2D boundary integral equations and the Nyström method

Alex Barnett[1] and **Fruzsina Agocs**[1]

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[1]Center for Computational Mathematics, Flatiron Institute, Simons Foundation

Integral equations on 1D interval

Given: i) function $\sigma(t)$ defined on interval $[0, 2\pi)$, periodic: $\sigma(2\pi) = \sigma(0)$, etc
ii) "kernel" function $k(t,s)$ defined on square $[0, 2\pi)^2$,

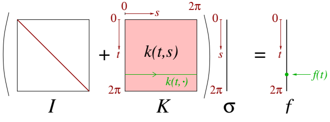
Integral *operator* K acts on σ to give another function $K\sigma$:

$$(K\sigma)(t) := \int_0^{2\pi} k(t,s)\sigma(s)ds, \quad t \in [0, 2\pi)$$

continuous analog of matrix-vector prod. $A\boldsymbol{x}$

Integral *equation*: $(I + K)\sigma = f$, ie

$$\sigma(t) + \int_0^{2\pi} k(t,s)\sigma(s)ds = f(t), \quad t \in [0, 2\pi)$$



analog of lin. sys. $A\boldsymbol{x} = \boldsymbol{b}$

Fredholm "second kind" (due to presence of I, otherwise called "first kind")

If kernel continuous, or "weakly" singular (integrable), K is *compact*:

- eigenvalues ($K\phi_k = \lambda_k \phi_k$) discrete, with $\lim_{k\to\infty} \lambda_k = 0$

  unless some $\lambda_k = -1$, 2nd kind IE has at most one soln: Nul $(I + K) = \{0\}$

- Nul $(I + K) = \{0\} \;\Rightarrow\;$ *existence* of solution for *any* f Fredholm Alternative

  like square matrix (finite-dim), recall: uniqueness $\Rightarrow$ consistent for any RHS

Contrast 1st kind IE $K\sigma = f$ is ill-posed problem (unstable)!

See references for lots of beautiful theory, precise statements

Nyström discretization of 2nd kind IE on interval

Simplest quadrature for periodic funcs: periodic trapezoid rule (PTR)

$$\int_0^{2\pi} f(t)dt \approx \sum_{j=1}^{N} \frac{2\pi}{N} f\left(\frac{2\pi j}{N}\right) = \sum_{j=1}^{N} w_j f(t_j)$$

w_j=weights, t_j=nodes

For f smooth, superalgebraically convergent ("spectral"): error $= \mathcal{O}(N^{-p})$ for any p

Apply quad to integral in 2nd kind IE: call the resulting approx soln $\tilde{\sigma}$

$$\tilde{\sigma}(t) + \sum_{j=1}^{N} k(t,t_j) w_j \tilde{\sigma}(t_j) = f(t), \quad t \in [0,2\pi) \qquad (*)$$

Holds for all t; in particular, holds at each t_i, $i = 1,\ldots,N$, giving:

$$\sigma_i + \sum_{j=1}^{N} k(t_i,t_j) w_j \sigma_j = f(t_i), \quad i = 1,\ldots,N$$

where $\sigma_j := \tilde{\sigma}(t_j)$

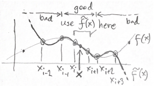
Write as: $A\boldsymbol{\sigma} = \boldsymbol{f}$ $\quad N \times N$ lin sys, entries $a_{ij} = \delta_{ij} + k(t_i,t_j)w_j$, and $f_j := f(t_j)$

solve? dense direct $\mathcal{O}(N^3)$; dense iter. $\mathcal{O}(N^2)$; fast iter. $\approx \mathcal{O}(N)$; fast direct $\approx \mathcal{O}(N^{(d+1)/2})$

Why 2nd kind? eigs(A) accumulate only at $+1$, iterative fast conv.

Sometimes for BIE (eg, far-field eval), node values $\{\sigma_j\}_{j=1}^N$ suffice.
If not, interpolate from them to any $t \in [0,2\pi)$. Two approaches:

- either: rearrange (*) to give $\tilde{\sigma}(t) = \ldots$, called "Nyström interpolant" (rare)
- or (common): use local high-order Lagrange (or global spectral) interpolation:



Demo Nyström on interval (1D)

day1/code/nyst1d_demos.m

```
kfun = @(s,t) exp(3*cos(t-s));               % smooth convolutional kernel, periodic domain [0,2pi)
ffun = @(t) cos(5*t+1);                      % smooth data (RHS) func
N = 30;                                      % number of unknowns: should study convergence as N grows...
tj = 2*pi/N*(1:N); wj = 2*pi/N*ones(1,N);    % PTR nodes and weights, row vecs
A = eye(N) + bsxfun(kfun,tj',tj)*diag(wj);   % identity plus fill k(t_i,t_j)w_j for i,j=1..N
rhs = ffun(tj');                             % col vec
sigmaj = A\rhs;                              % dense direct square solve (pivoted LU), gives col vec
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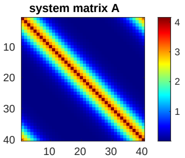

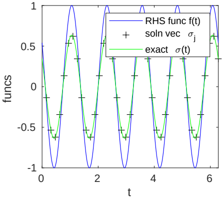

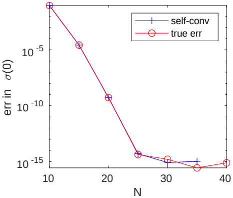

"self-convergence": use N=40 as "true"

f and k smooth

$\Rightarrow$ σ smooth

$\Rightarrow$ spectral conv?

Thm. (Anselone, Kress,. . .): error at node values (and Nyström interpolant) same order as that of quadrature rule applied to integrand $k(t, \cdot)\sigma$.

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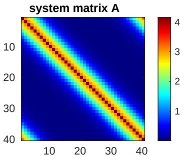

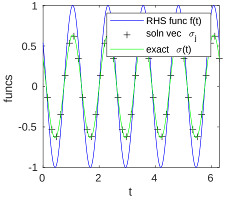

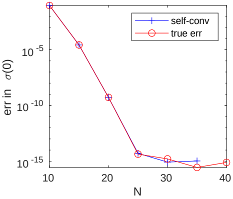

"self-convergence": use $N{=}40$ as "true"

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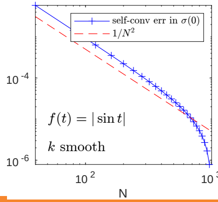
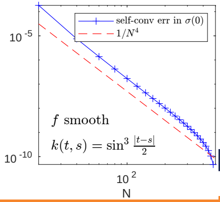
Thm. (Anselone, Kress,...): error at node values (and Nyström interpolant) same order as that of quadrature rule applied to integrand $k(t,\cdot)\sigma$.

- Then, f or k nonsmooth?

worse (here *algebraic*) convergence using plain PTR rule:

Qu: why does order appear to improve at end?





Fundamental solution in $\mathbb{R}^2$

Eg PDE: Poisson eqn $\Delta u = g$ $\qquad$ $\Delta := (\partial/\partial x_1)^2 + (\partial/\partial x_2)^2$ Laplacian

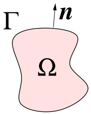
Notation: $\boldsymbol{x} := (x_1, x_2) \in \mathbb{R}^2$ is a point. This frees up $\boldsymbol{y} \in \mathbb{R}^2$ as another point (not y-coord!)

Not well-posed prob. unless add BC! BIEs are good for *homogeneous* PDEs (driving $g \equiv 0$)

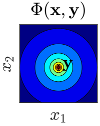
Eg well-posed* BVP:

$$\Delta u = 0 \text{ in } \Omega \qquad \text{PDE (u harmonic)}$$
$$u = f \text{ on } \Gamma \qquad \text{Dirichlet BC}$$

*exists, unique, continuous w.r.t. data



Laplace fundamental soln: $\Phi(\boldsymbol{x}, \boldsymbol{y}) = \frac{1}{2\pi} \log \frac{1}{r}$ $\quad$ where $r := \|\boldsymbol{x} - \boldsymbol{y}\|$



obeys $-\Delta_{\boldsymbol{x}} \Phi = -\Delta_{\boldsymbol{y}} \Phi = \delta_{\boldsymbol{x}}$ $\quad$ Φ harmonic except unit point-mass at $\mathbf{0}$

Normal $\boldsymbol{n}$ points outwards, $\|\boldsymbol{n}\| = 1$ $\quad$ normal deriv. notation $u_n := \boldsymbol{n} \cdot \nabla u$

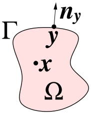
Green's 2nd identity: $\int_\Gamma v u_n - v_n u \, ds = \int_\Omega v \Delta u - (\Delta v) u \, d\boldsymbol{y}$ $\qquad$ calculus

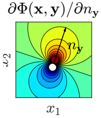
warm-up: set $u =$ BVP soln, $v \equiv 1$, G2I becomes $\int_\Gamma u_n ds - 0 = 0 - 0$: so u has zero *flux*

more fun: fix "target" $\boldsymbol{x} \in \Omega$, let $v = \Phi(\boldsymbol{x}, \cdot)$, G2I gives:

Green's representation formula:

$$\int_\Gamma \Phi(\boldsymbol{x}, \boldsymbol{y}) u_n(\boldsymbol{y}) - \frac{\partial \Phi(\boldsymbol{x}, \boldsymbol{y})}{\partial \boldsymbol{n}_{\boldsymbol{y}}} u(\boldsymbol{y}) \, ds_{\boldsymbol{y}} = u(\boldsymbol{x}) \qquad \text{for } \boldsymbol{x} \in \Omega$$





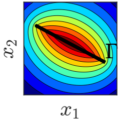
Gets soln from "Cauchy data" $(u, u_n)|_\Gamma$

also versions for Helmholtz, Stokes, Maxwell,...

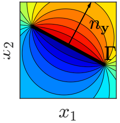
Layer potentials and their jump relations

Representations of harmonic functions off a curve Γ:

Single-layer potential $(\mathcal{S}\sigma)(\boldsymbol{x}) := \int_\Gamma \Phi(\boldsymbol{x},\boldsymbol{y})\sigma(\boldsymbol{y})ds_{\boldsymbol{y}}$ charge sheet



Double-layer potential $(\mathcal{D}\sigma)(\boldsymbol{x}) := \int_\Gamma \frac{\partial\Phi(\boldsymbol{x},\boldsymbol{y})}{\partial \boldsymbol{n}_{\boldsymbol{y}}}\sigma(\boldsymbol{y})ds_{\boldsymbol{y}}$ dipole sheet



interior (-) / exterior (+) limits:
$u^\pm(\boldsymbol{x}) := \lim_{h\to 0^+} u(\boldsymbol{x} \pm h\boldsymbol{n}_{\boldsymbol{x}})$
$u_n^\pm(\boldsymbol{x}) := \lim_{h\to 0^+} \boldsymbol{n}_{\boldsymbol{x}} \cdot \nabla u(\boldsymbol{x} \pm h\boldsymbol{n}_{\boldsymbol{x}})$

Jump relations:

$(\mathcal{S}\sigma)^\pm = S\sigma$ — S (Roman font) means *restriction* of $\mathcal{S}$ to Γ: a bdry int. op.

$(\mathcal{S}\sigma)_n^\pm = (D^T \mp I/2)\sigma$ — jump in normal derivative equal to the charge *density* σ

$(\mathcal{D}\sigma)^\pm = (D \pm I/2)\sigma$ — jump in value; D restriction in principal value sense

$(\mathcal{D}\sigma)_n^\pm = T\sigma$ — T hypersingular, kernel $\partial^2\Phi(\boldsymbol{x},\boldsymbol{y})/\partial\boldsymbol{n}_{\boldsymbol{x}}\partial\boldsymbol{n}_{\boldsymbol{y}}$

D smooth kernel on smooth Γ, while S always log-singular (weakly)

Recap GRF in LP notation: u harmonic in Ω $\Rightarrow$ $\mathcal{S}u_n^- - \mathcal{D}u^- = u$ in Ω

BIEs

Derive direct BIE via GRF

Derive indirect BIE via ansatz $u = \mathcal{D}\sigma$

Table: direct vs indirect pros/cons

we prefer

Parameterization to get a 1D IE

$\boldsymbol{x}(t)$

change of var

back to familiar IE on $[0, 2\pi)$ periodic: apply PTR + Nyström

see "speed weights"

Testing your codes

Test GRF first for a known soln

Exterior Laplace

subtlety of decay in 2D
mixed rep

Helmholtz

$(\Delta + \kappa^2)u = 0$ arises from scalar wave equation $u_{tt} - \Delta u = 0$
κ "wavenumber"; wavelength $\lambda = 2\pi/\kappa$
Also used for 2D Maxwell (z-invar); TE vs TM

Recap

TO DO

-
- Nyström discretization gets $\sigma(t_j)$ interpolate from them to other t
- Fancier quadratures needed for singular kernels and/or close eval
- Nyström is not the only discr. meth: Galerkin, collocation. but: simplest and no less accurate

Resources

Many numerical analysis (mathematical flavor), particularly:

- *Linear Integral Equations*, R. Kress, (1999, Springer). Ch. 6 & 12.
- *The Numerical Solution of Integral Equations of the Second Kind*, K. E. Atkinson, (1997, CUP).

Fewer on the practical/tutorial side:

- "High-order accurate methods for Nyström discretization of integral equations on smooth curves in the plane", S Hao, AH Barnett, PG Martinsson, P Young. *Adv. Comput. Math.* **40**, 245–272 (2014).

  goes beyond these slides for logarithmic singularities, eg SLP

- `https://users.flatironinstitute.org/~ahb/BIE/`
- `https://github.com/ahbarnett/BIEbook` in progress...