

Anisotropic Extensions of Nonlocal Interaction Models

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Abstract

In this project, we investigated pattern formations which arise from self-organizing mathematical models significant to physical, biological, and engineering systems. The mathematical problems examined in our research have applications in models of materials science and in collective behavior of multi-agent systems such as superconductor vortices, robotic swarms, and biological aggregations. These patterns can be described as minimal energy configurations N interacting particles.

Most nonlocal interaction models in the literature consider only *isotropic* interaction energies where the interaction kernel is radially symmetric. In this project, we first examined pattern formations induced by several *smooth and crystalline anisotropies*. The results of our numerical experiments showed significant quantitative differences in the properties of the energy-optimal configurations between smooth and crystalline anisotropies.

Introduction and Discussion

We have the Hamiltonian of a particle interaction system given by

$$E(x_1, \dots, x_N) = \sum_{\substack{i,j=1 \\ i \neq j}}^N K(\|x_i - x_j\|),$$

where x_1, \dots, x_N are vectors (x_1, x_2) in 2-dimensional Euclidean space \mathbb{R}^2 and $K : \mathbb{R} \rightarrow \mathbb{R}$.

To determine the ground state of this system, we solve the following ODE system

$$\frac{dx_i}{dt} = -\frac{1}{N} \sum_{\substack{i,j=1 \\ i \neq j}}^N \nabla K(\|x_i - x_j\|).$$

Our initial results involved a spherical kernel $K_s : \mathbb{R}^2 \rightarrow \mathbb{R}$ where for $p, q \in \mathbb{R}$ with $q > p > -2$:

$$K_s(x_i - x_j) = K_s(x_1, x_2) = \frac{(x_1^2 + x_2^2)^{\frac{q}{2}}}{q} - \frac{(x_1^2 + x_2^2)^{\frac{p}{2}}}{p}$$

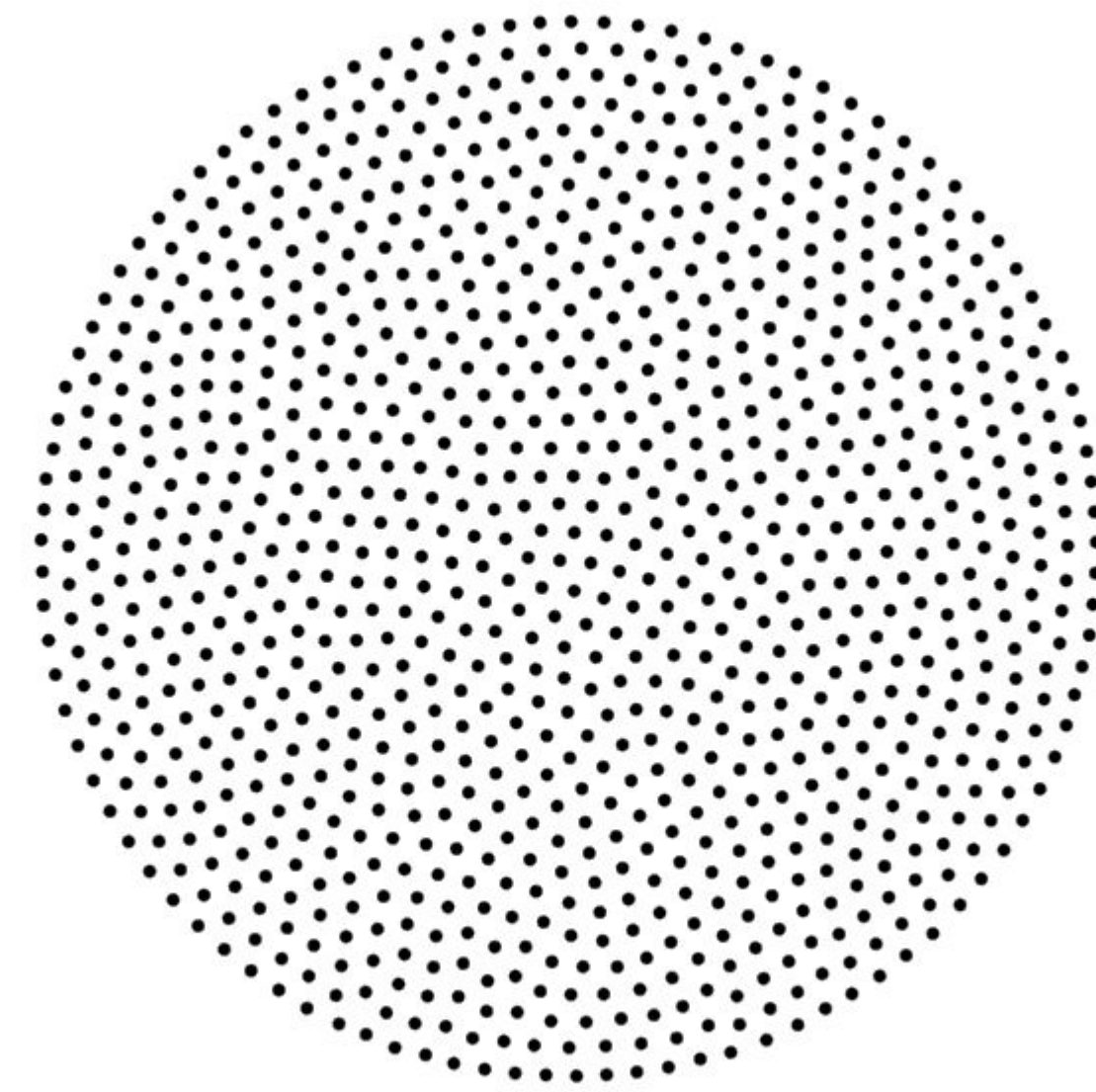
We later examined elliptical kernels $K_{e,c} : \mathbb{R}^2 \rightarrow \mathbb{R}$ with norm c and parameters $a, b \in \mathbb{R}$. For $c \geq 1$, we have

$$K_{e,c}(x_1, x_2) = \frac{(a^c|x_1|^c + b^c|x_2|^c)^{\frac{q}{c}}}{q} - \frac{(a^c|x_1|^c + b^c|x_2|^c)^{\frac{p}{c}}}{p}.$$

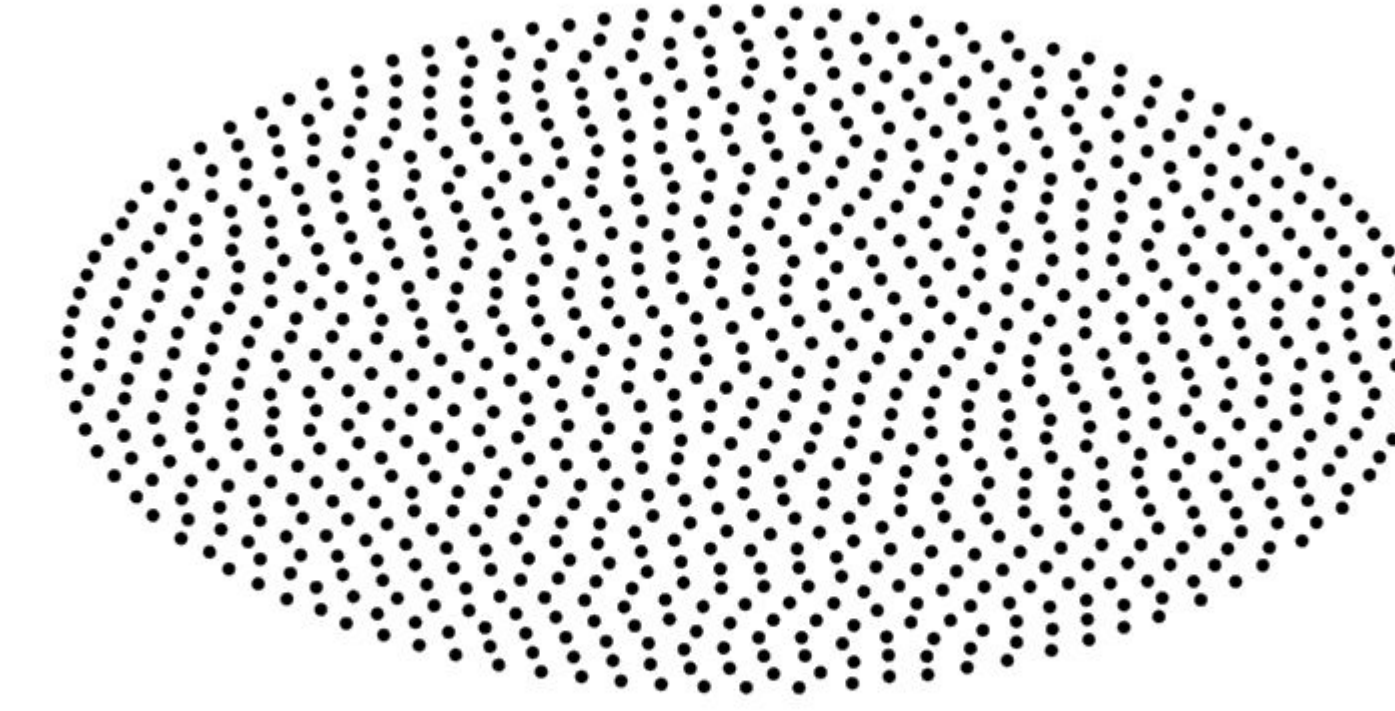
In our kernels, if $p = 0$ or $q = 0$, the p or q expressions were treated as logarithms. For the elliptical norm, this would therefore be

$$\frac{(a^c|x_1|^c + b^c|x_2|^c)^{\frac{q}{c}}}{0} = \log \left((a^c|x_1|^c + b^c|x_2|^c)^{\frac{1}{c}} \right)$$

Numerical Results

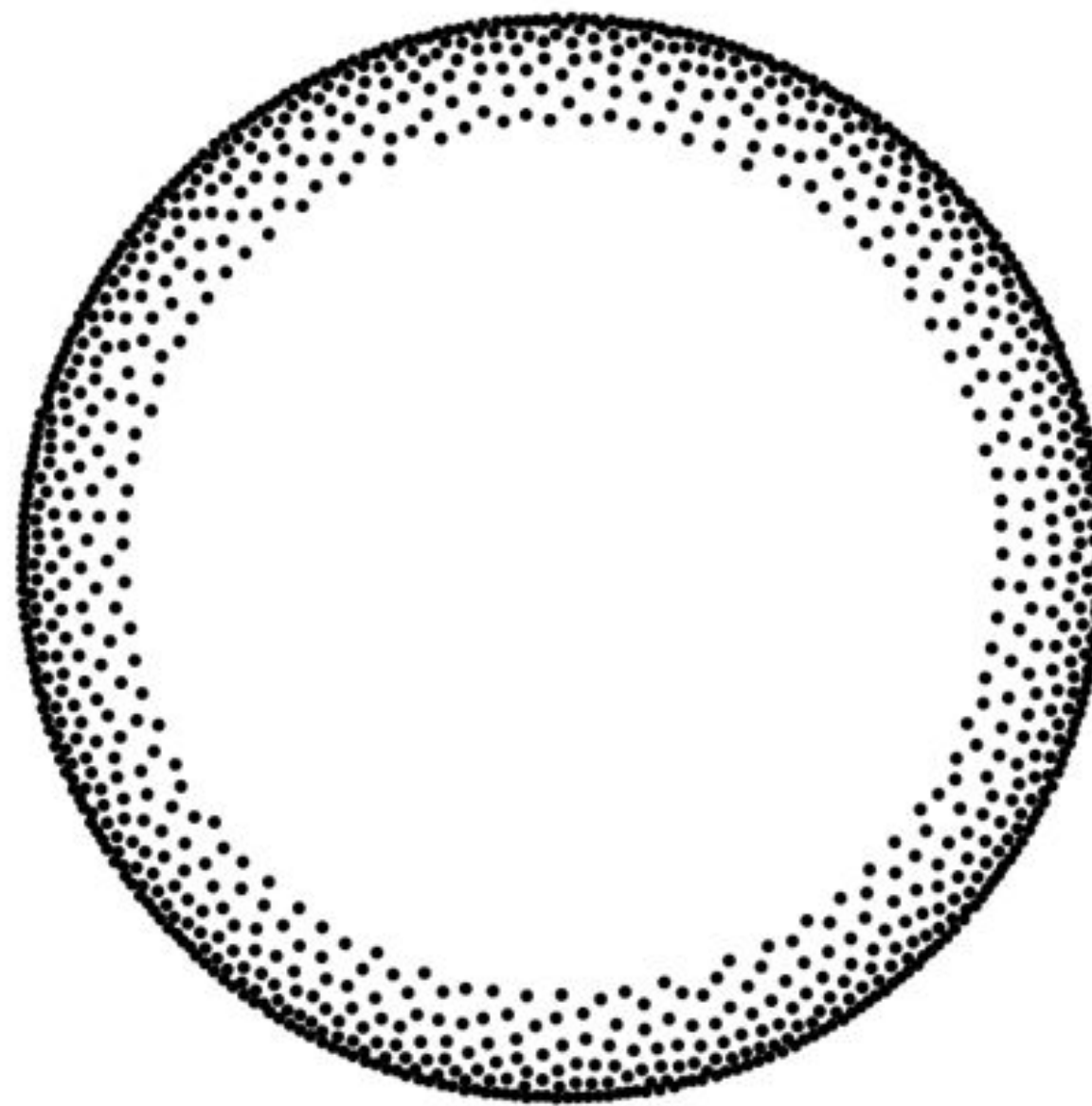


K_s with $p = 0, q = 2$,
and $N = 1000$

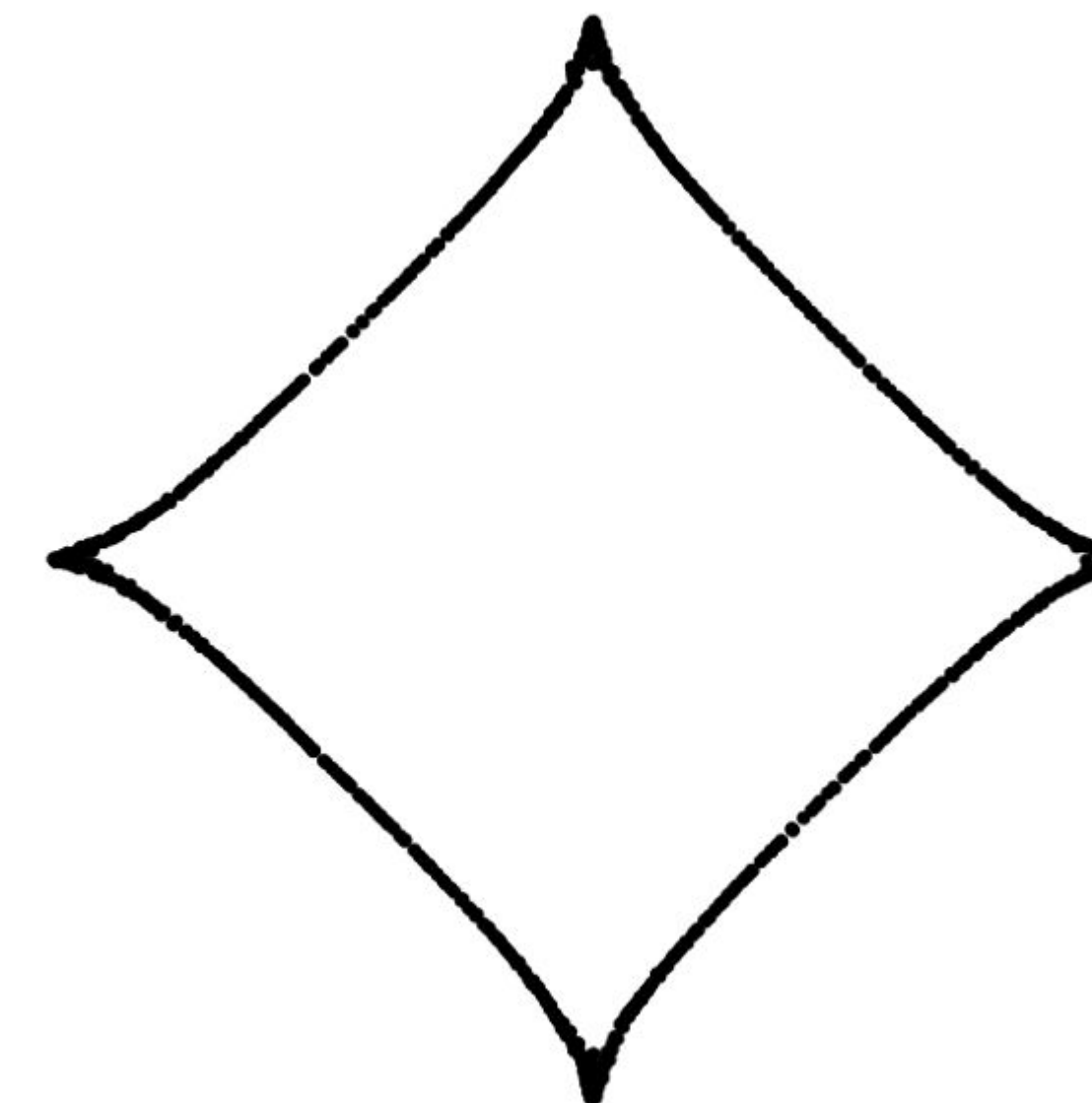


$K_{e,2}$ with $p = 0, q = 2, a = 0.5,$
 $b = 1$, and $N = 1000$

For $p = 0, q = 2$, K_s gives a uniform distribution of points throughout a circle. However, when using $K_{e,2}$, changing a and b away from 1 results in uniform distribution throughout an ellipse. As described in the Analytic Results section, we are able to describe this change through a change of variables.

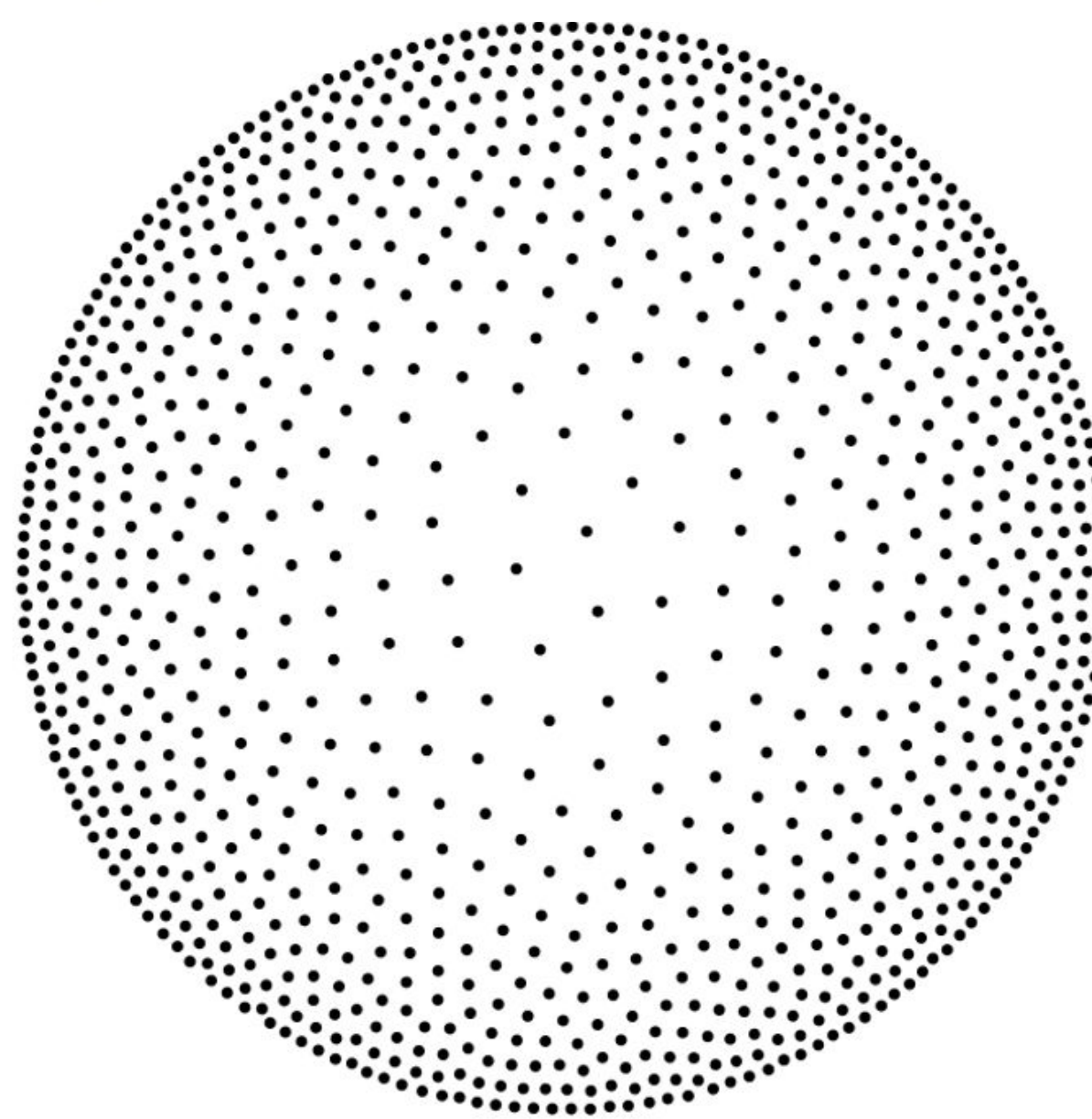


$K_{e,2}$ with $p = 1, q = 4,$
 $a = 1, b = 1$, and $N = 1000$

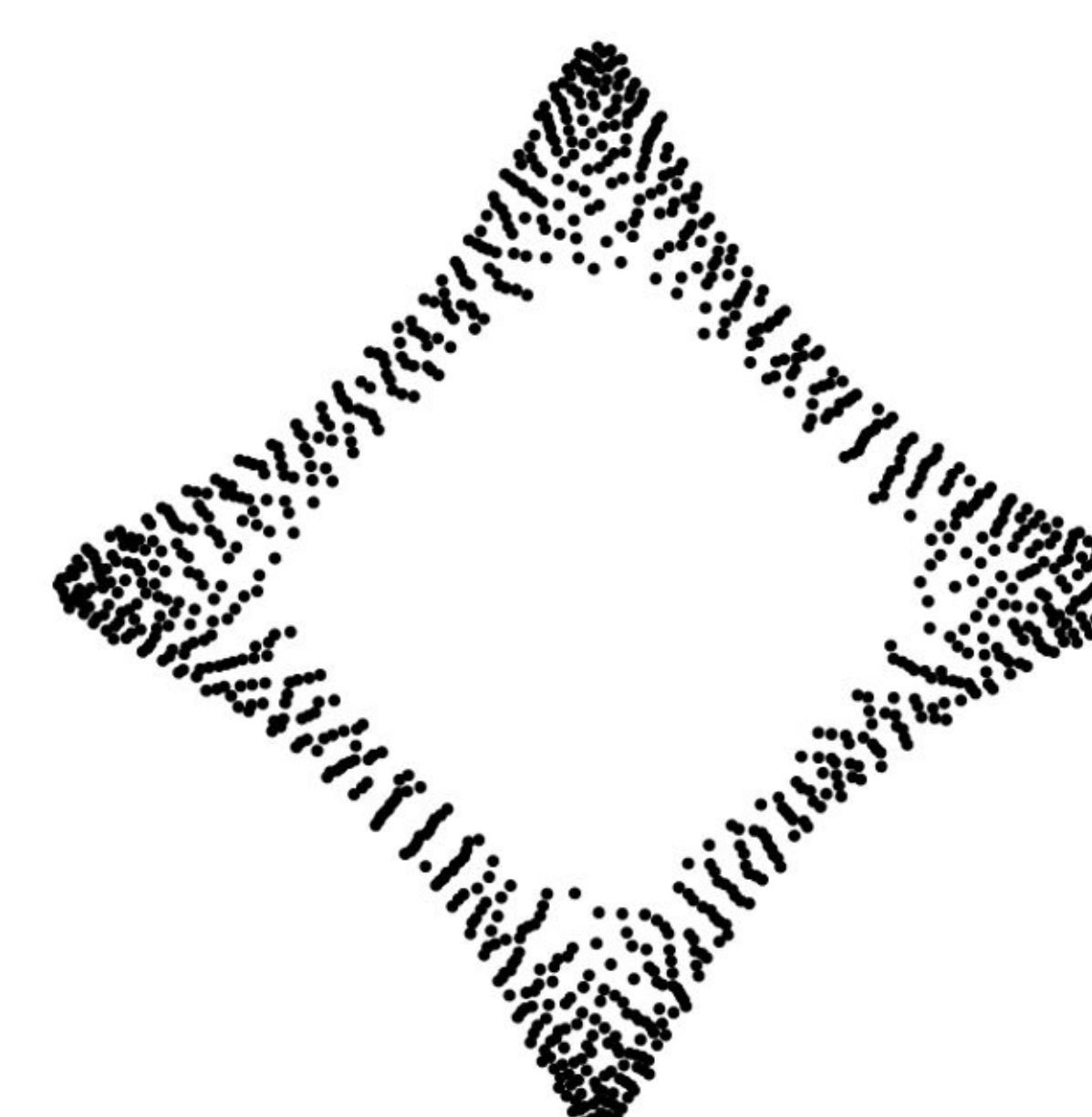


$K_{e,1}$ with $p = 1, q = 4,$
 $a = 1, b = 1$, and $N = 1000$

For $p = 1, q = 4$, $K_{e,1}$ gives a 1 dimensional curve (an astroid), while $K_{e,2}$ gives a 2 dimensional distribution. This difference in dimensionality reveals significant quantitative differences between these two types of kernels.



$K_{e,2}$ with $p = 0.5, q = 4,$
 $a = 1, b = 1$, and $N = 1000$



$K_{e,1}$ with $p = 0.5, q = 4,$
 $a = 1, b = 1$, and $N = 1000$

Even though $p = 1, q = 4$ led to differences in dimensionality, other parameters such as $p = 0.5, q = 4$ led to minimizers of the same dimension. We attempted to describe this behavior using an inequality described in the Analytic Results section.

Analytic Results

We can write K using K_s as

$$K_s(x - y) = K(\|x - y\|_2)$$

and we can represent the elliptical behavior using $f_{*,c} : \mathbb{R}^2 \rightarrow \mathbb{R}$ given by

$$f_{*,c}(x - y) = \left(\left(\frac{x_1 - y_1}{a} \right)^c + \left(\frac{x_2 - y_2}{b} \right)^c \right)^{1/c}$$

to write

$$K_{e,c}(x - y) = K(f_{*,c}(x - y)).$$

Consider a density function, ρ , which models the distribution of particles. For $\rho \geq 0$, we have

$$\int_{\mathbb{R}^2} \rho(x) dx = m.$$

For K_s , ρ is given by

$$\rho(x) = \frac{m}{\pi} \mathcal{X}_{B(0,1)}(x).$$

Our energy functional is given by

$$\mathcal{E}[\rho] = \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} K(\|x - y\|_2) \rho(x) \rho(y) dx dy.$$

Using

$$\tilde{x} = \left(\frac{x_1}{a}, \frac{x_2}{b} \right), \quad \tilde{y} = \left(\frac{y_1}{a}, \frac{y_2}{b} \right), \quad \tilde{\rho}(x) = ab\rho(x),$$

we were able to represent the minimizer of the $p = 0, q = 2$ case of K_s in terms of the $p = 0, q = 2$ case of $K_{e,2}$.

We also attempted to compare the energies of minimizers of $K_{e,1}$ and $K_{e,2}$ analytically. Our calculations revealed the following relation

$$\begin{aligned} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} K(\|x - y\|_1) \rho(x) \rho(y) dx dy &\leq \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} K(\|x - y\|_2) \rho(x) \rho(y) dx dy \\ &\leq 2^{\frac{q}{2}} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} K(\|x - y\|_1) \rho(x) \rho(y) dx dy \\ &\quad + \frac{2^{\frac{A}{2}}}{p} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} \|x - y\|_1^p \rho(x) \rho(y) dx dy, \end{aligned}$$

where

$$A = \frac{2 \log(2^{\frac{q}{2}} - 2^{\frac{p}{2}})}{\log(2)}.$$

Acknowledgements

We would like to thank UROP for funding our research. We would also like to thank the staff of the VCU High Performance Research Computing for allowing us to use the Teal Cluster in order to run the MATLAB code for our numerical results.