

I. VACUUM BOUNDARY CONDITION

A. Toroidal Coordinates

Right-handed orthogonal curvilinear coordinates: η, μ, ϕ . Here, ϕ is toroidal angle.

$$R = \frac{\sinh \mu}{\cosh \mu - \cos \eta}, \quad (1)$$

$$Z = \frac{\sin \eta}{\cosh \mu - \cos \eta}, \quad (2)$$

$$d_1 = [(R + 1)^2 + Z^2]^{1/2}, \quad (3)$$

$$d_2 = [(R - 1)^2 + Z^2]^{1/2}, \quad (4)$$

$$\cosh \mu = \frac{1}{2} \left(\frac{d_1}{d_2} + \frac{d_2}{d_1} \right), \quad (5)$$

$$\cos \eta = \left(\frac{d_1^2 + d_2^2 - 4}{2 d_1 d_2} \right). \quad (6)$$

Scale factors:

$$h_\eta = h_\mu = h = \frac{1}{\cosh \mu - \cos \eta}, \quad (7)$$

$$h_\phi = R = \frac{\sinh \mu}{\cosh \mu - \cos \eta}. \quad (8)$$

B. Magnetic Perturbation

$$\delta \mathbf{B} = \mathbf{i} \nabla V, \quad (9)$$

where

$$\nabla^2 V = 0. \quad (10)$$

Let

$$V(\eta, \mu, \phi) = \hat{V}(\eta, \mu) e^{-i n \phi}, \quad (11)$$

and

$$\delta \hat{\mathbf{B}} = \mathbf{i} \nabla \hat{V}. \quad (12)$$

So,

$$\delta \hat{B}_\eta = \frac{i}{h} \frac{\partial \hat{V}}{\partial \eta}, \quad (13)$$

$$\delta \hat{B}_\mu = \frac{i}{h} \frac{\partial \hat{V}}{\partial \mu}, \quad (14)$$

$$\delta \hat{B}_\phi = \frac{n}{R} \hat{V}. \quad (15)$$

C. Current Sheet

Suppose that there is a current sheet at $\mu = \mu_w$. Let

$$\delta \hat{\mathbf{J}} = \int_{\mu_{w-}}^{\mu_{w+}} \delta \hat{\mathbf{j}} h d\mu, \quad (16)$$

where

$$\delta \hat{\mathbf{j}} = \nabla \times \delta \hat{\mathbf{B}}. \quad (17)$$

Follows that

$$\delta \hat{J}_\eta = [\delta \hat{B}_\phi]_{\mu_{w-}}^{\mu_{w+}} = \frac{n}{R} [\hat{V}]_{\mu_{w-}}^{\mu_{w+}}, \quad (18)$$

$$\delta \hat{J}_\mu = 0, \quad (19)$$

$$\delta \hat{J}_\phi = -[\delta \hat{B}_\eta]_{\mu_{w-}}^{\mu_{w+}} = -\frac{i}{h} \left[\frac{\partial \hat{V}}{\partial \eta} \right]_{\mu_{w-}}^{\mu_{w+}}. \quad (20)$$

Of course,

$$\left[\frac{\partial \hat{V}}{\partial \mu} \right]_{\mu_{w-}}^{\mu_{w+}} = 0. \quad (21)$$

D. Toroidal Torque

Now,

$$\overline{(\delta \mathbf{J} \times \delta \mathbf{B})}_\phi = \overline{\delta J_\eta \delta B_\mu} = -\frac{i n}{2 R h} \left\{ [\hat{V}] \frac{\partial \hat{V}^*}{\partial \mu} - [\hat{V}^*] \frac{\partial \hat{V}}{\partial \mu} \right\}. \quad (22)$$

Thus, net toroidal torque acting on sheet is

$$\hat{T}_\phi = 2\pi \oint R \overline{(\delta \mathbf{J} \times \delta \mathbf{B})}_\phi h R d\eta = -i n \pi \oint R \left\{ [\hat{V}] \frac{\partial \hat{V}^*}{\partial \mu} - [\hat{V}^*] \frac{\partial \hat{V}}{\partial \mu} \right\} d\eta. \quad (23)$$

Let

$$z = \cosh \mu. \quad (24)$$

Follows that

$$\hat{T}_\phi = -i n \pi (z^2 - 1) \oint (z - \cos \eta)^{-1} \left\{ [\hat{V}] \frac{\partial \hat{V}^*}{\partial z} - [\hat{V}^*] \frac{\partial \hat{V}}{\partial z} \right\} d\eta. \quad (25)$$

Now,

$$[\hat{V}] = \hat{V}_+ - \hat{V}_-, \quad (26)$$

$$\frac{\partial \hat{V}_+}{\partial \mu} = \frac{\partial \hat{V}_-}{\partial \mu}. \quad (27)$$

So,

$$\hat{T}_\phi = -i n \pi (z^2 - 1) \oint (z - \cos \eta)^{-1} \left\{ \hat{V}_+ \frac{\partial \hat{V}_+^*}{\partial z} - \hat{V}_+^* \frac{\partial \hat{V}_+}{\partial z} - \hat{V}_- \frac{\partial \hat{V}_-^*}{\partial z} + \hat{V}_-^* \frac{\partial \hat{V}_-}{\partial z} \right\} d\eta. \quad (28)$$

Let

$$\hat{V}(z, \eta) = (z - \cos \eta)^{1/2} v(z, \eta). \quad (29)$$

It follows that

$$\hat{T}_\phi = -i n \pi (z^2 - 1) \oint \left\{ v_+ \frac{\partial v_+^*}{\partial z} - v_+^* \frac{\partial v_+}{\partial z} - v_- \frac{\partial v_-^*}{\partial z} + v_-^* \frac{\partial v_-}{\partial z} \right\} d\eta. \quad (30)$$

Let

$$v(z, \eta) = \sum_m v_m(z) \cos(m \eta). \quad (31)$$

It follows that

$$\hat{T}_\phi = -i n \pi^2 (z^2 - 1) \sum_m \left\{ v_{m+} \frac{dv_{m+}^*}{dz} - v_{m+}^* \frac{dv_{m+}}{dz} - v_{m-} \frac{dv_{m-}^*}{dz} + v_{m-}^* \frac{dv_{m-}}{dz} \right\}. \quad (32)$$

If

$$v_{m+} = p_m P_{m-1/2}^n(z) + q_m Q_{m-1/2}^n(z), \quad (33)$$

$$v_{m-} = r_m P_{m-1/2}^n(z), \quad (34)$$

then

$$\hat{T}_\phi = -i n \pi^2 (z^2 - 1) \sum_m (q_m p_m^* - q_m^* p_m) \left[\frac{dP_{m-1/2}^n}{dz} Q_{m-1/2}^n - P_{m-1/2}^n \frac{dQ_{m-1/2}^n}{dz} \right]. \quad (35)$$

However,

$$\frac{dP_{m-1/2}^n}{dz} Q_{m-1/2}^n - P_{m-1/2}^n \frac{dQ_{m-1/2}^n}{dz} = \frac{\Gamma(m+1/2+n)}{\Gamma(m+1/2-n)} \frac{(-1)^n}{1-z^2}, \quad (36)$$

so

$$\begin{aligned} \hat{T}_\phi &= i n \pi^2 (-1)^n \sum_m (q_m p_m^* - q_m^* p_m) \frac{\Gamma(m+1/2+n)}{\Gamma(m+1/2-n)} \\ &= -2 n \pi^2 (-1)^n \sum_m \text{Im}(q_m p_m^*) \frac{\Gamma(m+1/2+n)}{\Gamma(m+1/2-n)}. \end{aligned} \quad (37)$$

Thus, the torque on the plasma is $T_\phi = -\hat{T}_\phi$, giving

$$T_\phi = 2 n \pi^2 (-1)^n \sum_m \text{Im}(q_m p_m^*) \frac{\Gamma(m+1/2+n)}{\Gamma(m+1/2-n)}. \quad (38)$$

E. Matching to Vacuum Solution

In vacuum the solution vector $\tilde{\mathbf{Y}}_k^e$ has the expansion

$$\hat{V}_k^e = \sum_j p_{k,j}^e P_j(\mu, \eta), \quad (39)$$

where

$$P_j(\mu, \eta) = (\cosh \mu - \cos \eta)^{1/2} P_{m_j-1/2}^n(\cosh \mu) \cos(m_j \eta). \quad (40)$$

Likewise, the solution vector $\tilde{\mathbf{Y}}_k^o$ has the expansion

$$\hat{V}_k^o = \sum_j p_{k,j}^o P_j(\eta, \mu). \quad (41)$$

Finally, the $\tilde{\mathbf{Y}}^x$ solution vector has the expansion

$$\hat{V}^x = \sum_j [p_j^x P_j(\eta, \mu) + q_j^x Q_j(\eta, \mu)], \quad (42)$$

where

$$Q_j(\mu, \eta) = (\cosh \mu - \cos \eta)^{1/2} Q_{m_j-1/2}^n(\cosh \mu) \cos(m_j \eta). \quad (43)$$

The vacuum error-field has the expansion

$$\hat{V}^v = \sum_j q_j^x Q_j(\eta, \mu). \quad (44)$$

The general solution vector is

$$\mathbf{Y} = \sum_k (\Psi_k^e \mathbf{Y}_k^e + \Psi_k^o \mathbf{Y}_k^o) + \mathbf{Y}^x, \quad (45)$$

and has the expansion

$$\hat{V} = \sum_k \left(\Psi_k^e \hat{V}_k^e + \Psi_k^o \hat{V}_k^o \right) + \hat{V}_x. \quad (46)$$

It follows that

$$\hat{V} = \sum_j (p_j P_j + q_j Q_j), \quad (47)$$

where

$$p_j = p_j^x + \sum_k (\Psi_k^e p_{k,j}^e + \Psi_k^o p_{k,j}^o), \quad (48)$$

$$q_j = q_j^x. \quad (49)$$

Thus,

$$\text{Im}(q_j p_j^*) = \text{Im} \sum_k (\Psi_k^{e*} p_{k,j}^{e*} q_j^x + \Psi_k^{o*} p_{k,j}^{o*} q_j^x). \quad (50)$$

Here, we have assumed that

$$\text{Im}(p_j^{x*} q_j^x) = 0. \quad (51)$$

Thus,

$$T_\phi = 2 n \pi^2 (-1)^n \text{Im} \sum_k \sum_j (\Psi_k^{e*} p_{k,j}^{e*} q_j^x + \Psi_k^{o*} p_{k,j}^{o*} q_j^x) \frac{\Gamma(m_j + 1/2 + n)}{\Gamma(m_j + 1/2 - n)}. \quad (52)$$

However,

$$T_\phi = -2 n \pi^2 \text{Im} \sum_k (\Psi_k^{e*} \chi_k^e + \Psi_k^{o*} \chi_k^o). \quad (53)$$

It follows that

$$\chi_k^e = (-1)^{n+1} \sum_j p_{k,j}^{e*} q_j^x \frac{\Gamma(m_j + 1/2 + n)}{\Gamma(m_j + 1/2 - n)}, \quad (54)$$

$$\chi_k^o = (-1)^{n+1} \sum_j p_{k,j}^{o*} q_j^x \frac{\Gamma(m_j + 1/2 + n)}{\Gamma(m_j + 1/2 - n)}. \quad (55)$$

Finally, let

$$\hat{V}^v = \sum_j q_j^v (-1)^{n+1} \frac{\Gamma(m_j + 1/2 - n)}{\Gamma(m_j + 1/2 + n)} Q_j(\mu, \eta). \quad (56)$$

It follows that

$$\chi_k^e = \sum_j p_{k,j}^{e*} q_j^v, \quad (57)$$

$$\chi_k^o = \sum_j p_{k,j}^{o*} q_j^v. \quad (58)$$