## Advanced Numerical Methods for PDEs

Homework 3

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Consider the 2D problem

$$\partial_t u(x, y, t) + \partial_x f^x + \partial_y f^y = 0, \quad x, y \in [0, 1]$$
$$f^x = u - \partial_x u, \quad f^y = -\partial_y u$$

with boundary conditions

$$f^x(x = 0, y, t) = 0,$$
  $f^x(x = 1, y, t) = 0,$   
 $f^y(x, y = 0, t) = 0,$   $f^y(x, y = 1, t) = 0,$ 

and initial conditions

$$u(x, y, t = 0) = uI(x, y) = \delta(x - 0.5)\delta(y - 0.5)$$

### Problem 1

1. Write a code which solves the problem using the implicit Cranck - Nicholson scheme (the trapezoidal rule in time), using an  $M \times N$  grid with  $\Delta x = \Delta y$ . Explain how you solve the linear systems, using sparse Gauss elimination.

**Solution:** We can represent the PDE as

$$\frac{\partial u}{\partial t} + \frac{\partial u}{\partial x} - \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}\right) = 0$$

Next, we discretize the equation using Crank-Nicholson in time, central differences

in space:

$$\begin{split} \frac{u_{i,j}^{n+1} - u_{i,j}^n}{\Delta t} + \frac{1}{2} \left( \frac{u_{i+1,j}^{n+1} - u_{i-1,j}^{n+1}}{2\Delta x} + \frac{u_{i+1,j}^n - u_{i-1,j}^n}{2\Delta x} \right) \\ - \frac{1}{2} \left( \frac{u_{i+1,j}^{n+1} - 2u_{i,j}^{n+1} + u_{i-1,j}^{n+1}}{\Delta x^2} + \frac{u_{i+1,j}^n - 2u_{i,j}^n + u_{i-1,j}^n}{\Delta x^2} \right) \\ - \frac{1}{2} \left( \frac{u_{i,j+1}^{n+1} - 2u_{i,j}^{n+1} + u_{i,j-1}^{n+1}}{\Delta y^2} + \frac{u_{i,j+1}^n - 2u_{i,j}^n + u_{i,j-1}^n}{\Delta y^2} \right) = 0 \end{split}$$

Regrouping the n+1 terms on the left and the n terms on the right we obtaind

$$u_{i,j}^{n+1} + \frac{C}{4} \left( u_{i+1,j}^{n+1} - u_{i-1,j}^{n+1} \right) - \frac{\mu_x}{2} \left( u_{i+1,j}^{n+1} - 2u_{i,j}^{n+1} + u_{i-1,j}^{n+1} \right) - \frac{\mu_y}{2} \left( u_{i,j+1}^{n+1} - 2u_{i,j}^{n+1} + u_{i,j-1}^{n+1} \right)$$

$$= u_{i,j}^n - \frac{C}{4} \left( u_{i+1,j}^n - u_{i-1,j}^n \right) + \frac{\mu_x}{2} \left( u_{i+1,j}^n - 2u_{i,j}^n + u_{i-1,j}^n \right) + \frac{\mu_y}{2} \left( u_{i,j+1}^n - 2u_{i,j}^n + u_{i,j-1}^n \right),$$

with  $C = \Delta t/\Delta x$ ,  $\mu_x = \Delta t/\Delta x^2$ , and  $\mu_y = \Delta t/\Delta y^2$ .

Before we continue, let us define the grid with M and N horizontal and vertical cells, respectively. Note that the domain is then defined by (M+1)(N+1) grid nodes. Hence the coordinates  $x_i=(i-2)\Delta_x$  with  $i\in[1,M+3]$  and  $y_j=(j-2)\Delta_y$  with  $j\in[1,N+3]$  include the ghost nodes at indexes 1 and M+3, N+3 (chosen this way to match the MATLAB array indexing). We can collapse this 2D array into 1 column vector by stacking the columns one after another. We will obtain a vector of dimensions (M+3)(N+3), including ghost nodes. Thus, our solution vector is  $U_k=U_{i+(M+3)(j-1)}=u_{i,j}$ . It is easy to show that

$$u_{i\pm 1,j} = U_{i\pm 1+(M+3)(j-1)} = U_{k\pm 1},$$
  
 $u_{i,j\pm 1} = U_{i+(M+3)(j\pm 1-1)} = U_{k\pm (M+3)}.$ 

Then, we can turn our discretization to be in terms of the solution vector U,

$$U_{k}^{n+1} + \frac{C}{4} \left( U_{k+1}^{n+1} - U_{k-1}^{n+1} \right) - \frac{\mu_{x}}{2} \left( U_{k+1}^{n+1} - 2U_{k}^{n+1} + U_{k-1}^{n+1} \right) - \frac{\mu_{y}}{2} \left( U_{k+M+3}^{n+1} - 2U_{k}^{n+1} + U_{k-M-3}^{n+1} \right)$$

$$= U_{k}^{n} - \frac{C}{4} \left( U_{k+1}^{n} - U_{k-1}^{n} \right) + \frac{\mu_{x}}{2} \left( U_{k+1}^{n} - 2U_{k}^{n} + U_{k-1}^{n} \right) + \frac{\mu_{y}}{2} \left( U_{k+M+3}^{n} - 2U_{k}^{n} + U_{k-M-3}^{n} \right).$$

Reorganizing the equation we get

$$-\frac{\mu_y}{2}U_{k-M-3}^{n+1} - \left(\frac{C}{4} + \frac{\mu_x}{2}\right)U_{k-1}^{n+1} + \left(1 + \mu_x + \mu_y\right)U_k^{n+1} + \left(\frac{C}{4} - \frac{\mu_x}{2}\right)U_{k+1}^{n+1} - \frac{\mu_y}{2}U_{k+M+3}^{n+1}$$

$$= \frac{\mu_y}{2}U_{k-M-3}^n + \left(\frac{C}{4} + \frac{\mu_x}{2}\right)U_{k-1}^{n+1} + \left(1 - \mu_x - \mu_y\right)U_k^{n+1} + \left(\frac{\mu_x}{2} - \frac{C}{4}\right)U_{k+1}^{n+1} + \frac{\mu_y}{2}U_{k+M+3}^{n+1},$$

$$\begin{aligned} &a_{-M-3}U_{k-M-3}^{n+1} + a_{-1}U_{k-1}^{n+1} + a_{0}U_{k}^{n+1} + a_{+1}U_{k+1}^{n+1} + a_{M+3}U_{k+M+3}^{n+1} \\ &= b_{-M-3}U_{k-M-3}^{n} + b_{-1}U_{k-1}^{n+1} + b_{0}U_{k}^{n+1} + b_{+1}U_{k+1}^{n+1} + b_{M+3}U_{k+M+3}^{n+1}, \end{aligned}$$

This system of equations can be represented in matrix form,  $AU^{n+1} = BU^n$ , where A and B are sparse matrices defined by the 5 a diagonals and 5 b diagonals given in the equation above. To solve the system we use the MATLAB command backslash operator  $U^{n+1} = A \setminus BU^n$ . Lastly, we need to implement the boundary conditions after updating the solution. The boundary conditions in index form are:

• Bottom boundary:  $\partial_y u|_{y=0} = 0$ :

$$\frac{u_{i,j+1} - u_{i,j-1}}{2\Delta y} \Big|_{j=2} = 0$$

$$\frac{u_{i,3} - u_{i,1}}{2\Delta y} = 0$$

$$u_{i,1} = u_{i,3} \implies \boxed{U_i = U_{i+2(M+3)}}$$

• Top boundary:  $\partial_y u|_{y=1} = 0$ :

$$\frac{u_{i,j+1} - u_{i,j-1}}{2\Delta y} \Big|_{j=N+2} = 0$$

$$\frac{u_{i,N+3} - u_{i,N+1}}{2\Delta y} = 0$$

$$u_{i,N+3} = u_{i,N+1} \implies \boxed{U_{i+(N+2)(M+3)} = U_{i+N(M+3)}}$$

• Left boundary:  $u|_{x=0} - \partial_x u|_{x=0} = 0$ :

$$\begin{aligned} u_{i,j}\big|_{i=2} - \frac{u_{i+1,j} - u_{i-1,j}}{2\Delta x}\bigg|_{i=2} &= 0\\ u_{2,j} - \frac{u_{3,j} - u_{1,j}}{2\Delta x} &= 0\\ u_{1,j} &= u_{3,j} - 2\Delta x u_{2,j} \\ \hline U_{1+(M+3)(j-1)} &= U_{3+(M+3)(j-1)} - 2\Delta x U_{2+(M+3)(j-1)} \end{aligned}$$

• Right boundary:  $u|_{x=1} - \partial_x u|_{x=1} = 0$ :

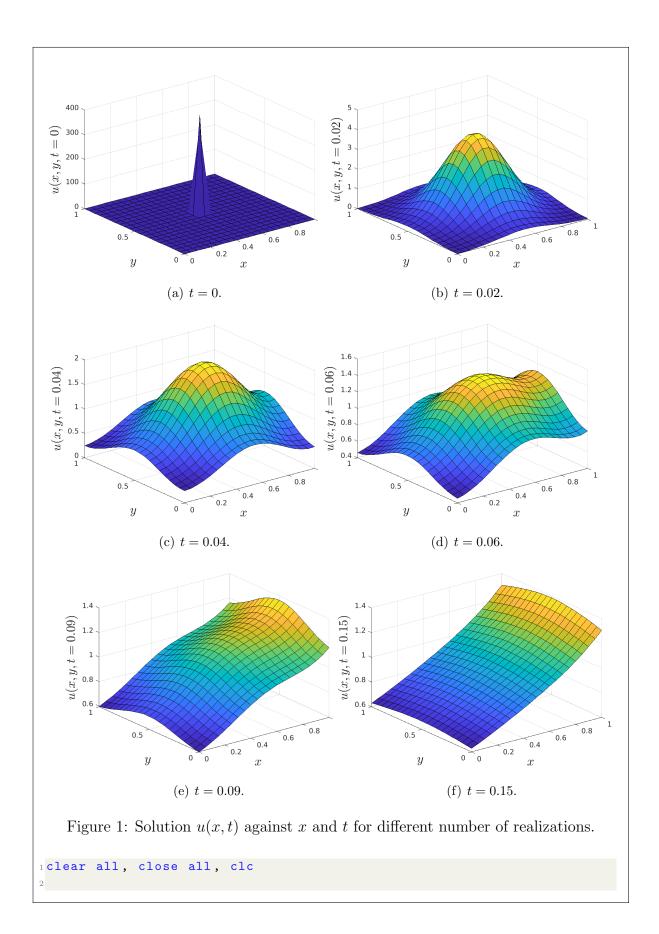
$$\begin{aligned} u_{i,j}|_{i=M+2} - \frac{u_{i+1,j} - u_{i-1,j}}{2\Delta x} \bigg|_{i=M+2} &= 0 \\ u_{M+2,j} - \frac{u_{M+3,j} - u_{M+1,j}}{2\Delta x} &= 0 \\ u_{M+3,j} &= u_{M+1,j} + 2\Delta x u_{M+2,j} \\ \hline U_{(M+3)j} &= U_{M+1+(M+3)(j-1)} + 2\Delta x U_{M+2+(M+3)(j-1)} \end{aligned}$$

The four boundary conditions are imposed after updating the solution by solving  $U^{n+1} = A \setminus BU^n$  every time-step.

# Problem 2

1. Using the Cranck - Nicholson scheme with a time and spatial step  $\Delta t = \Delta x = 0.01$  for 0 < t < 1. Solve the linear equations resulting from the implicit discretization in MATLAB sparse mode.

Solution:		



```
3 format short
5% Number of intervals
_{6}M = 20; % In x-direction
_{7}N = 20; \% In Y-direction
8
9% Define limits of domain
10 \times 0 = 0;
_{11}x_F = 1;
12 y_0 = 0;
y_F = 1;
_{14}T = 0.2;
15
16% Step sizes
_{17}dx = (x_F - x_0) / M;
18 dy = (y_F - y_0) / N;
_{19}dt = 0.0001;
20
21% Define grid (including ghost nodes)
_{22}x = linspace (x_0-dx, x_F+dx, M+3); % x(1) and x(M+3) are ghost
     nodes
_{23}y = linspace (y_0-dy, y_F+dy, N+3); % y(1) and y(N+3) are ghost
    nodes
25% Define useful constants
_{26}C = dt/dx;
27 \text{ mu}_x = \frac{dt}{(dx^2)};
28 \text{ mu}_y = \text{dt/(dy^2)};
29
30% Define initial condition
31U0 = kron((round(x,8) == 0.5),((round(y,8) == 0.5)))'/dx/dy;
32 plot_sol(U0,x,y,0)
34% Define big matrices size, including ghost nodes
_{35}S = (M + 3)*(N + 3);
36
37 %% Construct Matrix A
38% Diagonal component
39d = diag((1 + mu_x + mu_y) * ones(S,1));
40 % 1-lower diagonal component
4111 = diag((-C/4 - mu_x/2) * ones(S-1,1), -1);
42% 1-upper diagonal component
43u1 = diag((C/4 - mu_x/2) * ones(S-1,1), 1);
44% (M+3)-lower diagonal component
_{45}1M3 = diag(-mu_y/2 * ones(S-M-3,1), -M-3);
46% (M+3)-upper diagonal component
_{47}uM3 = diag(-mu_y/2 * ones(S-M-3,1), M+3);
48% Add them together
_{49}A = sparse(d + u1 + 11 + 1M3 + uM3);
51 %% Construct Matrix B
52% Diagonal component
```

```
53d = diag((1 - mu_x - mu_y) * ones(S,1));
54% 1-lower diagonal component
5511 = diag((C/4 + mu_x/2) * ones(S-1,1), -1);
56% 1-upper diagonal component
_{57}u1 = diag((-C/4 + mu_x/2) * ones(S-1,1), 1);
58% (M+3)-lower diagonal component
591M3 = diag(mu_y/2 * ones(S-M-3,1), -M-3);
60% (M+3)-upper diagonal component
g_1uM3 = diag(mu_y/2 * ones(S-M-3,1), M+3);
62 % Add them together
_{63}B = sparse(d + u1 + 11 + 1M3 + uM3);
65 %% Compute solution
67 t_plots = 0.01:0.01:T
69 t = 0;
70 \text{ tstep} = 0;
71 plot_idx = 1;
72% Initialize
_{73}U = U0;
74
75i = 1:M+3; % Horizontal array index
_{76}j = 1:N+3; \% Vertical array index
77 while t < T
78
79
     % Advance solution
     rhs = B*U;
     U = A \rhs;
81
82
     % Impose boundary conditions
     U(i) = U(i + 2*(M+3)); \% Bottom
84
85
     U(i + (N+2)*(M+3)) = U(i + N*(M+3)); % Top
     U(1 + (M+3)*(j-1)) = U(3 + (M+3)*(j-1)) - 2*dx*U(2 + (M+3)*(j-1))
86
     ); % Left
     U(M+3 + (M+3)*(j-1)) = U(M+1 + (M+3)*(j-1)) + 2*dx*U(M+2 + (M+3))
     *(j-1)); % Right
     % Advance time
     t = round(t + dt,4); % avoid floating point error to compare
90
     with ismember
91
     tstep = tstep + 1;
     if (mod(tstep, 100) == 0)
93
          plot_sol(U,x,y,t)
94
95
          pause (0.15)
     end
96
97
98 end
99
of unction plot_sol(U,x,y,t)
labelfontsize = 18;
```

```
102
      Lx = length(x);
103
     Ly = length(y);
104
      u = reshape(U, [Lx Ly]);
105
      figure(1)
106
07
      surf(x(2:Ly-1),y(2:Lx-1),u(2:Lx-1,2:Ly-1)')
108
109
      xlabel('$x$','interpreter','latex','fontsize',labelfontsize)
110
     ylabel('$y$','interpreter','latex','fontsize',labelfontsize)
      zlabel(append('$u(x, y, t = ',num2str(t),')$'),'interpreter','
112
     latex','fontsize',labelfontsize)
      figName = create_figName(t);
113
114
      saveas(gcf,figName,'png')
15 end
116
17 function figName = create_figName(t)
118
     if t == 0
          exponent = 0;
119
          base = 00;
120
121
     else
122
          exponent = -2;
123
          base = t*100;
     end
124
25
     path = '../figures/';
126
127
     if base > 9
128
          figName = append(path,'p2_u_t',num2str(base),'e',num2str(
129
     exponent));
     else
130
          figName = append(path, 'p2_u_t0', num2str(base), 'e', num2str(
131
     exponent));
      end
32
33 end
```

Consider the Schroedinger equation

$$i\hbar \frac{\partial \psi(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x,t) + V(x)\psi(x,t).$$

Here  $\hbar$  is Plancks constant, m is the mass of the particle, and V(x) is the external potential energy.  $\psi$  is the wave function. The particle density is given by  $u(x,t) = |\psi(x,t)|^2$ . Assume units for space, time and mass such that  $\hbar = m = 1$  holds.

### Problem 3

1. In these units derive the equations for a spectral Galerkin method on the interval  $x \in [\pi, pi]$  with periodic boundary conditions, using the basis functions

$$\phi_n(x) = e^{inx}, \quad n = -N:N,$$

and the weighted complex L2 scalar product  $\langle u, v \rangle = \frac{1}{2\pi} \int_{-\pi}^{\pi} u^* v dx$  for a potential energy given by

$$V(x) = \begin{cases} 0 & -\pi < x < 0 \\ 1 & 0 < x < \pi \end{cases}$$

**Solution:** Multiplying Schroedinger's equation by the complex constant i and making  $\hbar = m = 1$  we obtain

$$\frac{\partial \psi(x,t)}{\partial t} = \frac{i}{2} \frac{\partial^2}{\partial x^2} \psi(x,t) - iV(x)\psi(x,t).$$

Let us approximate the solution using the orthonormal basis functions  $\phi_n(x)$  given,

$$\psi(x,t) \approx \psi_N(x,t) = \sum_{n=-N}^{N} c_n(t)\phi_n(x) = \frac{1}{\sqrt{2\pi}} \sum_{n=-N}^{N} c_n(t)e^{inx}$$

Define

$$F(\psi_N) := \frac{i}{2} \frac{\partial^2}{\partial x^2} \psi_N(x, t) - iV(x)\psi_N(x, t)$$

$$= \frac{1}{\sqrt{2\pi}} \frac{i}{2} \frac{\partial^2}{\partial x^2} \sum_{n = -N}^{N} c_n(t) e^{inx} - \frac{1}{\sqrt{2\pi}} iV(x) \sum_{n = -N}^{N} c_n(t) e^{inx}$$

$$= -\frac{1}{\sqrt{2\pi}} \frac{i}{2} \sum_{n = -N}^{N} c_n(t) n^2 e^{inx} - \frac{1}{\sqrt{2\pi}} iV(x) \sum_{n = -N}^{N} c_n(t) e^{inx}$$

$$= \frac{-i}{\sqrt{2\pi}} \sum_{n = -N}^{N} \left[ c_n(t) \left( V(x) + \frac{n^2}{2} \right) e^{inx} \right].$$

Further,

$$\begin{split} b_{m} &= \langle \phi_{m}, F(\psi_{N}) \rangle \\ &= \frac{-i}{\sqrt{2\pi}} \frac{1}{\sqrt{2\pi}} \int_{-\pi}^{\pi} \left\{ e^{-imx} \sum_{n=-N}^{N} \left[ c_{n}(t) \left( V(x) + \frac{n^{2}}{2} \right) e^{inx} \right] \right\} dx \\ &= -i \sum_{n=-N}^{N} c_{n}(t) \int_{-\pi}^{\pi} \frac{1}{2\pi} \left( V(x) + \frac{n^{2}}{2} \right) e^{i(n-m)x} dx \\ &= -i \sum_{n=-N}^{N} c_{n}(t) \left[ \int_{-\pi}^{\pi} \frac{1}{2\pi} V(x) e^{i(n-m)x} dx + \int_{-\pi}^{\pi} \frac{1}{2\pi} \frac{n^{2}}{2} e^{i(n-m)x} dx \right] \\ &= -i \sum_{n=-N}^{N} c_{n}(t) \left[ I_{m,n}^{V} + \frac{n^{2}}{2} \delta_{m,n} \right], \end{split}$$

where  $I_{m,n}^V = \frac{1}{2\pi} \int_{-\pi}^{\pi} V(x) e^{i(n-m)x} dx$  and  $\delta_{m,n}$  is the Kronecker delta. Both  $I_{m,n}^V$ , and  $\delta_{m,n}$  can be expressed in matrix form. Hence,

$$\overrightarrow{b} = A\overrightarrow{c}(t)$$
.

where  $A = -i \left[ I_{m,n}^V + \frac{n^2}{2} \delta_{m,n} \right]$ . Since  $\overrightarrow{c}'(t) = \overrightarrow{b}(\overrightarrow{c})$ , we conclude

$$\overrightarrow{c}'(t) = A\overrightarrow{c}(t),$$

where the entries of A are determined by the entries of  $I_{m,n}^V$ , and  $\delta_{m,n}$ , which are detailed below:

$$I_{m,n}^{V} = \frac{1}{2\pi} \int_{-\pi}^{\pi} V(x)e^{i(n-m)x} dx$$
$$= \frac{1}{2\pi} \int_{0}^{\pi} e^{i(n-m)x} dx,$$

since V(x) = 0 from  $-\pi$  to 0. Hence,

• For n = m:

$$I_{m,n}^{V} = \frac{1}{2\pi} \int_{0}^{\pi} e^{i(n-m)x} dx$$
$$= \frac{1}{2\pi} \int_{0}^{\pi} dx$$
$$= \frac{1}{2}.$$

• For  $n \neq m$ :

$$I_{m,n}^{V} = \frac{1}{2\pi} \int_{0}^{\pi} e^{i(n-m)x} dx$$

$$= \frac{1}{2\pi} \frac{i \left(1 - e^{i(n-m)\pi}\right)}{(n-m)}$$

$$= \frac{i}{2\pi(n-m)} \left[1 - (-1)^{(n-m)}\right].$$

### Problem 4

1. Solve the SE for 0 < t < 1 for N = 3, 10, 50 with a time step  $\Delta t$  such that

$$\Delta t ||A||^2 = \Delta t \sqrt{\rho(A^H A)} = 1$$

holds. Here  $\rho$  is the spectral radius  $\rho(A^HA) = \max_n |\lambda_n|$  with  $\lambda_n$  the eigenvalues of  $A^HA$  and  $A^H$  the hermitian of the complex matrix A:  $A^H = (A^T)^*$ . Use the MATLAB function eig to compute the eigenvalues. Use

$$\psi^{I}(x) = \begin{cases} 1 & -\pi < x < -\frac{\pi}{2} \\ 0 & -\frac{\pi}{2} < x < \pi \end{cases}$$

as an initial condition, and the Cranck - Nicholson scheme for the time discretization.

**Solution:** From the previous problem, recover the equation

$$\overrightarrow{c}'(t) = A\overrightarrow{c}(t),$$

and apply the Crank-Nicholson scheme,

$$\frac{\overrightarrow{c}^{k+1} - \overrightarrow{c}^{k}}{\Delta t} = A \frac{1}{2} \left( \overrightarrow{c}^{k+1} + \overrightarrow{c}^{k} \right),$$

which yields

$$\overrightarrow{c}^{k+1} - A \frac{\Delta t}{2} \overrightarrow{c}^{k+1} = \overrightarrow{c}^k + A \frac{\Delta t}{2} \overrightarrow{c}^k.$$

Solving for  $\overrightarrow{c}^{k+1}$ ,

$$\overrightarrow{c}^{k+1} = \left(\mathbb{I} - A\frac{\Delta t}{2}\right)^{-1} \left(\mathbb{I} + A\frac{\Delta t}{2}\right) \overrightarrow{c}^{k}$$
$$= M\overrightarrow{c}^{k}.$$

Hence, we can find the solution at any time step k by

$$\overrightarrow{c}^k = M^k \overrightarrow{c}^0$$
.

Note that we use the Crank-Nicholson scheme on the *coefficients* of the expansion. Hence, we need to calculate the coefficients of the initial condition,

$$c_n(0) = \left\langle \phi_n, \psi^I(x) \right\rangle = \frac{1}{\sqrt{2\pi}} \int_{-\pi}^{\pi} e^{-inx} \psi^I(x) dx$$
$$= \frac{1}{\sqrt{2\pi}} \int_{-\pi}^{-\pi/2} e^{-inx} dx$$
$$= \frac{i}{n\sqrt{2\pi}} e^{in\pi/2} \left( 1 - e^{in\pi/2} \right), \text{ for } n \neq 0,$$

and,

$$c_0(0) = \langle \phi_0, \psi^I(x) \rangle = \frac{1}{\sqrt{2\pi}} \int_{-\pi}^{\pi} \psi^I(x) dx$$
$$= \frac{1}{\sqrt{2\pi}} \int_{-\pi}^{-\pi/2} dx$$
$$= \sqrt{\frac{\pi}{8}}.$$

In the following figures, we can see the effect of the number of terms, N, used in the expansion. In the first figure we can see the dramatic effect that N has on the accuracy of the approximation of the initial condition. The remaining figures show the evolution with time of the solution  $\psi$  and the probability density u for each choice of N. Find the MATLAB code at the end.

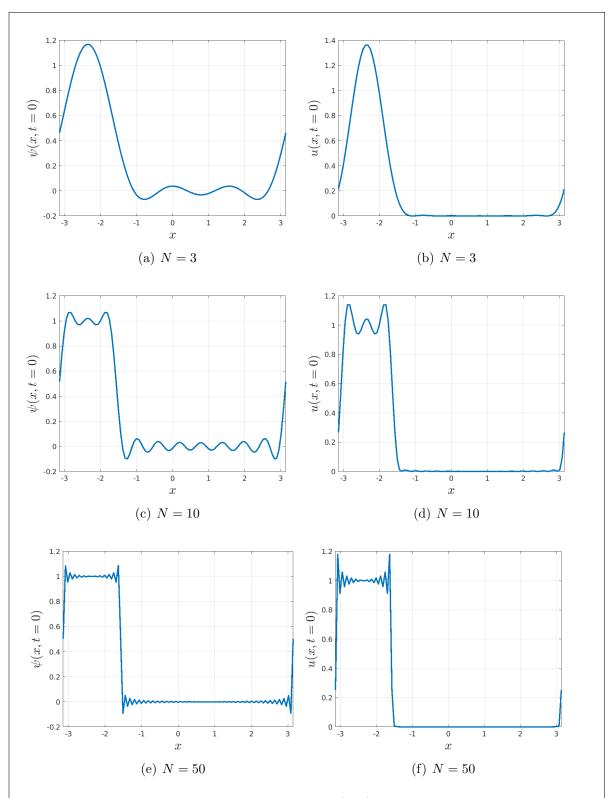


Figure 2: Initial condition of the wave function (left) and probability density function (right) for different number of terms in the approximation.

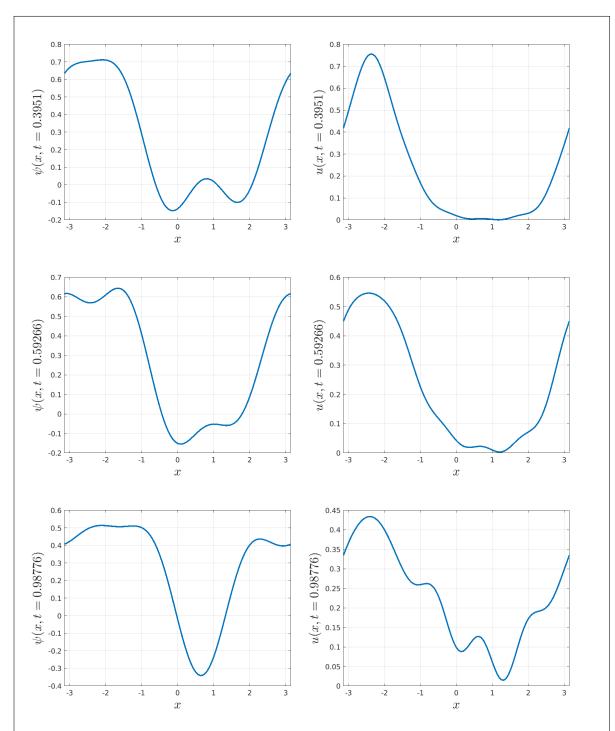


Figure 3: Evolution with time of the wave function  $\psi$  (left) and probability density function u (right) for N=3.

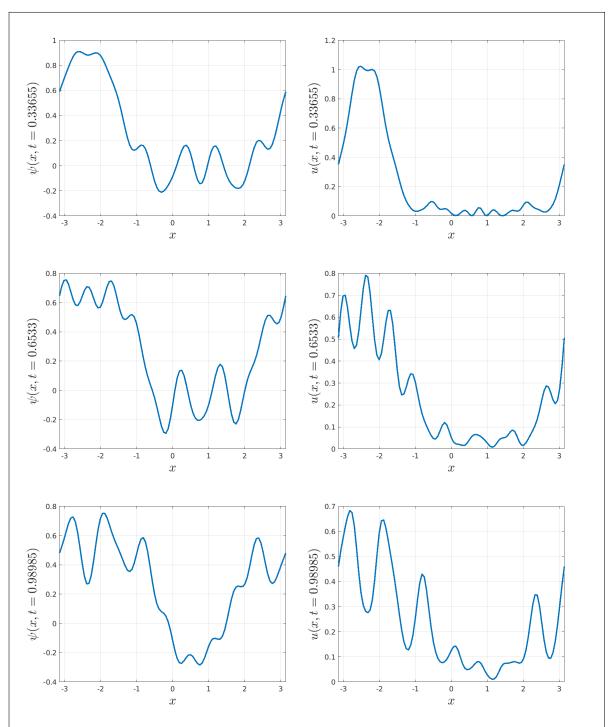


Figure 4: Evolution with time of the wave function  $\psi$  (left) and probability density function u (right) for N=10.

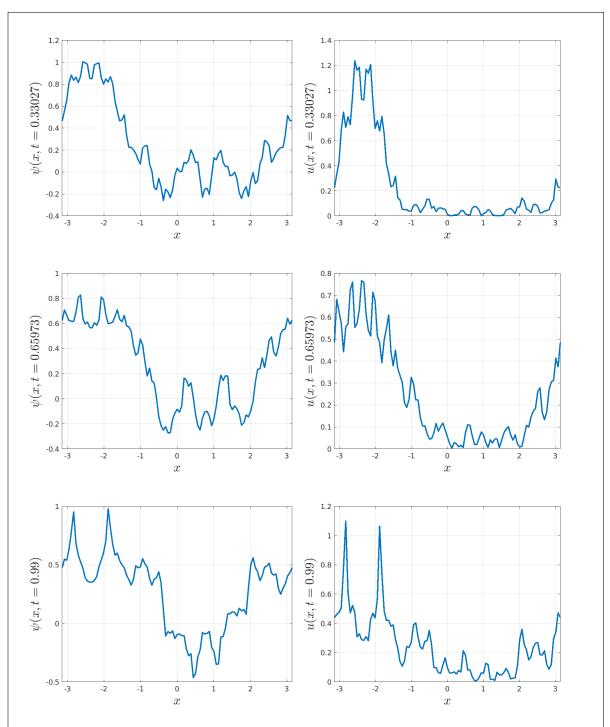


Figure 5: Evolution with time of the wave function  $\psi$  (left) and probability density function u (right) for N=50.

```
clear all, close all, clc
3N = 50; % Number of terms in expansion
_{4}n = (-N:N)'; \% (2N+1, 1)
6% Define vector x
7Nx = 100; % Number of grid steps
8x = linspace(-pi, pi, Nx+1); % (Nx+1, 1)
10% Construct matrix A
11 A = calculate_A(n,N);
12
13% Calculate dt
14 dt = calculate_dt(A);
16% Define useful matrices
_{17}M1 = eye(size(A)) - 0.5*dt*A;
_{18}M2 = eye(size(A)) + 0.5*dt*A;
_{19}M = M1\M2; \% As in c(k+1) = M c(k)
21% Define initial coefficients for initial condition
22 c0 = calculate_initial_condition(n);
23 psi = calculate_wave_function(c0,n,x);
24u = calculate_density(psi);
25 plot_psi(psi,x,0,0,N)
26 plot_u(u,x,0,0,N)
28%% Find the time-steps to plot
29 T = 1;
30 t = 0 : dt : T;
31t_{plots} = [0.33 \ 0.66 \ 0.99]; \% Times to plot
32 tstep_plots = zeros(3,1);
33 for i=1:length(t_plots)
      [minValue, minIdx] = min(abs(t-t_plots(i)));
      tstep_plots(i) = minIdx;
36 end
37 tstep_plots
38t(tstep_plots)
39 keyboard
40 %% Compute solutions and plot
41 for idx=1:length(tstep_plots)
    tstep = tstep_plots(idx);
    c = M^tstep*c0;
43
     t = (tstep-1)*dt;
44
    psi = calculate_wave_function(c,n,x);
    u = calculate_density(psi);
     plot_psi(psi,x,t,idx,N)
47
     plot_u(u,x,t,idx,N)
48
49 end
51 function A = calculate_A(n,N)
m = n;
```

```
I = 1i * (1 - (-1).^(n'-m)) ./ (2 * pi * (n'-m));
     I(1:2*N+2:(2*N+1)^2) = 1/2; \% Change diagonal values
     A = -1i * (I + diag(n.^2/2));
56 end
58 function dt = calculate_dt(A)
     spectral_radius = max(eig(A'*A));
     dt = 1/ sqrt(spectral_radius);
61 end
62
63 function c0 = calculate_initial_condition(n)
    c0 = 1i * exp(1i*n*pi/2) .* (1 - exp(1i*n*pi/2)) ./ (sqrt(2*pi
    ) * n);
     c0(n == 0) = sqrt(pi/8);
65
66 end
68 function psi = calculate_wave_function(c,n,x)
     psi = transpose(c) * exp(1i * n * x)/sqrt(2*pi);
69
70 end
72 function u = calculate_density(psi)
u = abs(psi).^2;
74 end
76 function plot_psi(psi,x,t,idx,N)
     labelfontsize = 18;
78
    figure(1)
79
    plot(x,psi,'linewidth',2)
    xlim([-pi pi])
81
    xlabel('$x$','interpreter','latex','fontsize',labelfontsize)
82
     ylabel(append('$\psi(x, t = ',num2str(t),')$'),'interpreter','
    latex','fontsize',labelfontsize)
     grid on
84
     figName = create_figName('psi',idx,N);
     saveas(gcf,figName,'png')
86
87 end
88
89 function plot_u(u,x,t,idx,N)
     labelfontsize = 18;
90
91
92
    figure(2)
     plot(x,u,'linewidth',2)
93
     xlim([-pi pi])
94
     xlabel('$x$','interpreter','latex','fontsize',labelfontsize)
95
     ylabel(append('$u(x, t = ',num2str(t),')$'),'interpreter','latex
96
     ','fontsize',labelfontsize)
97
     grid on
     figName = create_figName('u',idx,N);
98
     saveas(gcf,figName,'png')
99
00 end
01
```

```
function figName = create_figName(name,idx,N)
    path = '../figures/';
figName = append(path,'p4_',name,'_N',num2str(N),'_sample',num2str
    (idx));
for end
```