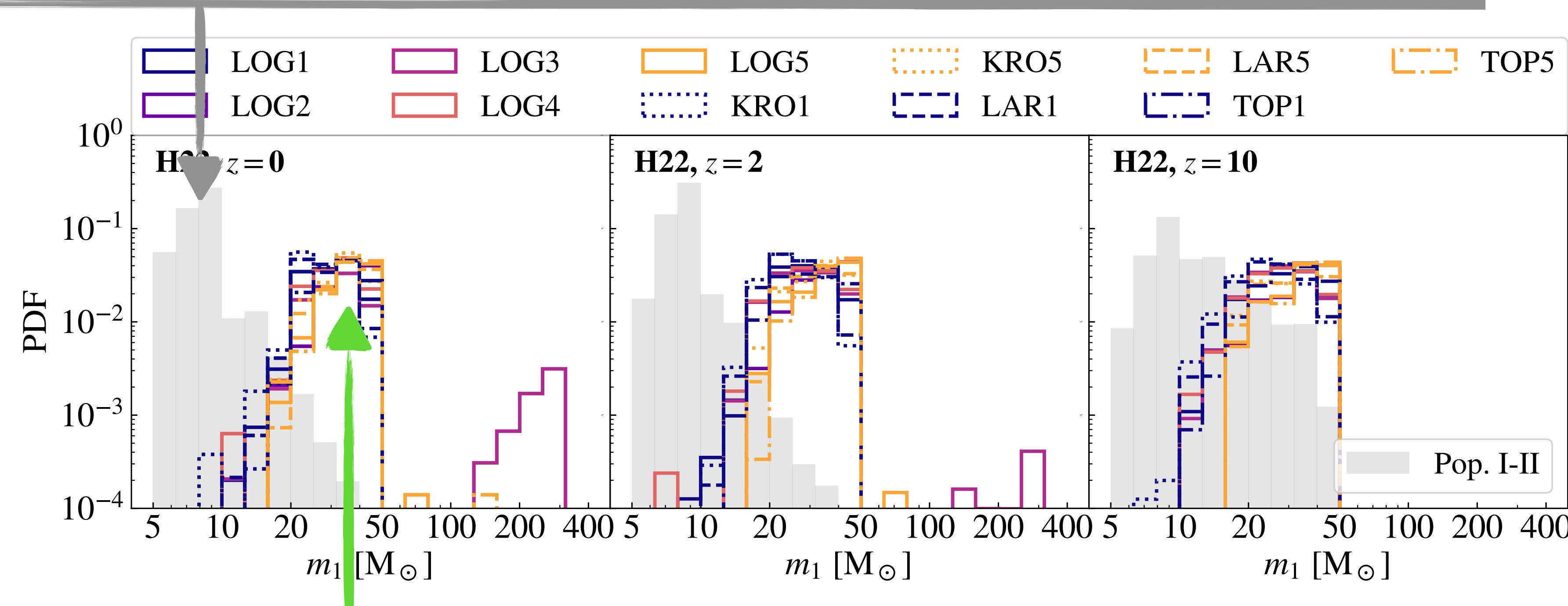


Primary mass



Primary mass

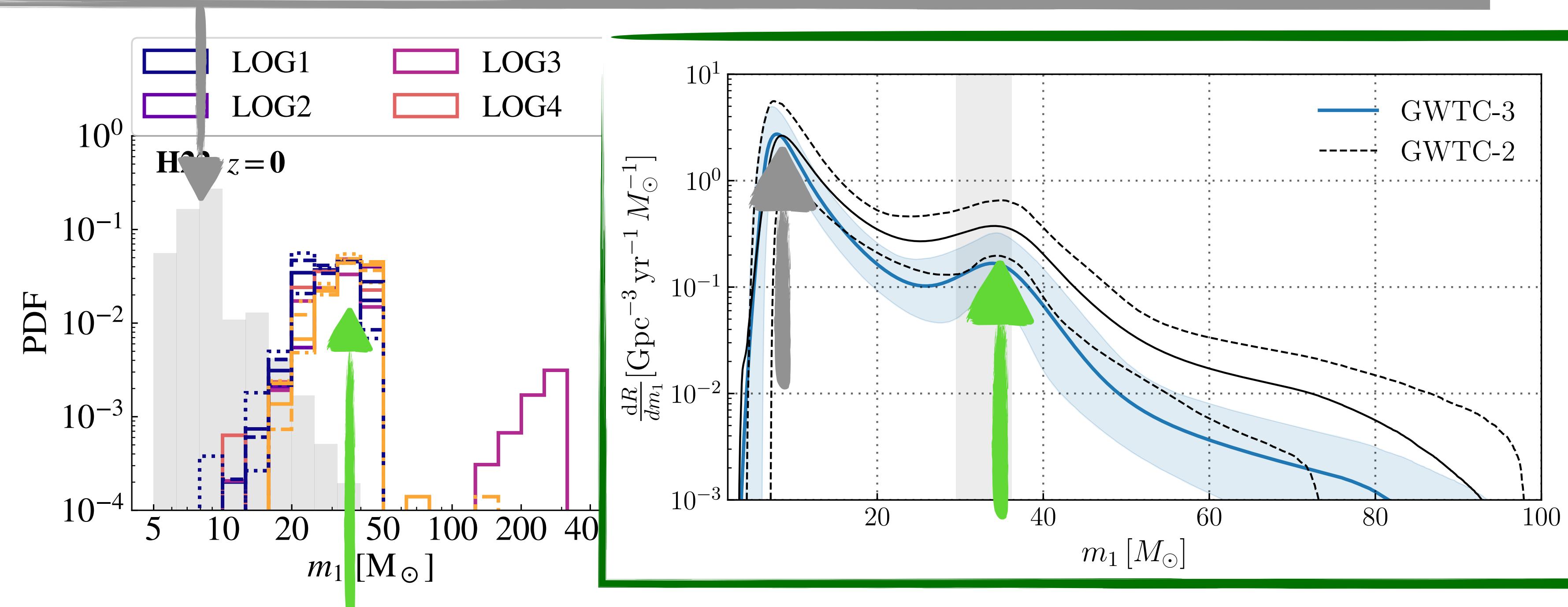
At $z = 0$, Pop. I-II BBHs show a **main peak at $8 - 10 M_{\odot}$**



Pop. III BBHs show a **main peak at $30 - 35 M_{\odot}$**

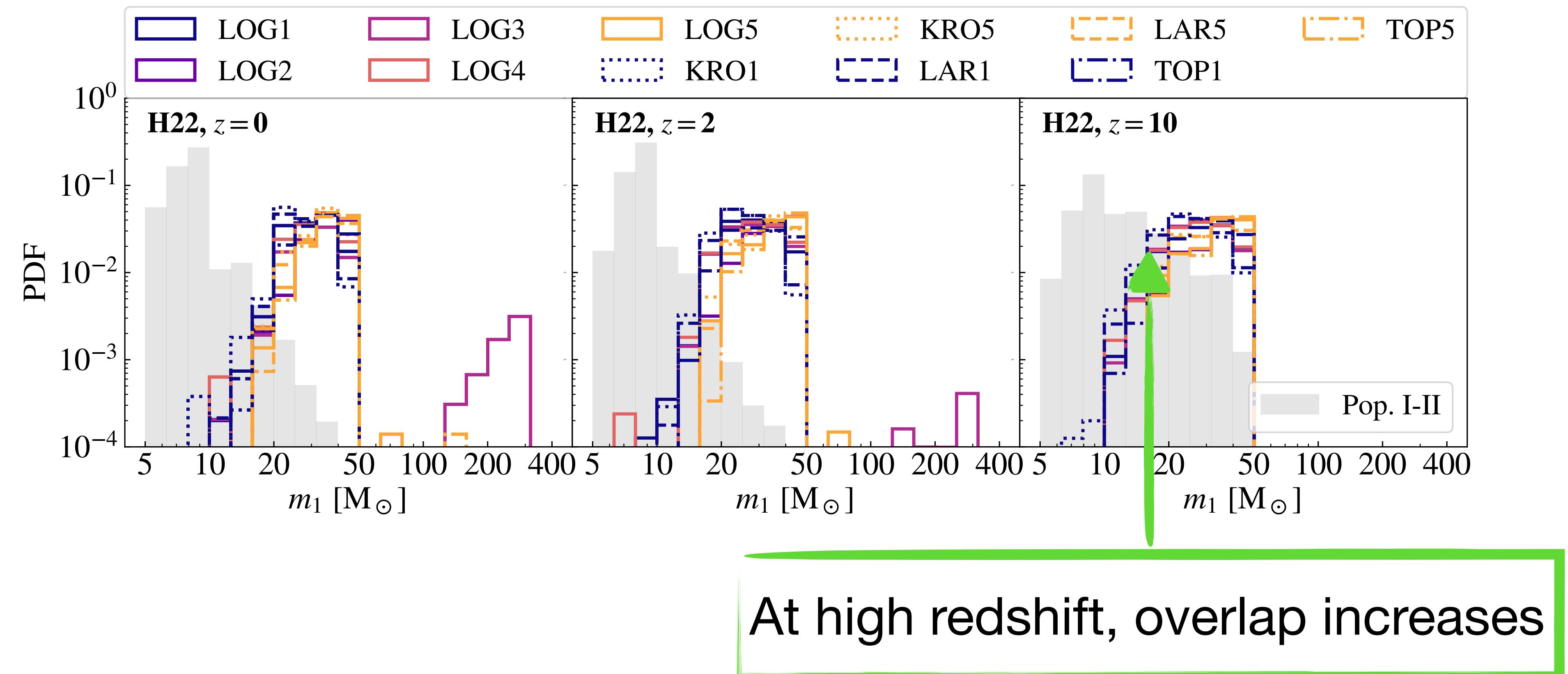
Primary mass

At $z = 0$, Pop. I-II BBHs show a **main peak at $8 - 10 M_{\odot}$**



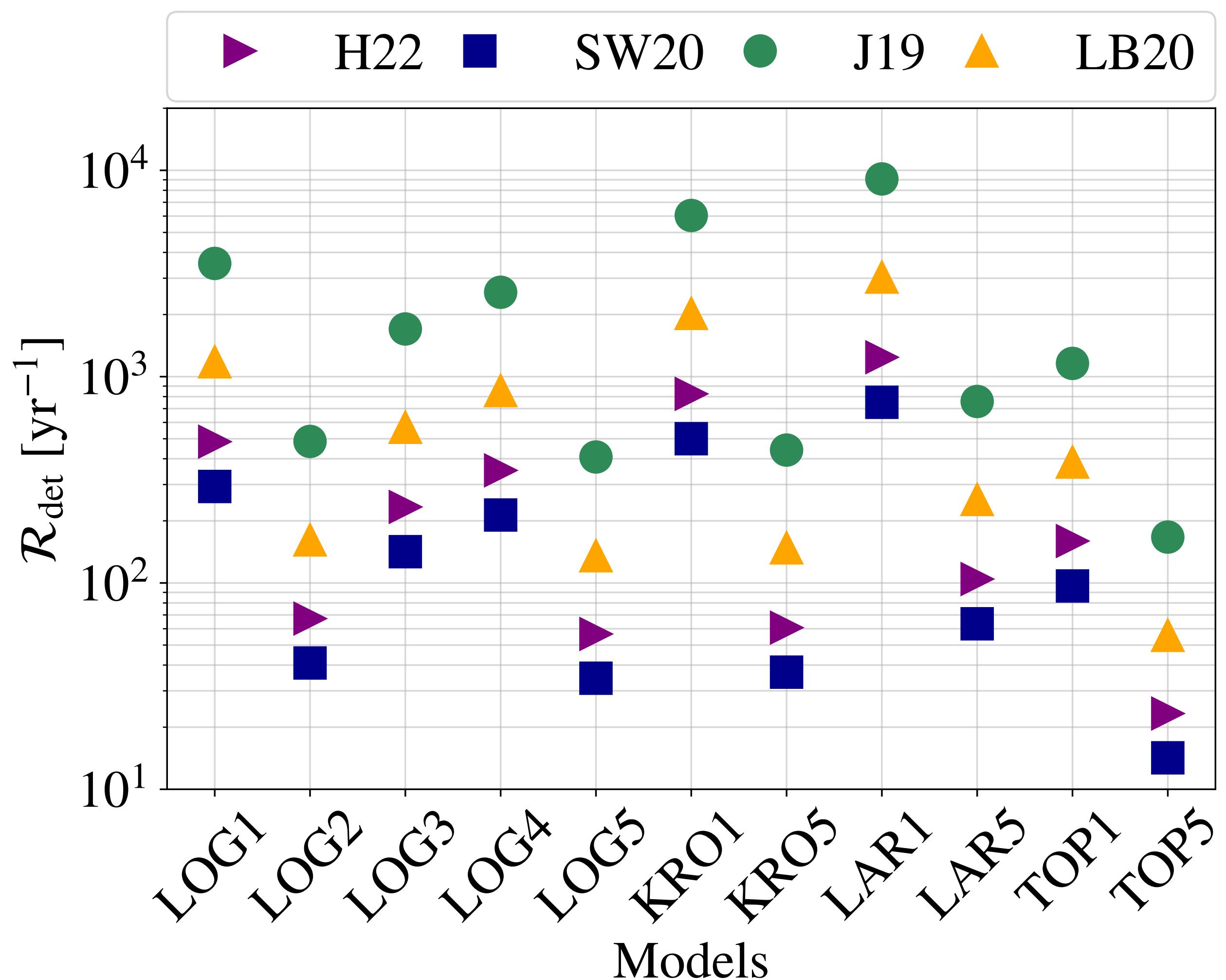
Pop. III BBHs show a **main peak at $30 - 35 M_{\odot}$**

Primary mass



Detection rate

- Einstein Telescope will detect $10 - 10^4$ Pop. III BBH mergers per year
- We expect between 23% and 73% of detections to occur at redshift $z > 8$

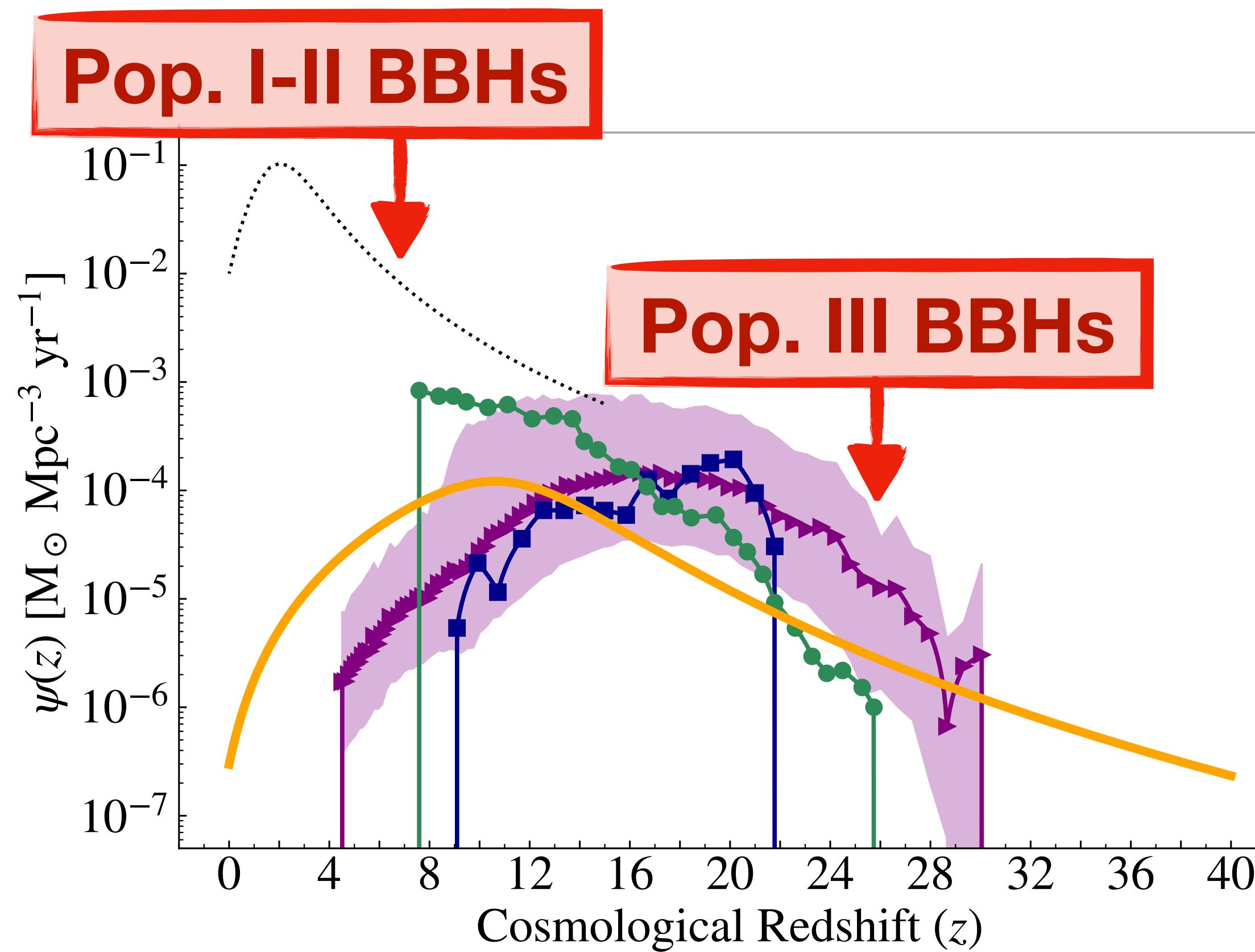


Can we distinguish Pop. III BBHs?

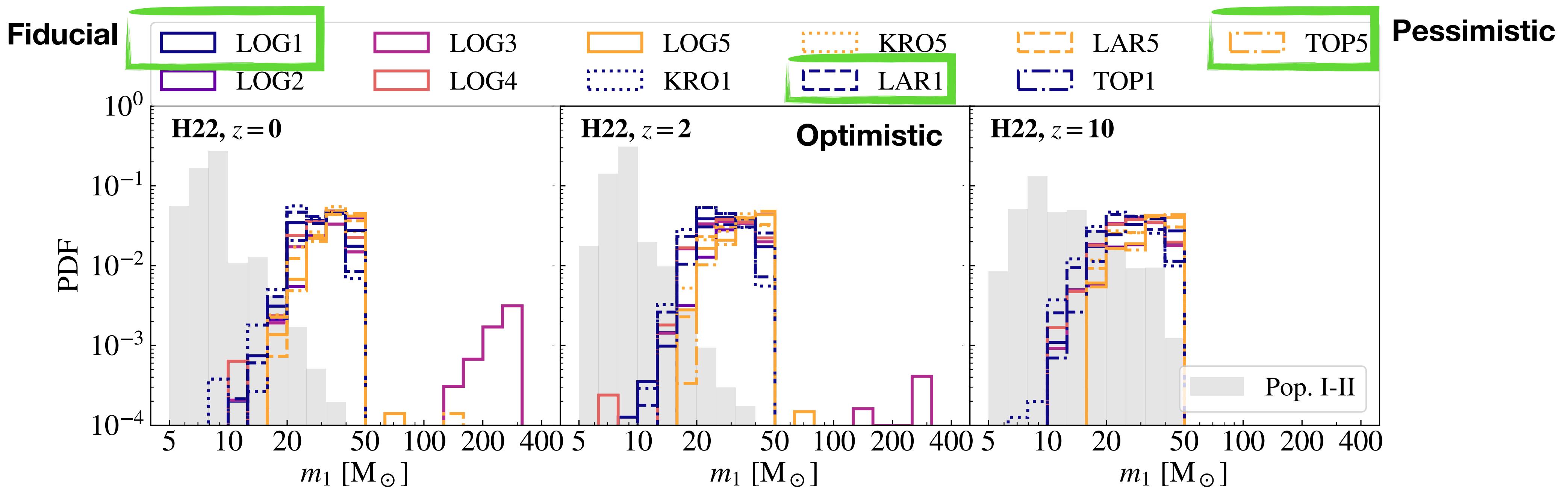
- **Einstein Telescope** will detect BBH mergers up to $z \sim 100$
- **high-redshift sources** with **low-SNR** and poor estimate of d_L
- **inferring the origin** of individual GW detections will not be granted

Ref. [Maggiore et al. 2020](#), [Ng et al. 2021](#), [2022](#), [Branchesi et al. 2023](#), [Mancarella et al. 2023](#)

Goal: classification



Simulation-based classification



Classification

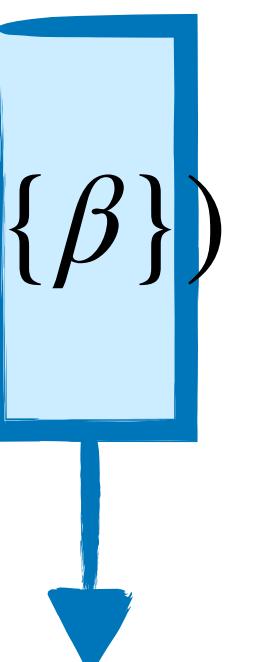
$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



This is the probability that the event j is a Pop. III BBH

Classification

Mixing fraction

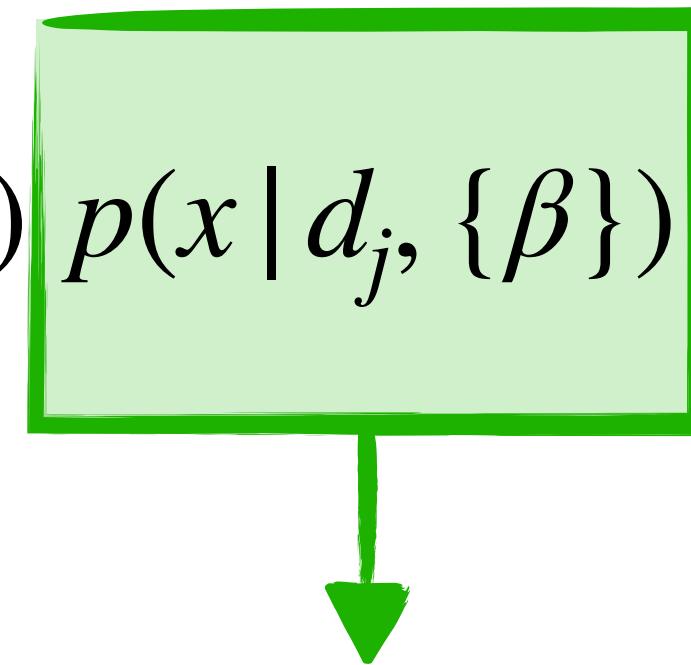
$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$


Fiducial ~4% $\beta_{III} \propto N_{\text{Pop. III}} \sim 400$

$\beta_{I-II} \propto N_{\text{Pop. I-II}} \sim 10^4$

Classification

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



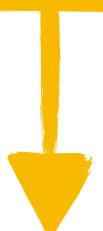
This is the posterior of waveform parameters →
parameter-estimation **performance of ET**

Fisher Information Matrix approximation



Classification

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



This is the probability that links waveform parameters to Pop. III BBHs

→ easiest approach is to consider a fix threshold

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

↓

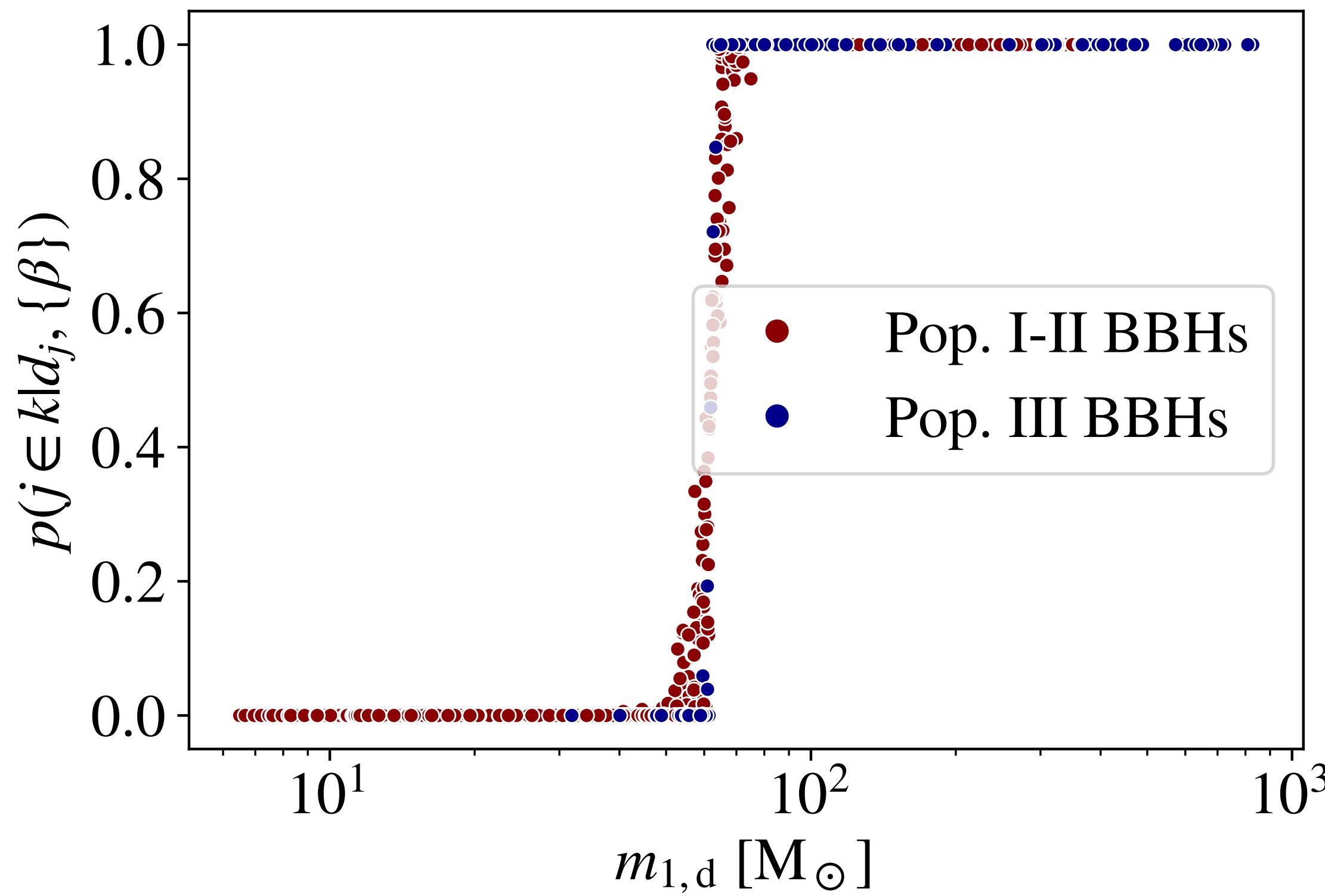
$$= 1 \quad \text{if } m_{1,d} \gtrsim 60 \text{ M}_\odot$$

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

↓

$$= 1 \quad \text{if } m_{1,d} \gtrsim 60 \text{ M}_\odot$$

manual classifier - fiducial



⚠️ low performances: precision is ~ 0.16

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



we can use Machine Learning

XGBoost

supervised ML based on decision trees

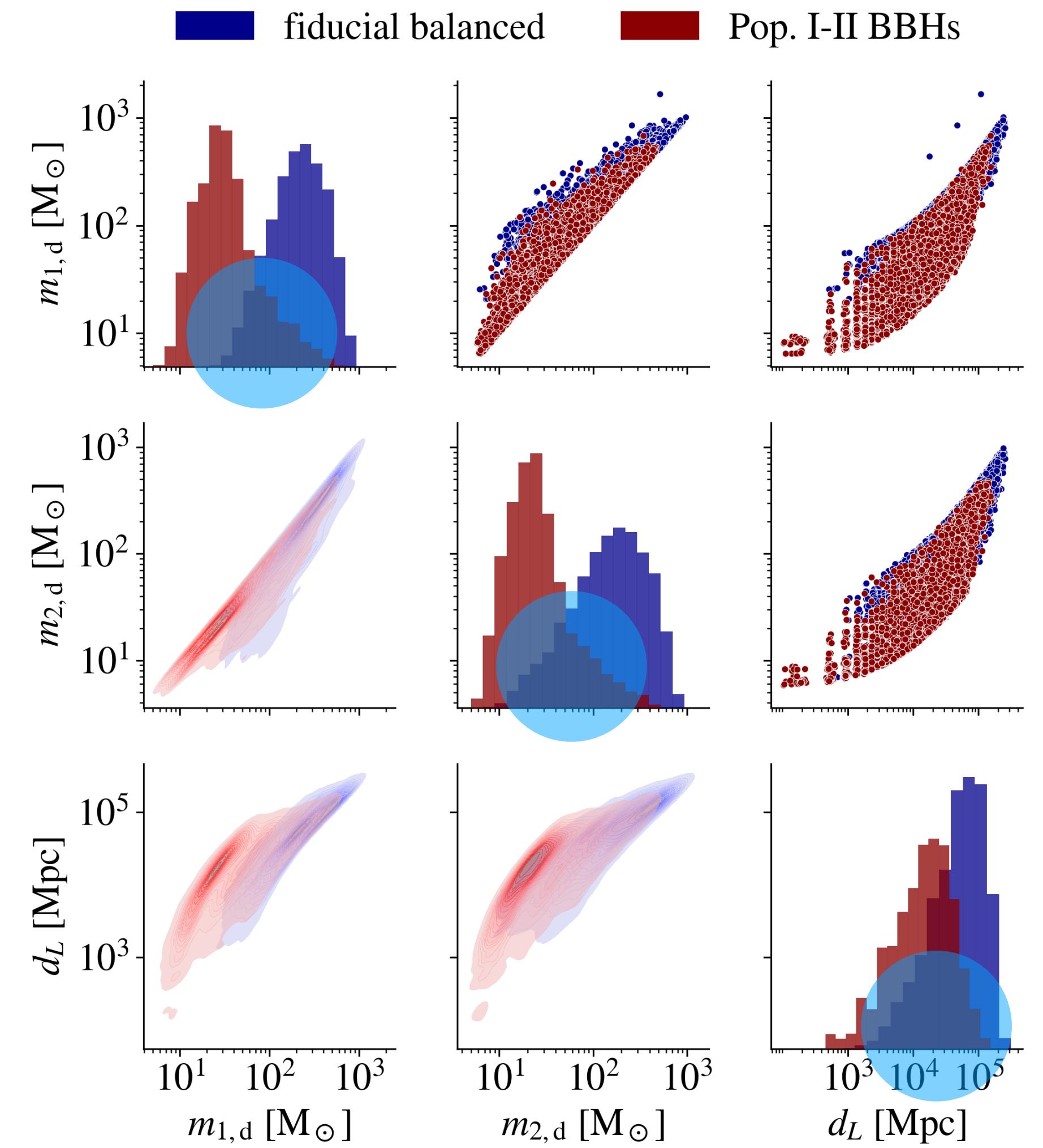
$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



Using Machine Learning



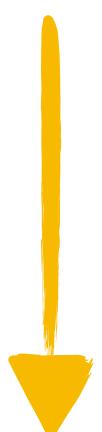
- * trained and tested on balanced classes + re-balancing
- * instances: $> 10^4$



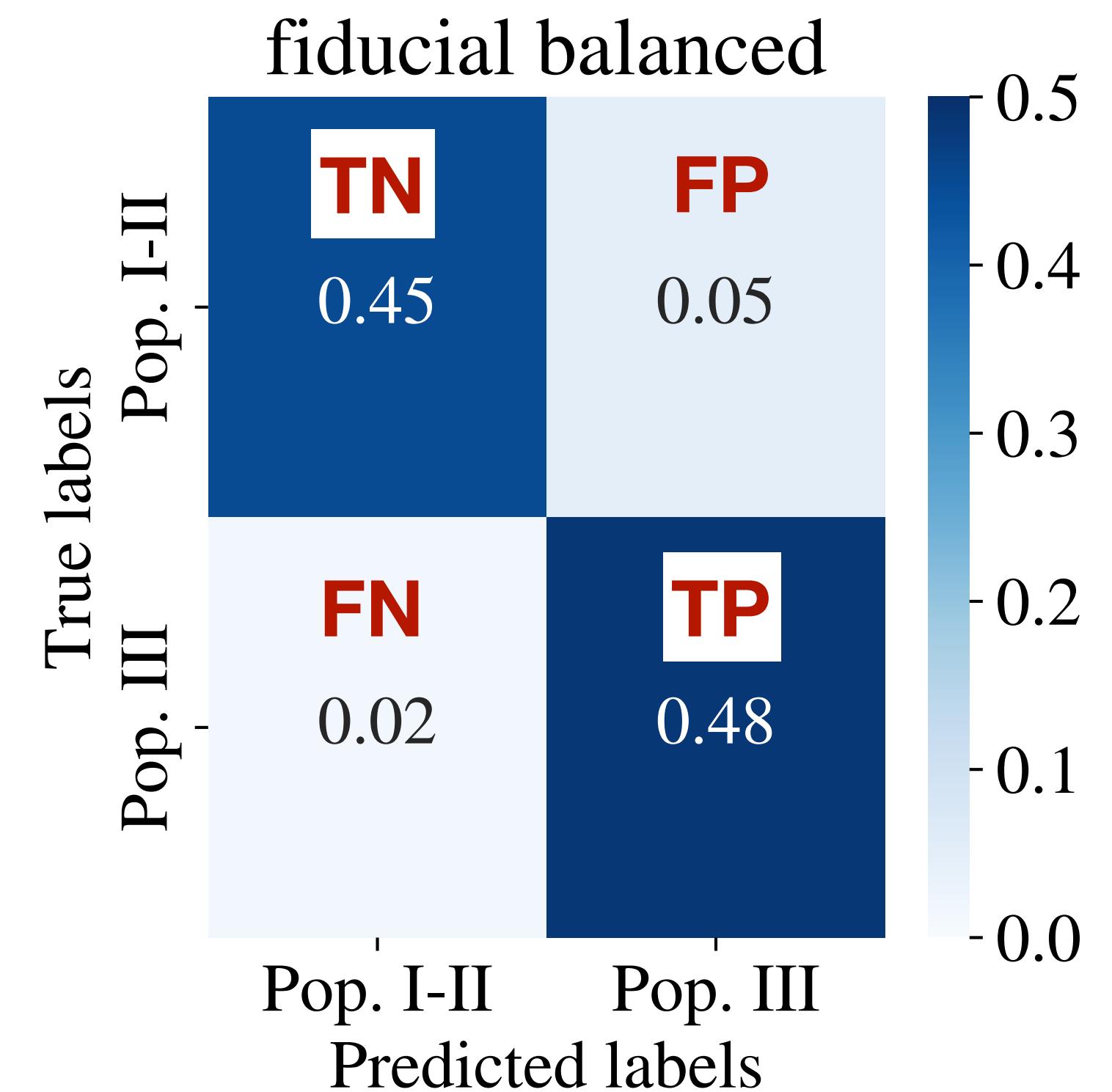
$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



Using Machine Learning

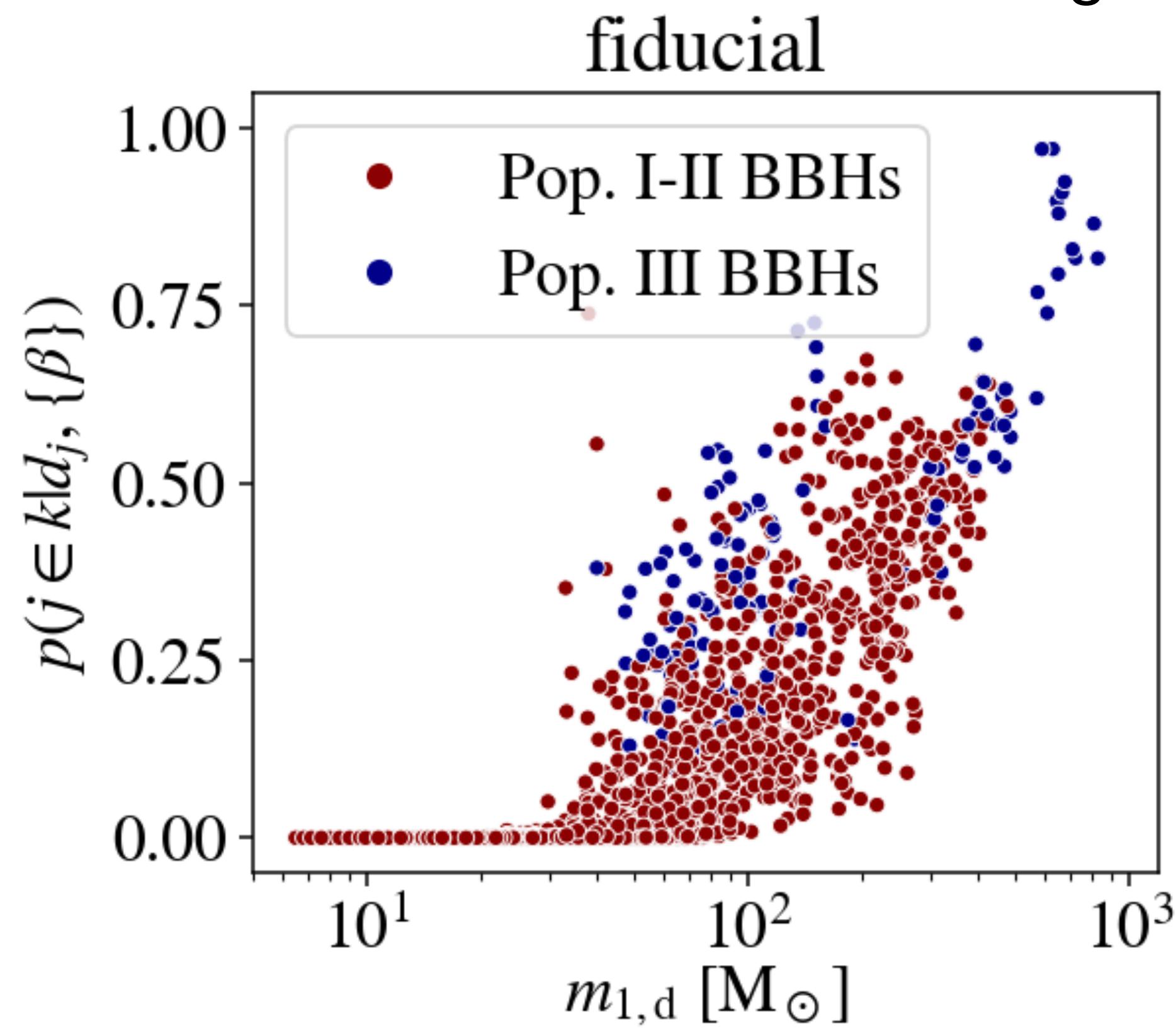


precision = TP/(TP+FP)
high precision > 0.90



$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$

Using Machine Learning

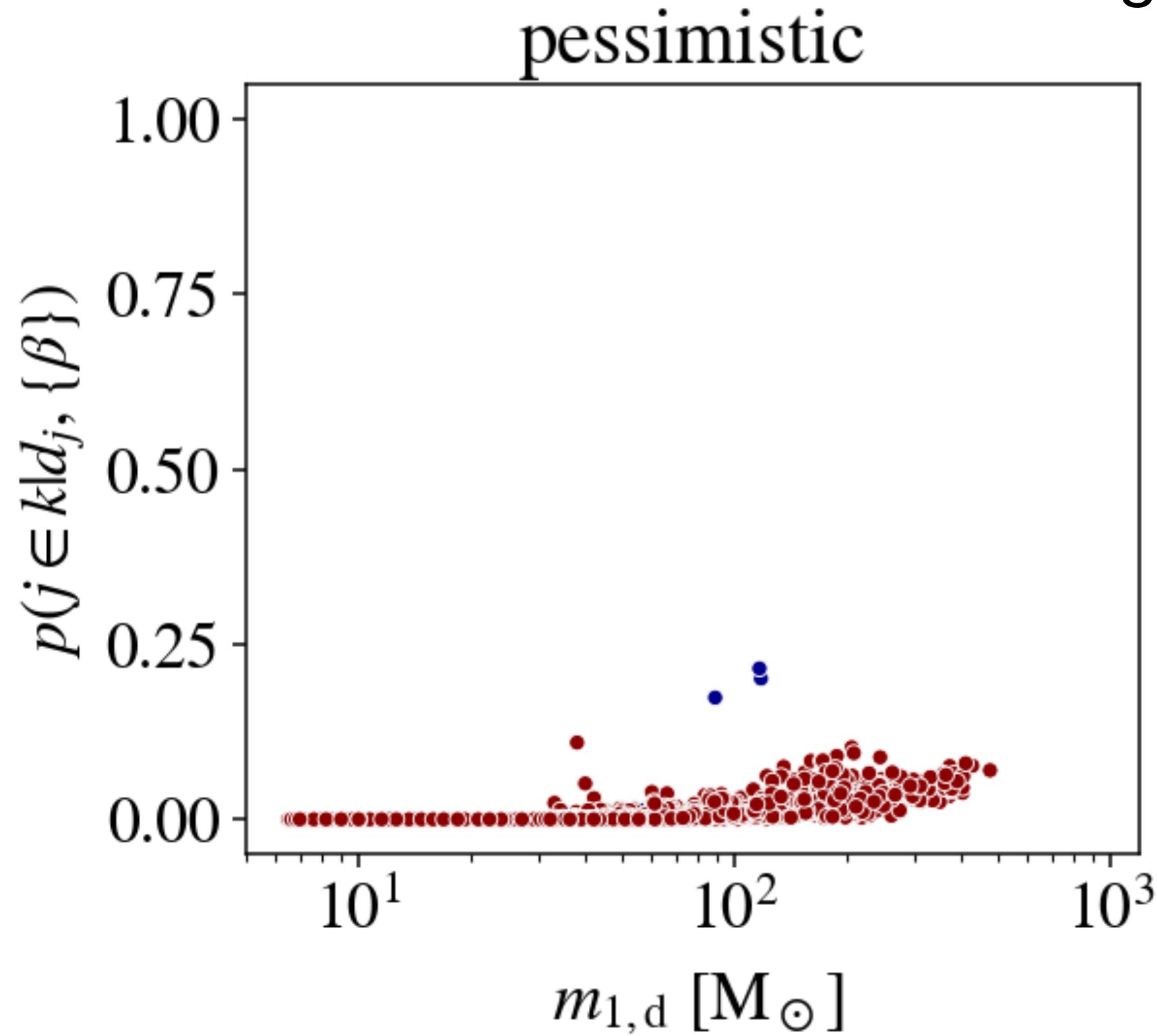


**~10% of detected sources are classified
with precision > 0.90**

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



Using Machine Learning

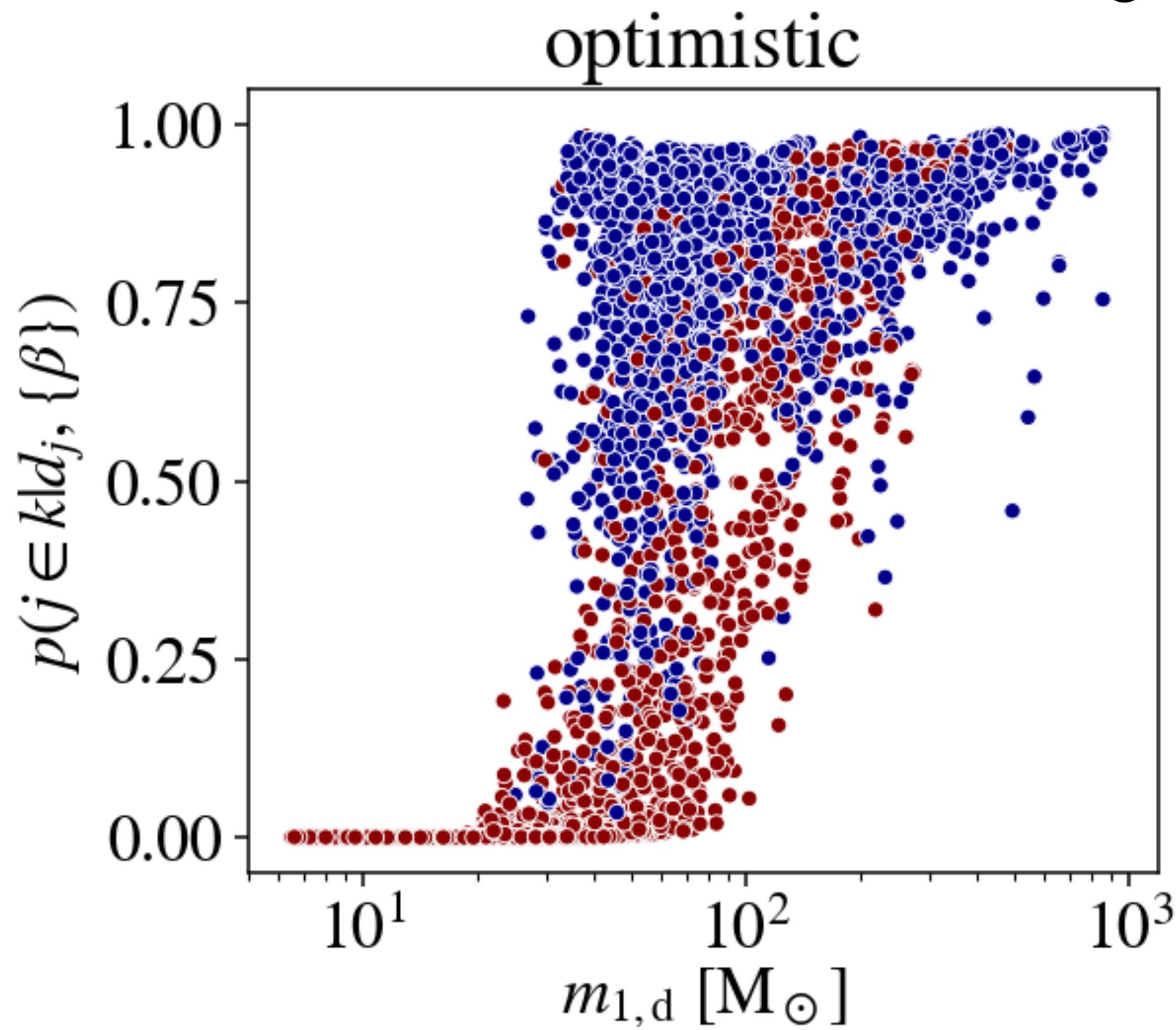


~30% of detected sources are classified with precision > 0.90

$$p(j \in k | d_i, \{\beta\}) = \int p(j \in k | x, d_j, \{\beta\}) p(x | d_j, \{\beta\})$$



Using Machine Learning



**~45% of detected sources are classified
with precision > 0.90**

Contributions

- First **large parameter exploration of Pop. III BHs**
 - **SFRD affects normalisation and shape of merger rate density**
 - primary mass of Pop. III BHs is **substantially massive**
- **ET will detect** these sources and **machine learning** boosts our ability to classify them

Backup slides

$\alpha\lambda$ formalism for modelling the common envelope

- $\Delta E = \alpha(E_{b,f} - E_{b,i}) = \alpha \frac{Gm_{c1}m_{c2}}{2} \left(\frac{1}{a_f} - \frac{1}{a_i} \right)$ This is the orbital energy before and after the common envelope phase
- $E_{\text{env}} = \frac{G}{\lambda} \left[\frac{m_{\text{env},1}m_1}{R_1} + \frac{m_{\text{env},2}m_2}{R_2} \right]$ This is the binding energy of the envelope
- By imposing $\Delta E = E_{\text{env}}$, $\frac{1}{a_f} = \frac{1}{\alpha\lambda} \frac{2}{m_{c1}m_{c2}} \left[\frac{m_{\text{env},1}m_1}{R_1} + \frac{m_{\text{env},2}m_2}{R_2} \right] + \frac{1}{a_i}$
- If α is larger, a_f is larger, following $a_f \sim \frac{\alpha}{1 + \alpha}$. Therefore larger α gets wider binaries
- Where λ is the parameter which measures the concentration of the envelope (the smaller λ is, the more concentrated is the envelope).
- The $\alpha\lambda$ formalism is a simplified prescription. When $\alpha > 1$, we account for other sources of energy that make the envelope less bind, for instance recombination energy. Recent works (e.g. [Fragos et al. 2019](#)) suggest that $\alpha > 1$ is necessary to reproduce the final orbital separation obtained with hydrodynamical simulations.

Initial conditions

Table 1. Initial conditions.

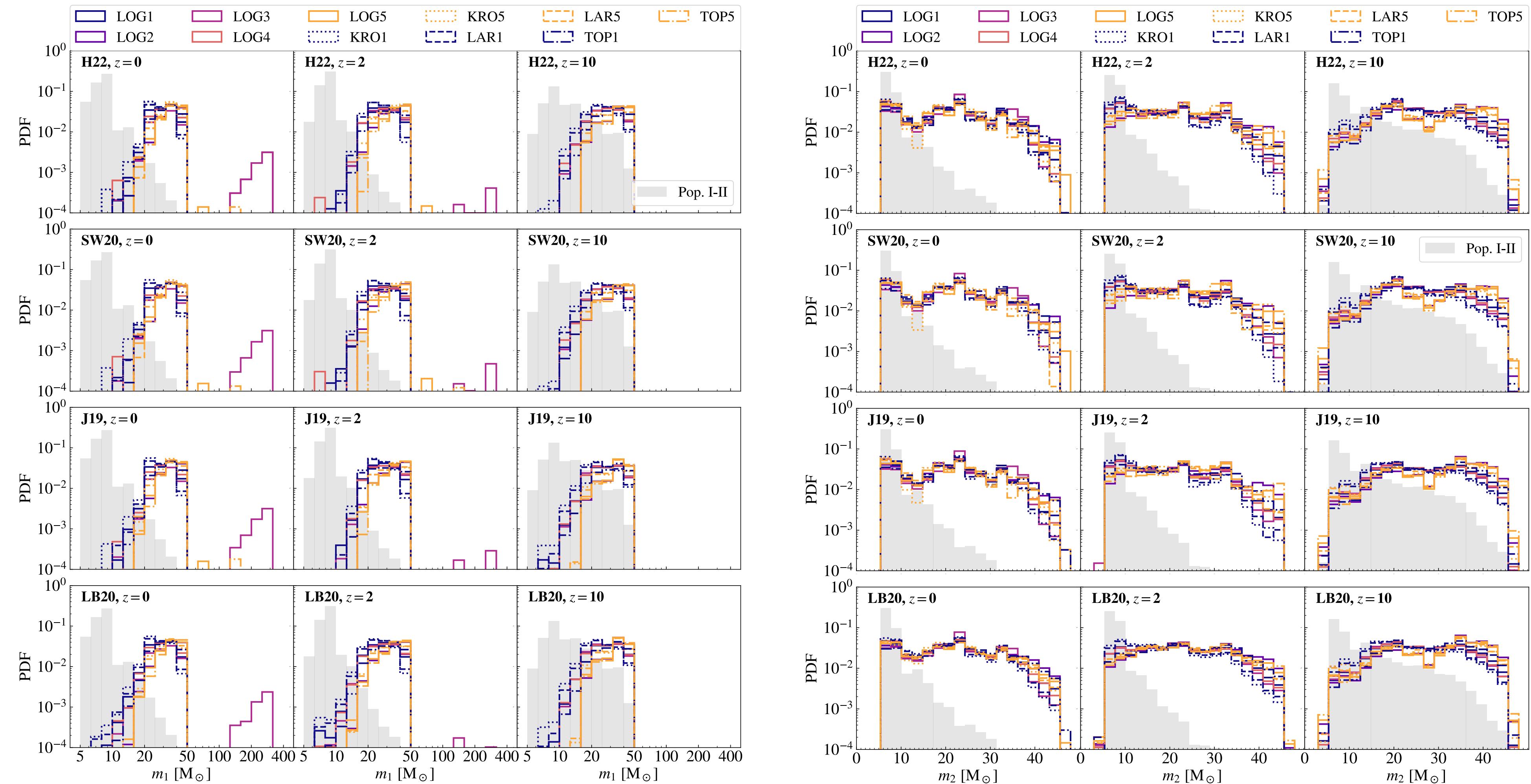
Model	$M_{\text{ZAMS},1}$	M_{ZAMS}	q	P	e
LOG1	Flat in log	–	S12	S12	S12
LOG2	Flat in log	–	S12	SB13	Thermal
LOG3	–	Flat in log	Sorted	S12	S12
LOG4	Flat in log	–	SB13	S12	Thermal
LOG5	Flat in log	–	SB13	SB13	Thermal
KRO1	K01	–	S12	S12	S12
KRO5	K01	–	SB13	SB13	Thermal
LAR1	L98	–	S12	S12	S12
LAR5	L98	–	SB13	SB13	Thermal
TOP1	Top heavy	–	S12	S12	S12
TOP5	Top heavy	–	SB13	SB13	Thermal

Column 1 reports the model name. Column 2 describes how we generate the ZAMS mass of the primary star (i.e., the most massive of the two members of the binary system). Column 3 describes how we generate the ZAMS mass of the overall stellar population (without differentiating between primary and secondary stars). We follow this procedure only for model LOG3 (see the text for details). Columns 4, 5, and 6 specify the distributions we used to generate the mass ratios q , the orbital periods P and the orbital eccentricity e . See Section 2.2 for a detailed description of these distributions.

Santoliquido et al. 2023:

<https://arxiv.org/pdf/2303.15515.pdf>

Pop. III BBHs: mass evolution



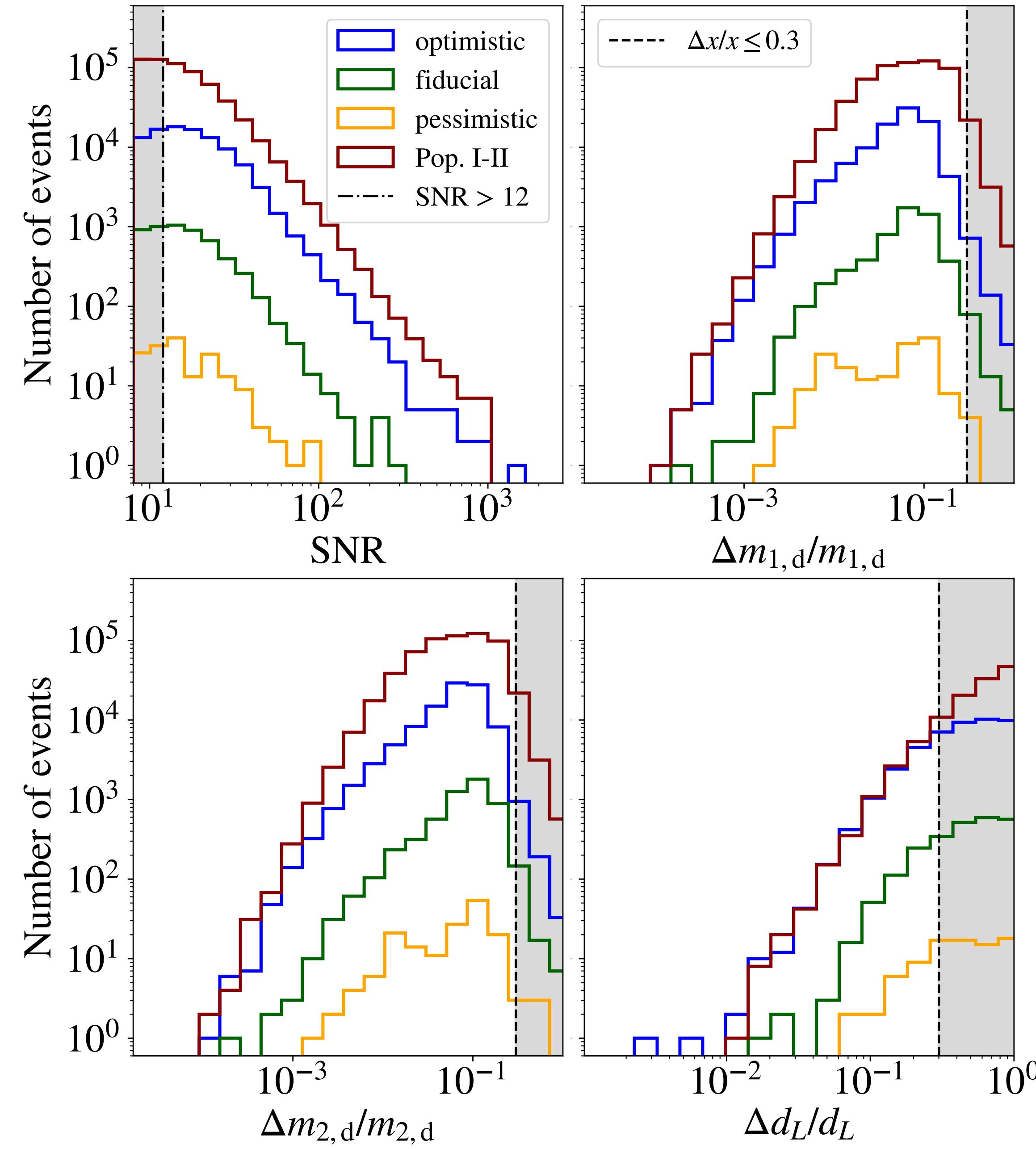
detection rate

$$\mathcal{R}_{\text{det}} = \int \frac{d^2\mathcal{R}(m_1, m_2, z)}{dm_1 dm_2} \frac{1}{(1+z)} \frac{dV_c}{dz} p_{\text{det}}(m_1, m_2, z) dm_1 dm_2 dz.$$

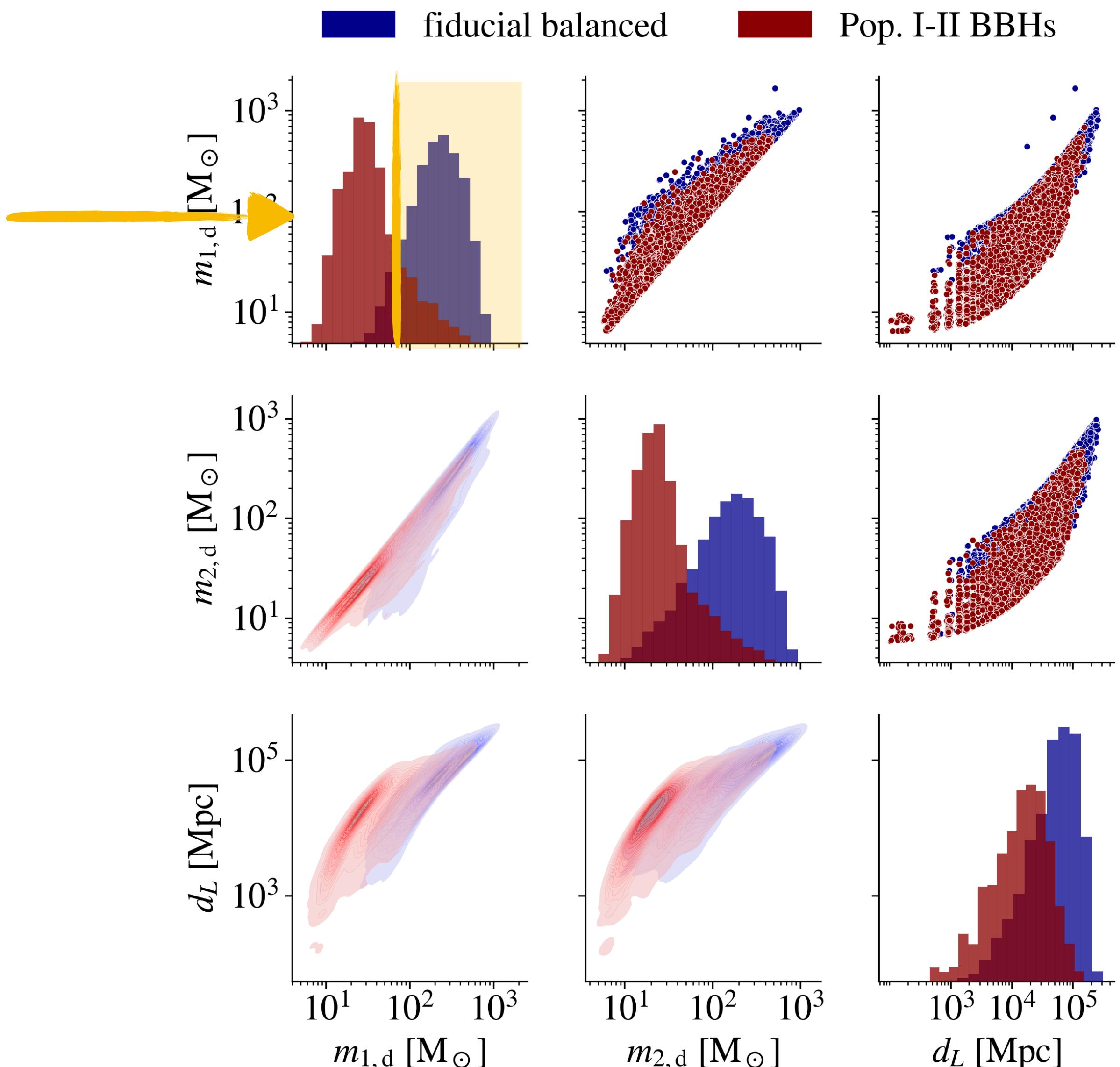
$$\frac{d^2\mathcal{R}(m_1, m_2, z)}{dm_1 dm_2} = \mathcal{R}(z) p(m_1, m_2 | z).$$

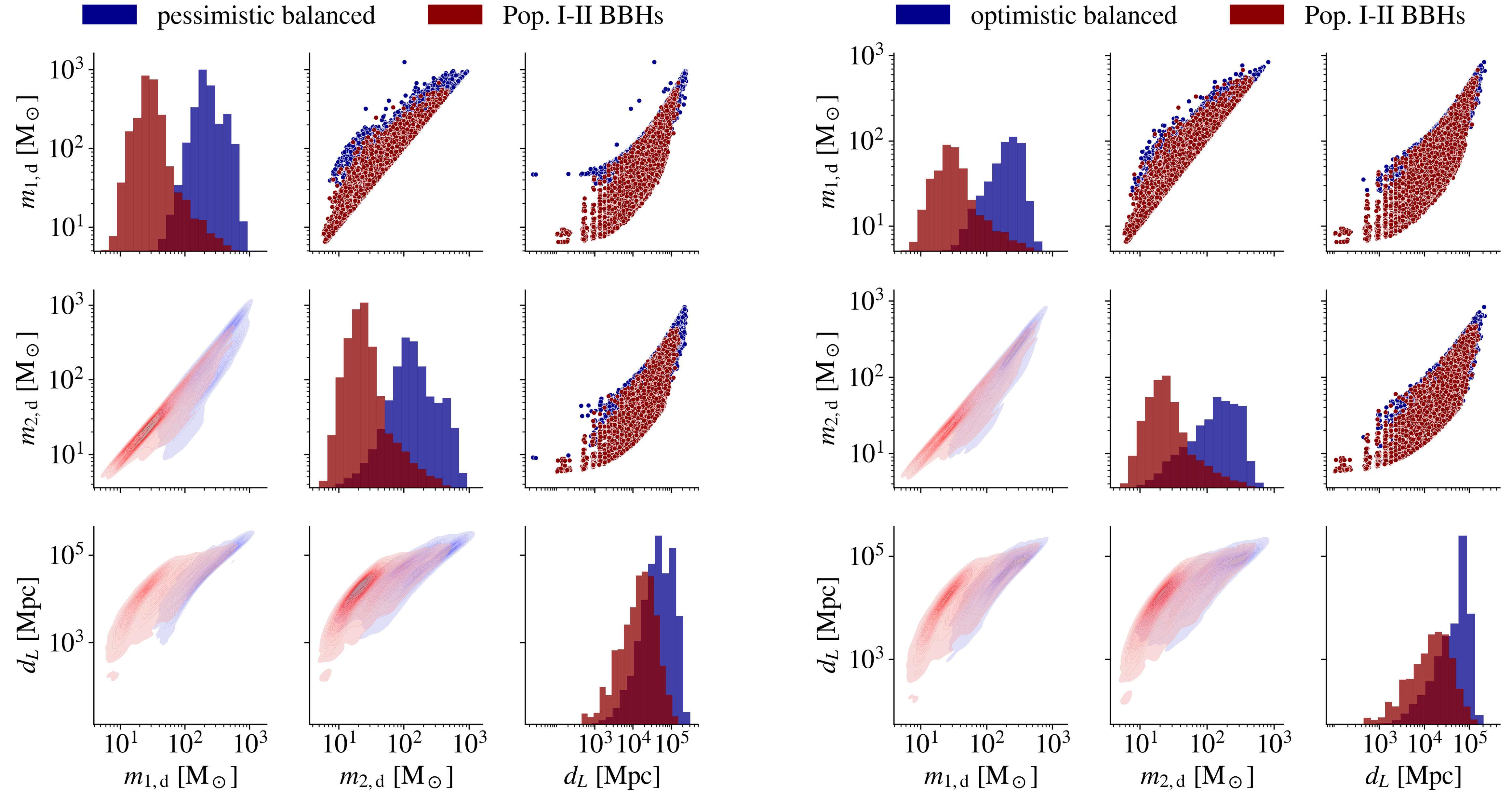
$$\rho = \rho_{\text{opt}} \sqrt{\omega_0^2 + \omega_1^2 + \omega_2^2}$$

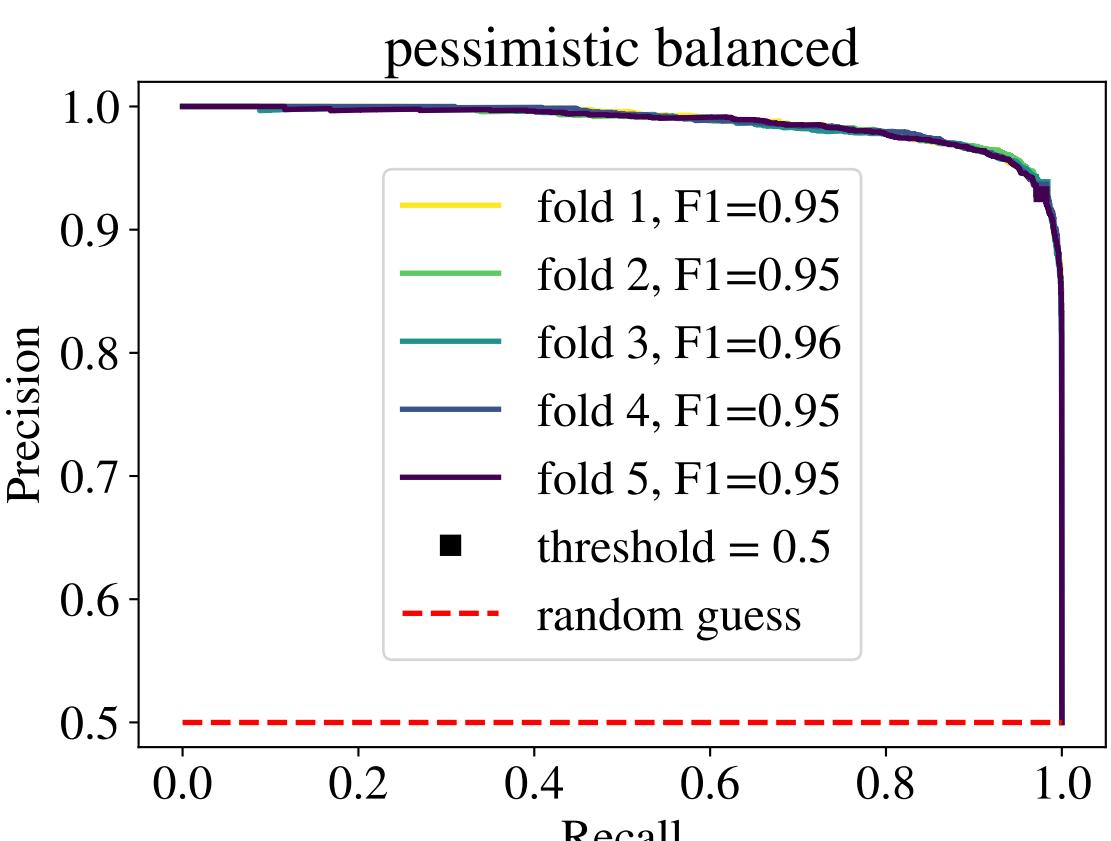
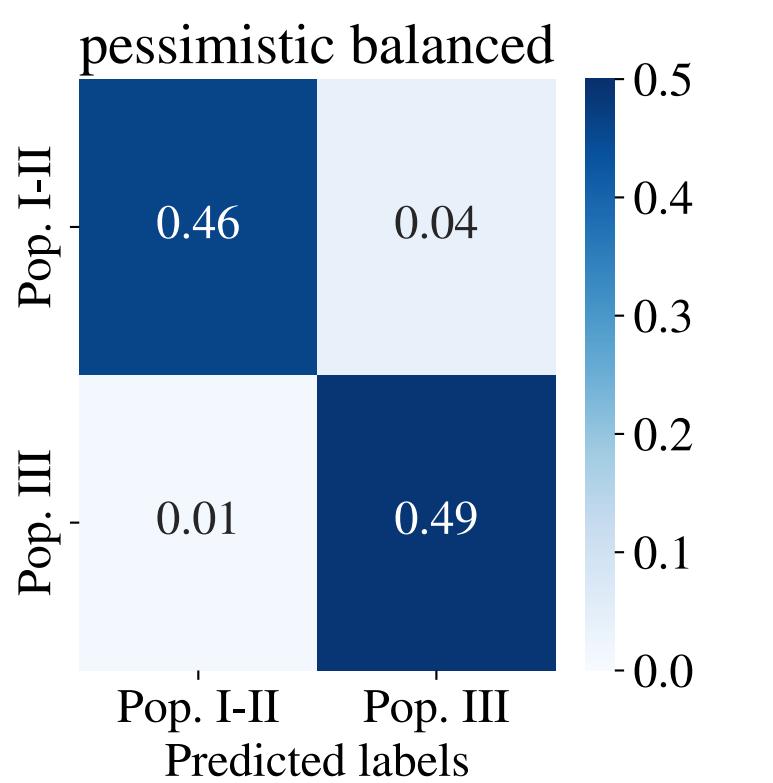
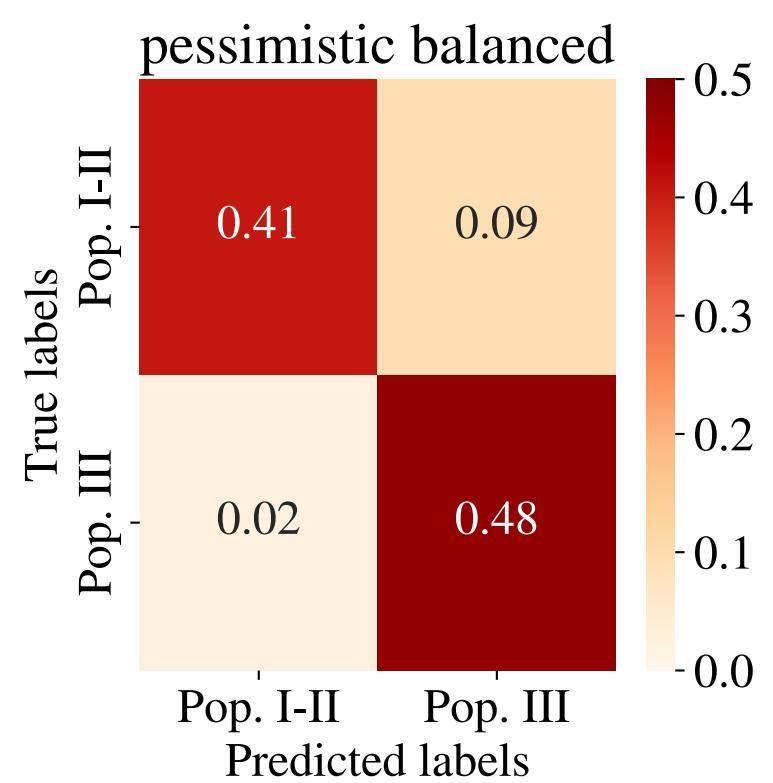
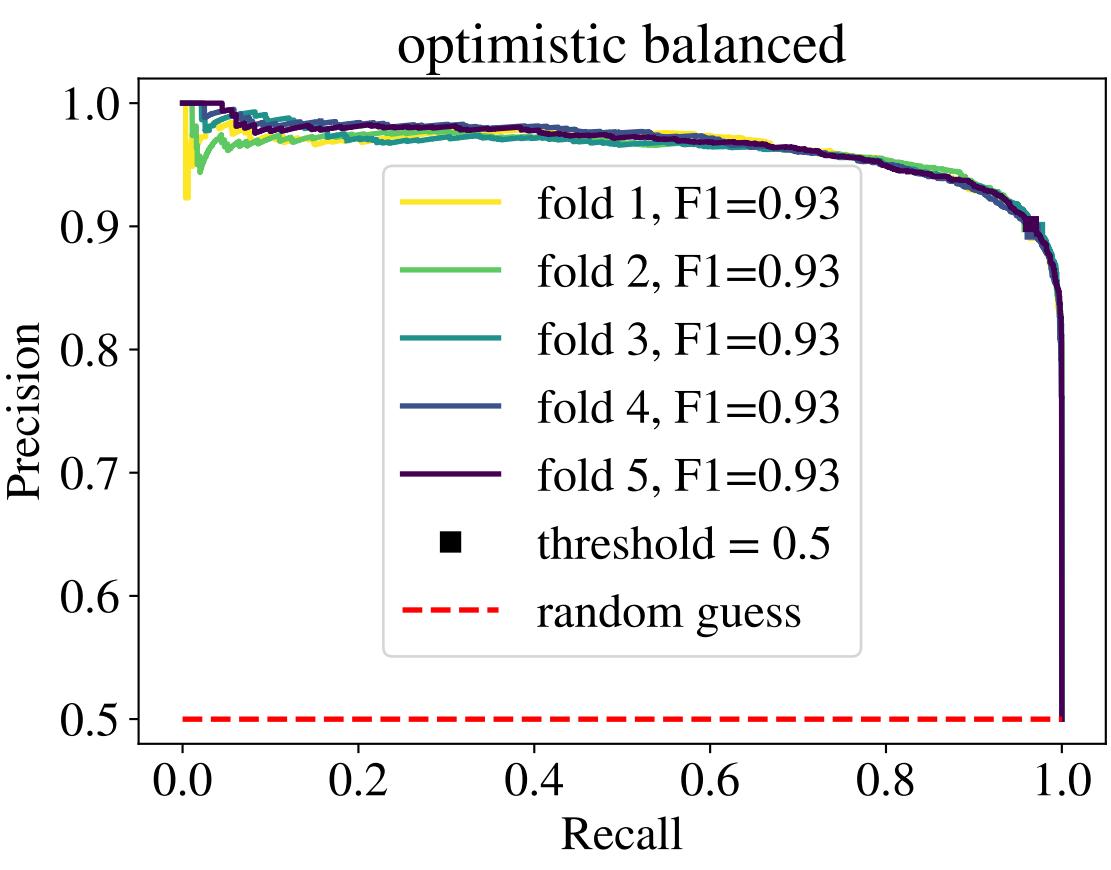
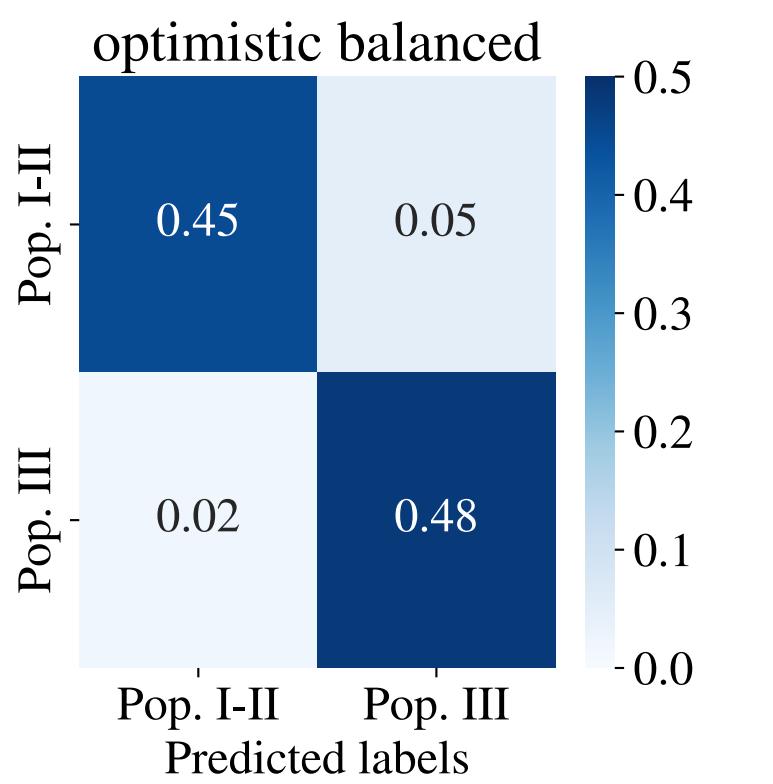
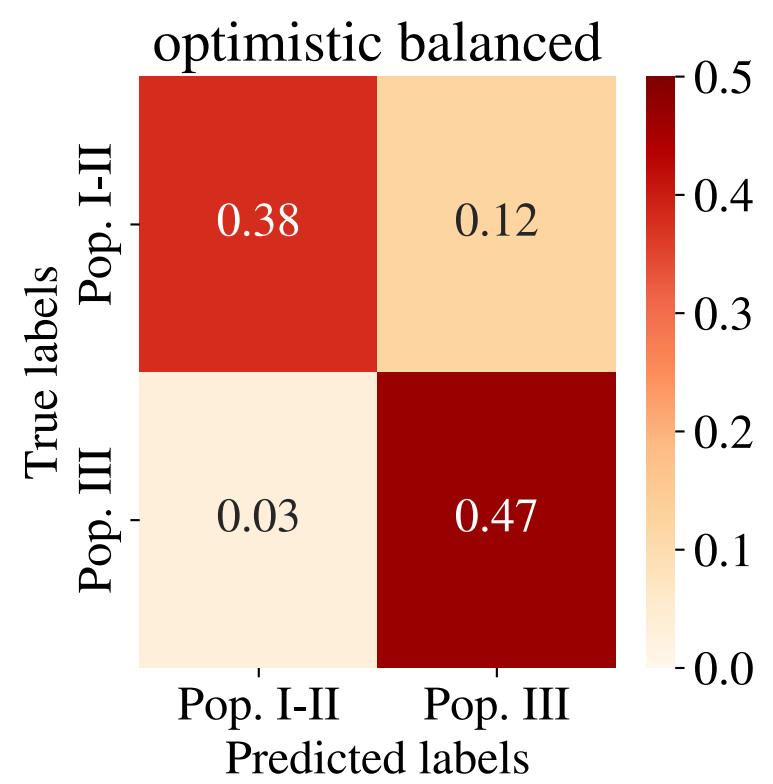
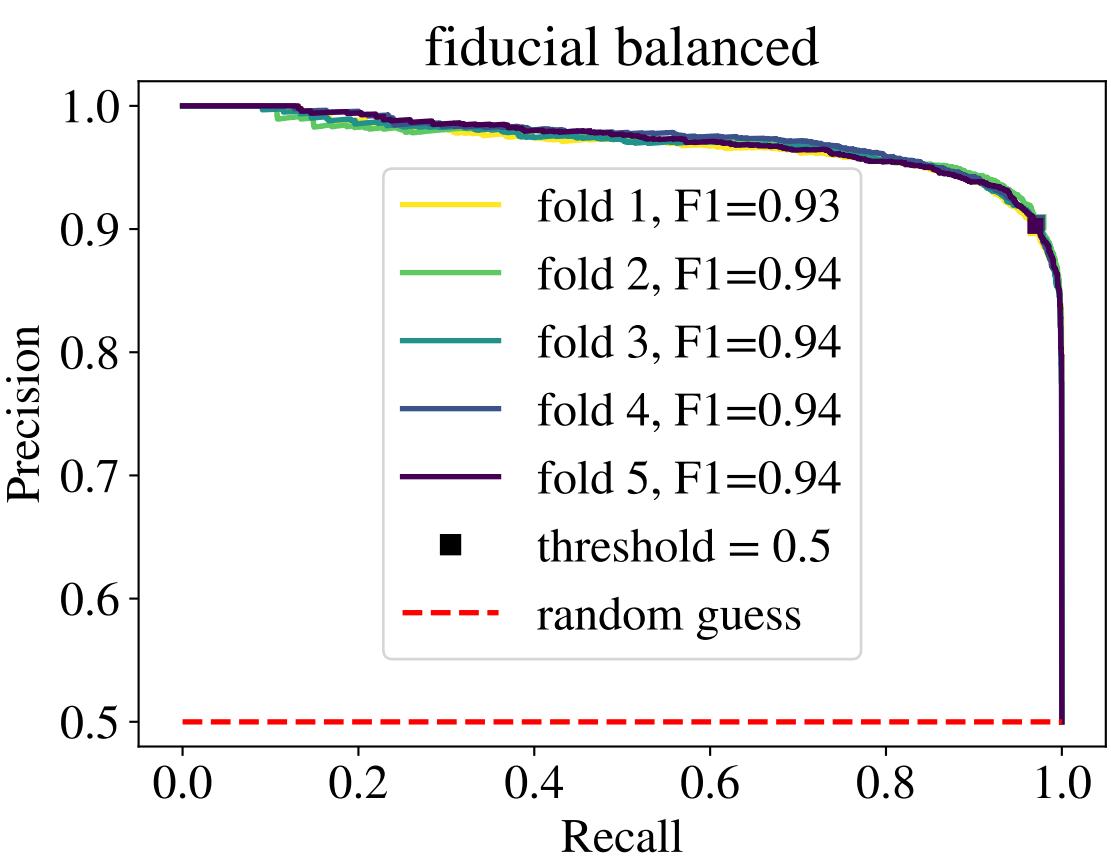
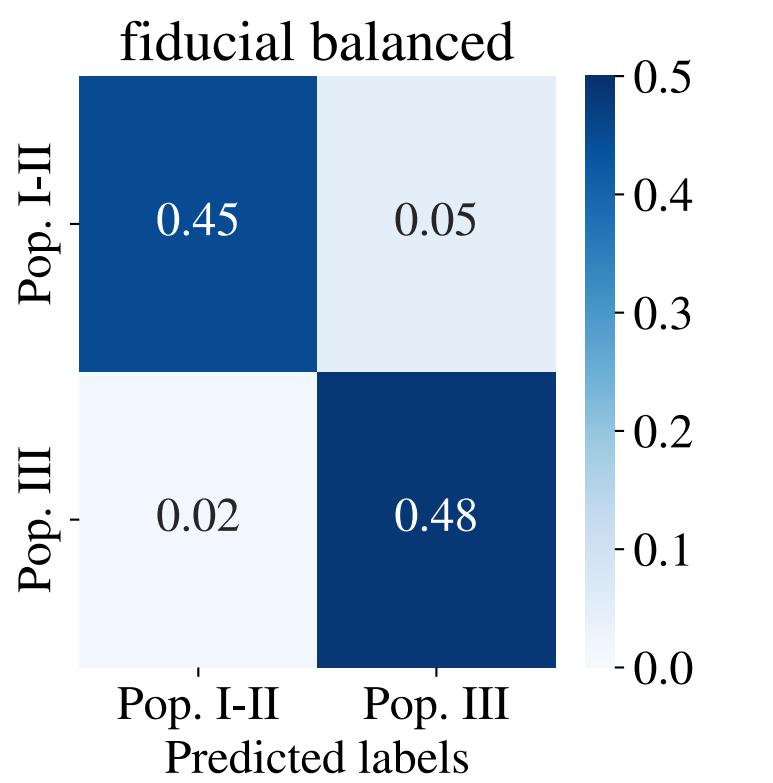
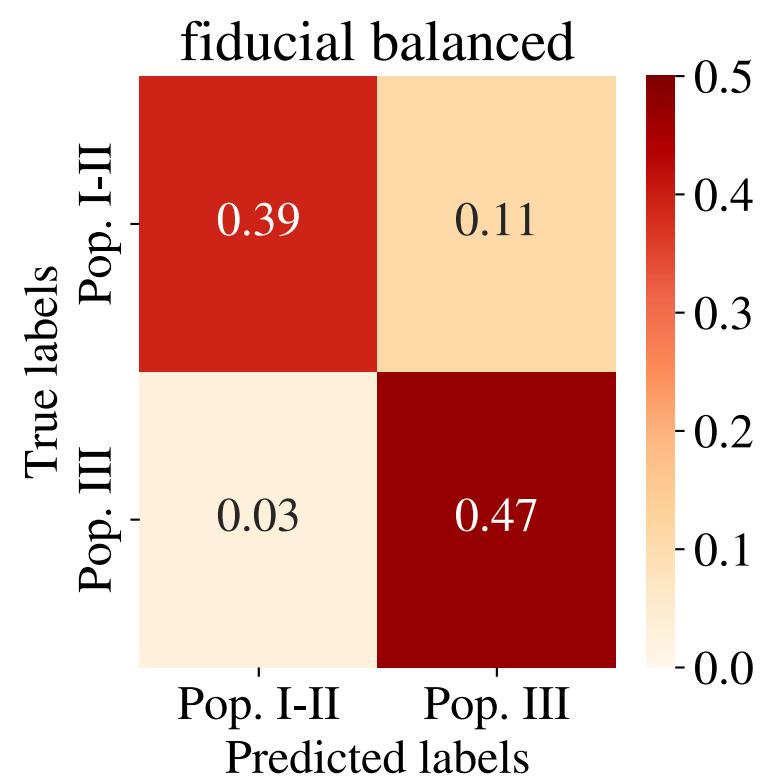
$$\rho_{\text{opt}}^2 = 4 \int_{f_{\text{low}}}^{f_{\text{high}}} df \frac{|\tilde{h}(f)|^2}{S_n(f)}$$

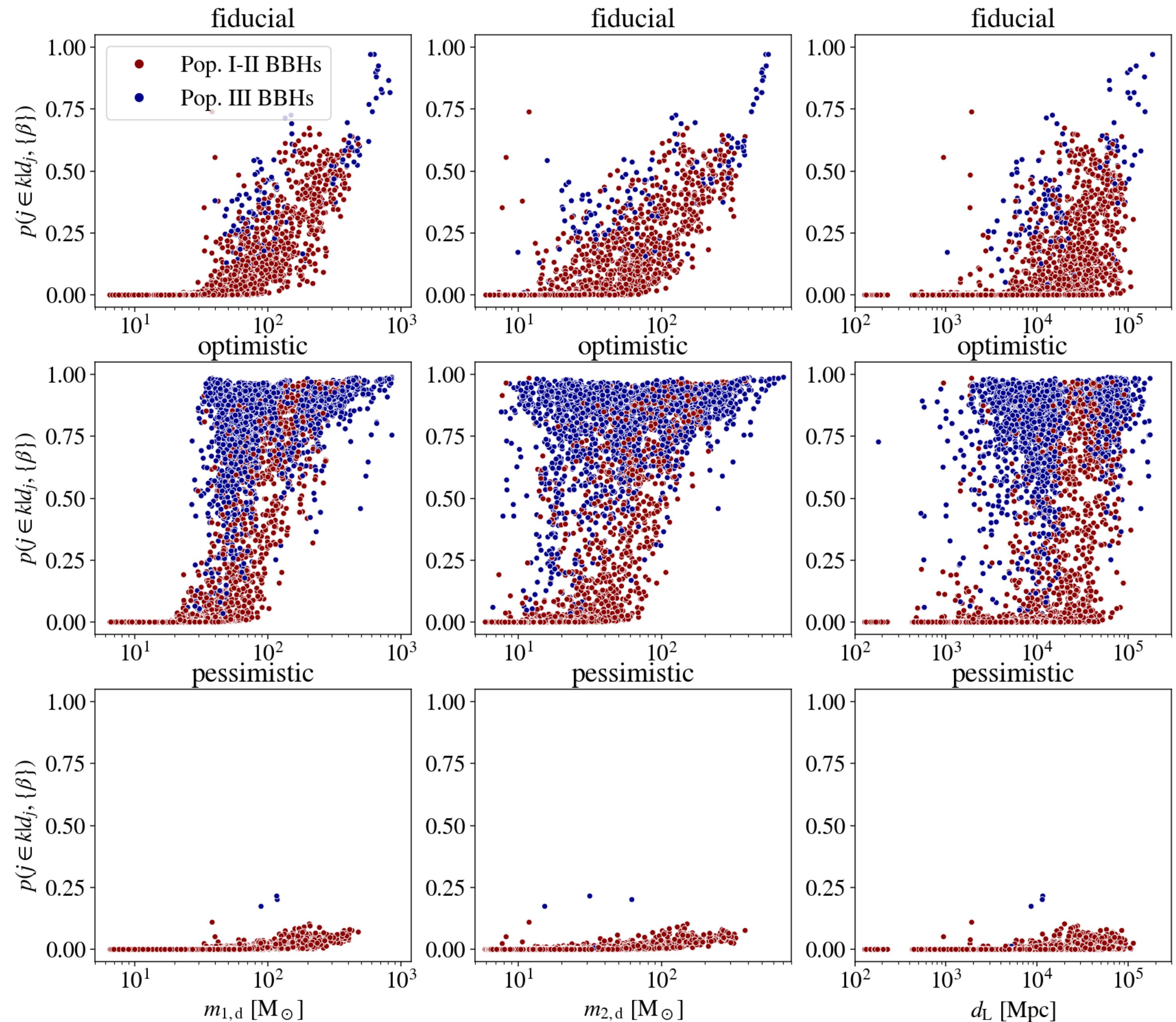


$$p(j \in k | x, d_j, \{\beta\}) = 1 \quad \text{if} \quad m_{1,d} \gtrsim 60 \text{ M}_\odot$$









Fiducial						
Thr.	%TP	%TN	%FP	%FN	Precision	Recall
0.1	96	85	15	4	0.20	0.96
0.2	86	90	10	14	0.26	0.86
0.5	33	98	2	67	0.43	0.33
0.7	11	100	0	89	0.94	0.11
0.9	3	100	0	97	1.00	0.03

Optimistic						
Thr.	%TP	%TN	%FP	%FN	Precision	Recall
0.1	100	77	23	0	0.80	1.00
0.2	99	80	20	1	0.81	0.99
0.5	95	85	15	5	0.85	0.95
0.7	87	89	11	13	0.88	0.87
0.9	46	96	4	54	0.91	0.46

Pessimistic						
Thr.	%TP	%TN	%FP	%FN	Precision	Recall
0.1	50	100	0	50	0.60	0.50
0.2	33	100	0	67	1.00	0.33
0.5	0	100	0	100	0	0
0.7	0	100	0	100	0	0
0.9	0	100	0	100	0	0