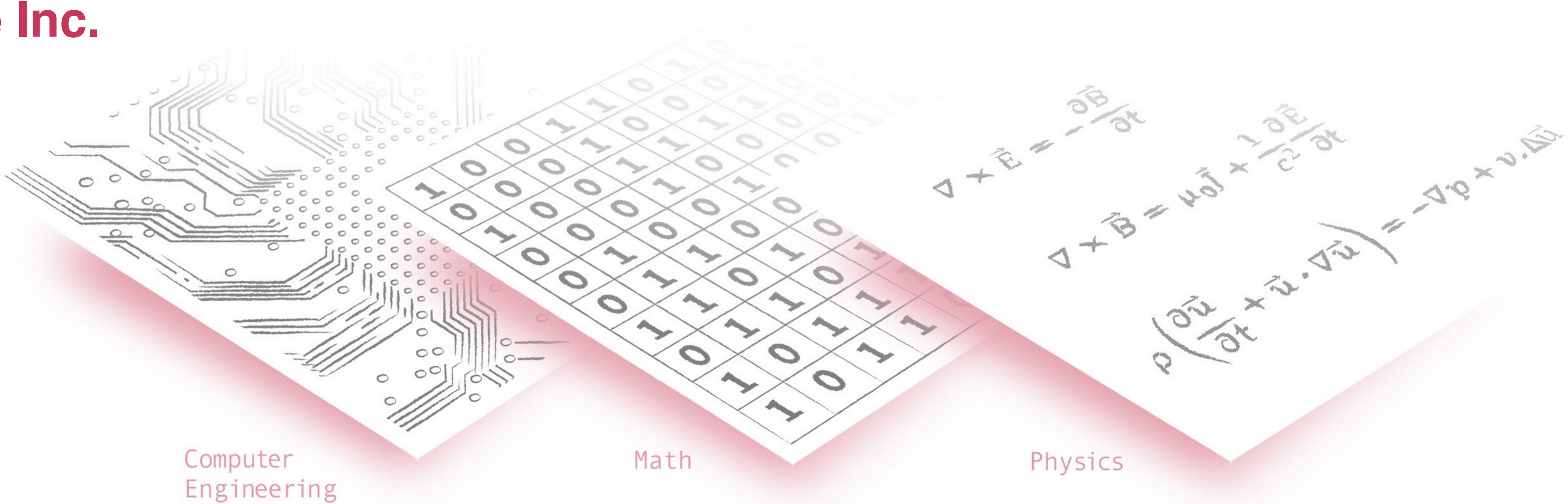




INTRO TO FDTD (7)

Flexcompute Inc.



Maxwell's equation in 1D:

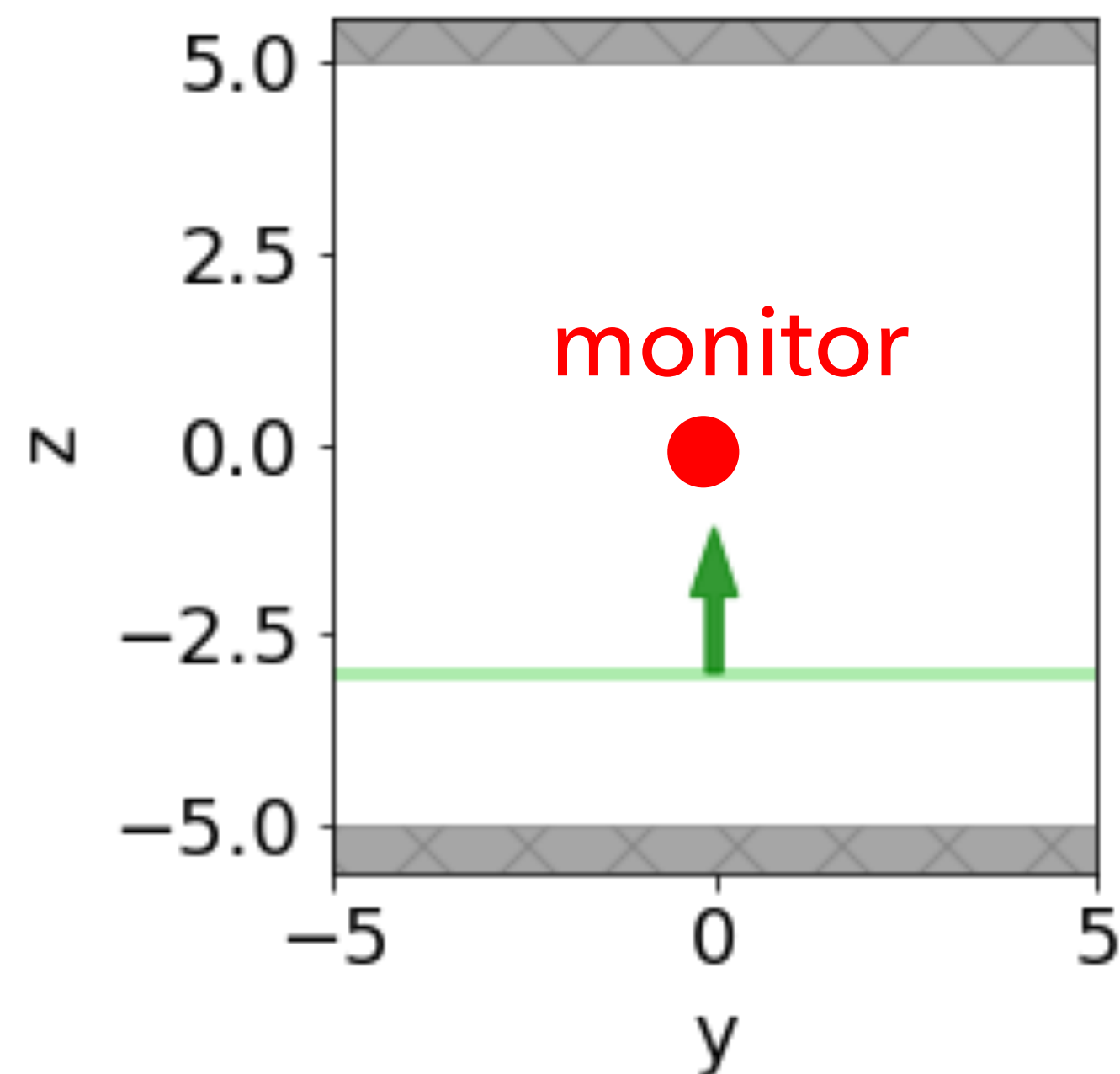
$$\epsilon \frac{\partial E_z}{\partial t} = \frac{\partial H_y}{\partial x}$$

$$\mu \frac{\partial H_y}{\partial t} = -\frac{\partial E_z}{\partial x}$$

Convert the differential operator to difference operator, for example

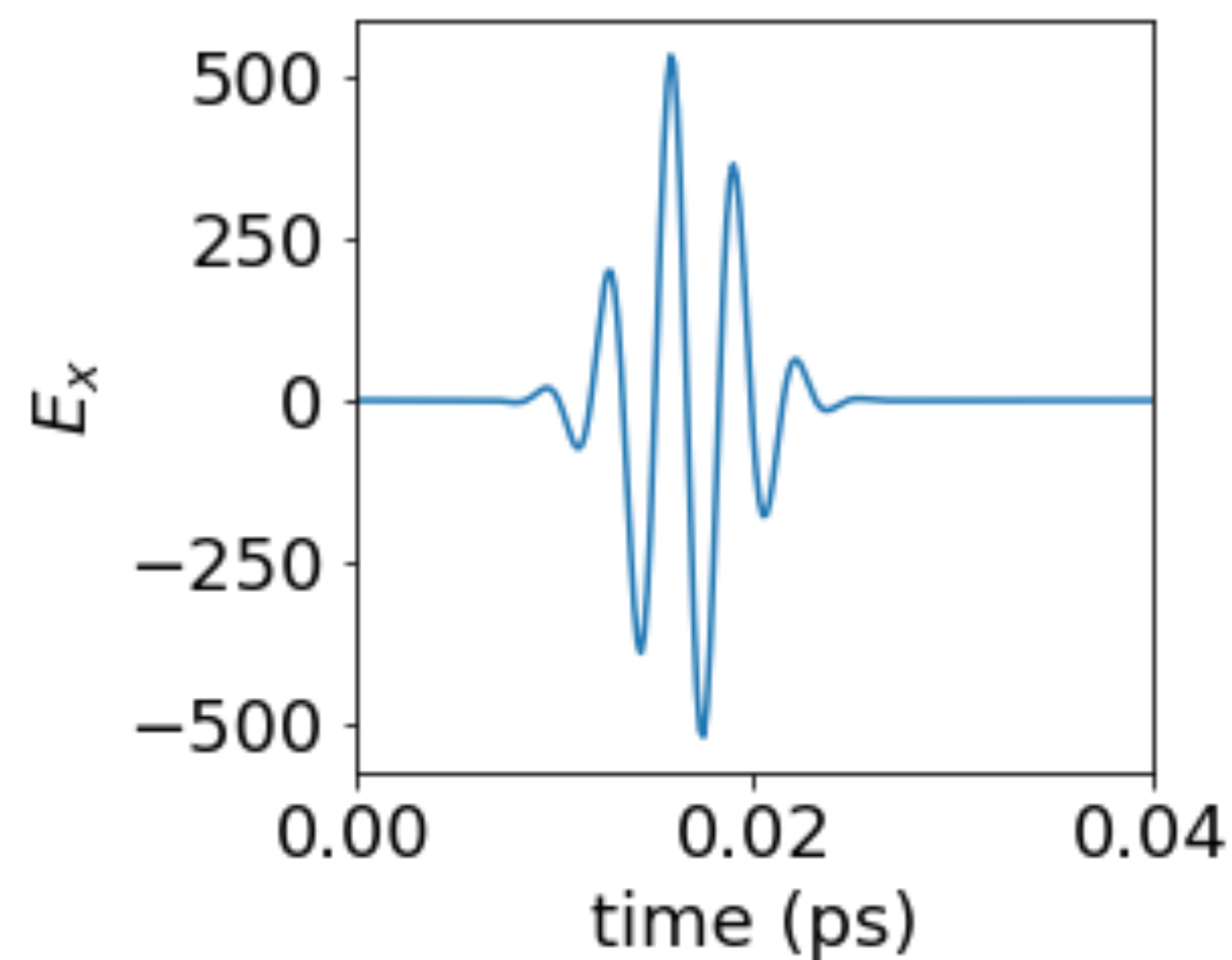
$$\frac{\partial E_z}{\partial t} \rightarrow \frac{E_z^{(\alpha+1)\Delta t} - E_z^{\alpha\Delta t}}{\Delta t}$$

Today's subject: The choice of time step Δt in FDTD

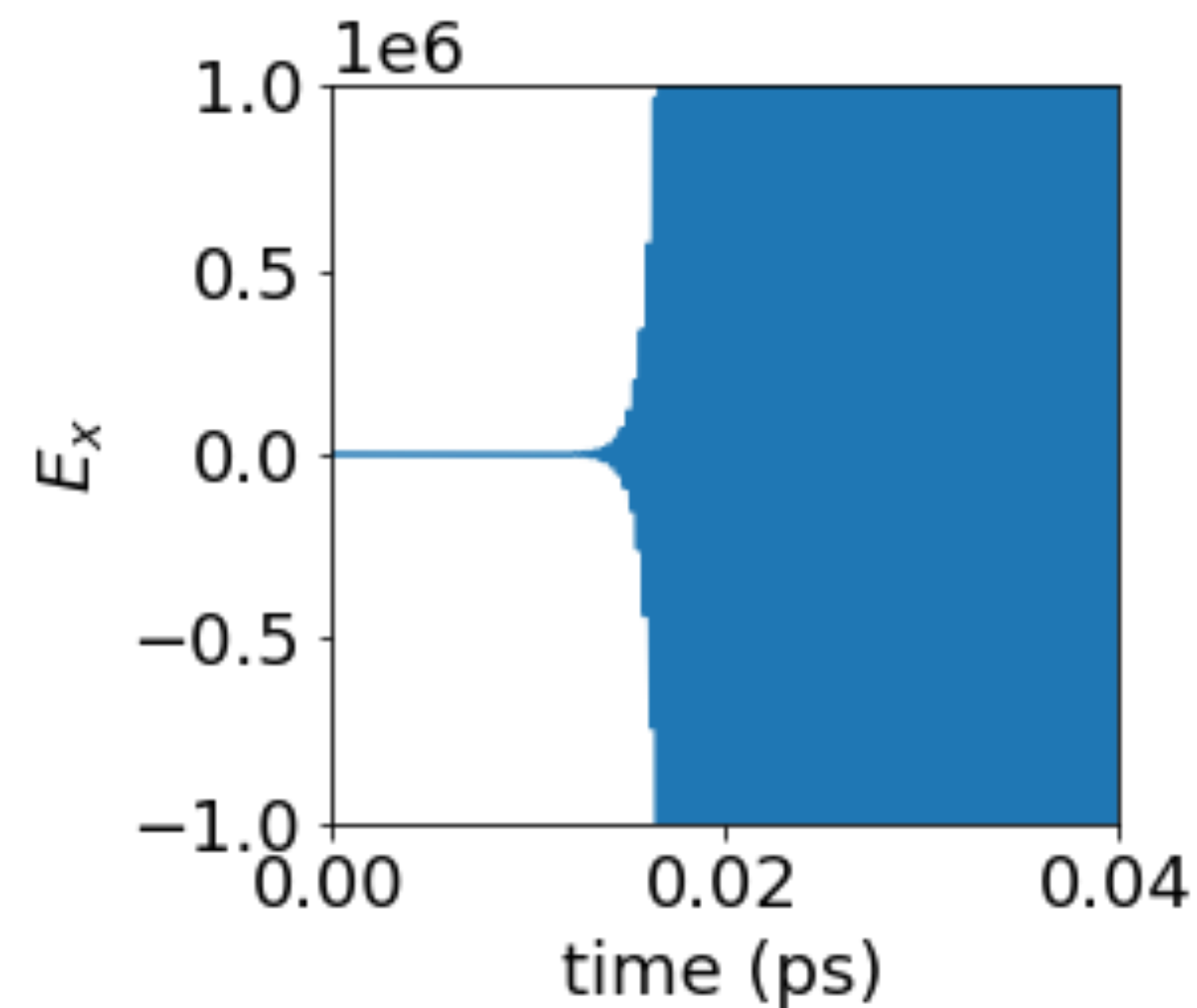


Planewave propagation
in vacuum

$$\Delta t = 0.99 \Delta x / c$$



$$\Delta t = 1.01 \Delta x / c$$



Only difference: increase time step size by 2%



Courant-Friedrichs-Lewy (CFL) condition

1D Yee grid in space and time

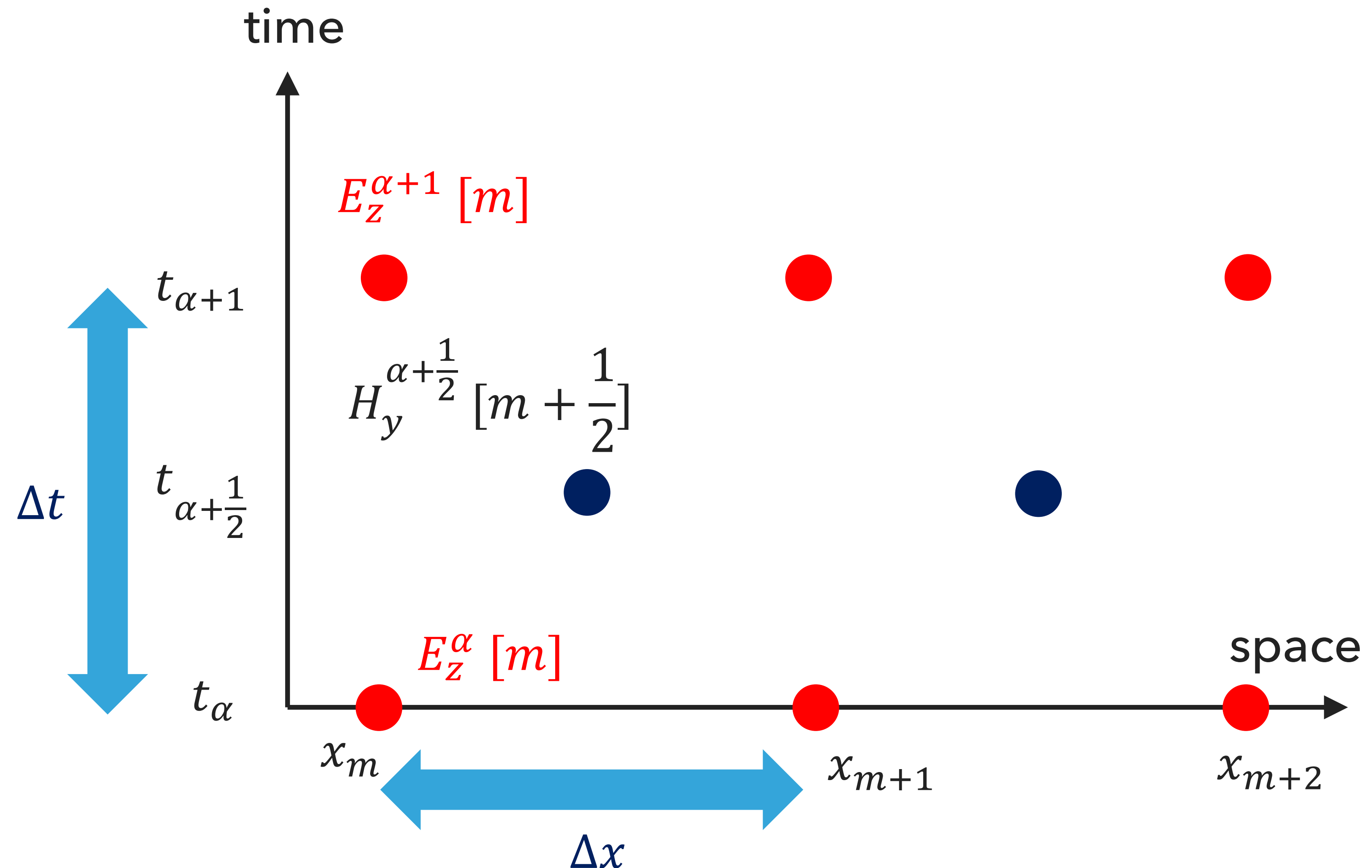
Maxwell's equation in 1D:

$$\varepsilon \frac{\partial E_z}{\partial t} = \frac{\partial H_y}{\partial x}$$

$$\mu \frac{\partial H_y}{\partial t} = -\frac{\partial E_z}{\partial x}$$

Discretize:

- Spatial step size: Δx
- Temporal step size: Δt
- Staggered E and H in space and time



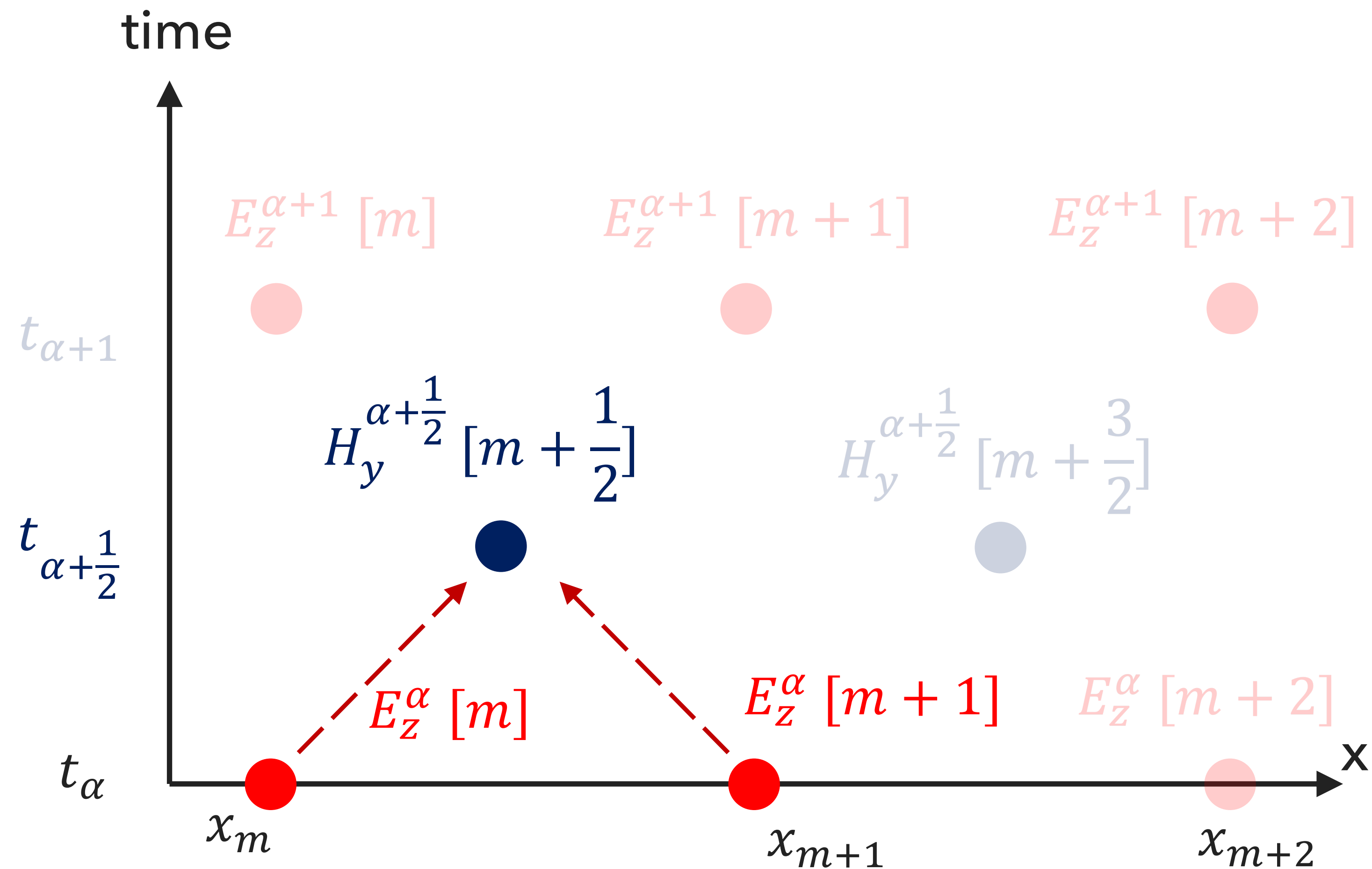
Time marching in 1D

Time update for H_y :

$$\mu \frac{\partial H_y}{\partial t} = \frac{\partial E_z}{\partial x}$$



$$H_y^{\alpha+\frac{1}{2}} \left[m + \frac{1}{2} \right] = H_y^{\alpha-\frac{1}{2}} \left[m + \frac{1}{2} \right] + \frac{\Delta t}{\mu \Delta x} (E_z^\alpha [m+1] - E_z^\alpha [m])$$



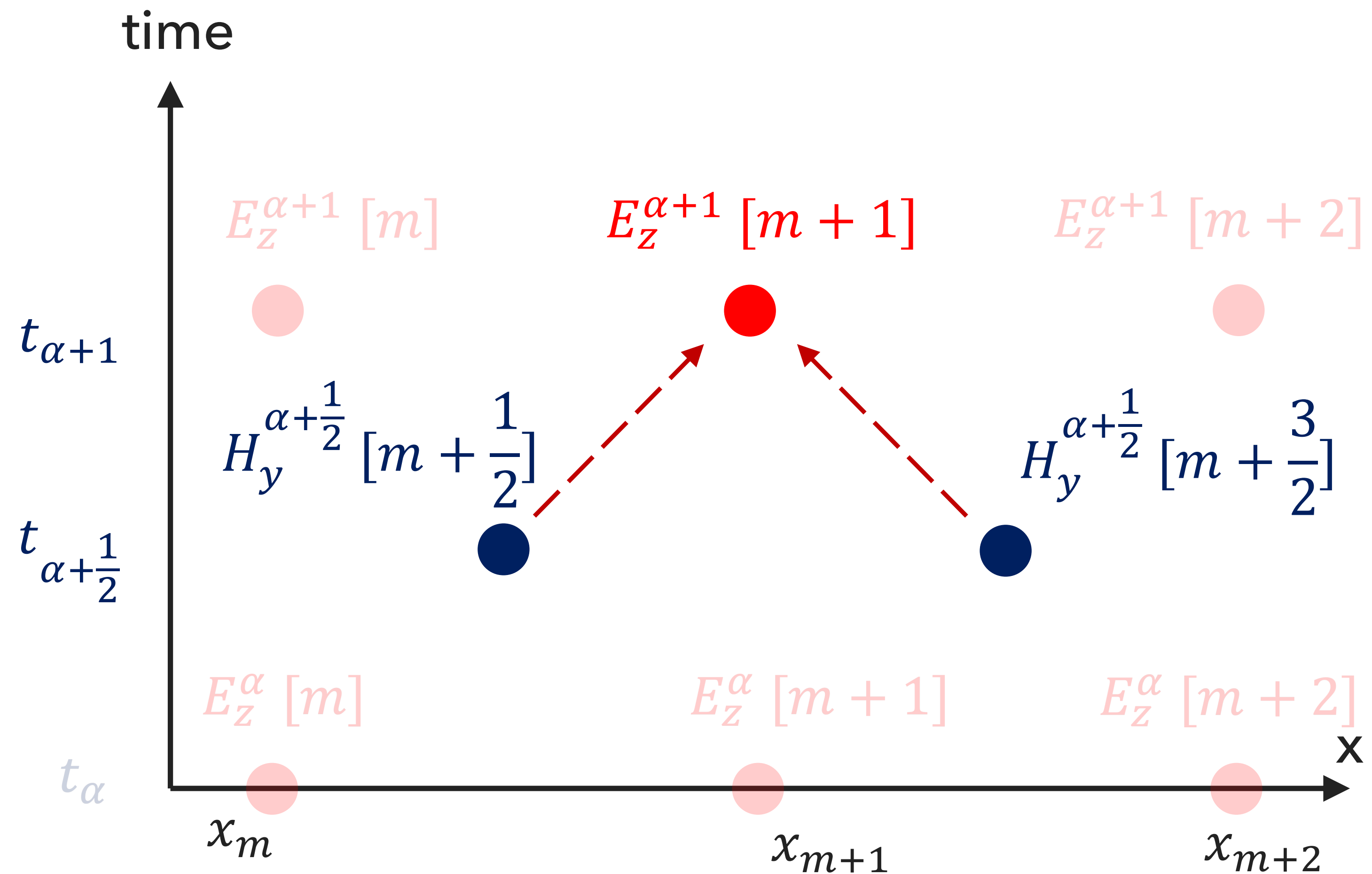
Time marching in 1D

Time update for E_z :

$$\varepsilon \frac{\partial E_z}{\partial t} = \frac{\partial H_y}{\partial x}$$



$$\begin{aligned}
 &E_z^{\alpha+1}[m+1] \\
 &= E_z^\alpha[m+1] + \frac{\Delta t}{\varepsilon \Delta x} \left(H_y^{\alpha+\frac{1}{2}} \left[m + \frac{3}{2} \right] \right. \\
 &\quad \left. - H_y^{\alpha+\frac{1}{2}} \left[m + \frac{1}{2} \right] \right)
 \end{aligned}$$



CFL condition in 1D

CFL condition:

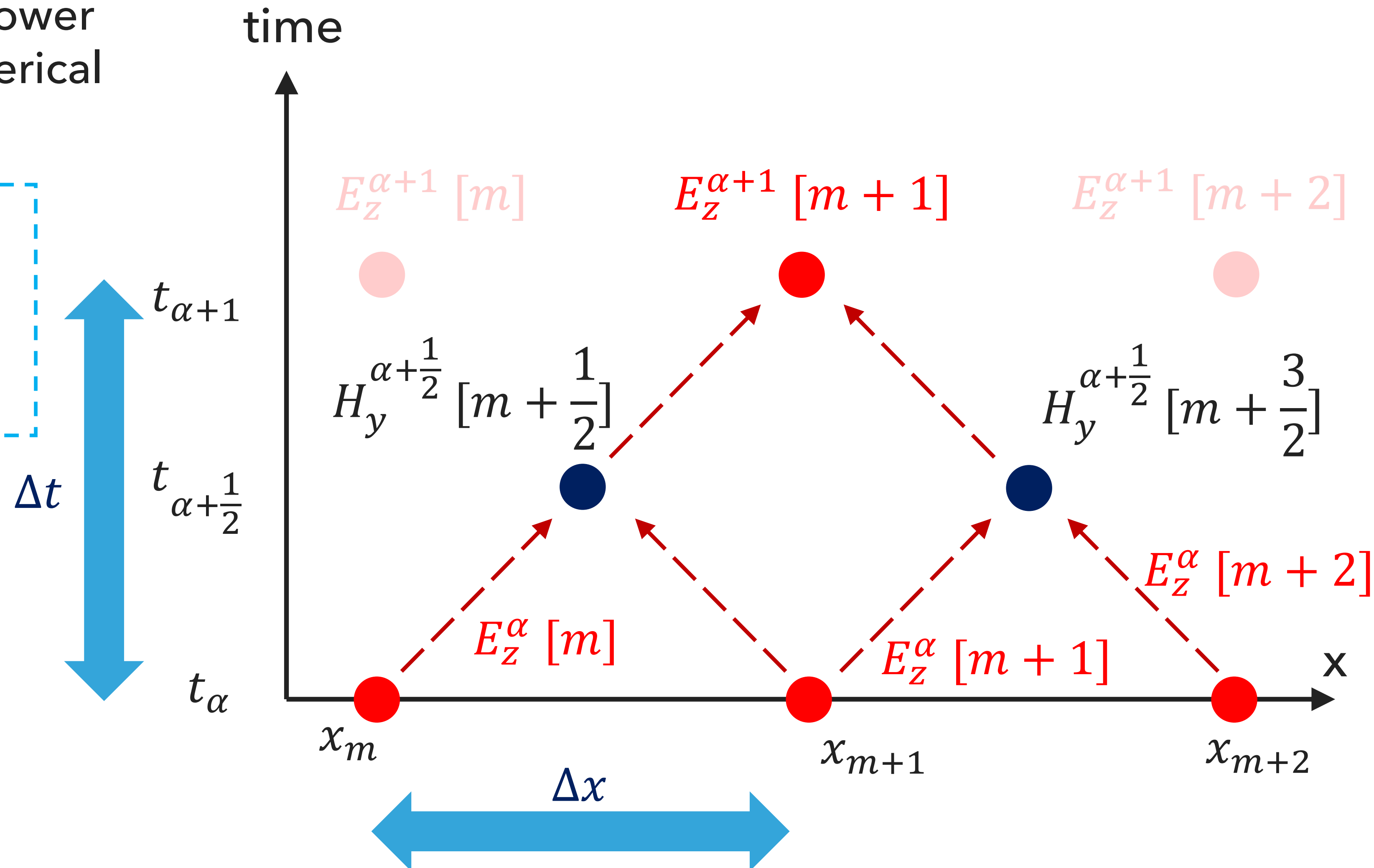
Physical wave should propagate slower than the speed of the slowest numerical dependence.

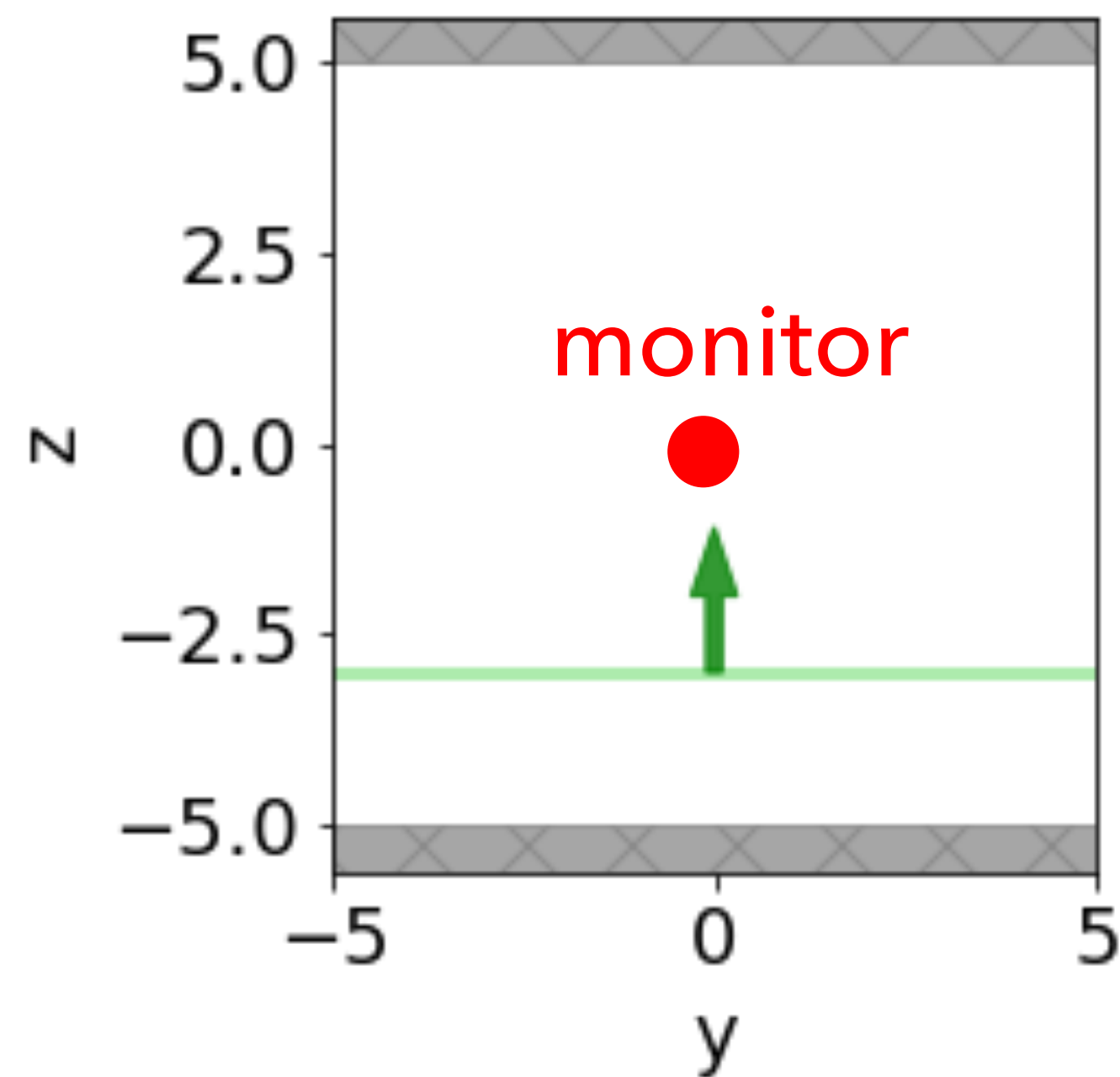
physical	numerical
c	$\frac{\Delta x}{\Delta t}$

$$c \leq \frac{\Delta x}{\Delta t}$$

$$\Delta t \leq \frac{\Delta x}{c}$$

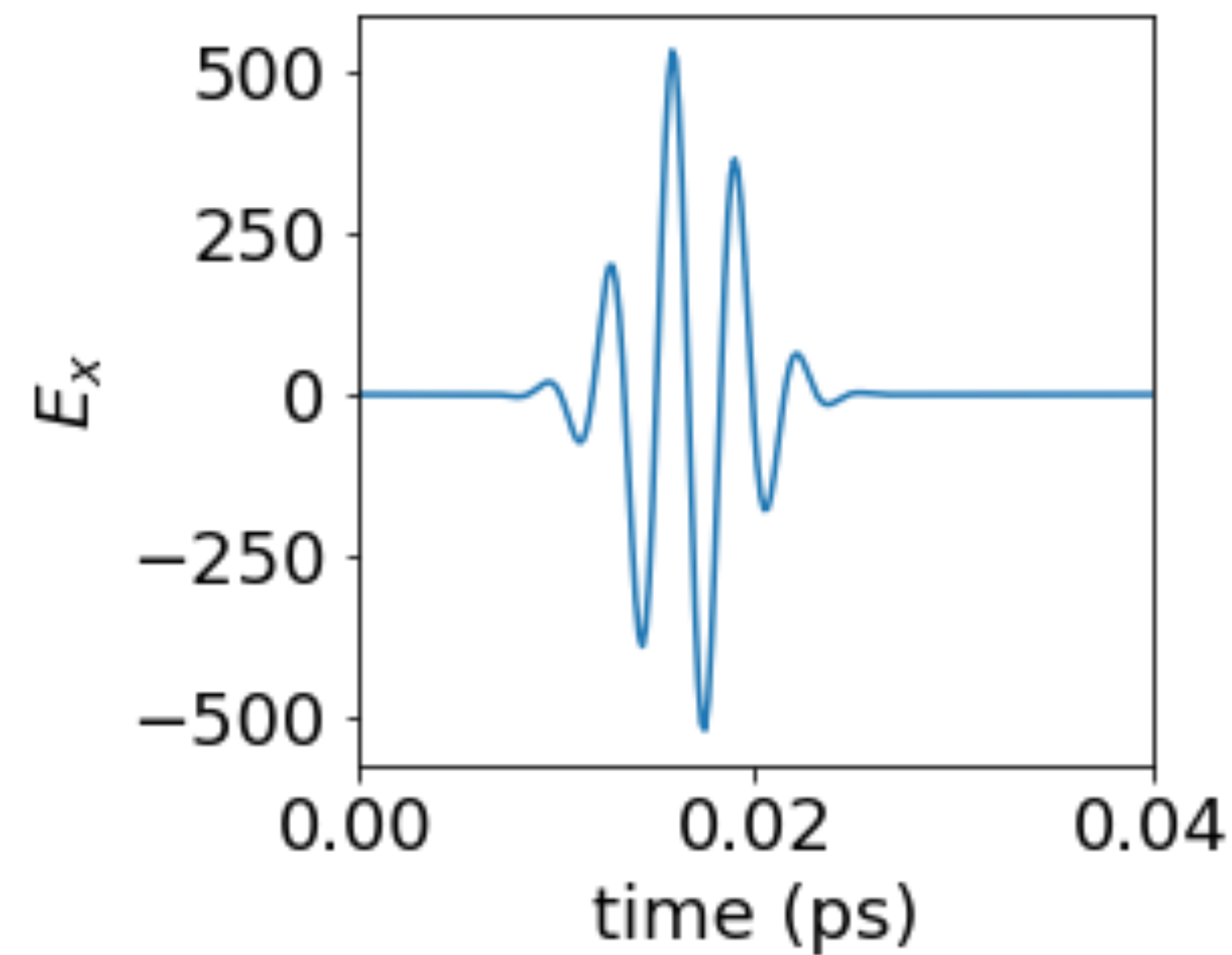
Courant number: $C = \frac{c\Delta t}{\Delta x} \leq 1$





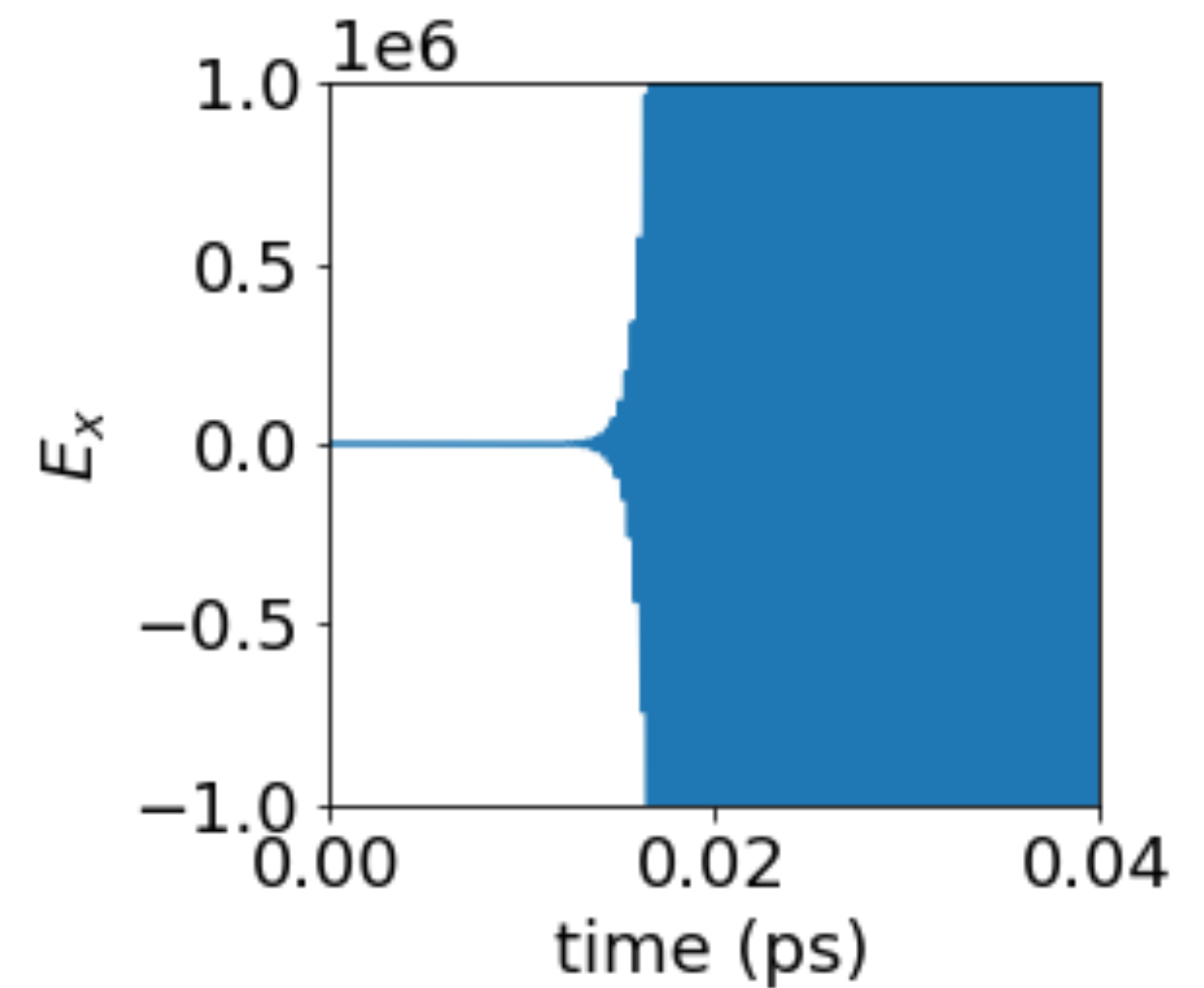
Planewave propagation
in vacuum

$$\Delta t = 0.99 \Delta x / c$$



Courant number: $C = 0.99$

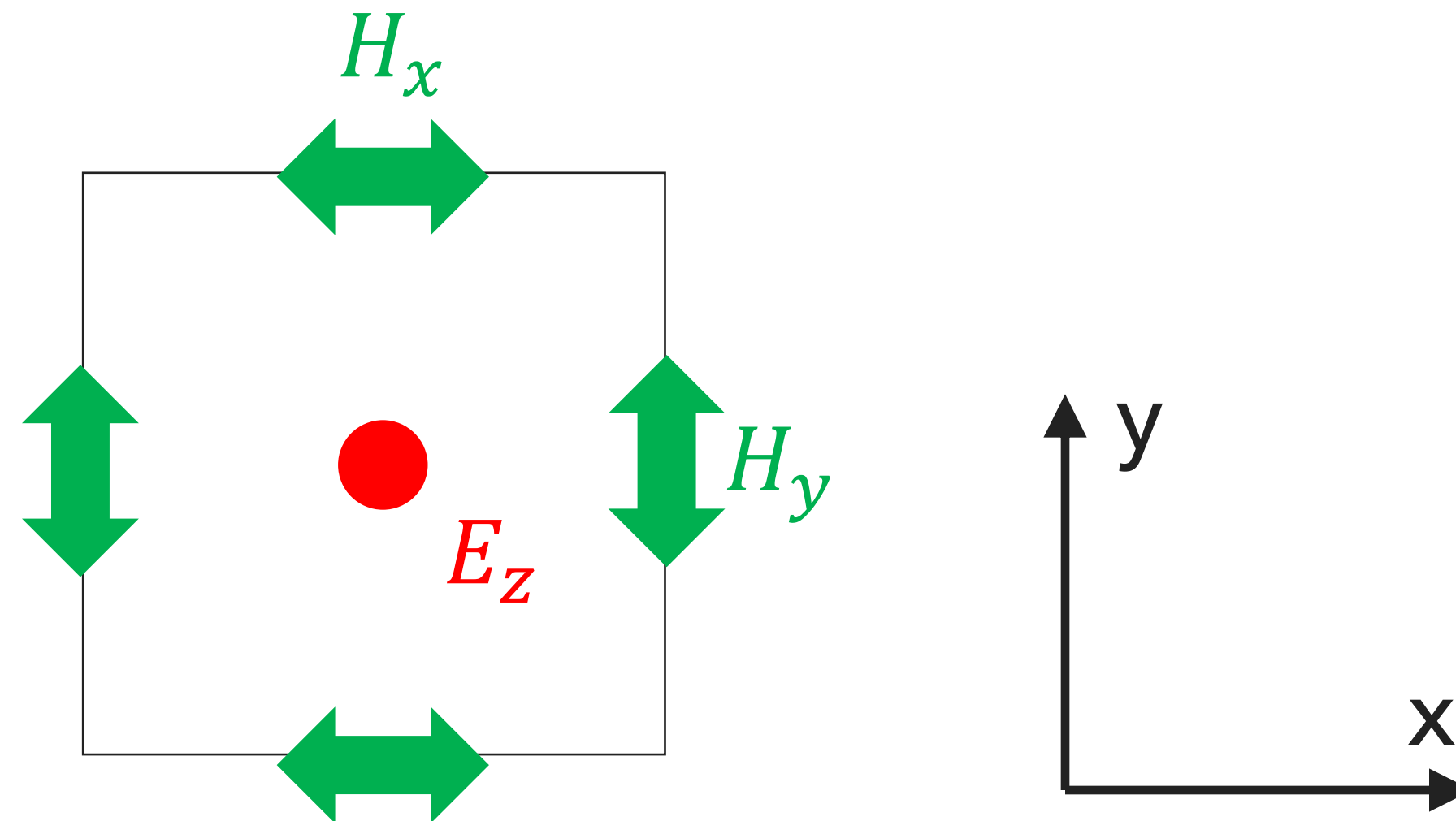
$$\Delta t = 1.01 \Delta x / c$$



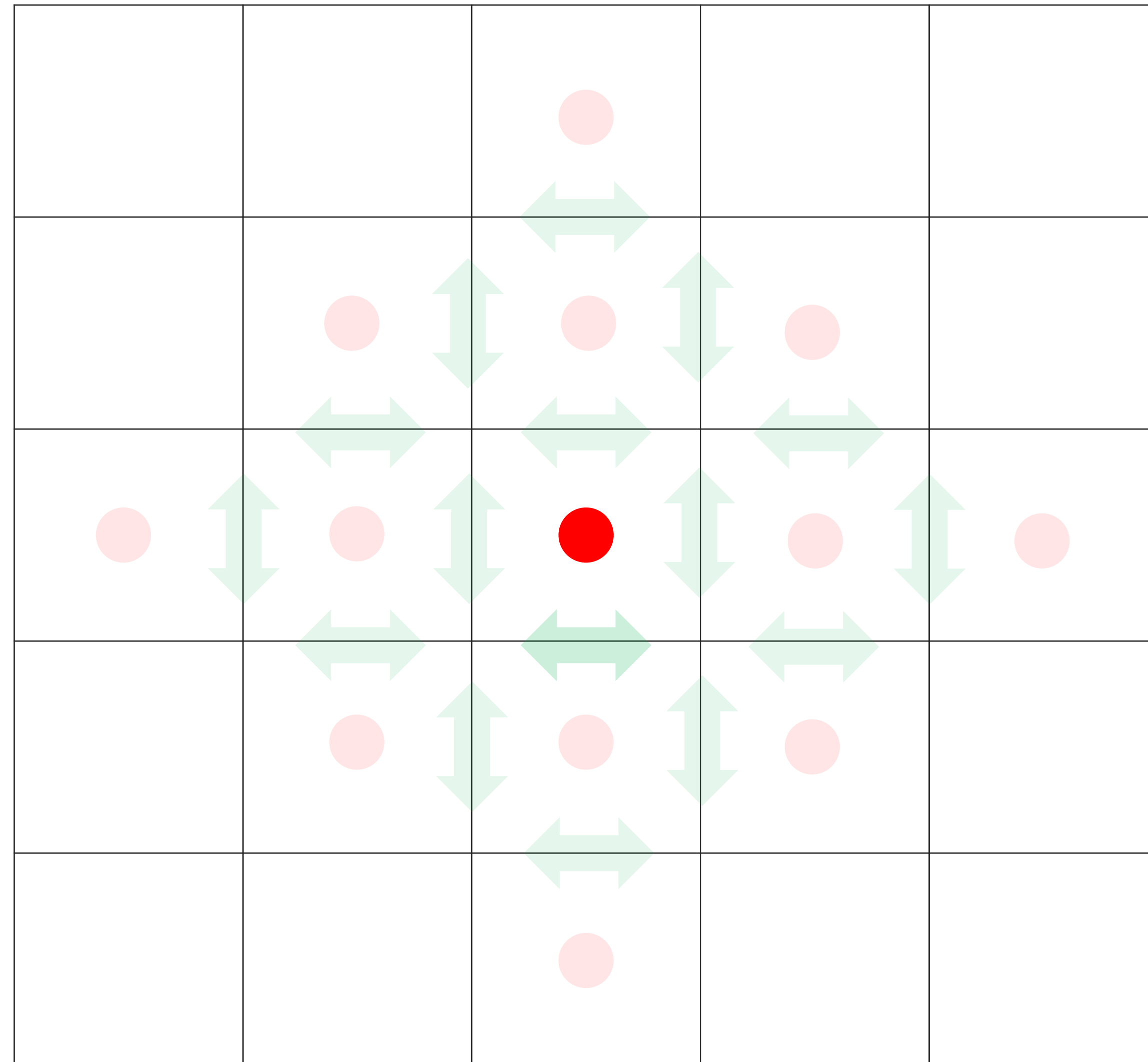
Courant number: $C = 1.01$

Extension to 2D

Yee-cell in 2D for TM
polarization

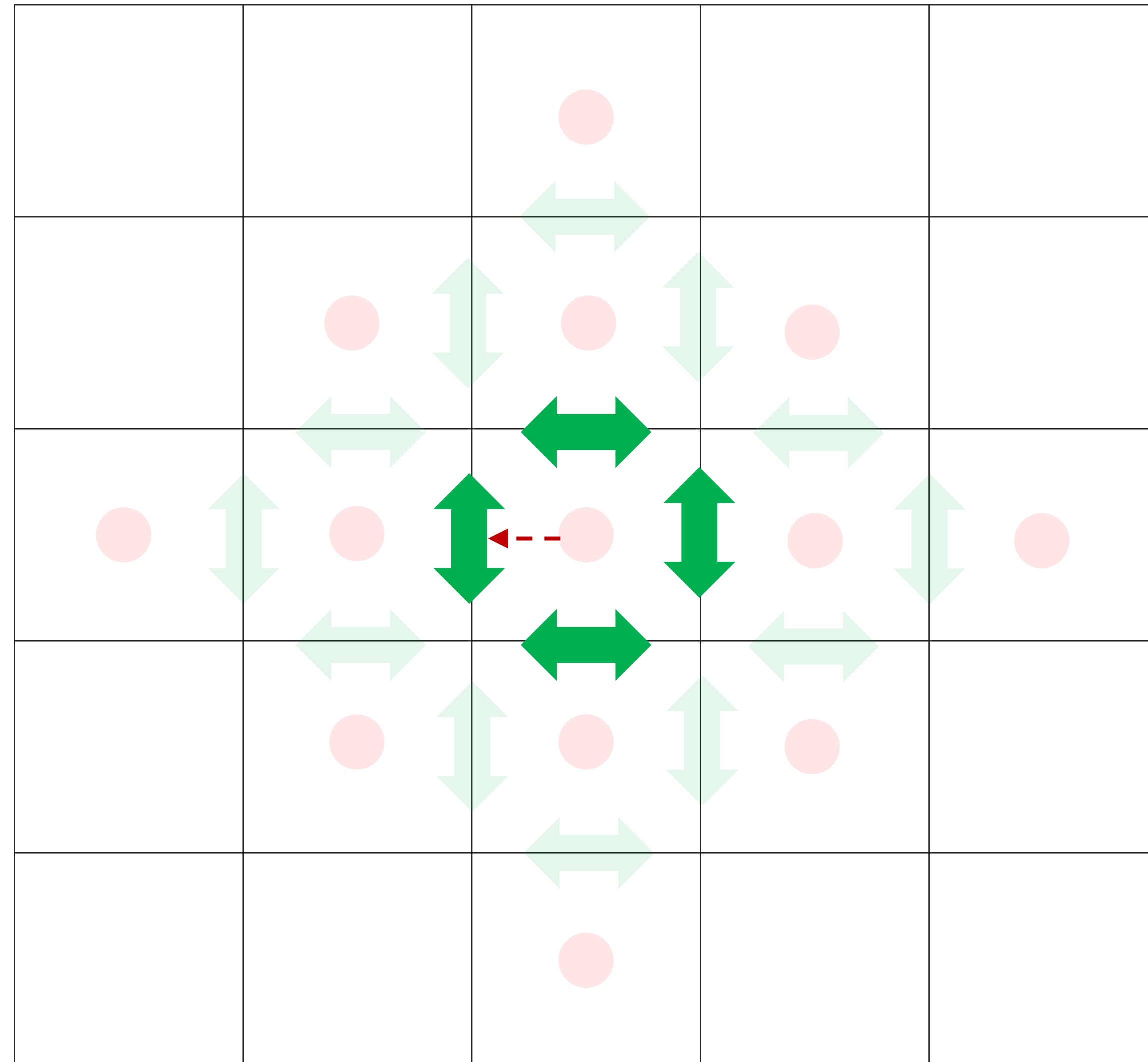


Time marching in 2D (isotropic)

 $t = 0$ 

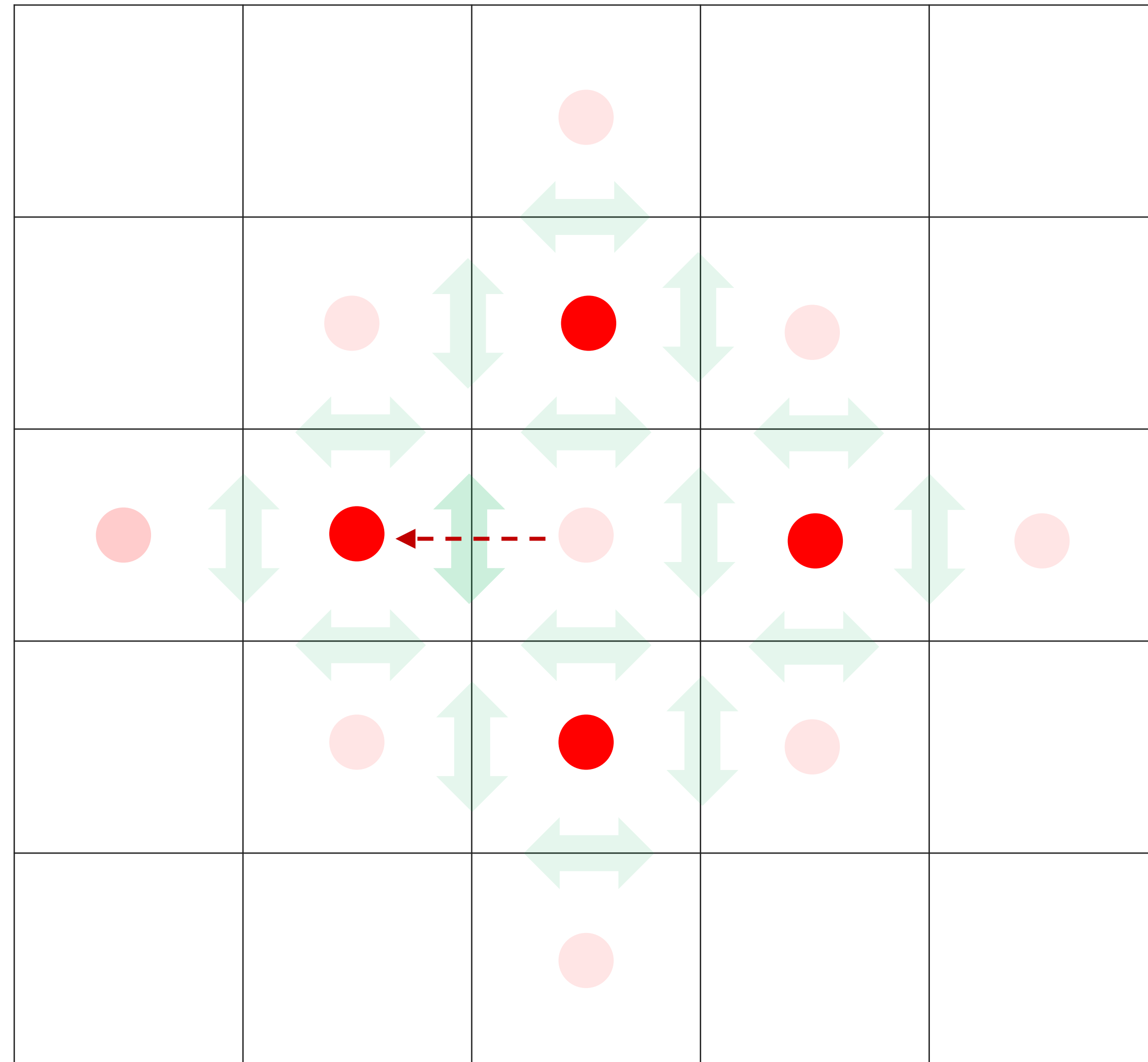
Time marching in 2D

$$t = \frac{\Delta t}{2}$$



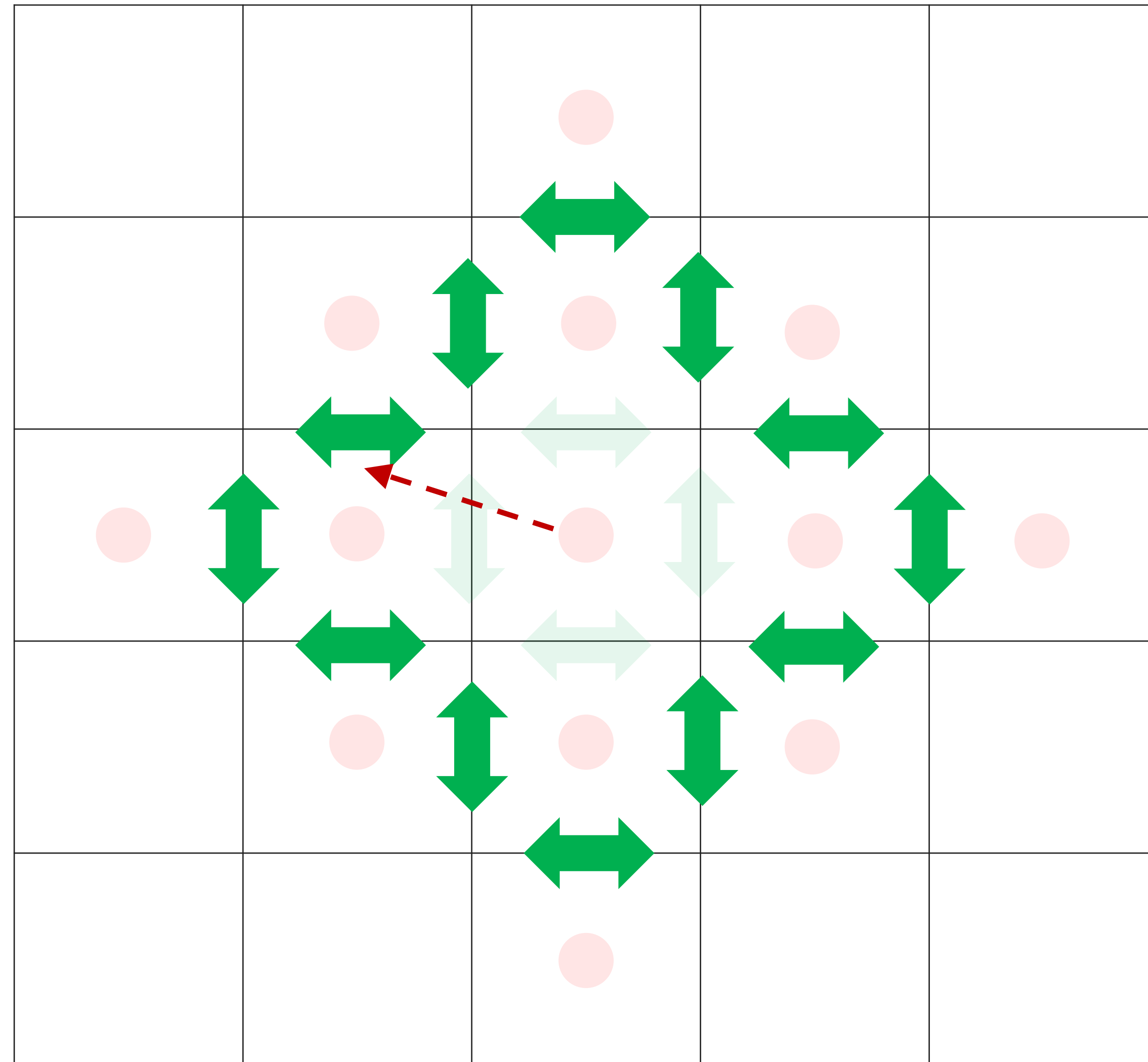
Time marching in 2D

$$t = \Delta t$$



Time marching in 2D

$$t = \frac{3\Delta t}{2}$$

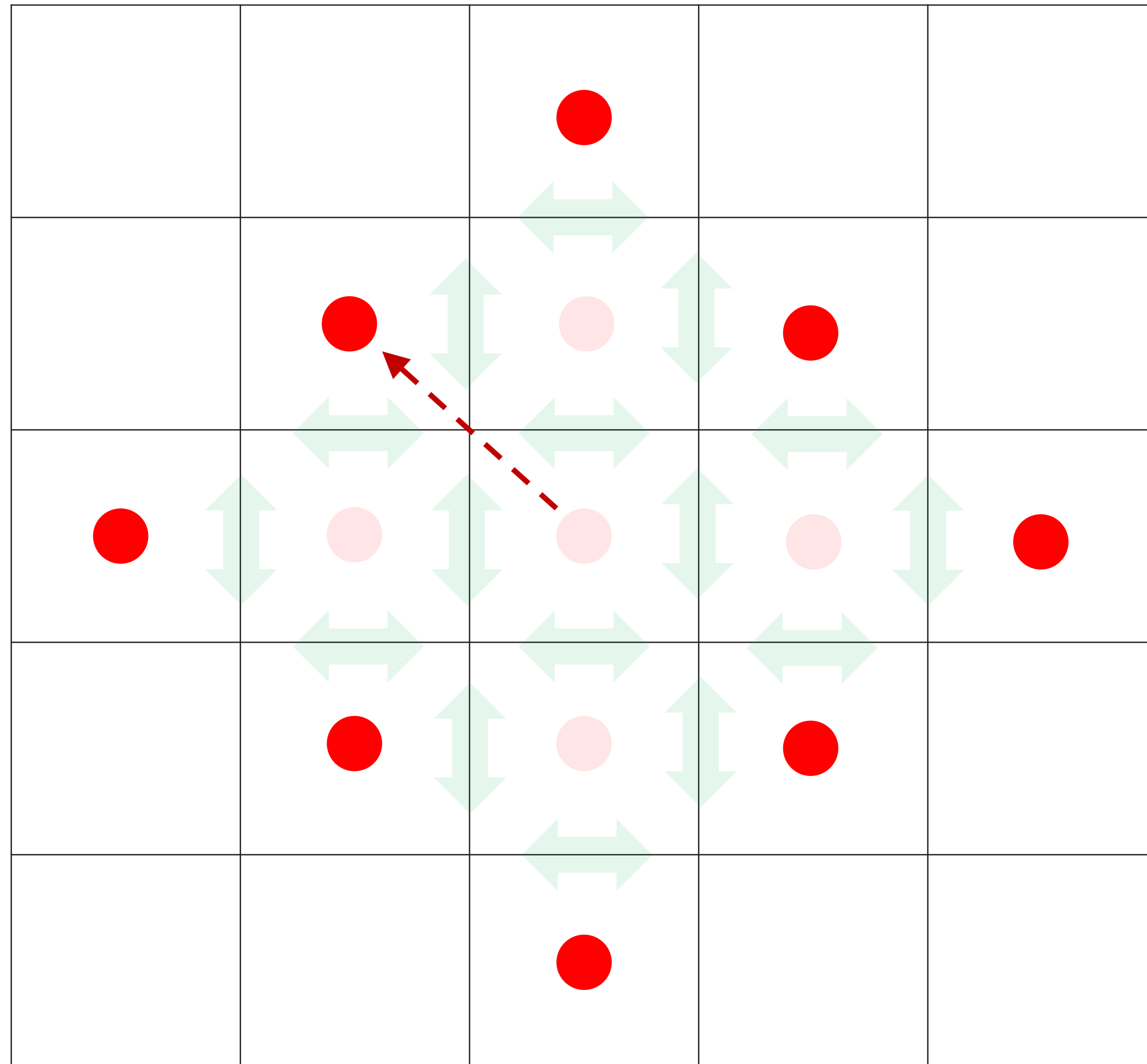


Time marching in 2D

$$t = 2\Delta t$$

Slowest numerical speed:

$$\frac{\sqrt{2}\Delta x}{2\Delta t}$$




CFL condition in 2D

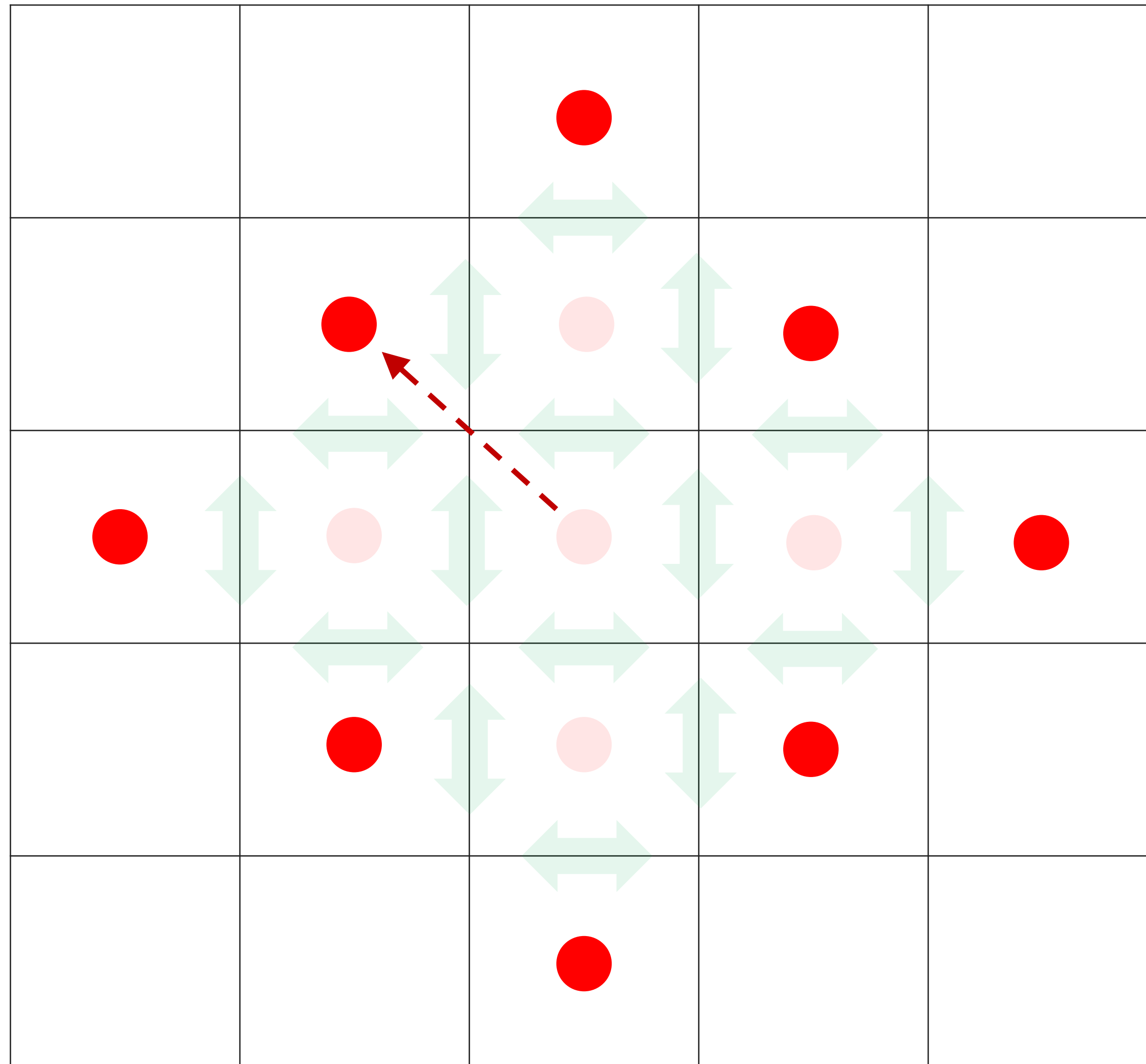
physical

Slowest numerical

$$c \leq \frac{\sqrt{2}\Delta x}{2\Delta t}$$

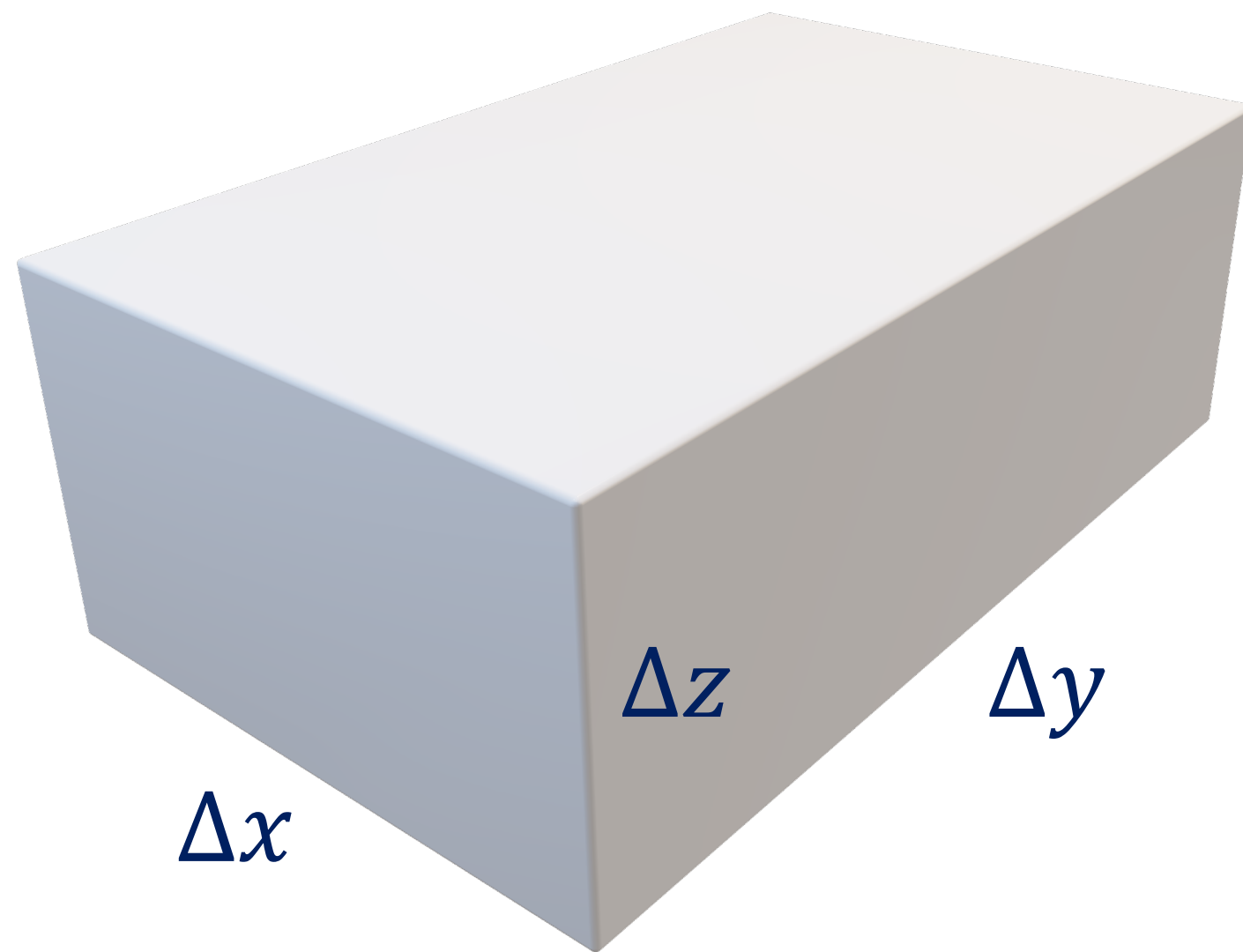

$$\Delta t \leq \frac{\Delta x}{\sqrt{2}c}$$

Courant number: $C = \frac{c\Delta t}{\Delta x} \leq \frac{1}{\sqrt{2}}$



General result for CFL condition in 3D

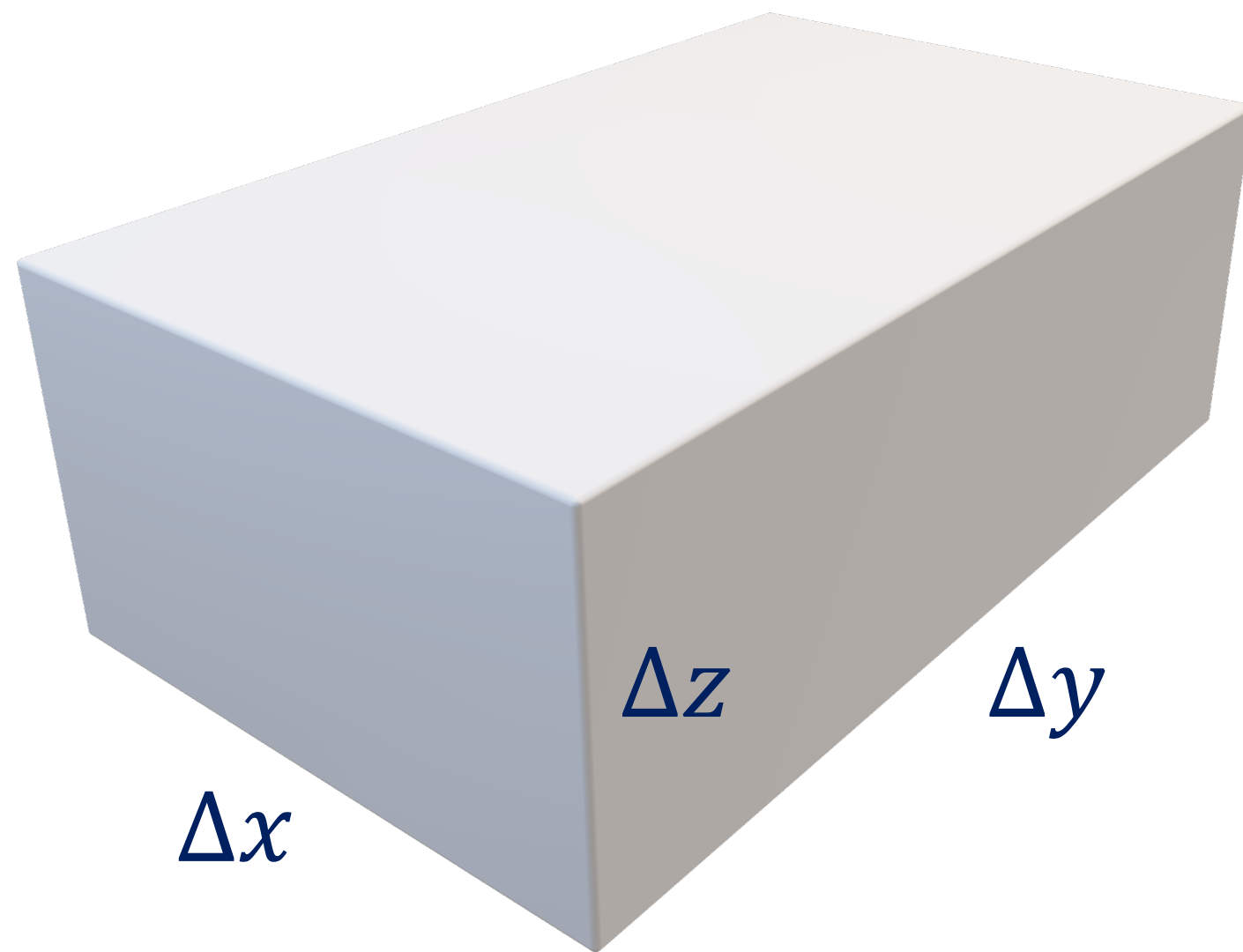
Index of refraction: n



$$\Delta t \leq \frac{n}{c \sqrt{\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2} + \frac{1}{\Delta z^2}}}$$

Practical implementations

Index of refraction: n



$$\Delta t \leq \frac{n}{c \sqrt{\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2} + \frac{1}{\Delta z^2}}}$$

- Nonuniform mesh: time step size is determined by the smallest grid size
- Computational cost for uniform isotropic mesh:
 - Number of grid: $\sim \frac{1}{\Delta x^3}$
 - Time marching steps: $\sim \frac{1}{\Delta x}$
 - In total: $\sim \frac{1}{\Delta x^4}$
 - As an example, reducing the grid size by half leads to increase of total cost by 16 times.