

Partial Differential Equations

TA Homework 4

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February 8, 2018

Problem 3.3.4

1. Solve the wave equation

$$\begin{aligned}\partial_t^2 u - \partial_x^2 u &= 0, & x, t \in \mathbb{R} \\ u(x, x) &= f(x), \quad \partial_t u(x, 0) = g(x), & x \in \mathbb{R},\end{aligned}$$

where $g, f : \mathbb{R} \rightarrow \mathbb{R}$.

State appropriate assumptions for f and g such that you really have a solution.

Notice that $c = 1$.

Solution: Given the previous PDE, the wave equation, we can express the solution $u(x, t)$ as

$$u(x, t) = F(x + t) + G(x - t),$$

where we have already taken into account that $c = 1$. Imposing the first initial condition,

$$u(x, x) = F(2x) + G(0) = f(x). \tag{1}$$

For the second initial condition, we calculate first $\partial_t u$,

$$\partial_t u(x, t) = F'(x + t) - G'(x - t),$$

and impose the initial condition,

$$\partial_t u(x, 0) = F'(x) - G'(x) = g(x).$$

Integrating the equation we get

$$F(x) - G(x) = F(0) - G(0) + \int_0^x g(y) dy,$$

or

$$F(2x) - G(2x) = F(0) - G(0) + \int_0^{2x} g(y) dy.$$

Subtract now equation (1) from the previous equation and get

$$\begin{aligned}-G(2x) - G(0) &= F(0) - G(0) - f(x) + \int_0^{2x} g(y) dy, \\ -G(2x) &= F(0) - f(x) + \int_0^{2x} g(y) dy, \quad G(2x) = f(x) - F(0) - \int_0^{2x} g(y) dy,\end{aligned}$$

or

$$G(x) = f\left(\frac{x}{2}\right) - F(0) - \int_0^x g(y)dy.$$

From the equation above we obtain $G(x - t)$ needed for our solution,

$$G(x - t) = f\left(\frac{x - t}{2}\right) - F(0) - \int_0^{x-t} g(y)dy.$$

Now, we retake equation (1) and rewrite it as

$$F(x + t) = f\left(\frac{x + t}{2}\right) - G(0).$$

Thus, the solution is

$$u(x, t) = f\left(\frac{x + t}{2}\right) + f\left(\frac{x - t}{2}\right) - (F(0) + G(0)) - \int_0^{x-t} g(y)dy,$$

and it is left to calculate $F(0) + G(0)$, which we do by evaluating equation (1) at $x = 0$,

$$F(0) + G(0) = f(0).$$

Hence,

$$u(x, t) = f\left(\frac{x + t}{2}\right) + f\left(\frac{x - t}{2}\right) - f(0) - \int_0^{x-t} g(y)dy,$$

which satisfies the initial conditions:

- $u(x, x) = f(x).$

$$u(x, x) = f\left(\frac{2x}{2}\right) + f(0) - f(0) - \int_0^0 g(y)dy = f(x).$$

- $\partial_t u(x, 0) = g(x).$

$$\partial_t u(x, 0) = \frac{1}{2}f\left(\frac{x}{2}\right) - \frac{1}{2}f\left(\frac{x}{2}\right) + g(x) = g(x),$$

where to evaluate the derivative of the integral we have used the Fundamental Theorem of Calculus.

Thus, since the $u(x, t)$ given satisfies the PDE and the initial conditions, it is the solution provided that f is twice differentiable and g is once differentiable.

Problem 3.3.7

1. Let $L > 0$. Solve the wave equation

$$\begin{aligned}\partial_t^2 u - c^2 \partial_x^2 u &= 0, & 0 \leq x \leq L, t \geq 0, \\ u(x, 0) &= f(x), & 0 \leq x \leq L, \\ \partial_t u(x, 0) &= g(x), & 0 \leq x \leq L, \\ \partial_x u(0, t) &= 0 = \partial_x u(L, t), & t \geq 0.\end{aligned}$$

Hint: Extend f, g in an even and $2L$ -periodic fashion.

Which assumptions do f and g have to satisfy to make u a solution?

Solution: Given the zero Neumann boundary condition at $x = 0$, we extend f and g to $[-L, L]$ in an even fashion:

$$\begin{aligned}f(-x) &= f(x), & x \in [0, L], \\ g(-x) &= g(x), & x \in [0, L],\end{aligned}$$

and noticing that the derivative of f is odd

$$f'(-x) = -f'(x), \quad x \in [0, L].$$

To take care of the zero Neumann boundary condition at $x = L$ we perform a $2L$ -periodic extension of f and g which gives us functions defined on all \mathbb{R} ,

$$\begin{aligned}f(x + 2kL) &:= f(x), & k \in \mathbb{Z}, x \in [-L, L], \\ g(x + 2kL) &:= g(x), & k \in \mathbb{Z}, x \in [-L, L].\end{aligned}$$

We can prove that the extended f is $2L$ periodic and even in a similar way as the *Lemma 3.13* is proved in the notes. Indeed, let $x \in \mathbb{R}$. Then $x = y + 2kL$ with $-L \leq y \leq L$ and $k \in \mathbb{Z}$. By extension,

$$f(x + 2L) = f(y + 2(k + 1)L) = f(y) = f(y + 2kL) = f(x).$$

Further

$$f(-x) = f(-y - 2kL) = f(-y) = f(y) = f(x),$$

so f is even around zero. Since f is even about zero and $2L$ -periodic,

$$f(L + x) = f(L + x - 2L) = f(-L + x) = f(-(L - x)) = f(L - x),$$

it is also even about L . We can prove the same for g . The conditions for f and g imply that their extensions to \mathbb{R} are twice and once differentiable, respectively. The D'Alembert's formula provides a solution to the PDE and the initial conditions,

$$u(x, t) = \frac{1}{2} [f(x + ct) + f(x - ct)] + \frac{1}{2c} \int_{x-ct}^{x+ct} g(s) ds.$$

We check now that the boundary conditions are satisfied. First we calculate

$$\partial_x u(x, t) = \frac{1}{2} [f'(x + ct) + f'(x - ct)] + \frac{1}{2c} [g(x + ct) - g(x - ct)],$$

where we have used the Fundamental Theorem of Calculus to differentiate the integral, and we make $x = 0$,

$$\begin{aligned}\partial_x u(0, t) &= \frac{1}{2} [f'(ct) + f'(-ct)] + \frac{1}{2c} [g(ct) - g(-ct)] \\ &= \frac{1}{2} [f'(ct) - f'(ct)] + \frac{1}{2c} [g(ct) - g(ct)] \\ &= 0.\end{aligned}$$

For the boundary condition at $x = L$,

$$\begin{aligned}\partial_x u(0, t) &= \frac{1}{2} [f'(L + ct) + f'(L - ct)] + \frac{1}{2c} [g(L + ct) - g(L - ct)] \\ &= \frac{1}{2} [f'(L + ct) - f'(L + ct)] + \frac{1}{2c} [g(L + ct) - g(L + ct)] \\ &= 0.\end{aligned}$$

since f' and g are odd and even around L , respectively. Thus, $u(x, t)$ is the solution provided that extended f and g are twice and once differentiable, respectively. The extended f is twice differentiable if and only if the original f is twice differentiable and

$$f(0) = f(L), \quad f'(0) = 0 = f'(L).$$

The extended g is once differentiable if and only if the original g is once differentiable and

$$g(0) = g(L).$$

Problem 3.3.9

1. Let u solve

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)u(x, t) &= \phi(x, t), \quad x, t \in \mathbb{R}. \\ u(x, 0) &= 0, \quad x \in \mathbb{R}, \\ \partial_t u(x, 0) &= 0, \quad x \in \mathbb{R}\end{aligned}$$

and \tilde{u} solve

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)\tilde{u}(x, t) &= 0, \quad x, t \in \mathbb{R}. \\ \tilde{u}(x, 0) &= f(x), \quad x \in \mathbb{R}, \\ \partial_t \tilde{u}(x, 0) &= g(x), \quad x \in \mathbb{R}\end{aligned}$$

Prove that $U = u + \tilde{u}$ solves

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)U(x, t) &= \phi(x, t), \quad x, t \in \mathbb{R}. \\ U(x, 0) &= f(x), \quad x \in \mathbb{R}, \\ \partial_t U(x, 0) &= g(x), \quad x \in \mathbb{R}\end{aligned}$$

This is a special case of the so-called principle of superposition. It works here because the problem is linear.

Solution: We start by proving that U satisfies the PDE,

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)U &= (\partial_t^2 - c^2 \partial_x^2)(u + \tilde{u}) \\&= \partial_t^2(u + \tilde{u}) - c^2 \partial_x^2(u + \tilde{u}) \\&= \partial_t^2 u + \partial_t^2 \tilde{u} - c^2 \partial_x^2 u - c^2 \partial_x^2 \tilde{u} &= (\partial_t^2 - c^2 \partial_x^2)u + (\partial_t^2 - c^2 \partial_x^2)\tilde{u} \\&= \phi(x, t) + 0 &= \phi(x, t)\end{aligned}$$

Now we check that it satisfies the first initial condition,

$$\begin{aligned}U(x, 0) &= u(x, 0) + \tilde{u}(x, 0) \\&= 0 + f(x) \\&= f(x).\end{aligned}$$

To check the second initial condition we first calculate

$$\partial_t U(x, t) = \partial_t(u + \tilde{u}) = \partial_t u(x, t) + \partial_t \tilde{u}(x, t),$$

and make $t = 0$,

$$\begin{aligned}\partial_t U(x, 0) &= \partial_t u(x, 0) + \partial_t \tilde{u}(x, 0) \\&= 0 + g(x) \\&= g(x).\end{aligned}$$

Hence, since $U = u + \tilde{u}$ satisfies the PDE and the initial conditions, it is the solution.

Problem 3.3.10

1. Solve the vibrating string equation with external force,

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)u(x, t) &= t \sin(x), & t \geq 0, 0 \leq x \leq \pi, \\u(x, 0) &= \sin(x), & x \in [0, \pi], \\\partial_t u(x, 0) &= \sin(x), & x \in [0, \pi], \\u(0, t) &= 0 = u(\pi, t), & t \geq 0.\end{aligned}$$

Show that the solution is of the form $u(x, t) = \psi(t) \sin(x)$. Determine $\psi(t)$ using d'Alembert. **Do not assume that the solution is of this form.**

Solution: We start performing a 2π -periodic extension of $f(x) = \sin x$ in an odd and 2π -periodic fashion. We express the PDE as

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)u(x, t) &= t \sin(x), & x \in \mathbb{R}, t \in \mathbb{R} \\u(x, 0) &= \sin(x), & x \in \mathbb{R} \\\partial_t u(x, 0) &= \sin(x), & x \in \mathbb{R} \\u(0, t) = 0 &= u(\pi, t), & t \in \mathbb{R}\end{aligned}$$

Like in the previous problem, we can separate the PDE in two and, by the superposition principle, add the solutions of the two following PDEs:

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)u_1(x, t) &= t \sin(x), & x \in \mathbb{R}, t \in \mathbb{R} \\ u_1(x, 0) &= 0, & x \in \mathbb{R} \\ \partial_t u_1(x, 0) &= 0, & x \in \mathbb{R} \\ u_1(0, t) &= 0 = u_1(\pi, t), & t \in \mathbb{R}\end{aligned}$$

and

$$\begin{aligned}(\partial_t^2 - c^2 \partial_x^2)u_2(x, t) &= 0, & x \in \mathbb{R}, t \in \mathbb{R} \\ u_2(x, 0) &= \sin(x), & x \in \mathbb{R} \\ \partial_t u_2(x, 0) &= \sin(x), & x \in \mathbb{R} \\ u_2(0, t) &= 0 = u_2(\pi, t), & t \in \mathbb{R}\end{aligned}$$

We start with the first PDE. As it is proved in the notes, we can solve this inhomogeneous wave equation and the solution is

$$\begin{aligned}u_1(x, t) &= \frac{1}{2c} \int_0^t \int_{x-c(t-r)}^{x+c(t-r)} r \sin \rho \, d\rho dr \\ &= \frac{1}{2c} \int_0^t r [-\cos \rho]_{x-c(t-r)}^{x+c(t-r)} dr \\ &= \frac{1}{2c} \int_0^t r [\cos(x - c(t - r)) - \cos(x + c(t - r))] dr.\end{aligned}$$

By the trigonometric identities, $\cos(x - c(t - r)) - \cos(x + c(t - r)) = 2 \sin(x) \cos(ct - cr)$. Therefore,

$$\begin{aligned}u_1(x, t) &= \frac{1}{2c} \int_0^t 2r \sin(x) \cos(ct - cr) dr \\ &= \frac{\sin(x)}{c} \int_0^t r \cos(ct - cr) dr \\ &= \frac{\sin(x)}{c} \left(\left[\frac{r}{c} \cos(ct - cr) \right]_0^t - \frac{1}{c} \int_0^t \cos(ct - cr) dr \right) \\ &= \sin(x) \left(\frac{t}{c^2} - \frac{\sin(ct)}{c^3} \right),\end{aligned}$$

where we have used integration by parts in the last steps. For the second PDE, d'Alembert formula gives us the solution

$$\begin{aligned}u_2(x, t) &= \frac{1}{2} (f(x + ct) + f(x - ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} g(s) ds \\ &= \frac{1}{2} (\sin(x + ct) + \sin(x - ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} \sin(s) ds \\ &= \frac{1}{2} (\sin(x) \cos(ct) + \sin(ct) \cos(x) + \sin(x) \cos(ct) - \sin(ct) \cos(x)) + \frac{1}{2c} [-\cos(s)]_{x-ct}^{x+ct} \\ &= \sin(x) \cos(ct) + \frac{1}{2c} [\cos(x - ct) - \cos(x + ct)] \\ &= \sin(x) \cos(ct) + \frac{1}{2c} 2 \sin(x) \sin(ct) \\ &= \sin(x) \left[\cos(ct) + \frac{1}{c} \sin(ct) \right].\end{aligned}$$

Finally, by superposition principle, we obtain

$$u(x, t) = \sin(x) \left[\frac{t}{c^2} + \cos(ct) + \frac{1}{c} \sin(ct) - \frac{\sin(ct)}{c^3} \right].$$

We check now if the solution satisfies the initial condtions

$$u(x, 0) = \sin(x),$$

and

$$\begin{aligned} \partial_t u(x, t) &= \sin(x) \left[\frac{1}{c^2} - c \sin(ct) + \cos(ct) - \frac{\cos(ct)}{c^2} \right], \\ \partial_t u(x, 0) &= \sin(x) \left[\frac{1}{c^2} + 1 - \frac{1}{c^2} \right] = \sin(x). \end{aligned}$$

To finish, we check that satisfies the boundry conditions

$$u(0, t) = 0 = u(\pi, t),$$

since the sine mutlplying the whole expression is zero at those points of x . Thus, $u(x, t)$ given is the solution to our problem.