Calculus, review

Intro to ODE:

Solution using a Taylor expansion

higher-orde

first order ODEs



Ordinary differential equations CHEM-E7190 (was E7140), 2022

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A brief review of calculus Ordinary differential equation

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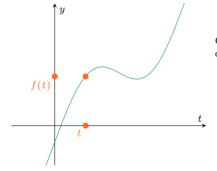
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Functions and their derivatives

A function y = f(t) encodes the relation between two quantities or variables, y and t



Consider the rate of change of quantity y corresponding to a change in t

ullet It is the ratio between the differential change in y and the corresponding differential change in variable t

We conventionally call the ratio of differential changes the derivative of function f

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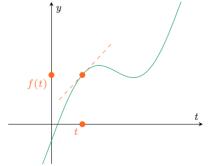
From high to first order ODE

Functions and their derivatives (cont.)

The derivative of a function f(t) is the rate of change of the function, it is a number

- \rightarrow The derivative is defined with respect to the independent variable (here, t)
- \rightarrow It can be computed at any point t of the domain of the function

We are given some function f(t), we are interested in its derivative at some point t



Derivative of function f with respect to t

$$\Rightarrow \frac{\mathrm{d}f(t)}{\mathrm{d}t}$$

The rate of change is understood as the slope of the tangent line to the function,

• ... at that specific point t

Functions and their derivatives (cont.)

The value of the derivative can be approximated by using small changes in t and f(t)

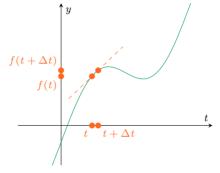
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Consider the small change $t \to t + \Delta t$ and the associated $f(t) \to f(t + \Delta t)$

$$\begin{split} \frac{\mathrm{d}f(t)}{\mathrm{d}t} &\approx \frac{f(t+\Delta t) - f(t)}{\Delta t} \\ &\leadsto \frac{\Delta y}{\Delta t} \end{split}$$

The tangent line will be approximated

- By the secant line to the function
- Its slope is the approximation

Remember the equation of a line y = mx + c through two points (x_1, y_1) and (x_2, y_2)

$$y = \frac{y_2 - y_1}{x_2 - x_1}(x - x_1) + y_1 = \underbrace{\left(\frac{y_2 - y_1}{x_2 - x_1}\right)}_{\Delta y / \Delta x} x + \underbrace{\frac{y_2 - y_1}{x_2 - x_1}x_1 + y_1}_{\text{constant}}$$

Functions and their derivatives (cont.)

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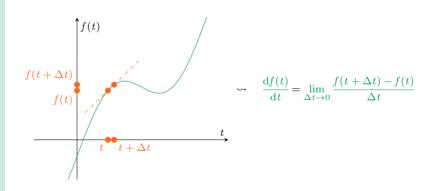
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We can improve the quality of this approximation, by letting Δt become smaller

- As $\Delta t \to 0$, the approximation will converge to the true derivative
- (Because the secant line will get closer to the tangent line)



Functions and their derivatives (cont.)

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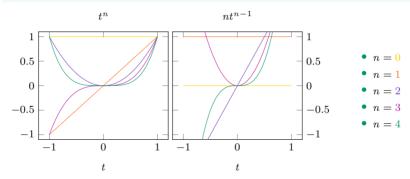
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Power law

Consider the function $f(t) = t^n$ (the power law) and its derivative $\frac{\mathrm{d}f(t)}{\mathrm{d}t} = nt^{n-1}$



- The derivative is commonly known (remembered), but we can derive it
- We will be using the approximation of derivative that we defined

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Functions and their derivatives (cont.)

$$f(t) = t^n$$

By definition of derivative, we have

$$\begin{split} \frac{\mathrm{d}f(t)}{\mathrm{d}t} &\approx \frac{f(t+\Delta t)-f(t)}{\Delta t} = \frac{1}{\Delta t} \Big[\underbrace{(t+\Delta t)^n}_{\text{Powers of a binomial}} - t^n \Big] \\ &= \frac{1}{\Delta t} \Big[\underbrace{t^n + nt^{n-1}(\Delta t) + \frac{n(n-1)}{2} t^{n-2}(\Delta t)^2 + \cdots - t^n}_{\text{Powers of a binomial}} \\ &= \frac{1}{\Delta t} \Big[\underbrace{t^n + nt^{n-1}(\Delta t) + \frac{n(n-1)}{2} t^{n-2}(\Delta t)^2 + \cdots - t^n}_{\text{Powers of a binomial}} \Big] \\ &= \frac{1}{\Delta t} \Big[nt^{n-1}(\Delta t) + \frac{n(n-1)}{2} t^{n-2}(\Delta t)^2 + \underbrace{\mathcal{O}((\Delta t)^3)}_{\text{H.O. terms}} \Big] \\ &= nt^{n-1} + \frac{n(n-1)}{2} t^{n-2}(\Delta t) + \mathcal{O}((\Delta t)^2) \\ &= nt^{n-1} + \mathcal{O}(\Delta t) \\ &\approx nt^{n-1} \end{split}$$

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Functions and their derivatives (cont.)

The first order derivative df(t)/dt is the ratio of two distinct quantities df(t) and dtThe ratio can be manipulated by conventional algebraic procedures

Thus, we can have multiplication by some quantity

$$\mathrm{d}t \, \frac{\mathrm{d}f}{\mathrm{d}t} = \mathrm{d}f$$

And, multiplication and division by some quantity

$$\frac{\mathrm{d}f}{\mathrm{d}t}\frac{\mathrm{d}t}{\mathrm{d}z} = \frac{\mathrm{d}f}{\mathrm{d}z}$$

As an application, we get the chain law of derivation

$$\frac{\mathrm{d}f(g(t))}{\mathrm{d}t} = \underbrace{\frac{\mathrm{d}f(g(t))}{\mathrm{d}g(t)}}_{f'(g(t))} \underbrace{\frac{\mathrm{d}g(t)}{\mathrm{d}t}}_{g'(t)} = f'(t)$$

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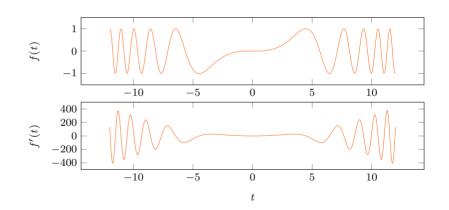
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Functions and their derivatives (cont.)

Example

Consider the function $f(t) = \sin(t^3)$, compute its first derivative with respect to t



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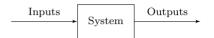
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Introduction to ODEs

Ordinary differential equations (ODEs) are probably our most useful modelling tool

- → (Together with probability)
- → (Not used in this course)



First some motivating and yet simple examples of ODEs (understood as system models)

Systems for which the input is identically null over time

- Non-zero initial conditions
- Force-free response
- y(t), when u(t) = 0



Systems for which the input is not identically null over time

- Zero initial conditions
- Forced response
- y(t), when $u(t) \neq 0$



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Example

Consider the problem of modelling the number of bacteria in some bacterial colony

- We assume that each bacterium in the colony gives rise to new individuals
- We also assume that we know the birth-rate, let us denote it by $\lambda>0$

We assume that, on average, each bacterium will produce λ offsprings per unit time

- \rightarrow The size y of the colony varies (grows) in time proportionally to its size
- → (That is, the larger the population, the larger the rate of growth)

$$\frac{\mathrm{d}y(t)}{\mathrm{d}t} = \lambda y(t) \qquad \text{(This identity is an ODE)}$$

We are interested in knowing (determining) the size y(t) of the population, over time

- The function y(t) is the solution to the ordinary differential equation
- This is the function that satisfies the model $\dot{y}(t) = \lambda y(t)$

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Introduction to ODEs (cont.)

The solution to the ODE is a (family of) function(s) y(t) that satisfies the identity

$$\frac{\mathrm{d}y(t)}{\mathrm{d}t} = \dot{y}(t) = \lambda y(t)$$

There are many techniques that can be used to solve ordinary differential equations

 \bullet For the simple growth model we can separate the variables, then integrate

$$\frac{\mathrm{d}y(t)}{\mathrm{d}t} = \lambda y(t) \quad \rightsquigarrow \quad \int_{y_0}^{y} \frac{1}{y(t)} \mathrm{d}y = \int_{t_0}^{t} \lambda \mathrm{d}t$$

- \bullet Move all terms in y to one side
- \bigcirc Move all terms in t to the other side
- 3 Integrate both sides over appropriate intervals
- 4 The intervals are set in terms of initial conditions
- **6** (The initial, at time t_0 , size of the population, $y_0 = y(t = t_0)$)

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Introduction to ODEs (cont.)

We have,

Taking the exponential of both sides, we have

$$y(t) = \underbrace{\frac{e^{(\lambda t + \text{constant})}}{e^{(\alpha+\beta)} = e^{\alpha}e^{\beta}}}_{e^{(\alpha+\beta)} = e^{\lambda t} \cdot \text{constant}}$$
$$= e^{\lambda t} \cdot \text{constant}$$

The bacteria population y(t) evolves in time as an exponential function, it grows

- The exponential grow $(\lambda > 0)$ is weighted by some constant
- The constant must be determined, we use initial conditions

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Introduction to ODEs (cont.)

$$y(t) = e^{\lambda t} \cdot \text{constant}$$

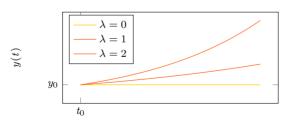
Suppose that at time t = 0, the population size is known to be $y(t = 0) = y_0 = y(0)$

$$y_0 = \underbrace{e^{\lambda \cdot 0}}_1 \cdot \text{constant} \quad \leadsto \quad \text{constant} = y_0$$

That is, the solution to the ordinary differential equation is given by $y(t) = (e^{\lambda t})y_0$

- We can solve this ODE analytically (We have a closed-form solution)
- Function $e^{\lambda t}$ is very important (The state transition function)

The system evolution, starting from an initial bacterial population size y_0 at time t_0



$$y(t) = e^{\lambda t} y_0$$

$$y_0 = 10$$

For $\lambda = 0$ the population size remains constant

• Zero birth-rate

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Introduction to ODEs (cont.)

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```
v0 = 10:
                                                              % Set initial condition
   lambda = 2:
                                                              % Set model parameter
   tMin = 00: tMax = 01: deltaT=0.1:
                                                              % Define time range
   tRange = tMin:deltaT:tMax;
                                                              % Min, max, delta
   v_clf = Q(t) exp(lambda*t)*v0;
                                                              % Set analytical solution
   [timeR, v_num] = ode45(@(t,v) lambda*v, tRange, v0);
                                                              % Compute the numerical
                                                              % solution using ODE45
12 figure (1);
                                                              % Plotting stuff
13 hold on
   fplot(v_clf,[tMin,tMax],'k');
                                                               Analytical
16 plot(timeR.v_num.'.-k');
                                                               Numerical
   stairs (timeR, y_num, '--r');
                                                                Numerical
   hold off
   xlabel('Time', 'FontSize', 24)
   vlabel('N. of bacteria', 'FontSize', 24)
24 xlim([tMin,tMax]);
   vlim([0.max(v_num)]):
                                                              % Could set a legend. ...
```

Introduction to ODEs (cont.)

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Example

Reconsider the problem of modelling the number of bacteria in some bacterial colony

- We assume that bacteria procreate, at rate λ_1
- We assume that bacteria die, at rate λ_2

$$\dot{y}(t) = \lambda_1 y(t) - \lambda_2 y(t)$$

$$= \underbrace{(\lambda_1 - \lambda_2)}_{\lambda} y(t)$$

$$= \lambda y(t)$$

Formally, the resulting model (ODE) has not changed

• We know the solution for some initial condition

$$y(t) = (e^{\lambda t})y_0$$

• λ is no longer restricted to be non-negative

Introduction to ODEs (cont.)

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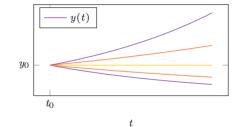
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Suppose that at time t = 0, the population size is known to be $y(t = 0) = y_0 = y(0)$



$$\lambda = \{-2, -1, 0, 1, 2\}$$

 $y_0 = 10$

We cannot discriminate between the effect of birth λ_1 and death λ_2 any longer (!)

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Solution by Taylor expansion

Ordinary differential equations

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Solution by Taylor series expansion

Consider ODE $\dot{y}(t) = \lambda y(t)$, but suppose that we want approximate the solution y(t)

• Suppose we express the solution y(t) by its Taylor series expansion

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + \cdots$$

- \leadsto This is a parametric representation of function y(t)
- \rightarrow The parameters $\{c_0, c_1, c_2, c_3, \dots\}$ are constants

We are interested in determining the actual solution y(t), from this approximation

- To characterise a specific y(t) we must set the parameters
- (We must determine the constants in the expansion)

In general, the Taylor series expansion of some function f(x) around some point x_0

$$f(x) = f(x_0) + \frac{\mathrm{d}f}{\mathrm{d}x}\Big|_{x_0} \frac{(x - x_0)}{1!} + \frac{\mathrm{d}^2f}{\mathrm{d}x^2}\Big|_{x_0} \frac{(x - x_0)^2}{2!} + \frac{\mathrm{d}^3f}{\mathrm{d}x^3}\Big|_{x_0} \frac{(x - x_0)^3}{3!} + \cdots$$

Solution by Taylor series expansion (cont.)

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Solution using a Taylor expansion

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Consider the ODE $\dot{y}(t) = \lambda y(t)$, we could compute its solution by variable separation

- \rightarrow We considered some value of λ and some initial condition y(t=0)=y(0)
- \rightarrow Then, we calculated the closed-form solution $y(t) = e^{\lambda t}y(0)$

By expressing the solution y(t) in terms of its Taylor series expansion, we have

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + c_4 t^4 + \mathcal{O}(t^5)$$

Given this expression of y(t), we could also calculate its first derivative $\dot{y}(t)$

$$\dot{y}(t) = 0 + c_1 + 2c_2t + 3t^2 + 4c_4t^3 + \mathcal{O}(t^4)$$

We substitute $\dot{y}(t)$ and y(t) into the ordinary differential equation, $\dot{y}(t) = \lambda y(t)$

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Solution by Taylor series expansion (cont.)

We substitute $\dot{y}(t)$ and y(t) into the given ordinary differential equation, $\dot{y}(t) = \lambda y(t)$

$$\underbrace{\frac{0+c_1+2c_2t+3c_3t^2+4c_4t^3+\mathcal{O}(t^4)}_{\dot{y}(t)}}_{=\underbrace{\lambda c_0+\lambda c_1t+\lambda c_2t^2+\lambda c_3t^3+\mathcal{O}(t^4)}_{\lambda y(t)}$$

The identity is satisfied when the coefficients of the powers of t in both sides match

$$\begin{array}{lll}
 & \sim (t^0) & 1c_1 = \lambda c_0 \\
 & \sim (t^1) & 2c_2 = \lambda c_1 \\
 & \sim (t^2) & 3c_3 = \lambda c_2 \\
 & \sim (t^3) & 4c_4 = \lambda c_3
\end{array}$$

 \rightarrow c_0 can be determined

If we knew c_0 , we could calculate c_1 , then given c_1

we could calculate c_2 , from c_2 we could calculate c_2 , then given c_2 we could calculate c_3 , ...

Using the initial condition y_0 , we know that at t=0, we have

$$y_0 = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + \cdots$$

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Solution by Taylor series expansion (cont.)

$$\begin{array}{lll}
 & c_0 = y_0 \\
 & c_1 = \lambda c_0 = \lambda y_0 \\
 & c_2 = \frac{1}{2}\lambda c_1 = \frac{1}{2}\lambda^2 y_0 \\
 & c_3 = \frac{1}{3}\lambda c_2 = \frac{1}{3!}\lambda^3 y_0 \\
 & c_4 = \frac{1}{4}\lambda c_3 = \frac{1}{4!}\lambda^4 y_0
\end{array}$$

By substituting the coefficients $\{c_0, c_1, c_2, \dots\}$ in the series expansion of y(t), we obtain

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + c_4 t^4 + \cdots$$

$$= y_0 + \lambda y_0 t + \frac{1}{2} \lambda^2 y_0 t^2 + \frac{1}{3!} \lambda^3 y_0 t^3 + \frac{1}{4!} \lambda^4 y_0 t^4 + \mathcal{O}(t^5)$$

$$= \left[\underbrace{1 + \lambda t + \frac{1}{2} \lambda^2 t^2 + \frac{1}{3!} \lambda^3 t^3 + \frac{1}{4!} \lambda^4 t^4 + \mathcal{O}(t^5)}_{} \right] y_0 = e^{\lambda t} y_0$$

Exponential function $e^{\lambda t}$

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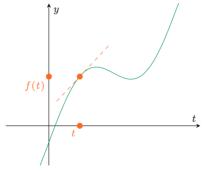
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Solution by Taylor series expansion (cont.)

Taylor series expansion

Any smooth function $f(t + \Delta t)$ can be expanded as a Taylor series at some point t

$$f(t \pm \Delta t) = f(t) + \frac{\mathrm{d}f(t)}{\mathrm{d}t} \Delta t + \frac{\mathrm{d}^2 f(t)}{\mathrm{d}t^2} \frac{(\Delta t)^2}{2!} + \frac{\mathrm{d}^3 f(t)}{\mathrm{d}t^3} \frac{(\Delta t)^3}{3!} + \dots + \frac{\mathrm{d}^n f(t)}{\mathrm{d}t^n} \frac{(\Delta t)^n}{n!} + \dots$$



If we know the function, its first derivative, its second order derivative, ... at point t, we can approximate f near that point

The more derivatives we add, the more accurate the approximation

$$\Rightarrow \text{Also, } f(t) = f(t_0) + \frac{\mathrm{d}f(t_0)}{\mathrm{d}t}(t - t_0) + \frac{\mathrm{d}^2f(t_0)}{\mathrm{d}t^2} \frac{(t - t_0)^2}{2!} + \frac{\mathrm{d}^3f(t_0)}{\mathrm{d}t^3} \frac{(t - t_0)^3}{3!} + \cdots$$

Solution by Taylor series expansion (cont.)

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Example

Consider functions $f(t) = \sin(t)$ and $f(t) = \cos(t)$, compute the Taylor expansions

• Expand them about point $t_0 = 0$ (MacLaurin series expansions)

In general, we can write the expansion

$$f(t) = f(t_0) + \frac{\mathrm{d}f(t_0)}{\mathrm{d}t}(t - t_0) + \frac{\mathrm{d}^2f(t_0)}{\mathrm{d}t^2}(t - t_0)^2 + \frac{\mathrm{d}^3f(t_0)}{\mathrm{d}t^3}(t - t_0)^3 + \cdots$$

For the sine function, we have

$$\sin(t) = \sin(0) + \cos(0)t - \frac{1}{2!}\sin(0)t^2 - \frac{1}{3!}\cos(0)t^3 + \frac{1}{4!}\sin(0)t^4 - \frac{1}{5!}\cos(0)t^5 + \cdots$$

$$= \sin(0) + \cos(0)t - \frac{1}{2!}\sin(0)t^2 - \frac{1}{3!}\cos(0)t^3 + \frac{1}{4!}\sin(0)t^4 - \frac{1}{5!}\cos(0)t^5 + \cdots$$

$$= t - \frac{t^3}{3!} + \frac{t^5}{5!} - \frac{t^7}{7!} + \cdots = \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k+1)}t^{2k+1}$$

Solution by Taylor series expansion (cont.)

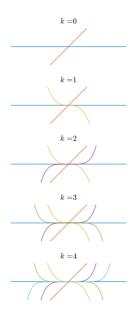
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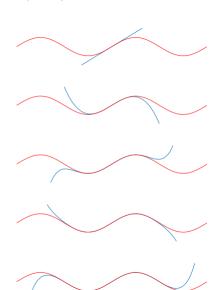
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Solution by Taylor series expansion (cont.)

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$$f(t) = f(t_0) + \frac{\mathrm{d}f(t_0)}{\mathrm{d}t}(t - t_0) + \frac{\mathrm{d}^2f(t_0)}{\mathrm{d}t^2}(t - t_0)^2 + \frac{\mathrm{d}^3f(t_0)}{\mathrm{d}t^3}(t - t_0)^3 + \cdots$$

For the cosine function, we have

$$\begin{split} \cos{(t)} &= \cos{(0)} - \sin{(0)t} - \frac{1}{2!}\cos{(0)}t^2 + \frac{1}{3!}\sin{(0)}t^3 + \frac{1}{4!}\cos{(0)}t^4 - \frac{1}{5!}\sin{(0)}t^5 + \cdots \\ &= 1 - \frac{t^2}{2!} + \frac{t^4}{4!} - \frac{t^6}{6!} + \cdots \\ &= \sum_{k=0}^{\infty} \frac{(-1)^k}{(2k)!}t^{2k} \end{split}$$

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Solution by Taylor series expansion (cont.)

$$\cos(t) = 1 + 0t - \frac{t^2}{2!} + 0t^3 + \frac{t^4}{4!} + 0t^5 - \frac{t^6}{6!} + \dots = \sum_{k=0}^{\infty} \frac{(-1)^k}{(2n)!} t^{2n}$$

```
tRange = -2*pi:0.01:+2*pi;
                                                          # Define the t-range
  F\cos = Q(t) \cos(t):
                                                          # Define functional variable
                                                          # Ecos of t
   C\cos 1 = [0 1]:
                                                          # Set coefficients of a 1st
       Tcos_1 = polyval(Ccos_1,xRange);
                                                          # order expansion, evaluation
 8
   Ccos 3 = [-1/factorial(2) \ 0 \ 1]:
                                                          # Set coefficients of a 2nd
       Tcos 3 = polyval(Ccos 3.xRange):
                                                          # order expansion, evaluation
   Ccos_5 = [1/factorial(4) \ 0 \ -1/factorial(2) \ 0 \ 1];
                                                          # Set coefficients of a 3rd
       Tcos 5 = polyval(Ccos 5.xRange):
                                                          # order expansion, evaluation
14
15 figure(1); hold on
                                                          # Some plotting
17 fplot(Fcos):
                                                          # Plots function Fcos
1.8
19 plot(xRange, Tcos 1):
                                                          # Plots 1st approximation
20 plot(xRange, Tcos_3);
                                                          # Plots 2nd approximation
21 plot(xRange, Tcos 5):
                                                          # Plots 3rd approximation
23 hold off
```

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Solution by Taylor series expansion (cont.)

We can consider the Taylor series expansion of the exponential function e^t (important)

$$\rightarrow$$
 $e^t = 1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \frac{t^4}{4!} + \cdots$

We can consider the Taylor series expansion of the function $e^{\lambda t}$ (this is also important)

• By replacing t with λt , we obtain

$$\Rightarrow$$
 $e^{\lambda t} = 1 + (\lambda t) + \frac{(\lambda t)^2}{2!} + \frac{(\lambda t)^3}{3!} + \frac{(\lambda t)^4}{4!} + \cdots$

We may want to write the Taylor series expansion of function e^{it} , with $i = \sqrt{(-1)}$

$$e^{it} = 1 + it + \frac{(it)^2}{2!} + \frac{(it)^3}{3!} + \frac{(it)^4}{4!} + \cdots$$

$$= 1 + it - \frac{t^2}{2!} - i\frac{t^3}{3!} + \frac{t^4}{4!} + i\frac{t^5}{5!} + \cdots$$

$$= \underbrace{1 - \frac{t^2}{2!} + \frac{t^4}{4!} + \cdots + i\left(t - \frac{t^3}{3!} + \frac{t^5}{5!} + \cdots\right)}_{\sin(t)}$$

$$= \cos(t) + i\sin(t)$$
Again, by replacing t with it

$$\xrightarrow{} (\text{Euler's formula})$$

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Second- and higher-order systems Ordinary differential equation

Second- and higher-order systems

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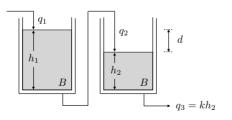
From high to first order ODE

Example

Two tanks (SS to IO representation and return)

Consider a system consisting of two cylindric liquid tanks, same cross section $B\ [\mathrm{m}^2]$

- A main inflow to tank 1, a main outflow from tank 2
- The outflow from tank 1 is the inflow to tank 2



First liquid tank

- Inflow, rate q_1 [m³s⁻¹]
- Outflow, rate q_2 [m³s⁻¹]
- h_1 is the liquid level [m]

Second liquid tank

- Inflow, rate q_2 [m³s⁻¹]
- Outflow, rate q_3 [m³s⁻¹]
- h_2 is the liquid level [m]

Calculus, review

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Second- and higher-order systems (cont.)

Suppose that flow-rates q_1 and q_2 can be set to some desired value (pumps)

Also, suppose that q_3 depends linearly on the liquid level in the tank, h_2

• $q_3 = k \cdot h_2$ [m³s⁻¹], with k [m²s⁻¹] some appropriate constant

$$\underbrace{u(t) = [q_1(t), q_2(t)]'}_{\text{Two tanks}} \text{Two tanks} \underbrace{y(t) = d(t)}_{\text{Two tanks}}$$

Inputs, q_1 and q_2 , both measurable and manipulable

 \leadsto They influence the liquid levels in the tanks

Output, $d = h_1 - h_2$, measurable, not manipulable

 \leadsto It is influenced by the inputs

State variables, V_1 and V_2 , not measurable and not manipulable

- → They evolve according to own dynamics
- → They are also influenced by the inputs

Calculus, review

Solution using a Taylor expansion

Second- and higher-order

first order ODEs

Second- and higher-order systems (cont.)

For an incompressible fluid, by mass conservation

$$\begin{cases} \frac{\mathrm{d}V_1(t)}{\mathrm{d}t} = q_1(t) - q_2(t) \\ \frac{\mathrm{d}V_2(t)}{\mathrm{d}t} = q_2(t) - q_3(t) \end{cases} \longrightarrow \begin{cases} \dot{h}_1(t) = \frac{1}{B}q_1(t) - \frac{1}{B}q_2(t) \\ \dot{h}_2(t) = \frac{1}{B}q_2(t) - \frac{k}{B}h_2(t) \end{cases}$$

By taking the first derivative of $y(t) = h_1(t) - h_2(t)$ and rearranging, we obtained

$$\dot{y}(t) = \dot{h}_1(t) - \dot{h}_2(t) = \frac{1}{B}u_1(t) - \frac{2}{B}u_2(t) + \frac{k}{B}[h_1(t) - y(t)]$$

By taking the second derivative of y(t) and rearranging, we obtained

$$\ddot{y}(t) = \frac{1}{B}\dot{u}_1(t) - \frac{2}{B}\dot{u}_2(t) + \frac{k}{B}\dot{h}_1(t) - \frac{k}{B}\dot{y}(t)$$

$$= \frac{1}{B}\dot{u}_1(t) - \frac{2}{B}\dot{u}_2(t) + \underbrace{\frac{k}{B^2}u_1(t) - \frac{k}{B^2}u_2(t)}_{\frac{k}{B}\dot{h}_1(t)} - \frac{k}{B}\dot{y}(t)$$

Calculus, review

Solution using a

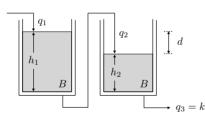
Second- and higher-order

first order ODEs

Second- and higher-order systems (cont.)

Rearranging terms, the IO system's representation is an ordinary differential equation

$$\label{eq:control_eq} \leadsto \quad \ddot{y}(t) + \frac{k}{B} \dot{y}(t) - \frac{1}{B} \dot{u}_1(t) + \frac{2}{B} \dot{u}_2(t) - \frac{k}{B^2} u_1(t) + \frac{k}{B} u_2(t) = 0$$



Suppose that the inputs are zero

$$u_1(t) = q_1(t) = 0$$

 $u_2(t) = q_2(t) = 0$

Also their derivatives are zero

$$\ddot{y}(t) + \frac{k}{B}\dot{y}(t) = +\frac{1}{B}\dot{u}_1(t) - \frac{2}{B}\dot{u}_2(t) + \frac{k}{B^2}\dot{u}_1(t) - \frac{k}{B}\dot{u}_2(t) \quad \leadsto \quad \ddot{y}(t) + \frac{k}{B}\dot{y}(t) = 0$$

What's y(t), for some y(0) and $\dot{y}(0)$?

Second- and higher-order systems (cont.)

Calculus, review

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODE

Homogeneous equation

Consider the ordinary differential equation of a IO model (linear and time invariant)

$$\alpha_n \frac{\mathrm{d}^n y(t)}{\mathrm{d}t^n} + \dots + \alpha_1 \frac{\mathrm{d}y(t)}{\mathrm{d}t} + \alpha_0 y(t) = \beta_m \frac{\mathrm{d}^m u(t)}{\mathrm{d}t^m} + \dots + \beta_1 \frac{\mathrm{d}u(t)}{\mathrm{d}t} + \beta_0 u(t)$$

Let the RHS of be zero, define the homogenous equation associated to the model

$$\Rightarrow a_n \frac{\mathrm{d}^n y(t)}{\mathrm{d}t^n} + \dots + a_1 \frac{\mathrm{d}y(t)}{\mathrm{d}t} + a_0 y(t) = 0$$

The solution y(t) to the homogeneous equation can be defined as the system response (the output) for an input u(t) that is null for $t \ge t_0$ and for given initial conditions

Input- or force-free response

• We may denote it as h(t)



Calculus, review

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From high to first order ODE

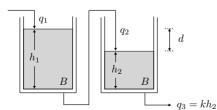
Second- and higher-order systems (cont.)

Example

$$\ddot{y}(t) + \frac{k}{B}\dot{y}(t) = 0$$

Interest is in y(t) for some given initial conditions y(0) and $\dot{y}(0)$ (assuming u(t)=0)

- We want to use the Taylor expansion of the solution y(t)
- For simplicity, let k/B = 1
- \rightarrow We solve $\ddot{y}(t) + \dot{y}(t) = 0$



The differential equation is second-order

→ Initial position

$$y(t=0) = y(0)$$

 \leadsto Initial velocity

$$\dot{y}(t=0) = \dot{y}(0)$$

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Calculus, review

Solution using a Taylor expansion

Second- and higher-order

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Second- and higher-order systems (cont.)

$$\ddot{y}(t) + \dot{y}(t) = 0$$

• We assume that y(t) can be expressed by using a Taylor series expansion

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + c_4 t^4 + \cdots$$

• We can compute the first derivative of the assumed solution y(t), $\dot{y}(t)$

$$\dot{y}(t) = c_1 + 2c_2t + 3c_3t^2 + 4c_4t^3 + 5c_5t^4 + \cdots$$

• We compute the second derivative of the assumed solution y(t), $\ddot{y}(t)$

$$\ddot{y}(t) = 2c_2 + 2 \cdot 3c_3t + 3 \cdot 4c_4t^2 + 4 \cdot 5c_5t^3 + \cdots$$

• Then, proceed by substituting function and derivatives into the ODEs

Calculus, review

Solution using a

Second- and higher-order

first order ODE

Second- and higher-order systems (cont.)

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + c_4 t^4 + \cdots$$

After considering initial conditions y(t=0) = y(0) and $\dot{y}(t=0) = \dot{y}(0)$, we have

$$y(t = 0) = c_0 + \cot t + \cot^2 + \cot^3 + \cot^4 + \dots = y(0)$$

$$\sim c_0 = y(0)$$

$$\dot{y}(t = 0) = c_1 + 2\cot t + 3\cot^2 + 4\cot^3 + 5\cot^4 + \dots = \dot{y}(0)$$

$$\sim c_1 = \dot{y}(0)$$

• Then, from the ordinary differential equation $\ddot{y}(t) + \dot{y}(t) = 0$ we have

$$\underbrace{2c_2 + 2 \cdot 3c_3t + 3 \cdot 4c_4t^2 + 4 \cdot 5c_5t^3 + \cdots}_{\ddot{y}(t)}$$

$$= \underbrace{-(c_1 + 2c_2t + 3c_3t^2 + 4c_4t^3 + 5c_5t^4 + \cdots)}_{\dot{y}(t)}$$

Calculus, review

Solution using a

Second- and higher-order

first order ODEs

Second- and higher-order systems (cont.)

$$\underbrace{\frac{2c_2 + 2 \cdot 3c_3t + 3 \cdot 4c_4t^2 + 4 \cdot 5c_5t^3 + \cdots}_{\ddot{y}(t)}}_{\ddot{y}(t)} = \underbrace{-c_1 - 2c_2t - 3c_3t^2 - 4c_4t^3 - 5c_5t^4 - \cdots}_{\dot{y}(t)}$$

By equating the coefficients to satisfy the identity and rearranging, we obtain

•
$$c_0 = y(0)$$

• $c_1 = \dot{y}(0)$
• $c_2 = -\frac{1}{2}\dot{y}(0)$
• $c_3 = +\frac{1}{3!}\dot{y}(0)$
• $c_4 = -\frac{1}{4!}\dot{y}(0)$
• $c_5 = +\frac{1}{5!}\dot{y}(0)$
• ...

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Taylor expansion

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first order ODE

Second- and higher-order systems (cont.)

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + c_4 t^4 + \cdots$$

Substituting the coefficients in the assumed (Taylor's) solution form, we obtain

$$y(t) = y(0) + \dot{y}(0)t - \frac{1}{2}\dot{y}(0)t^2 + \frac{1}{3!}\dot{y}(0)t^3 - \frac{1}{4!}\dot{y}(0)t^4 + \frac{1}{5!}\dot{y}(0)t^5 - \cdots$$

$$= y(0) - \dot{y}(0)\left(\underbrace{-t + \frac{1}{2}t^2 - \frac{1}{3!}t^3 - \frac{1}{4!}t^4 + \frac{1}{5!}t^5 - \cdots}_{-1+e^{-t}}\right)$$

$$= \underbrace{y(0) + \dot{y}(0)}_{-}\underbrace{-\dot{y}(0)}_{-}e^{-t}$$

$$= \underbrace{y(0) + \dot{y}(0)}_{k_1} \underbrace{-\dot{y}(0)}_{k_2} e^{-t}$$
$$= k_1 + k_2 e^{-t}$$

With k_1 and k_2 constant values depending on the initial conditions

$$\rightarrow k_1 = y(0) + \dot{y}(0)$$

$$\rightarrow k_2 = -\dot{y}(0)$$

We used
$$e^{-t} = 1 + (-t) + \frac{(-t)^2}{2} + \frac{(-t)^3}{3!} + \cdots$$

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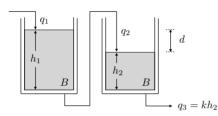
Solution using a

Second- and higher-order

first order ODEs

Second- and higher-order systems (cont.)

Example



$$\ddot{y}(t) + \dot{y}(t) = 0$$

For simplicity, we let $\frac{k}{B} = 1$ and obtained the system evolution by solving the ODE

$$y(t) = y(0) + \dot{y}(0) - \dot{y}(0)e^{-t}$$

Calculus, review

Solution using a

Second- and higher-order

first order ODE

Second- and higher-order systems (cont.)

Higher-order systems

Consider the general linear time-invariant system and homogeneous (with no inputs)

$$\alpha_n \frac{\mathrm{d}^n y}{\mathrm{d}t^n} + \alpha_{n-1} \frac{\mathrm{d}^{n-1} y}{\mathrm{d}t^{n-1}} + \dots + \alpha_2 \frac{\mathrm{d}^2 y}{\mathrm{d}t^2} + \alpha_1 \frac{\mathrm{d}y}{\mathrm{d}t} + \alpha_0 y = 0$$

Or, equivalently

$$\alpha_n y^{(n)} + \alpha_{n-1} y^{(n-1)} + \dots + \alpha_1 \ddot{y} + \alpha_0 y = 0$$

We consider an alternative to assuming that the solution is written as Taylor expansion

Instead of using Taylor expansions, we assume that the solution is given by $y(t) = e^{\lambda t}$

• (Which is not very different, in practice)

$$y(t) = c_0 + c_1 t + c_2 t^2 + c_3 t^3 + \cdots$$
$$e^{\lambda t} = 1 + \lambda t + \frac{(\lambda t)^2}{2!} + \frac{(\lambda t)^3}{3!} + \cdots$$

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Second- and higher-order systems (cont.)

If we set the solution to be $y(t) = e^{\lambda t}$, then we can easily compute its derivatives

$$\rightsquigarrow \dot{y}(t) = \frac{\lambda}{\lambda} e^{\lambda t}$$

$$\rightsquigarrow \ddot{y}(t) = \frac{\lambda^2}{\lambda^2} e^{\lambda t}$$

~→ ··

$$\rightsquigarrow y^{(n)}(t) = \frac{\lambda^n e^{\lambda t}}{}$$

These functions can be substituted into homogeneous linear time-invariant ODEs

$$\alpha_n y^{(n)} + \alpha_{n-1} y^{(n-1)} + \dots + \alpha_1 \dot{y} + \alpha_0 y = 0$$

By substituting the assumed solution and derivatives into the differential equation

$$\Rightarrow \left[\alpha_n \lambda^n + \alpha_{n-1} \lambda^{n-1} + \dots + \alpha_2 \lambda^2 + \alpha_1 \lambda + \alpha_0 \right] e^{\lambda t} = 0$$

The identity is verified for all n values of λ solving the characteristic equation

$$\underbrace{\alpha_n \lambda^n + \alpha_{n-1} \lambda^{n-1} + \dots + \alpha_2 \lambda^2 + \alpha_1 \lambda + \alpha_0}_{\text{Characteristic polynomial}} = 0$$

Characteristic equation

Second- and higher-order systems (cont.)

Calculus, review

 $\alpha_n \lambda^n + \alpha_{n-1} \lambda^{n-1} + \dots + \alpha_2 \lambda^2 + \alpha_1 \lambda + \alpha_0 = 0$

$$a_n \times + a_{n-1} \times + \cdots + a_{2N} + a_{1N} + a_0 = 0$$

The characteristic equation has n solutions, or roots, collected in set $\{\lambda_1, \lambda_2, \dots, \lambda_n\}$

- They can be real and/or complex (and associated complex-conjugate) numbers
- They can be positive and/or negative, distinct and repeated (multiplicity)

$$\alpha_n y^{(n)} + \alpha_{n-1} y^{(n-1)} + \dots + \alpha_1 \dot{y} + \alpha_0 y = 0$$

For distinct (real and complex) roots, the ODE solution has the simple form

$$y(t) = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} + \dots + c_n e^{\lambda_n t}$$
$$= \sum_{i=1}^n c_i e^{\lambda_i t}$$

The solution is a sum of exponential functions, each weighted by coefficients

 \bullet The coefficients are determined from the n initial conditions

Second- and higher-order systems (cont.)

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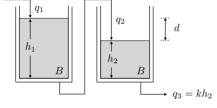
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Example



$$\ddot{y}(t) + \frac{k}{B}\dot{y}(t) = 0$$

For simplicity, let $\frac{k}{B}=1$ and obtained the system evolution of $\ddot{y}(t)+\dot{y}(t)=0$

$$y(t) = y(0) + \dot{y}(0) - \dot{y}(0)e^{-t}$$

Start by assuming a solution $y(t) = e^{\lambda t}$ and computing its derivatives $\dot{y}(t)$ and $\ddot{y}(t)$

- \leadsto Substitute then in the original system ODE
- ightarrow Compute the characteristic equation
- → Solve the characteristic equation

Calculus, review

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Second- and higher-order systems (cont.)

Definition

Characteristic polynomial

Consider the homogeneous part of the linear and time-invariant differential equation

$$\alpha_n \frac{\mathrm{d}^n y(t)}{\mathrm{d}t^n} + \dots + \alpha_1 \frac{\mathrm{d}y(t)}{\mathrm{d}t} + \alpha_0 y(t) = 0$$

The characteristic polynomial is a *n*-order polynomial in the variable λ whose coefficients correspond to the coefficients $\{\alpha_0, \alpha_1, \dots, \alpha_n\}$ of the homogeneous equation

$$P(\lambda) = \alpha_n \lambda^n + \alpha_{n-1} \lambda^{n-1} + \dots + \alpha_1 \lambda + \alpha_0$$
$$= \sum_{i=0}^n \alpha_i \lambda^i$$

Any polynomial of order n with real coefficients has n real or complex-conjugate roots

• The roots are solutions of the characteristic equation

$$P(\lambda) = \sum_{i=0}^{n} \alpha_i \lambda^i = 0$$

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Second- and higher-order

Second- and higher-order systems (cont.)

In general, there are $r \leq n$ distinct roots p_i , each with multiplicity ν_i



$$\leadsto$$
 If $i \neq j$, then $p_i \neq p_j$

$$\rightsquigarrow \sum_{i=1}^r \nu_i = n$$

Consider the case in which all roots have multiplicity equal one (no repetitions)

$$\rightsquigarrow \quad \overbrace{p_1 \quad p_2 \quad \cdots \quad p_{n-1} \quad p_n}^n$$

$$\rightsquigarrow$$
 If $i \neq j$, then $p_i \neq p_j$

$$\rightsquigarrow \nu_i = 1$$
, for every i

Second- and higher-order systems (cont.)

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Definition

Modes

Let p be one of the roots with multiplicity ν of the characteristic polynomial

The modes associated to that root are the ν functions of time

$$\Rightarrow$$
 $e^{pt}, te^{pt}, t^2 e^{pt}, \cdots, t^{\nu-1} e^{pt}$

A system with a n-order characteristic polynomial has n modes

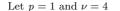
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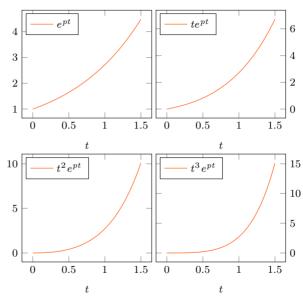
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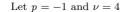


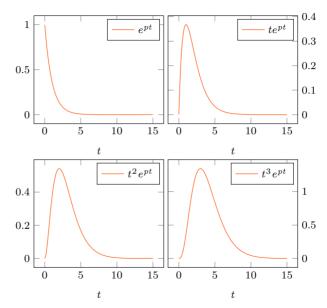
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Second- and higher-order systems (cont.)

The modes from the characteristic polynomial, the mixing coefficients are parameters

$$h(t) = \sum_{i=1}^{r} \left(\sum_{k=0}^{\nu_i - 1} A_{i,k} t^k e^{p_i t} \right)$$

The coefficients determine the force-free evolution, from every possible initial condition

Theorem

Solution of the homogeneous equation

Consider the homogeneous equation

$$a_n \frac{d^n y(t)}{dt^n} + \dots + a_1 \frac{dy(t)}{dt} + a_0 y(t) = 0$$

A real function h(t) is the solution of a homogeneous linear time-invariant differential equation if and only if h(t) can be written as a linear combination of the modes

$$\rightarrow h(t) = \sum_{i=1}^{r} \left(\sum_{k=0}^{\nu_i - 1} A_{i,k} t^k e^{p_i t} \right)$$

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Second- and higher-order systems (cont.)

Modes are functions of time, their linear combinations are a family of functions of time

- The family is parameterised by the coefficients of the combination
- (Different coeffcients correspond to different family members)

Definition

Linear combinations of modes

A linear combination of the n modes is a function h(t), a weighted sum of the modes

• Each mode is weighted by some coefficient

Each individual root p_i with multiplicity ν_i is associated to a combination of ν_i terms

$$A_{i,0}e^{p_it} + A_{i,1}te^{p_it} + \dots + A_{i,\nu_i-1}t^{\nu_i-1}e^{p_it} = \underbrace{\sum_{k=0}^{\nu_i-1} A_{i,k}t^k e^{p_it}}_{\text{root } p_i}$$

There is a total of r distinct roots, $i = 1, \ldots, r$

Second- and higher-order systems (cont.)

Calculus, review

From high to first order ODE

$$A_{i,0}e^{p_it} + A_{i,1}te^{p_it} + \dots + A_{i,\nu_i-1}t^{\nu_i-1}e^{p_it} = \underbrace{\sum_{k=0}^{\nu_i-1} A_{i,k}t^k e^{p_it}}_{\text{root } p_i}$$

As there are r distinct roots, i = 1, ..., r, the complete linear combination of modes

$$h(t) = \underbrace{\sum_{k=0}^{\nu_1 - 1} A_{1,k} t^k e^{p_1 t}}_{\text{root } p_1} + \underbrace{\sum_{k=0}^{\nu_2 - 1} A_{2,k} t^k e^{p_2 t}}_{\text{root } p_2} + \dots + \underbrace{\sum_{k=0}^{\nu_r - 1} A_{r,k} t^k e^{p_r t}}_{\text{root } p_r}$$

$$\Rightarrow = \sum_{i=1}^{r} \left(\sum_{k=0}^{\nu_i - 1} A_{i,k} t^k e^{p_i t} \right)$$

Consider the case in which all roots (n) have multiplicity equal to one (no repetitions)

$$\rightarrow$$
 $h(t) = A_1 e^{p_1 t} + A_2 e^{p_2 t} + \dots + A_n e^{p_n t} = \sum_{i=1}^n A_i e^{p_i t}$

(We have omitted the second subscript of coefficients A)

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Second- and higher-order systems (cont.)

Example

Consider the following homogenous differential equation

$$3\frac{d^4y(t)}{dt^4} + 21\frac{d^3y(t)}{dt^3} + 45\frac{d^2y(t)}{dt^2} + 39\frac{dy(t)}{dt} + 12y(t) = 0$$

The associated characteristic polynomial

$$P(\lambda) = 3\lambda^4 + 21\lambda^3 + 45\lambda^2 + 39\lambda + 12 = 3(\lambda + 1)^3(\lambda + 4)$$

The characteristic equation has four roots

→ The system has four modes

$$p_1 = -1, \quad (\nu_1 = 3) \quad \leadsto \quad \begin{cases} e^{-t} \\ te^{-t} \\ t^2 e^{-t} \end{cases}$$
 $p_2 = -4, \quad (\nu_2 = 1) \quad \leadsto \quad \begin{cases} e^{-4t} \\ e^{-4t} \end{cases}$

The family of functions h(t) is given as a linear combination of the modes

$$h(t) = \underbrace{A_{1,0}e^{-t} + A_{1,1}te^{-t} + A_{1,2}t^{2}e^{-t}}_{\text{root } p_{1}} + \underbrace{A_{2}e^{-4t}}_{\text{root } p_{2}}$$

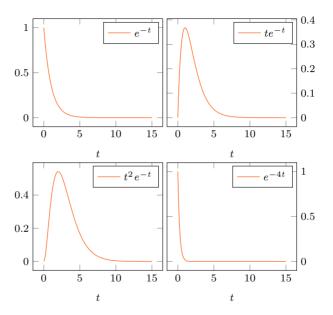
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Second- and higher-order systems (cont.)

Complex and conjugate roots

A characteristic polynomial P(s) with complex roots will have complex signal modes

$$h(t) = \sum_{i=1}^{r} \left(\sum_{k=0}^{\nu_i - 1} A_{i,k} t^k e^{p_i t} \right)$$
 (Yet, their combination must be a real function)

Let P(s) be a characteristic polynomial with roots $p_i = \alpha_i + j\omega_i$ of multiplicity ν_i

• Let $p_i' = \alpha_i - j\omega_i$ with multiplicity $\nu_i' = \nu_i$ be the conjugate complex root

The contribution of each pair (p_i, p'_i) to the linear combination can be re-written

$$\sum_{k=0}^{\nu_i-1} M_{i,k} t^k e^{\alpha_i t} \cos(\omega_i t + \phi_{i,k}) \quad \text{(Coefficients } M_{i,k} \text{ and } \phi_{i,k})$$

Or, equivalently

$$\sum_{k=0}^{\nu_i-1} \left[B_{i,k} \, t^k e^{\alpha_i t} \cos(\omega_i t) + C_{i,k} \, t^k e^{\alpha_i t} \sin(\omega_i t) \right] \quad \text{(Coefficients $B_{i,k}$ and $C_{i,k}$)}$$

Calculus, review

Solution using a

Second- and higher-order

first order ODI

Second- and higher-order systems (cont.)

The solution equations

$$h(t) = \sum_{i=1}^{R} \sum_{k=0}^{\nu_i - 1} A_{i,k} t^k e^{p_i t} + \sum_{i=R+1}^{R+S} \sum_{k=0}^{\nu_i - 1} M_{i,k} t^k e^{\alpha_i t} \cos(\omega_i t + \phi_{i,k})$$

$$\left(\leadsto \sum_{i=1}^{R} A_i e^{p_i t} + \sum_{i=R+1}^{R+S} M_i e^{\alpha_i t} \cos(\omega_i t + \phi_i) \right)$$

The solution equations

$$\begin{split} h(t) &= \sum_{i=1}^{R} \sum_{k=0}^{\nu_{i}-1} A_{i,k} t^{k} e^{p_{i}t} + \sum_{i=R+1}^{R+S} \sum_{k=0}^{\nu_{i}-1} \left[B_{i,k} t^{k} e^{\alpha_{i}t} \cos(\omega_{i}t) + C_{i,k} t^{k} e^{\alpha_{i}t} \sin(\omega_{i}t) \right] \\ &\left(\leadsto \sum_{i=1}^{R} A_{i} e^{p_{i}t} + \sum_{i=R+1}^{R+S} \left[B_{i} e^{\alpha_{i}t} \cos(\omega_{i}t) + C_{i} e^{\alpha_{i}t} \sin(\omega_{i}t) \right] \right) \end{split}$$

They provide the parametric structure of the linear combination and are all equivalent

Second- and higher-order systems (cont.)

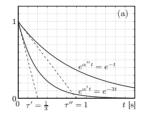
Calculus, review

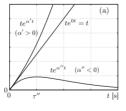
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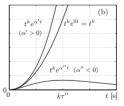
Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs







Second- and higher-order systems (cont.)

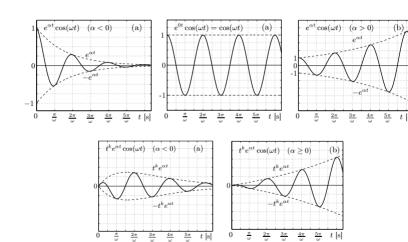
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From high to first order ODEs

From high to first order ODEs Ordinary differential equation

From high-order ODEs to systems of first-order ODEs

Calculus, review

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs

Consider the general linear time-invariant n-order system and homogeneous (no inputs)

$$\alpha_n \frac{\mathrm{d}^n y}{\mathrm{d}t^n} + \alpha_{n-1} \frac{\mathrm{d}^{n-1} y}{\mathrm{d}t^{n-1}} + \dots + \alpha_2 \frac{\mathrm{d}^2 y}{\mathrm{d}t^2} + \alpha_1 \frac{\mathrm{d}y}{\mathrm{d}t} + \alpha_0 y = 0$$

Or, equivalently

$$\alpha_n y^{(n)} + \alpha_{n-1} y^{(n-1)} + \dots + \alpha_1 \ddot{y} + \alpha_0 y = 0$$

We can convert the n-order equation into a set of n first order equations, and solve it

As a preprocessing step, we start by dividing all the coefficients by α_n

$$y^{(n)}(t) + \underbrace{\mathbf{a}_{n-1}}_{\alpha_{n-1}/\alpha_n} y^{(n-1)}(t) + \dots + \underbrace{\mathbf{a}_{2}}_{\alpha_{2}/\alpha_n} \ddot{y}(t) + \underbrace{\mathbf{a}_{1}}_{\alpha_{1}/\alpha_n} \dot{y}(t) + \underbrace{\mathbf{a}_{0}y}_{\alpha_{0}/\alpha_n} = 0$$

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

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Second- and higher-order

From high to first order ODEs

$$y^{(n)}(t) + \mathbf{a_{n-1}}y^{(n-1)}(t) + \dots + \mathbf{a_{2}}\ddot{y}(t) + \mathbf{a_{1}}\dot{y}(t) + \mathbf{a_{0}}y = 0$$

Firstly, we introduce a set of n new variables $x(t) = [x_1(t), x_2(t), \dots, x_n(t)]'$

$$x_1(t) = y(t)$$

$$x_2(t) = \dot{y}(t)$$

$$x_3(t) = \ddot{y}(t)$$

$$\cdots = \cdots$$

$$x_{n-1}(t) = y^{(n-2)}(t)$$

$$x_n(t) = y^{(n-1)}(t)$$

Then, we introduce their first-order derivatives $\dot{x}(t) = [\dot{x}_1(t), \dot{x}_2(t), \dots, \dot{x}_n(t)]'$

$$\dot{x}_1(t) = \dot{y}(t) = x_2(t)
\dot{x}_2(t) = \ddot{y}(t) = x_3(t)
\dot{x}_3(t) = \dddot{y}(t) = x_4(t)
\dots = \dots
\dot{x}_{n-1}(t) = y^{(n-1)}(t) = x_n(t)
\dot{x}_n(t) = y^{(n)}(t) = -a_{n-1}x_n(t) - a_{n-1}x_{n-1}(t) - \dots - a_2x_3(t) - a_1x_2(t) - a_0x_1(t)$$

From high-order ODEs to first-order ODEs (cont.)

 $\dot{x}_1(t) = \dot{y}(t) = x_2(t)$

Calculus, review

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Taylor expansion

Second- and higher-order

From high to first order ODEs

$$\dot{x}_2(t) = \ddot{y}(t) = x_3(t)
\dot{x}_3(t) = \ddot{y}(t) = x_4(t)
\dots = \dots
\dot{x}_{n-1}(t) = y^{(n-1)}(t) = x_n(t)
\dot{x}_n(t) = y^{(n)}(t) = -a_{n-1}x_n(t) - a_{n-1}x_{n-1}(t) - \dots - a_2x_3(t) - a_1x_2(t) - a_0x_1(t)$$

That is, we get the set of linear equations with explicit dependences between terms

$$\begin{aligned} \dot{x}_1 &= 0x_1 + \frac{1x_2}{2} + 0x_3 + 0x_4 + \dots + 0x_n \\ \dot{x}_2 &= 0x_1 + 0x_2 + \frac{1x_3}{2} + 0x_4 + \dots + 0x_n \\ \dot{x}_3 &= 0x_1 + 0x_2 + 0x_3 + \frac{1x_4}{2} + \dots + 0x_n \\ \dots &= \dots \\ \dot{x}_{n-1} &= 0x_1 + 0x_2 + 0x_3 + 0x_4 + \dots + 1x_n \\ \dot{x}_n &= -a_0x_1 - a_1x_2 - a_2x_3 - a_3x_4 - \dots - a_{n-1}x_n \end{aligned}$$

From high-order ODEs to first-order ODEs (cont.)

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From high to first order ODEs

$$\dot{x}_1 = 0x_1 + \frac{1}{2} + 0x_3 + 0x_4 + \dots + 0x_n
\dot{x}_2 = 0x_1 + 0x_2 + \frac{1}{2} + 0x_4 + \dots + 0x_n
\dot{x}_3 = 0x_1 + 0x_2 + 0x_3 + \frac{1}{2} + \dots + 0x_n
\dots = \dots
\dot{x}_{n-1} = 0x_1 + 0x_2 + 0x_3 + 0x_4 + \dots + 1x_n
\dot{x}_n = -\frac{a_0}{a_0} x_1 - \frac{a_1}{a_1} x_2 - \frac{a_2}{a_2} x_3 - \frac{a_3}{a_3} x_4 - \dots - \frac{a_{n-1}}{a_{n-1}} x_n$$

We can write a $\dot{x}(t)$ as a matrix-vector multiplication Ax(t), a system of equations

$$\underbrace{ \begin{bmatrix} \dot{x_1} \\ \dot{x_2} \\ \dot{x_3} \\ \vdots \\ x_{n-1} \\ \dot{x_n} \end{bmatrix}}_{\dot{x}(t)} = \underbrace{ \begin{bmatrix} 0 & 1 & 0 & 0 & \cdots & 0 \\ 0 & 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & 0 & 1 & \cdots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & 0 & \cdots & 1 \\ -a_0 & -a_1 & -a_2 & -a_3 & \cdots & -a_{n-1} \end{bmatrix} }_{A} \underbrace{ \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ \vdots \\ x_{n-1} \\ x_n \end{bmatrix}}_{x(t)}$$

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

Intro to ODE

Solution using a Taylor expansion

Second- and higher-orde

From high to first order ODEs

Example

Consider a linear and time-invariant homogeneous system representation

$$\ddot{y} + a_2 \ddot{y} + a_1 \dot{y} + a_0 y = 0$$

The system in IO representation is a third-order ODE

We are interested in formulating the system as a matrix system

- The system is third-order (max derivative of y)
- A system of 3 first-order ODEs

We first introduce three dummy variables $x(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$

Then, we get

$$x_1(t) = y(t)$$

$$x_2(t) = \dot{y}(t)$$

$$x_3(t) = \ddot{y}(t)$$

From high-order ODEs to first-order ODEs (cont.)

$$\ddot{y} + a_2 \ddot{y} + a_1 \dot{y} + a_0 y = 0$$

We compute the derivatives of the x(t) variables with respect to time $\dot{x}(t) = \begin{pmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \\ \dot{x}_3(t) \end{pmatrix}$ • Remember that we defined them as $x(t) = \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_2(t) \end{bmatrix} = \begin{bmatrix} y(t) \\ \dot{y}(t) \\ \dot{y}(t) \\ \ddot{y}(t) \end{bmatrix}$

That is,

$$\begin{split} \dot{x_1}(t) &= \dot{y}(t) \\ &= x_2(t) \\ \dot{x_2}(t) &= \ddot{y}(t) \\ &= x_3(t) \\ \dot{x_3}(t) &= \dddot{y}(t) \\ &= -a_2 x_3(t) - a_1 x_2(t) - a_0 x_1(t) \end{split}$$

In matrix form, we can write

$$\begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \\ \dot{x}_3(t) \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 \end{bmatrix} \begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}$$

From high to first order ODEs

Calculus, review

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs

From high-order ODEs to first-order ODEs (cont.)

Example

Consider the following linear and time-invariant homogeneous system

$$3\frac{d^4y(t)}{dt^4} + 21\frac{d^3y(t)}{dt^3} + 45\frac{d^2y(t)}{dt^2} + 39\frac{dy(t)}{dt} + 12y(t) = 0$$

The system in IO representation is a forth-order ODE

→ A system of 4 first-order ODEs

We first divide by the leading coefficient $(a_4 = 3)$

$$\frac{\mathrm{d}^4 y(t)}{\mathrm{d}t^4} + \underbrace{7}_{a_3} \frac{\mathrm{d}^3 y(t)}{\mathrm{d}t^3} + \underbrace{15}_{a_2} \frac{\mathrm{d}^2 y(t)}{\mathrm{d}t^2} + \underbrace{13}_{a_1} \frac{\mathrm{d}y(t)}{\mathrm{d}t} + \underbrace{4}_{a_0} y(t) = 0$$

By using the general expression derived earlier,

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \\ \dot{x}_4 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 & -a_3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix}$$

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

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\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \\ \dot{x}_4 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 & -a_3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix}
= \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -4 & -13 & -15 & -7 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix}
```

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

Solution using a

Second- and higher-order

From high to first order ODEs

Example

Consider the second-order linear and homogeneous differential equation

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

The initial conditions (at t = 0),

$$\begin{cases} y(0) = 2\\ \dot{y}(0) = -3 \end{cases}$$

We want to determine in its solution

We start by assuming that its solution is given by function $y(t) = e^{\lambda t}$

Then, we compute the derivatives of the assumed solution

• Up to order n=2

Then, we have

$$\dot{y}(t) = \lambda e^{\lambda t}$$
$$\ddot{y}(t) = \lambda^2 e^{\lambda t}$$

Calculus, review

Solution using a

Second- and higher-order

From high to first order ODEs

From high-order ODEs to first-order ODEs (cont.)

Now we substitute the solution and its derivatives into the original equation, to get

$$\underbrace{\ddot{y}(t)}_{\lambda^2 e^{\lambda t}} + 3 \underbrace{\dot{y}(t)}_{\lambda e^{\lambda t}} + 2 \underbrace{y(t)}_{e^{\lambda t}} = 0$$

Rearranging, we have

$$\lambda^{2} e^{\lambda t} + 3\lambda e^{\lambda t} + 2e^{\lambda t} = 0$$
$$e^{\lambda t} (\lambda^{2} + 3\lambda + 2) = 0$$

The roots of the characteristic polynomial $\lambda^2 + 3\lambda + 2 = (\lambda + 2)(\lambda + 1) = 0$,

$$\begin{cases} \lambda_1 = -1 \\ \lambda_2 = -2 \end{cases}$$
 (Real, with negative real part)

We formulate the general solution,

$$y(t) = C_1 e^{\lambda_1 t} + C_2 e^{\lambda_2 t}$$
$$= C_1 e^{(-1)t} + C_2 e^{(-2)t}$$

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

Intro to ODEs

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs

$$y(t) = C_1 e^{(-1)t} + C_2 e^{(-2)t}$$

By using the initial conditions, we determine the unknown coefficients,

$$\begin{cases} y(t=0) = C_1 \underbrace{e^{-t}}_{=1} + C_2 \underbrace{e^{-2t}}_{=1} = 2 \\ \dot{y}(t=0) = -C_1 \underbrace{e^{-t}}_{=1} - 2C_2 \underbrace{e^{-2t}}_{=1} = -3 \end{cases}$$

We can then solve for C_1 and C_2 , to get the pair of coefficients

$$\rightsquigarrow C_1 = 1$$

$$\sim$$
 $C_2 = 1$

The solution is stable, as it is the sum of stable exponentials

$$y(t) = e^{-t} + e^{-2t}$$

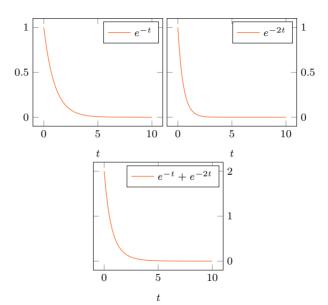
Calculus, revie

Intro to ODEs

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs



Second- and higher-order

From high to first order ODEs We can reformulate $\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$ as a system of 2 first-order equations

1 We start by introducing two dummy variables

$$\begin{cases} x_1 = y^{(0)} \\ x_2 = y^{(1)} \end{cases}$$

2 We compute the time derivatives

$$\begin{cases} \dot{x_1} = y^{(1)} = x_2 \\ \dot{x_2} = y^{(2)} = -3x_2 - 2x_1 \end{cases}$$

3 Rewriting in matrix form

$$\underbrace{\begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \end{bmatrix}}_{\dot{x}(t)} = \underbrace{\begin{bmatrix} 0 & 1 \\ -2 & -3 \end{bmatrix}}_{A} \underbrace{\begin{bmatrix} x_1(t) \\ x_2(t) \end{bmatrix}}_{x(t)}$$

Note how the two eigenvalues of A equal the roots of the characteristic polynomial

Calculus, review

Solution using a

Taylor expansion

Second- and higher-order

From high to first order ODEs

From high-order ODEs to first-order ODEs (cont.)

Example

Consider the second-order linear and homogeneous differential equation

$$\ddot{y}(t) \underbrace{-3\dot{y}(t)}_{\text{flipped sign}} +2y(t) = 0$$

The same initial conditions (at t = 0),

$$\begin{cases} y(0) = 2\\ \dot{y}(0) = -3 \end{cases}$$

We want to determine in its solution

We start by assuming that its solution is given by function $y(t) = e^{\lambda t}$

Then, we compute the derivatives up to order n = 2, to get

$$\dot{y}(t) = \lambda e^{\lambda t}$$

$$\ddot{y}(t) = \lambda^2 e^{\lambda}$$

Calculus, review

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Taylor expansion

Second- and higher-order

From high to first order ODEs

From high-order ODEs to first-order ODEs (cont.)

Now we substitute the solution and its derivatives into the original equation, to get

$$\underbrace{\ddot{y}(t)}_{\lambda^2 e^{\lambda t}} - 3 \underbrace{\dot{y}(t)}_{\lambda e^{\lambda t}} + 2 \underbrace{y(t)}_{e^{\lambda t}} = 0$$

Rearranging, we have

$$\lambda^{2} e^{\lambda t} - 3\lambda e^{\lambda t} + 2e^{\lambda t} = 0$$
$$e^{\lambda t} (\lambda^{2} - 3\lambda + 2) = 0$$

The roots of the characteristic polynomial $\lambda^2 - 3\lambda + 2 = (\lambda - 2)(\lambda - 1) = 0$,

$$\begin{cases} \lambda_1 = 1 \\ \lambda_2 = 2 \end{cases}$$
 (Real, with positive real part)

We formulate the general solution,

$$y(t) = C_1 e^{\lambda_1 t} + C_2 e^{\lambda_2 t}$$
$$= C_1 e^{(+1)t} + C_2 e^{(+2)t}$$

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

Intro to ODE

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs

$$y(t) = C_1 e^t + C_2 e^{2t}$$

The solution is unstable because at least one of the exponentials in the sum is unstable By using the initial conditions, we can still determine the unknown coefficients,

$$\begin{cases} y(t=0) = C_1 \underbrace{e^t}_{=1} + C_2 \underbrace{e^{2t}}_{=1} = 2 \\ \dot{y}(t=0) = C_1 \underbrace{e^t}_{=1} + 2C_2 \underbrace{e^{2t}}_{=1} = -3 \end{cases}$$

We can then solve for C_1 and C_2 , to get the pair of coefficients

$$\sim$$
 $C_1 = 7$

$$\sim$$
 $C_2 = -5$

The solution,

$$y(t) = 7e^t + 5e^{2t}$$

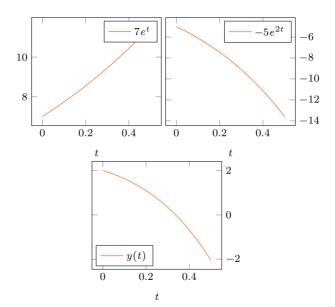
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Second- and higher-order

From high to first order ODEs



From high-order ODEs to first-order ODEs (cont.)

Calculus, review

Intro to ODEs

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Second- and higher-order

From high to first order ODEs

Example

Consider the second-order linear and homogeneous differential equation

$$\ddot{y}(t) + 1\dot{y}(t) - 2y(t) = 0$$

The initial conditions (at t = 0),

$$\begin{cases} y(0) = 2\\ \dot{y}(0) = -3 \end{cases}$$

We want to determine in its solution

Calculus, review

Solution using a Taylor expansion

Second- and higher-order

From high to first order ODEs

From high-order ODEs to first-order ODEs (cont.)

The roots of the characteristics polynomial and the eigenvalues of the state matrix

• The two are closely connected

This fact can be easily checked for small-size systems

Consider the general linear and homogeneous equation $y^{(3)} + {\color{red}a_2}y^{(2)} + {\color{red}a_1}y^{(1)} + {\color{red}a_0}y = 0$

- A third-order ordinary differential equation
- Its characteristic polynomial $P(\lambda)$

$$\lambda^3 + \frac{a_2}{a_1}\lambda^2 + \frac{a_1}{a_1}\lambda + \frac{a_0}{a_0}$$

The system as three first-order differential equations

$$\underbrace{\begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \\ \dot{x}_3(t) \end{bmatrix}}_{\dot{x}(t)} = \underbrace{\begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 \end{bmatrix}}_{A} \underbrace{\begin{bmatrix} x_1(t) \\ x_2(t) \\ x_3(t) \end{bmatrix}}_{x(t)},$$

(After we defined the dummy variables $x_1(t) = y(t)$, $x_2(t) = \dot{y}(t)$, and $x_3(t) = \ddot{y}(t)$)

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

We have,

Solution using a Taylor expansion

first order ODEs

$$A - \lambda I = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 \end{bmatrix} - \begin{bmatrix} \lambda & 0 & 0 \\ 0 & \lambda & 0 \\ 0 & 0 & \lambda \end{bmatrix}$$
$$= \begin{bmatrix} -\lambda & 1 & 0 \\ 0 & -\lambda & 1 \\ -a_0 & -a_1 & -a_2 - \lambda \end{bmatrix}$$

The eigenvalues of matrix A are given by the values of λ such that $\det(A - \lambda I) = 0$

From high to

The determinant,

$$\det \begin{bmatrix} -\lambda & 1 & 0 \\ 0 & -\lambda & 1 \\ -a_0 & -a_1 & -a_2 - \lambda \end{bmatrix} = -\lambda \left[\lambda(\lambda + a_2) + a_1 \right] - a_0$$
$$= -\lambda \left[\lambda^2 + a_2 \lambda + a_1 \right] - a_0$$
$$= -\lambda^3 - a_2 \lambda^2 - a_1 \lambda - a_0$$

The determinant is zero for values of λ that are roots of the characteristic polynomial The eigenvalues of A correspond to the roots of the characteristic polynomial $P(\lambda)$ \Rightarrow This is because $\det(A - \lambda I) = 0$ equals $P(\lambda)$

From high-order ODEs to first-order ODEs (cont.)

Calculus, review

Intro to ODEs

Taylor expansion

Second- and higher-order

From high to first order ODEs

Example

Consider a linear and time-invariant homogeneous ODE $\ddot{y} + a_2\ddot{y} + a_1\dot{y} + a_0y = 0$

$$\begin{bmatrix} \dot{x_1} \\ \dot{x_2} \\ \dot{x_3} \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix}$$

$$A - \lambda I = \underbrace{\begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -a_0 & -a_1 & -a_2 \end{bmatrix}}_{A} - \lambda \underbrace{\begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{I} = \begin{bmatrix} -\lambda & 1 & 0 \\ 0 & -\lambda & 1 \\ -a_0 & -a_1 & -a_2 - \lambda \end{bmatrix}$$
$$= -\lambda \Big[-\lambda (-\lambda - a_2) + a_1 \Big] + 1 \Big[-a_0 \Big] + 0$$
$$= -\lambda^3 - \lambda^2 a_2 - \lambda_1 - a_0$$
$$\Rightarrow \lambda^3 + a_2 \lambda^2 + a_1 \lambda a_0 = 0$$

Calculus, review

Solution using a Taylor expansion

higher-orde

From high to first order ODEs

From high-order ODEs to first-order ODEs (cont.)

When we convert a N_x -order ordinary differential equation that is linear and homogeneous, we obtain a system of n first-order ordinary differential equations (unforced)

The general form of the system,

$$\dot{x}(t) = Ax(t)$$

We will look more closely at it

Case 1: The dynamics of the state variables are decoupled

Let $x(t) = (x_1(t), x_2(t), \dots, x_{N_n}(t))'$ be the set of state variables

- The evolution of state variable x_i is not affected by x_j
- For all pairs (i, j), with $i, j \in \{1, \ldots, N_x\}$
- We say, the dynamics are decoupled

This condition corresponds to a special structure of matrix A

• Matrix A is diagonal

$$\underbrace{ \begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \\ \vdots \\ \dot{x}_n(t) \end{bmatrix}}_{\dot{x}(t)} = \underbrace{ \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \lambda_n \end{bmatrix}}_{A} \underbrace{ \begin{bmatrix} x_1(t) \\ x_2(t) \\ \vdots \\ x_n(t) \end{bmatrix}}_{x(t)}$$

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$$\underbrace{\begin{bmatrix} \dot{x}_1(t) \\ \dot{x}_2(t) \\ \vdots \\ \dot{x}_n(t) \end{bmatrix}}_{\dot{x}(t)} = \underbrace{\begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \lambda_n \end{bmatrix}}_{A} \underbrace{\begin{bmatrix} x_1(t) \\ x_2(t) \\ \vdots \\ x_n(t) \end{bmatrix}}_{x(t)}$$

The dynamics of the individual state variables have this simple structure

$$\begin{cases} \dot{x_1} = \frac{\lambda_1}{1} x_1(t) \\ \dot{x_2} = \frac{\lambda_2}{2} x_2(t) \\ \vdots \\ \dot{x_n} = \frac{\lambda_n}{1} x_n(t) \end{cases}$$

We can solve them for a initial condition $x(0) = (x_1(0), x_2(0), ..., x_{N_x}(0))'$

$$\begin{cases} x_1(t) = e^{\lambda_1 t} x_1(0) \\ x_2(t) = e^{\lambda_2 t} x_2(0) \\ \vdots \\ x_n(t) = e^{\lambda_n t} x_n(0) \end{cases}$$

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When we re-write the system of solutions in vector form, we obtain the general solution

$$\underbrace{\begin{bmatrix} x_1(t) \\ x_2(t) \\ \vdots \\ x_n(t) \end{bmatrix}}_{x(t)} = \underbrace{\begin{bmatrix} e^{\lambda_1 t} & 0 & \cdots & 0 \\ 0 & e^{\lambda_2 t} & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & e^{\lambda_n t} \end{bmatrix}}_{e^{At}} \underbrace{\begin{bmatrix} x_1(0) \\ x_2(0) \\ \vdots \\ x_n(0) \end{bmatrix}}_{x(0)}$$

The matrix exponential of the state matrix A is a matrix, e^{At} , and it has size $N_x \times N_x$

- ullet Its computation is generally difficult for an arbitrary matrix A
- ullet But, it is very easy to compute when matrix A is diagonal

Matrix e^{At} is called the state transition matrix

- It makes state variables transition in time
- From an initial condition x(0), to x(t)
- According to $x(t) = e^{\mathbf{A}t}x(0)$
- (Remember $y(t) = e^{\lambda t} y_0$)

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From high-order ODEs to first-order ODEs (cont.)

Case 2: The dynamics of the state variables are not decoupled

The standard form of the state-space model $\dot{x}(t) = Ax(t)$ characterises the (unforced) dynamics in a coordinate system whose components are physically meaningful, typically

- The components $x = (x_1, x_2, \dots, x_{N_x})'$ often correspond to physical variables
- Because our state-space models are usually derived from conservation laws

Though interpretable from a process viewpoint, this representation is however arbitrary and not necessarily convenient in terms of solving for the time evolution of the system

- The solution to systems with decoupled dynamics is much easier to compute
- Simplicity is merely due to the difficulty to compute matrix exponentials

However, the vast majority of process systems do now present decoupled dynamics

- Composition of compounds affect each other
- Temperature affects compositions
- . . .