

Zeitharmonische Maxwell-Gleichungen

für anisotrope, homogene Medien

$$\nabla \times \underline{\vec{H}} = \underline{\vec{J}} + \mathrm{j}\omega\varepsilon \underline{\vec{E}} \qquad \Longleftrightarrow \qquad \oint_{\partial A} \underline{\vec{H}} \cdot \mathrm{d}\vec{s} = \iint_A \left(\underline{\vec{J}} + \mathrm{j}\omega\varepsilon \underline{\vec{E}} \right) \cdot \mathrm{d}\vec{A}$$
$$-\nabla \times \underline{\vec{E}} = \underline{\vec{M}} + \mathrm{j}\omega\mu \underline{\vec{H}} \qquad \Longleftrightarrow \qquad -\oint_{\partial A} \underline{\vec{E}} \cdot \mathrm{d}\vec{s} = \iint_A \left(\underline{\vec{M}} + \mathrm{j}\omega\mu \underline{\vec{H}} \right) \cdot \mathrm{d}\vec{A}$$
$$\nabla \cdot \underline{\vec{H}} = \rho_m / \mu \qquad \Longleftrightarrow \qquad \oiint_{\partial V} \underline{\vec{H}} \cdot \mathrm{d}\vec{A} = \frac{1}{\mu} \iiint_V \rho_m \mathrm{d}V$$
$$\nabla \cdot \underline{\vec{E}} = \rho_e / \varepsilon \qquad \Longleftrightarrow \qquad \oiint_{\partial V} \underline{\vec{E}} \cdot \mathrm{d}\vec{A} = \frac{1}{\varepsilon} \iiint_V \rho_e \mathrm{d}V$$

für Freiraumausbreitung und allgemein Vakuum: $\rho_e = \rho_m = 0$ und $\underline{\vec{J}} = \underline{\vec{M}} = \vec{0}$
Poynting-Vektor; Wirkleistung

$$\underline{\vec{S}} = \frac{1}{2} \underline{\vec{E}} \times \underline{\vec{H}}^* \quad ; \quad P = \oiint \Re \left\{ \underline{\vec{S}} \right\} \mathrm{d}\vec{A} = \frac{1}{2} \oiint \Re \left\{ \underline{\vec{E}} \times \underline{\vec{H}}^* \right\} \mathrm{d}\vec{A}$$

EM-Wellen: allgemeine Zusammenhänge

Wellenleiter (konstanter Querschnitt)

$$\lambda = \frac{\lambda_0}{\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2}} = \frac{\lambda_0}{\sqrt{1 - \left(\frac{f_c}{f}\right)^2}} = \frac{\lambda_0}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}$$

Phasengeschwindigkeit

$$v_\mathrm{p} = \frac{c}{\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2}} = \frac{c}{\sqrt{1 - \left(\frac{f_c}{f}\right)^2}} = \frac{c}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}$$

Gruppengeschwindigkeit

$$v_\mathrm{g} = c\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2} = c\sqrt{1 - \left(\frac{f_c}{f}\right)^2} = c\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}$$

Zusammenhang Geschwindigkeiten

$$c_0^2 = v_\mathrm{g} \cdot v_\mathrm{p}$$

Kritische Frequenzen ($a > b$)

$$f_{c,m,n} = \frac{c}{2} \cdot \sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}$$

Zweileitersysteme

Gruppen- und Phasengeschwindigkeit

$$v_\mathrm{g} = v_\mathrm{p} = \frac{1}{\sqrt{\varepsilon\mu}} = \frac{1}{\sqrt{L'C'}} = c$$

Leitungsbeläge Koaxialleitung

$$C' = \frac{C}{l} = \frac{2\pi\varepsilon}{\ln\left(\frac{D}{d}\right)} \qquad L' = \frac{\varepsilon\mu}{C'} = \frac{\mu \ln\left(\frac{D}{d}\right)}{2\pi}$$

Feldwellenwiderstand

TEM-Wellen

$$\frac{\underline{\vec{E}}_x}{\underline{\vec{H}}_y} = -\frac{\underline{\vec{E}}_y}{\underline{\vec{H}}_x} = \pm \frac{\omega\mu}{\beta} = \pm \sqrt{\frac{\mu}{\varepsilon}} = \pm Z_F$$

$$Z_{F0} = \sqrt{\frac{\mu_0}{\varepsilon_0}} = 120\pi\Omega = 377\Omega$$

TE-Wellen

$$Z_{FH} = \pm \frac{\underline{\vec{E}}_x}{\underline{\vec{H}}_y} = \mp \frac{\underline{\vec{E}}_y}{\underline{\vec{H}}_x} = \frac{\omega\mu}{\beta} = \frac{\omega\mu}{\beta_0\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2}} = \frac{Z_F}{\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2}} = \frac{Z_F}{\sqrt{1 - \left(\frac{f_c}{f_0}\right)^2}} = \frac{Z_F}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}$$

TM-Wellen

$$Z_{FE} = \pm \frac{\underline{\vec{E}}_x}{\underline{\vec{H}}_y} = \mp \frac{\underline{\vec{E}}_y}{\underline{\vec{H}}_x} = \frac{\beta}{\omega\varepsilon} = \frac{\beta_0\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2}}{\omega\varepsilon}$$
$$= Z_F\sqrt{1 - \left(\frac{\beta_c}{\beta_0}\right)^2} = Z_F\sqrt{1 - \left(\frac{f_c}{f_0}\right)^2} = Z_F\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}$$