# Geodesic X-ray transform and streaking artifacts on simple surfaces or on spaces of constant curvature

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## X-ray transform on the plane

• All the planar lines are parametrized by  $(\theta, t) \in [0, \pi] \times \mathbb{R}$ :

$$\ell = \{ (-s\sin\theta + t\cos\theta, s\cos\theta + t\sin\theta) : s \in \mathbb{R} \}.$$

The X-ray transform of f(x,y) on  $\mathbb{R}^2$  is defined by

$$\mathcal{R}f(\theta,t) := \int_{\ell} f = \int_{-\infty}^{\infty} f(-s\sin\theta + t\cos\theta, s\cos\theta + t\sin\theta) ds.$$

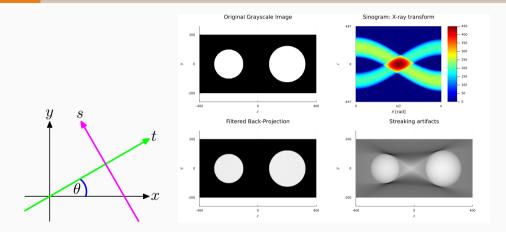
This is considered to be the measurements of CT scanners for normal tissue. The FBP formula  $f = (-\partial_x^2 - \partial_y^2)^{1/2} \circ \mathcal{R}^T \circ \mathcal{R} f$  is well-known.

We consider a model of human body f containing a metal region D such as dental
implants, stents in blood vessels, and etc. We observe that the metal streaking artifacts
caused by beam hardening effect in the energy level of X-ray. The main term is the filtered
back-projection of nonlinear term

$$(-\partial_x^2 - \partial_y^2)^{1/2} \circ \mathcal{R}^T [(\mathcal{R}1_D)^2],$$

This is a conormal distribution whose singular support is the streaking artifact.

# Figures: metal streaking artifacts



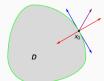
The main part of artifacts:  $(-\partial_x^2 - \partial_y^2)^{1/2} \mathcal{R}^T \left[ (\mathcal{R} \mathbf{1}_D)^2 \right]$ .

#### **Conormal distributions**

## **Definition 1 (Conormal distributions)**

Let X be an N-dim manifold, and let Y be a closed submanifold of X. We say that  $u \in \mathscr{D}'(X)$  is conormal with respect to Y of degree m if  $L_1 \cdots L_{\mu} u \in {}^{\infty}H^{loc}_{(-m-N/4)}(X)$  for all  $\mu = 0, 1, 2, \ldots$  and all vector fields  $L_1, \ldots, L_{\mu}$  tangential to Y. Denote by  $I^m(N^*(Y))$  the set of all distributions on X conormal wrt Y of degree m. Note that  $WF(u) \subset N^*(Y) \setminus 0$ .

• The characteristic function of a domain:  $1_D \in I^{-1/2-n/4}(N^*(\partial D))$  for  $D \subset \mathbb{R}^n$ , which is a domain with smooth boundary.



• The Schwartz kernel of a PsDO:  $\int_{\mathbb{R}^n} e^{i(x-y)\cdot\xi} a(x,\xi) d\xi \in I^m(N^*(\Delta)),$   $\Delta = \{(x,x)\} \text{ for } a(x,\xi) \in S^m(\mathbb{R}^n \times \mathbb{R}^n).$ 



# Geodesic X-ray transform 1

Suppose that (M,g) is a compact nontrapping simple Riemannian manifold with strictly convex smooth boundary. A map  $\pi:S(M)\to M$  is the natural projection. Denote by  $\partial_-S(M)$  the set of all unit incoming tangent vectors on the boundary  $\partial M$ :

$$\partial_{-}S(M) = \{ w \in S(M) : \pi(w) \in \partial M, \langle \nu, w \rangle < 0 \},$$

where  $\nu(x)$  is the unit outer normal vector at  $x \in \partial M$ . Note that the nontrapping condition ensures that  $\partial_- S(M)$  is identified with the manifold of all the normal geodesics on (M,g):

$$\partial_{-}S(M)\simeq\mathcal{G}:=\{\gamma_{v}:\ \nabla_{\dot{\gamma}_{v}(t)}\dot{\gamma}_{v}(t)=0,\ \dot{\gamma}_{v}(0)=v\in S(M)\}.$$

The geodesic X-ray transform of a function (more precisely a half-density) f on M is defined by

$$\mathcal{X}f(w) := \int_0^{\tau(w)} f(\gamma_w(s)) ds, \quad w \in \partial_- S(M),$$

where  $\tau(w)$  is the exit time of  $\gamma_w$ .

## Geodesic X-ray transform 2

Set  $n = \dim(M)$ . Then  $\dim(S(M)) = 2n - 1$  and  $\dim(\partial_{-}S(M)) = 2n - 2$ .

Let  $F: S(M) \to \partial_- S(M)$  be the submersion defined by  $F(\dot{\gamma}_w(t)) = w$  for  $w \in \partial_- S(M)$  and  $t \in [0, \tau(w)]$ . Then we have  $\mathcal{X} = F_* \circ \pi^*$  and  $\mathcal{X}^T = \pi_* \circ F^*$ . See Holman-Uhlmann (2018).

## **Proposition 2**

 ${\mathcal X}$  is an elliptic Fourier integral operator, and its Schwartz kernel belongs to

$$I^{-n/4}(\partial_{-}S(M)\times M^{int}, C'_{\mathcal{X}}; \Omega^{1/2}_{\partial_{-}S(M)\times M^{int}}),$$

where  $C_{\mathcal{X}}$  is the canonical relation of  $\mathcal{X}$ : we say that  $(\xi, \eta) \in C_{\mathcal{X}}$  if  $\exists v \in S(M^{int})$  such that

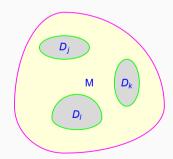
$$\xi \in T^*_{F(v)}\big(\partial_-S(M)\big)\setminus\{0\}, \quad \eta \in T^*_{\pi(v)}(M^{int})\setminus\{0\}, \quad DF|_v^T\xi = D\pi|_v^T\eta,$$

## Assumption 1

• Assume that dim(M) = 2 or (M, g) is a space of constant curvature.

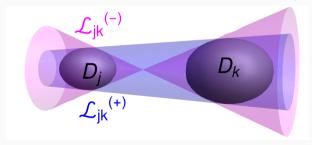
This ensures that all the Jacobi fields are of the form scalar function  $\times$  parallel transport.

• Suppose that the metal region  $D \subset M^{\text{int}}$  is a disjoint union of  $D_j$   $(j=1,\ldots,J)$  which are simply connected, strictly convex and bounded with smooth boundaries  $\partial D_j$ .



# A hypersurface $\mathscr L$ surrounding the metal region D

- For any j and  $x \in \partial D_j$ , denote by  $v_j(x)$  the unit outer normal vector at x. Consider the tangent hyperplane  $\exp_x v_j(x)^{\perp} \cap M^{\text{int}}$  at  $x \in \partial D_j$ .
- There are some common tangent hyperplanes of  $\partial D_j$  and  $\partial D_k$  for  $j \neq k$ . In this case there is common tangent geodesics in such hyperplanes. The union of all these geodesics forms a conical or cylindrical hypersurface denoted by  $\mathscr{L}_{jk}^{(\pm)}$ . Set  $\mathscr{L} := \bigcup \left( \mathscr{L}_{jk}^{(+)} \cup \mathscr{L}_{jk}^{(-)} \right)$ .



# Assumption 2 (The simple model of beam hardening effect)

Let  $E \ge 0$  be a parameter describing the energy level of the X-ray beam, and let  $E_0$  be the fixed standard level for the normal tissue. The measurement P is of the form:

$$P = -\log \left\{ \int_0^\infty 
ho(E) \exp(-\mathcal{X} f_E) dE, \, 
ight\},$$

where  $\rho(E)$  is a probability density function on  $[0, \infty)$  and is called the spectral function. Let  $f_{CT}$  be the FBP of P. We employ the simple model of the form

$$f_E(x) = f_{E_0}(x) + \alpha(E - E_0)1_D(x), \quad \rho(E) = \frac{1}{2\varepsilon}1_{[E_0 - \varepsilon, E_0 + \varepsilon]}(E)$$

with small parameters  $\alpha > 0$  and  $\varepsilon > 0$ .

#### Main Theorem

Then the nonlinear effect  $f_{MA}$  in the CT image becomes

$$f_{\mathsf{MA}} := f_{\mathsf{CT}} - f_{\mathsf{E}_0} = \sum_{k=1}^{\infty} (\alpha \varepsilon)^{2k} A_k Q \mathcal{X}^{\mathsf{T}} [(\mathcal{X}1_D)^{2k}] \mod C^{\infty}(M^{\mathsf{int}}), \quad \{A_k\} \subset \mathbb{R},$$

where Q is a parametrix of  $\mathcal{X}^T \circ \mathcal{X}$ :  $Q\mathcal{X}^T \mathcal{X} = Id$  modulo smoothing operators locally.

Our main result is as follows:

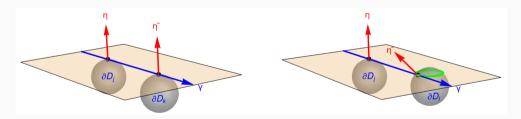
#### Theorem 3

$$f_{MA} \in I^{-3n/4-1/2}ig(N^*(\mathscr{L})ig)$$
 away from  $\partial D$ , and  $\sigma_{prin}ig(Q\mathcal{X}^T[(\mathcal{X}1_D)^2]ig) 
eq 0$ .

- Park-Choi-Seo (2017) proved that WF $(f_{MA}) \subset N^*(\mathscr{L})$  for  $M = \mathbb{R}^2$ .
- Palacios-Uhlmann-Wang (2018) proved Theorem 3 for  $M = \mathbb{R}^2$ .
- C (2022) proved Theorem 3 for the d-plane transform on  $\mathbb{R}^n$ . We could NOT understand the meaning in many parts of this paper.

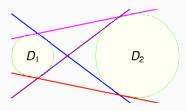
## What does Theorem 3 say?

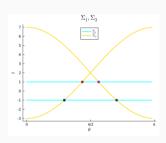
- If  $\partial D_j$  and  $\partial D_k$  have a common tangent hyperplane, then the conormal singularities propagate along the common tangent geodesic. See the left figure.
- Suppose  $n \ge 3$ . If  $\partial D_j$  and  $\partial D_k$  have a common tangent geodesic, but the conormal directions at the tangent points are different, then the conormal singularities do not propagate along the common tangent geodesic. See the right figure.



# Outline of the proof of Theorem 3

- $1_{D_j} \in I^{-1/2-n/4}(N^*(\partial D_j) \setminus 0).$
- $\mathcal{X}1_{D_j} \in I^{-(n+1)/2}(N^*(\Sigma_j) \setminus 0)$  with some hypersurface  $\Sigma_j$  in  $\partial_-S(M)$ .
- For  $j \neq k$ ,  $\Sigma_j$  is transversal to  $\Sigma_k$ .





ullet Set  $\Sigma_{jk}:=\Sigma_j\cap\Sigma_k$  for short. For  $j{
eq}k$ ,

$$\mathcal{X}1_{D_j}\cdot\mathcal{X}1_{D_k}\in \begin{cases} I^{-(n+1)/2-1}\big(N^*(\Sigma_{jk})\setminus 0\big) & \text{at }\Sigma_{jk},\\ I^{-(n+1)/2}\big(N^*(\Sigma_j)\setminus 0\big)+I^{-(n+1)/2}\big(N^*(\Sigma_k)\setminus 0\big) & \text{away from }\Sigma_{jk}. \end{cases}$$

**Key:** 
$$C_{\mathcal{X}^T} \circ \Sigma_{jk} \setminus 0 = N^*(\mathcal{L}_{jk}) \setminus 0$$

- Fix arbitrary geodesic  $\gamma_w \simeq w \in \Sigma_{jk}$ .
- If  $\xi, \tilde{\xi} \in T_w^*(\partial_- S(M))$ ,  $w = F(v) = F(\tilde{v})$ ,  $\pi(v) \in \partial D_j$ ,  $\pi(\tilde{v}) \in \partial D_k$ ,  $DF|_v^T \xi = D\pi|_v^T \eta, \quad \eta \in N_v^*(\partial D_j) \setminus \{0\}, \quad DF|_{\tilde{v}}^T \tilde{\xi} = D\pi|_{\tilde{v}}^T \tilde{\eta}, \quad \tilde{\eta} \in N_{\tilde{v}}^*(\partial D_k) \setminus \{0\},$

then  $\xi$  and  $\tilde{\xi}$  are linearly independent, and the nonlinear effect on the geodesic  $\gamma_w$  creates two-dimensional singularity span $\langle \xi, \tilde{\xi} \rangle$  in  $T_w^*(\partial_- S(M))$  due to the simplicity condition.

- WLOG WMA  $\eta$  and  $\tilde{\eta}$  are unit covectors.
- WLOG WMA  $\eta$  is the parallel transport of  $\tilde{\eta}$  if  $\eta \parallel \tilde{\eta}$ .
- We shall show that if  $\tilde{\eta}$  is the parallel transport of  $\eta$ , then

$$C_{\mathcal{X}}^{T} \circ \operatorname{span} \langle \xi, \tilde{\xi} \rangle = \bigcup_{\mathbf{a} \in \mathbb{R}} (\text{the parallel transport of } \mathbf{a} \boldsymbol{\eta} \text{ along } \gamma_{w}) = \bigcup_{t \in [0, \tau(w)]} N_{\gamma_{w}(t)}^{*}(\mathcal{L}_{jk}),$$

otherwise, 
$$C_{\mathcal{X}}^T \circ \operatorname{span}\langle \xi, \tilde{\xi} \rangle = N_{\pi(v)}^*(\partial D_j) \bigcup N_{\pi(\tilde{v})}^*(\partial D_k).$$

# When $\tilde{\eta}$ is the parallel transport of $\eta$

- Set  $\gamma_w(t_0) = \pi(\eta) \in \partial D_j$  and  $\gamma_w(\tilde{t}_0) = \pi(\tilde{\eta}) \in \partial D_k$ , and suppose  $\tilde{\eta} = P(\tilde{t}_0, t_0; \gamma_w)^T \eta$ , where  $P(t_0, \tilde{t}_0; \gamma_w)$  is the parallel transport of  $T_{\gamma_w(\tilde{t}_0)}(M)$  onto  $T_{\gamma_w(t_0)}(M)$  along  $\gamma_w$ . Set  $\eta(s) := P(s, t_0; \gamma_w)^T \eta \in T^*_{\gamma_w(s)}(M^{\text{int}})$  for  $s \in (0, \tau(w))$ . Then  $\eta(\tilde{t}_0) = \tilde{\eta}$ .
- Let k(x) be a sectional curvature at  $x \in M$ , which is a constant when  $n \ge 3$ .
- Let  $a(t;s), b(t;s) \in C^{\infty}(0,\tau(w))$  be solutions to

$$a_{tt}(t;s) + k(\gamma_w(t))a(t;s) = 0, \quad a(s;s) = 1, \quad a_t(s;s) = 0,$$
  $b_{tt}(t;s) + k(\gamma_w(t))b(t;s) = 0, \quad b(s;s) = 0, \quad b_t(s;s) = 1.$ 

•  $\Delta(t_0, \tilde{t}_0; s) := \det \begin{bmatrix} a(t_0; s) & a(\tilde{t}_0; s) \\ b(t_0; s) & b(\tilde{t}_0; s) \end{bmatrix}$  never vanish due to the simplicity. If we set

$$\xi(s) := \frac{b(\tilde{t}_0; s)}{\Delta(t_0, \tilde{t}_0; s)} \xi - \frac{b(t_0; s)}{\Delta(t_0, \tilde{t}_0; s)} \tilde{\xi} \in \operatorname{span}\langle \xi, \tilde{\xi} \rangle, \quad s \in (0, \tau(w)),$$

then we have  $DF|_{\dot{\gamma}_{W}(s)}^{T}\xi(s)=D\pi|_{\dot{\gamma}_{W}(s)}^{T}\eta(s)$  in  $T_{\dot{\gamma}_{W}(s)}^{*}\big(S(M^{\mathrm{int}})\big)$  for  $s\in(0,\tau(w)).$ 

## References

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