

Complex-scaling method for the complex plasmonic resonances of particles with corners

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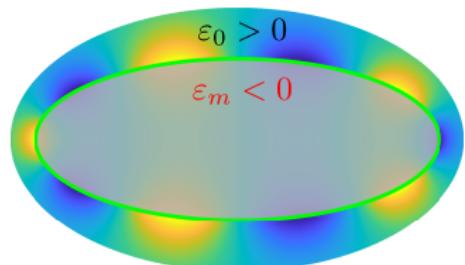
4 Numerical results using corner perturbations

5 Conclusion

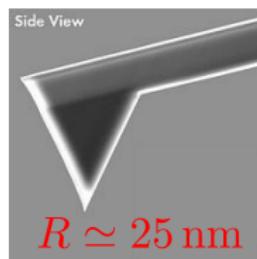
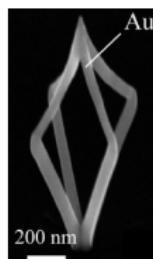
Motivation & objective

Motivation: Light concentration using “surface plasmons”.

Surface plasmon



Photonic devices



(Lalanne et al. 2018) (NanoWorld 2022)

Computational challenges

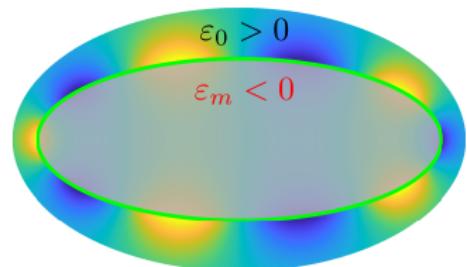
► **Interface geometry**

► Nonlinear materials

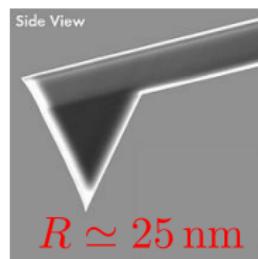
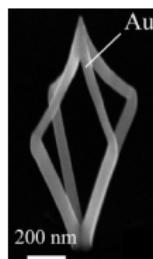
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Objective: Evidence of **complex resonances** associated with a sign-changing corner.

1 – What are “complex resonances”? (Zworski 2017)

In scattering, complex resonances model **energy leaking at infinity**.

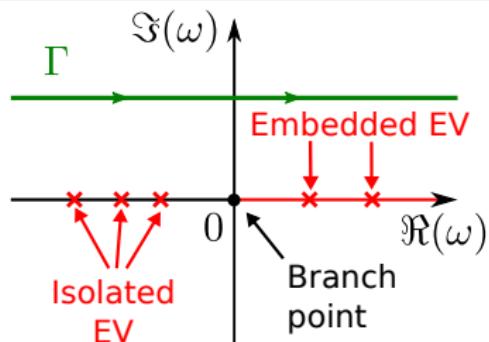
$$i\partial_t \psi(t, \mathbf{x}) = H\psi(t, \mathbf{x}) + f(\mathbf{x}), \quad \psi(0, \mathbf{x}) = 0 \quad (x \in \mathbb{R}^3).$$

The wave function is formally given by

$$\psi(t, \mathbf{x}) = \frac{1}{2\pi} \int_{\Gamma} R(\omega) f(\mathbf{x}) e^{-i\omega t} d\omega \quad (t > 0),$$

where the outgoing resolvent is $R(\omega) = (H - \omega I)^{-1}$ for $\Im(\omega) > 0$.

$\sigma(H) \neq \sigma_p(H)$, bound states and quasi-normal modes



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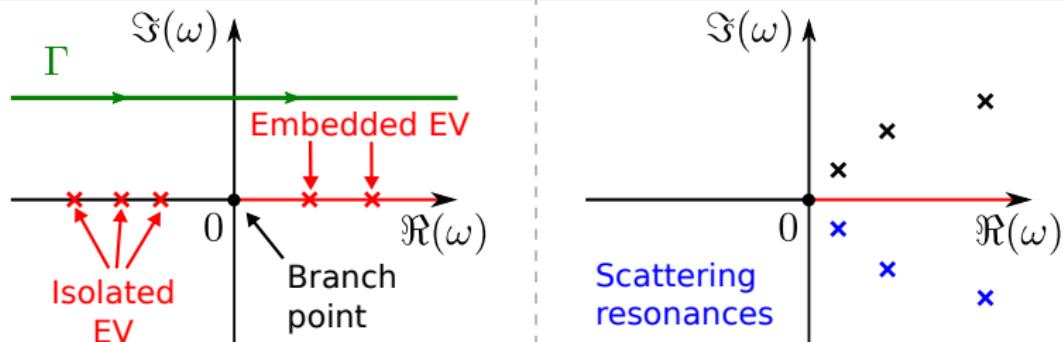
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⇒ This work investigates a **plasmonic analogue** of scattering resonances.

2 – Corners as a source of essential spectrum (Bonnet-Ben Dhia et al. 2013)

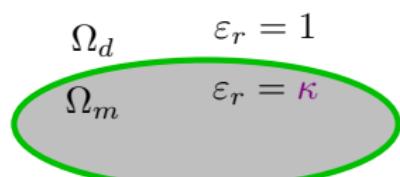
Plasmonic Eigenvalue Problem (PEP)

Find $(u, \kappa) \in H_0^1(\Omega) \times \mathbb{C}$ such that

$$\operatorname{div} [\varepsilon_r(\kappa)^{-1} \nabla u] = 0$$

with piecewise-constant permittivity:

$$\varepsilon_r(\kappa) = \kappa \mathbb{1}_{\Omega_m} + \mathbb{1}_{\Omega_d}$$



⚠ Spectral parameter is **contrast** κ .

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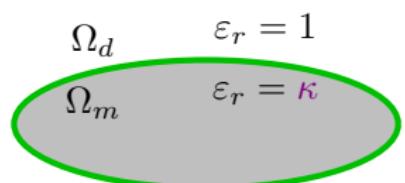
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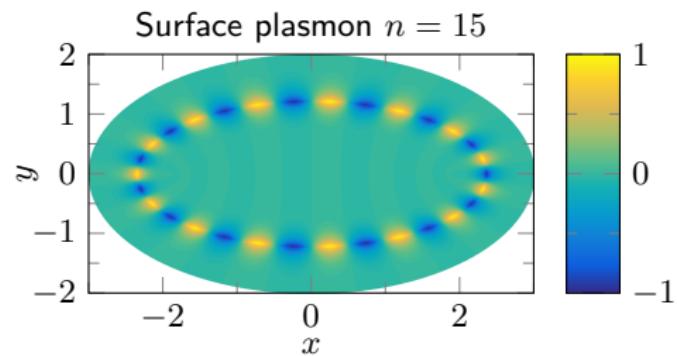
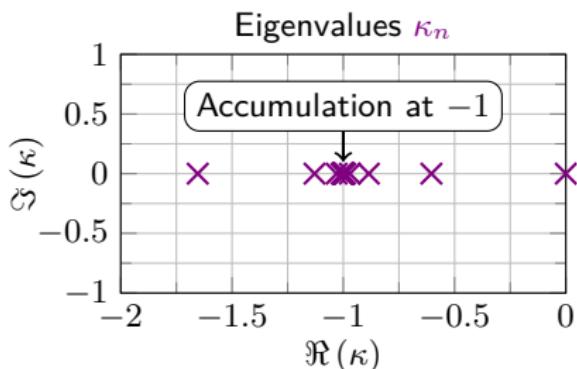
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⚠ Spectral parameter is contrast κ .

Point spectrum in H_{loc}^1 : $\kappa_n < 0$, $\kappa_n \rightarrow -1$ (Grieser 2014, Thm. 1).



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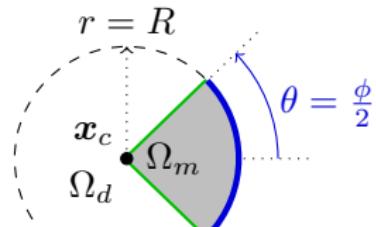
Local problem around corner x_c :

$$\operatorname{div} [\varepsilon_r(\kappa)^{-1} \nabla u] = 0 \quad (\star).$$

Countable family of solutions:

$$u_\eta(r, \theta) = r^{i\eta} \times \Phi_\eta(\theta) \quad (\eta \in \mathbb{H}(\kappa, \phi)),$$

where $\Phi_\eta \in H^1_{\text{per}}(-\pi, \pi)$.



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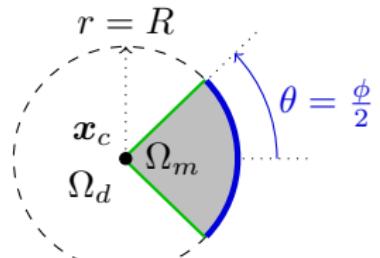
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There is a **critical interval** I_c such that

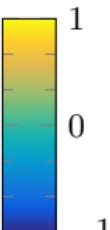
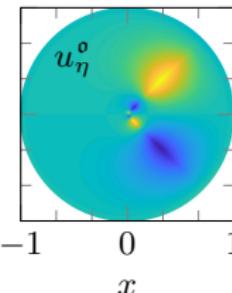
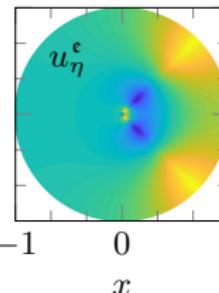
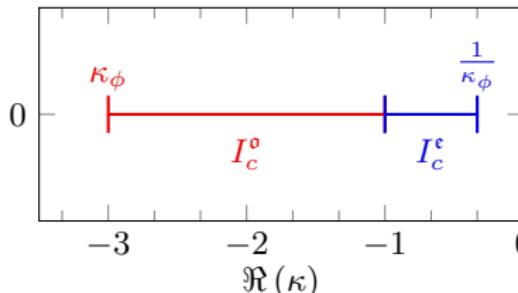
$$\kappa \in I_c \iff \exists \eta_{bh} \in \mathbb{R} : u_{\eta_{bh}} \text{ solves } (\star).$$

⚠ $u_{\eta_{bh}} \in L^2_{\text{loc}} \setminus H^1_{\text{loc}}$ is a strongly-oscillating “black-hole” wave.

Critical interval $I_c = I_c^o \cup I_c^e$

$\kappa \in I_c^e, \eta = 2$

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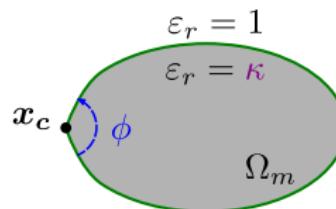


Objectives and outline

Plasmonic Eigenvalue Problem (PEP)

Find $(u, \kappa) \in U \times \mathbb{C}$ such that

$$\operatorname{div} \left[\varepsilon_r(\kappa)^{-1} \nabla u \right] = 0.$$



Objective: Numerical evidence of complex resonances for a piecewise-smooth negative particle.

Outline

2 Definition of complex plasmonic resonances

What are they?

3 Applicability of corner complex scaling

How to compute them?

4 Numerical results using corner perturbations

Does this actually work?

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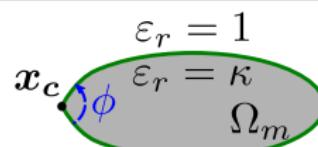
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- 2 Definition of complex plasmonic resonances
 - Definition
 - Sketch of construction
 - Characterization
- 3 Applicability of corner complex scaling
- 4 Numerical results using corner perturbations
- 5 Conclusion

Definition of complex plasmonic resonances

Let $\Omega_m \Subset \Omega \Subset \mathbb{R}^2$, $\partial\Omega_m$ smooth except for one corner.

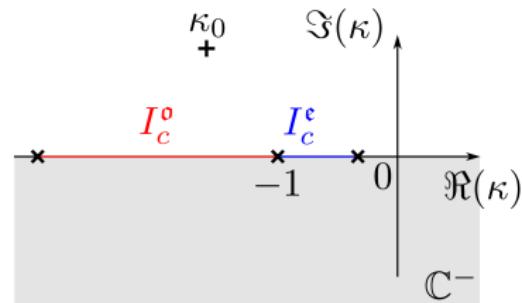
Operators: $\mathcal{A}(\kappa)u := \operatorname{div} [\varepsilon_r(\kappa)^{-1} \nabla u]$

$R(\kappa)f := \mathcal{A}(\kappa)^{-1}f$



When $\Im(\kappa) > 0$, $R(\kappa) : H^{-1}(\Omega) \rightarrow H_0^1(\Omega)$ is bounded.

When $\kappa \rightarrow I_c^o \cup I_c^e$, $\|R(\kappa)f\|_{H^1(\Omega)} \rightarrow \infty$ (Bonnet-Ben Dhia et al. 2013).



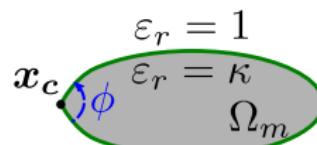
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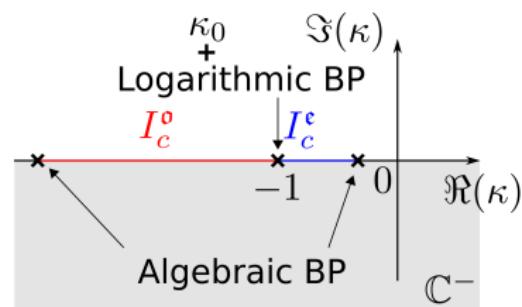
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Crossing \mathbb{R} once yields 3 continuations of $\kappa \mapsto R(\kappa)$:



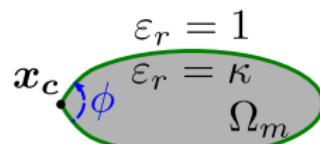
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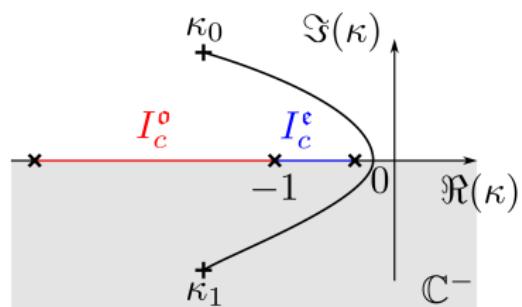
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Crossing \mathbb{R} once yields **3 continuations** of $\kappa \mapsto R(\kappa)$:

$$\tilde{R}(\kappa) := \overline{R(\bar{\kappa})} \quad (\kappa \in \mathbb{C}^- \cup \mathbb{R} \setminus \overline{I_c})$$



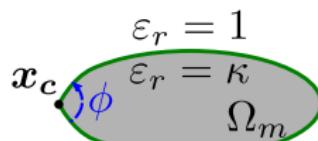
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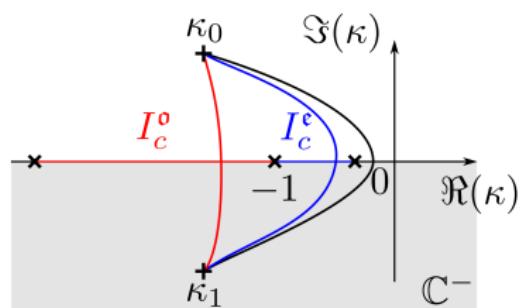


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$$R^{|\textcolor{blue}{e}(\textcolor{red}{o})}(\kappa) \quad (\kappa \in \mathbb{C}^- \cup I_c^{|\textcolor{blue}{e}(\textcolor{red}{o})})$$



Definition. A *complex plasmonic (CP) resonance* is a pole of $\kappa \mapsto R^{|\textcolor{blue}{e}}(\kappa)$ or $\kappa \mapsto R^{|\textcolor{red}{o}}(\kappa)$ in \mathbb{C}^- .

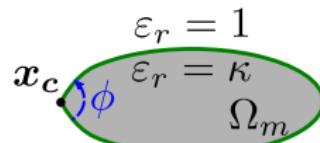
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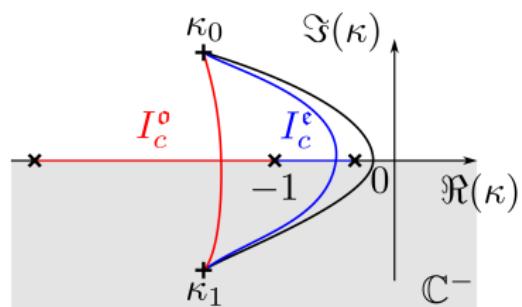


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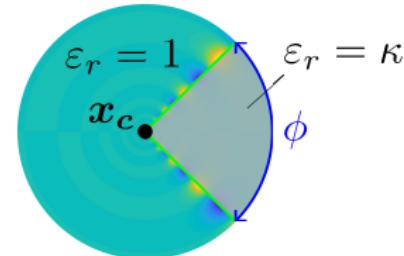
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Next: characterization of resonance functions as $x \rightarrow x_c$.

Sketch of Mellin analysis (1) (Dauge and Texier 1997)

Local problem for $\kappa \in \mathbb{C}^+$:

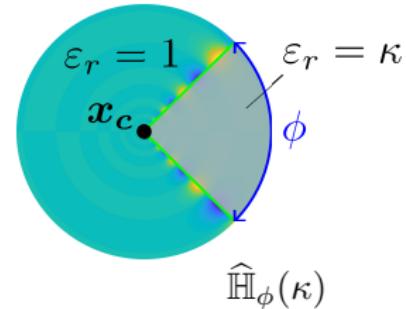
$$\operatorname{div} \left[\varepsilon_r(\kappa)^{-1} \nabla u \right] = 0 \quad (\mathbf{x} \in D).$$



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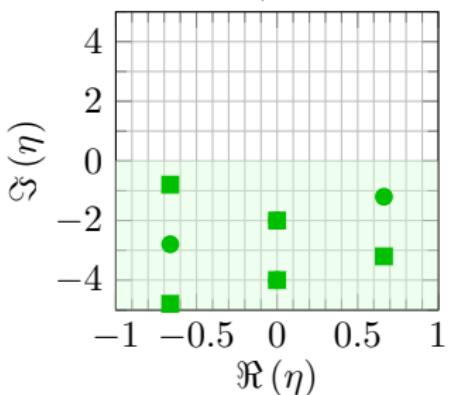


Expansion: If $u \in H^1(D)$ then $\forall \eta_* < 0$,

$$u \underset{r \rightarrow 0}{=} c_0 + \sum_{\substack{\eta \in \widehat{\mathbb{H}}_\phi(\kappa) \\ \Im(\eta) > \eta_*}} c_\eta r^{i\eta} \Phi_\eta(\theta) + \mathcal{O}(r^{-\eta_*})$$

with $\Phi_\eta \in H_{\text{per}}^1(-\pi, \pi)$ and

$$\widehat{\mathbb{H}}_\phi(\kappa) := \{ \eta \mid f_\phi(\eta, \kappa) = 0, \Im(\eta) < 0 \}.$$



Strategy: Characterize resonance functions by studying the continuation to \mathbb{C}^- of the map

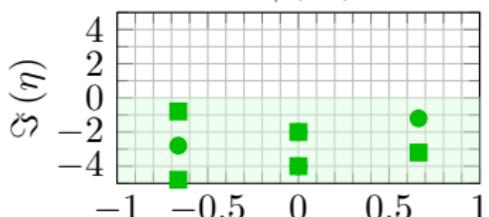
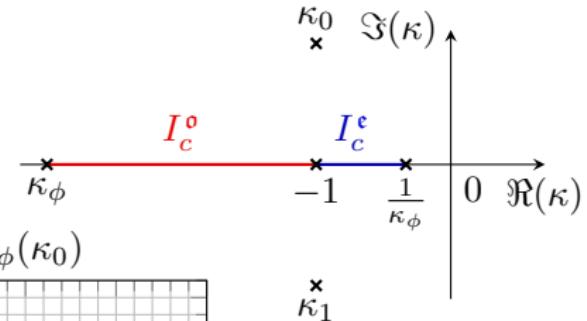
$$\mathbb{C}^+ \ni \kappa \mapsto \widehat{\mathbb{H}}_\phi(\kappa).$$

Sketch of Mellin analysis (2)

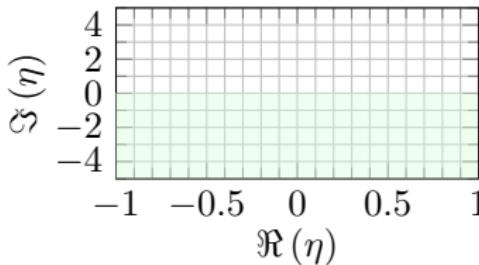
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$$\Gamma : (0, 1) \rightarrow \mathbb{C}, \quad \Gamma(0) = \kappa_0, \quad \Gamma(1) = \kappa_1$$

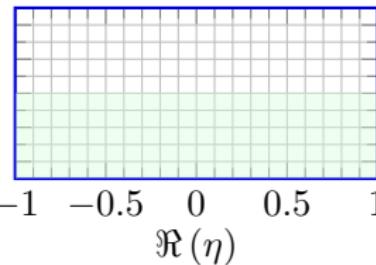
We track $\widehat{\mathbb{H}}_\phi(\kappa)$ as κ moves.



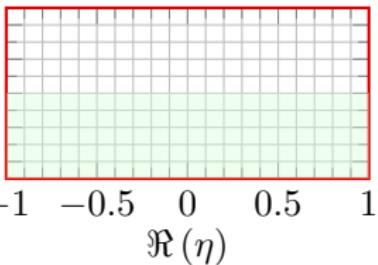
$\widehat{\mathbb{H}}_\phi(\kappa_1)$



$\widehat{\mathbb{H}}_\phi^{|\epsilon}()$



$\widehat{\mathbb{H}}_\phi^{|\sigma}()$

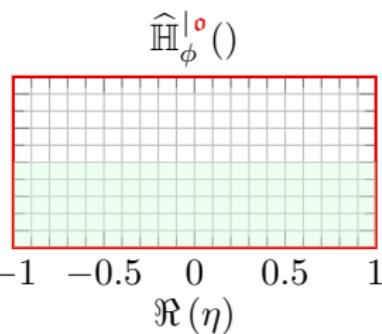
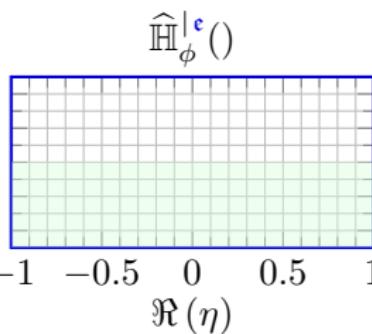
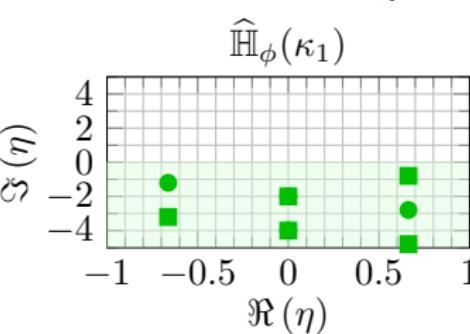
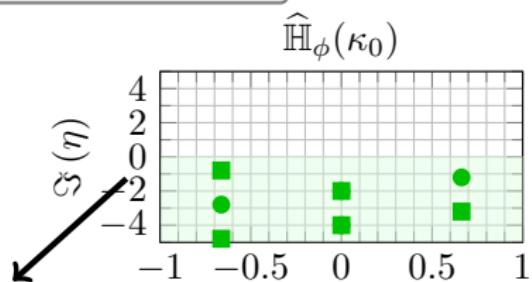
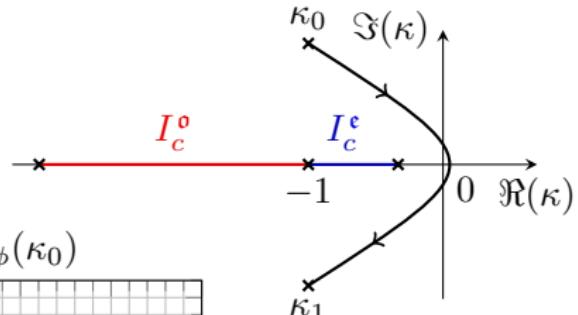


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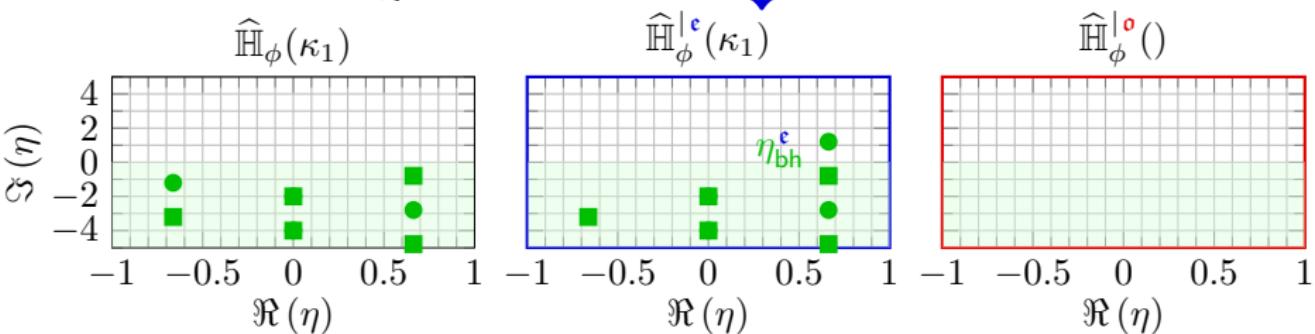
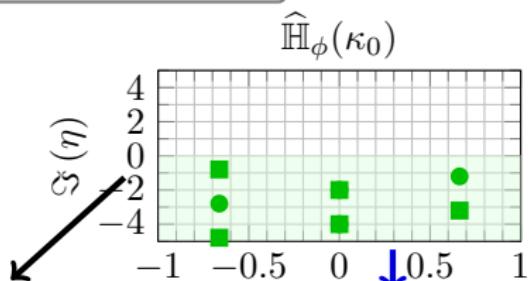
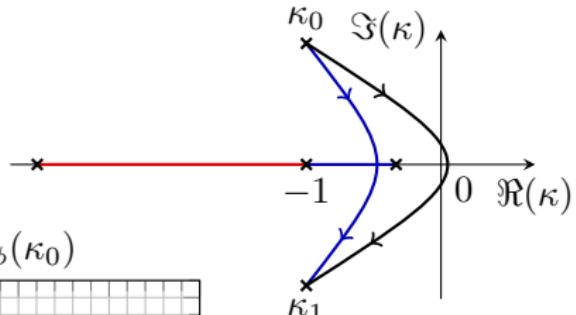


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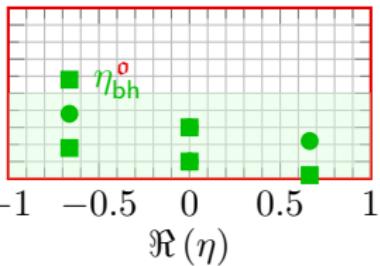
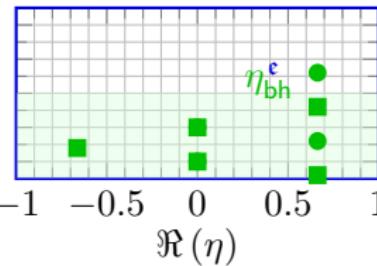
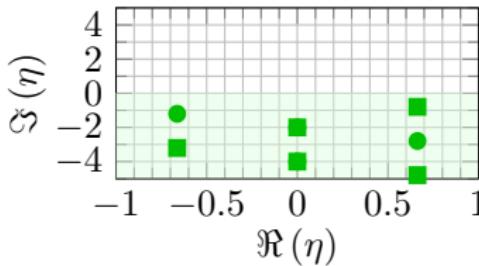
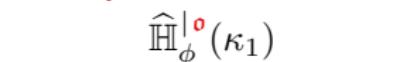
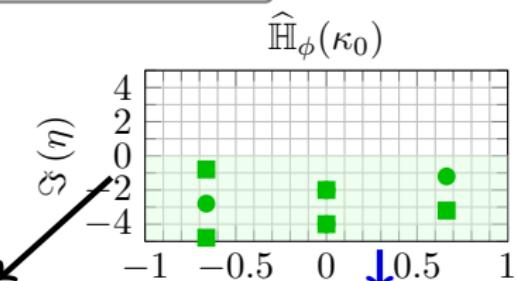
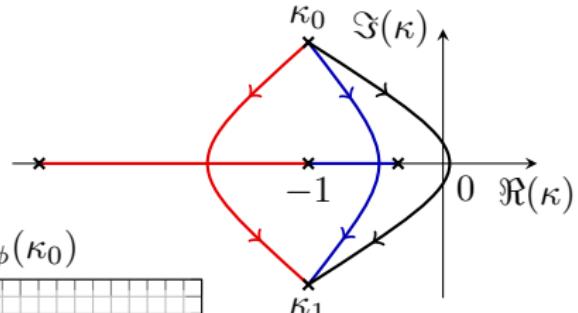


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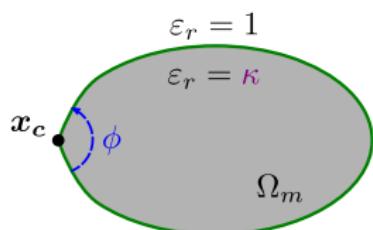
$$\Gamma : (0, 1) \rightarrow \mathbb{C}, \quad \Gamma(0) = \kappa_0, \quad \Gamma(1) = \kappa_1$$

We track $\widehat{\mathbb{H}}_\phi(\kappa)$ as κ moves.

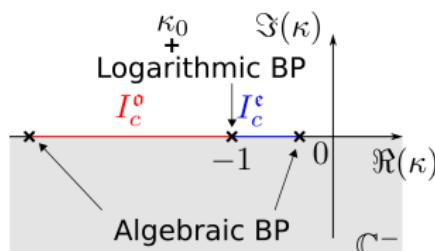


Characterization of resonance functions

Particle:



Spectrum:



Characterization. $\kappa \in \mathbb{C}^-$ is a resonance $\Leftrightarrow \exists u \notin H_{\text{loc}}^1(\Omega)$:

$$\operatorname{div} [\epsilon_r(\kappa)^{-1} \nabla u](x) = 0 \quad (x \neq x_c), \quad u|_{\partial\Omega} = 0,$$

$$u(r, \theta) \underset{r \rightarrow 0}{\sim} c_1 r^{i\eta_{bh}} \Phi(\theta) + c_0,$$

where $\eta_{bh} = \eta_{bh}(\kappa)$ and $\Im(\eta_{bh}) > 0$.

Next: applicability of corner complex scaling

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- 2 Definition of complex plasmonic resonances
- 3 Applicability of corner complex scaling
 - Corner complex scaling
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- 5 Conclusion

Corner complex scaling: formulation

Principle. Let $\alpha \in \mathbb{C}$. Define a non self-adjoint “ $\text{PEP}\alpha$ ” such that:
 κ complex resonance of $\text{PEP} \iff \kappa$ eigenvalue of $\text{PEP}\alpha$.

Intuitively, we would like

$$(\text{PEP}) \quad u_{\text{res}} \underset{r \rightarrow 0}{\sim} e^{i\eta \ln r} \Phi_\eta(\theta) + c_0 \quad (\Im(\eta) > 0)$$

\downarrow

$$(\text{PEP}\alpha) \quad u_{\text{res},\alpha} \underset{r \rightarrow 0}{\sim} e^{i\frac{\eta}{\alpha} \ln r} \Phi_\eta(\theta) + c_0 \quad \left(\Im\left(\frac{\eta}{\alpha}\right) < 0 \right)$$

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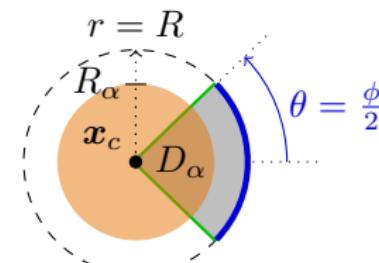
$$(\text{PEP}\alpha) \quad u_{\text{res},\alpha} \underset{r \rightarrow 0}{\sim} e^{i\frac{\eta}{\alpha} \ln r} \Phi_\eta(\theta) + c_0 \quad \left(\Im\left(\frac{\eta}{\alpha}\right) < 0 \right)$$

Definition of $\text{PEP}\alpha$. Substitution

$$r\partial_r \rightarrow \alpha r\partial_r$$

around the corner.

(Bonnet-Ben Dhia, Carvalho, Chesnel, and Ciarlet
2016)



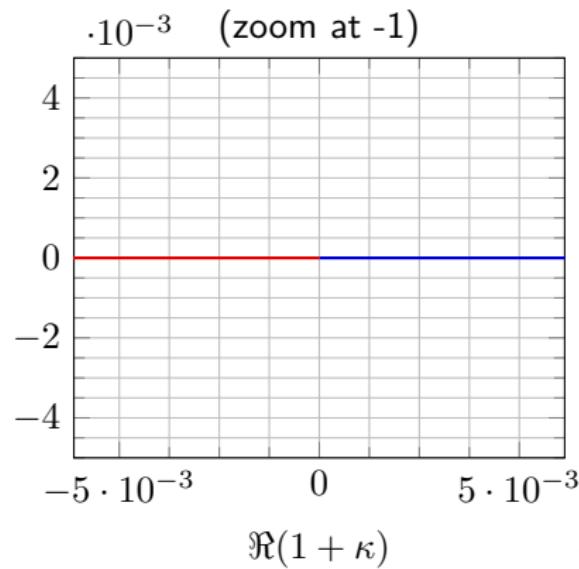
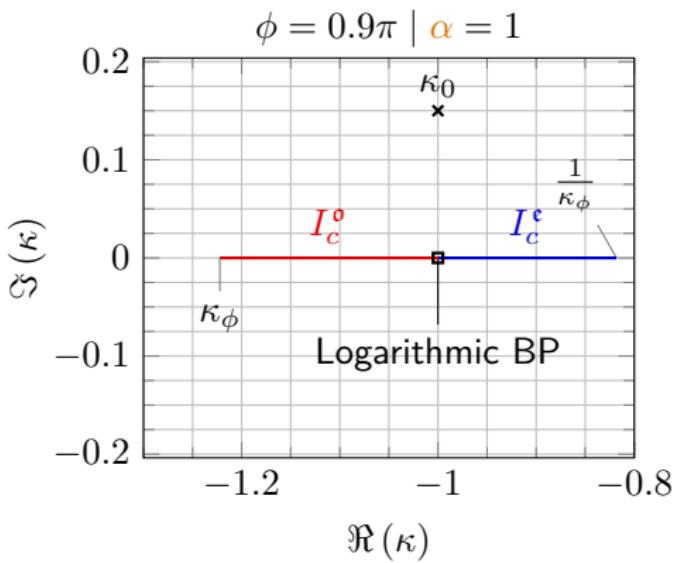
Next: Domain of validity?

Corner complex scaling: uncovered region

Definition of PEP α . Substitution $r\partial_r \rightarrow \alpha r\partial_r$ around the corner.
(Bonnet-Ben Dhia, Carvalho, Chesnel, and Ciarlet 2016)

Proposition. Let κ be an eigenvalue of PEP α with $\alpha \in \mathbb{C} \setminus \mathbb{R}$. Then,

$$\kappa \in U_\phi^\alpha \Rightarrow \kappa \text{ is a complex resonance.}$$

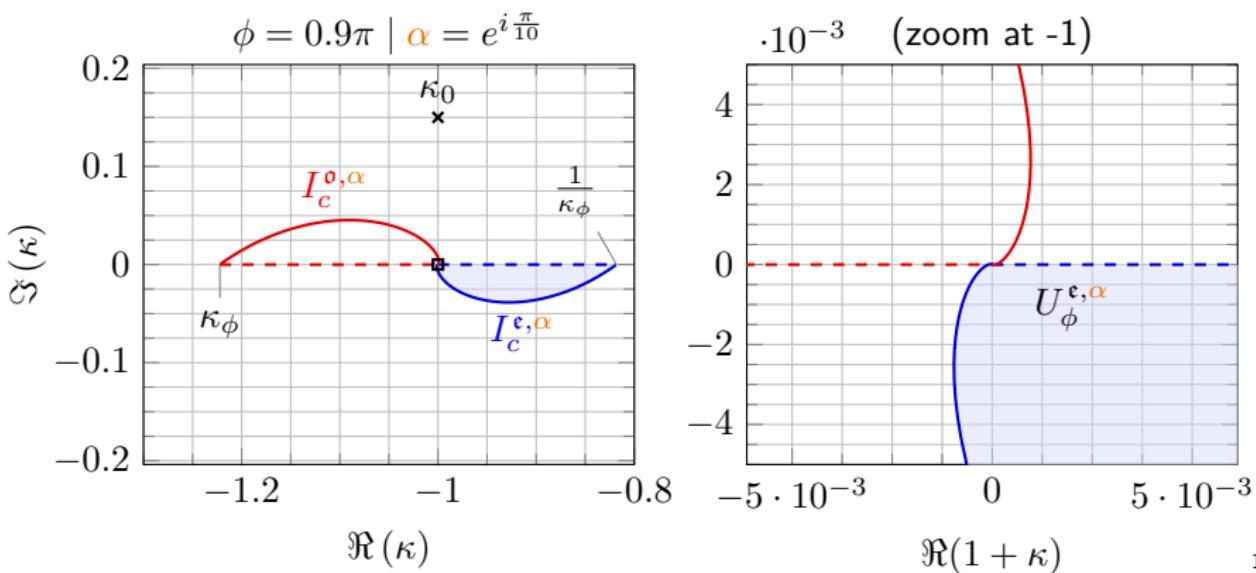


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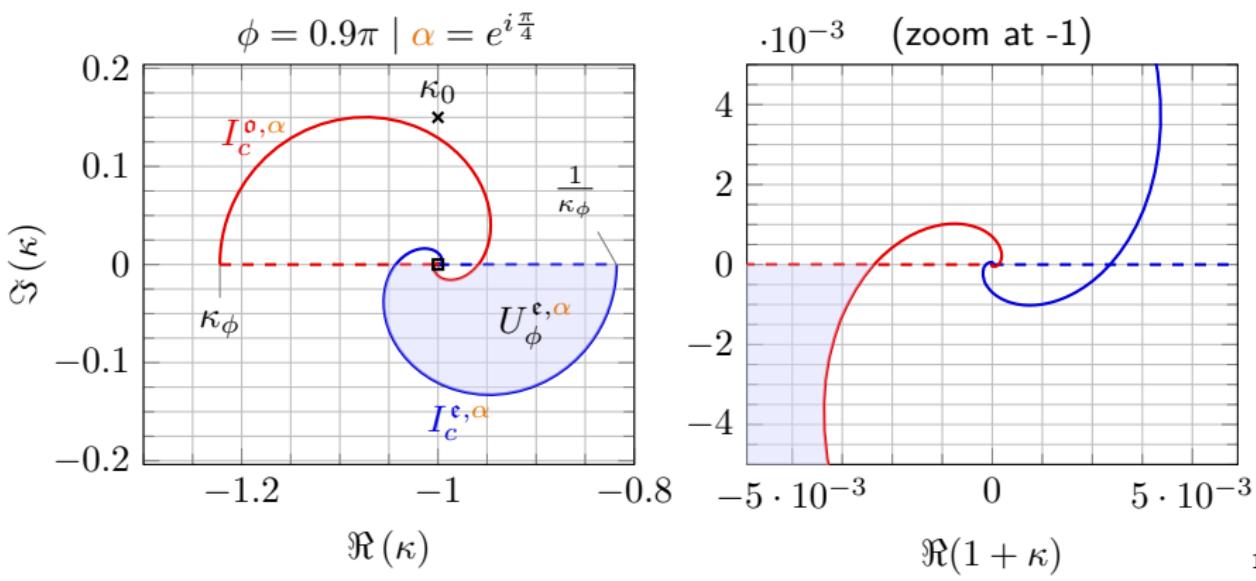


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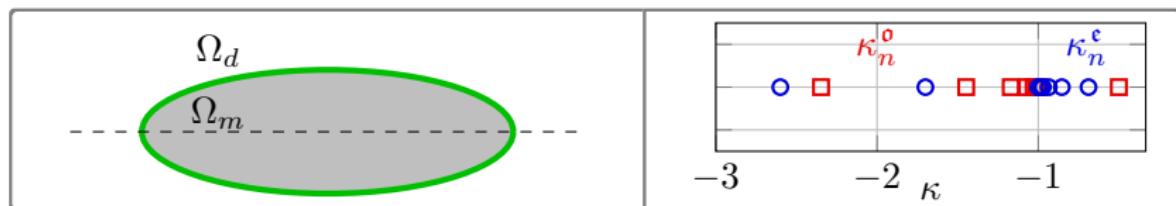
How to obtain complex resonances?

Perturbation of elliptical Ω_m by corner along major axis.

Embedded eigenvalues

→ Existence proof (Li and Shipman 2019, § 5.2)

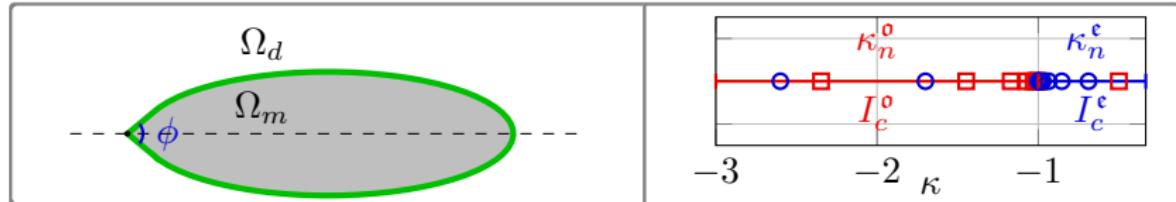
→ Numerical evidence (Helsing, Kang, and Lim 2017)



+



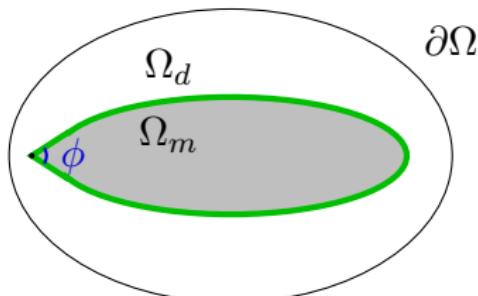
=



Discretization without scaling

Geometry. Piecewise-smooth $\partial\Omega_m$.

- ▶ Ellipse perturbed by a straight corner of angle $\phi \in (0, \pi)$.
- ▶ \mathcal{C}^1 junction.



Weak Formulation: Find $(u, \kappa) \in H_0^1(\Omega) \times \mathbb{C}$ s.t.

$$\forall v \in H_0^1(\Omega), \int_{\Omega_m} \nabla u(\mathbf{x}) \cdot \nabla v(\mathbf{x}) \, d\mathbf{x} = -\kappa \int_{\Omega_d} \nabla u(\mathbf{x}) \cdot \nabla v(\mathbf{x}) \, d\mathbf{x}.$$

Discretization:

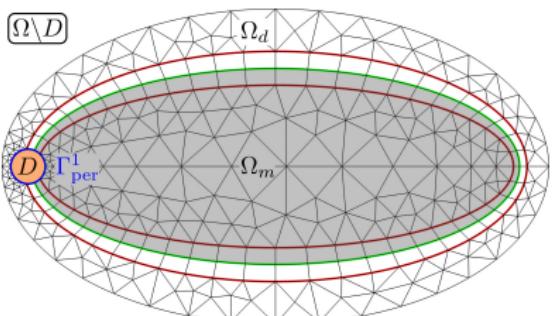
$$A_{\Omega_m} U = -\kappa A_{\Omega_d} U,$$

where $A_{\Omega_m}, A_{\Omega_d}$ are real symmetric and positive (but not definite).

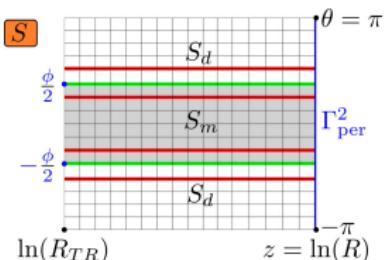
Next: addition of a complex scaling region around \mathbf{x}_c .

Discretization with scaling

$\partial\Omega_m = \text{ellipse perturbed by a corner of angle } \phi \in (0, \pi), \mathcal{C}^1 \text{ junction.}$



$$H_e := \{u \in H^1(\Omega \setminus \bar{D}) \mid u|_{\partial\Omega} = 0\}$$



Euler coordinates ($z = \ln(r), \theta$).

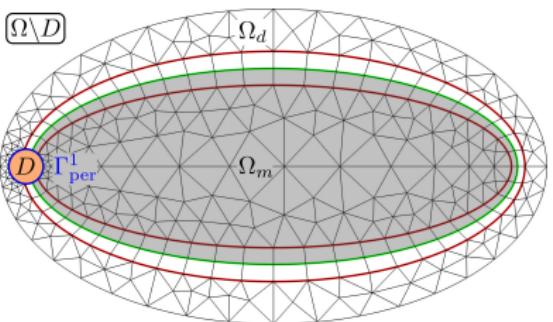
$$H_c := \{\check{u} \in H^1(S) \mid \check{u}(\cdot, \pi) = \check{u}(\cdot, -\pi)\}$$

Solution space:

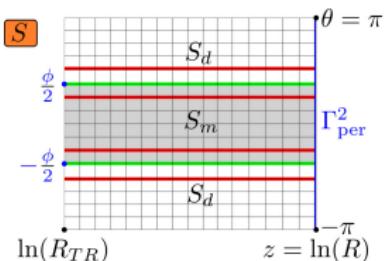
$$V = \left\{ (u, \check{u}) \in H_e \times H_c \mid u|_{\Gamma_{\text{per}}^1} = \check{u}|_{\Gamma_{\text{per}}^2} \right\}.$$

Discretization with scaling

$\partial\Omega_m = \text{ellipse perturbed by a corner of angle } \phi \in (0, \pi), \mathcal{C}^1 \text{ junction.}$



$$H_e := \{u \in H^1(\Omega \setminus \bar{D}) \mid u|_{\partial\Omega} = 0\}$$



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Discretization with H^1 -conforming elements (isoparametric P^2/Q^2).

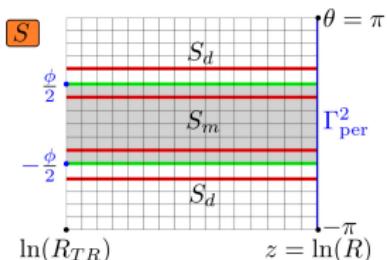
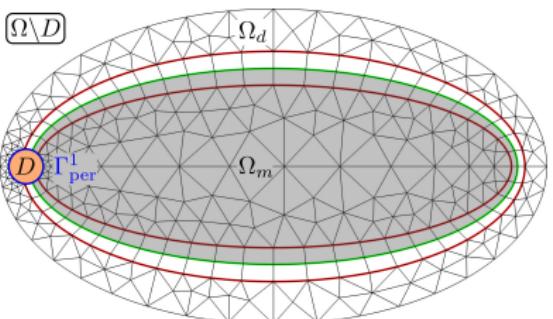
Find $(\kappa, U) \in \mathbb{C} \times \mathbb{C}^N$:

$$\left[A_{\Omega_m \setminus D}^{(x,y)} + \alpha A_{S_m}^{(z)} + \frac{1}{\alpha} A_{S_m}^{(\theta)} \right] U = -\kappa \left[A_{\Omega_d \setminus D}^{(x,y)} + \alpha A_{S_d}^{(z)} + \frac{1}{\alpha} A_{S_d}^{(\theta)} \right] U,$$

where all matrices are real.

Discretization with scaling

$\partial\Omega_m = \text{ellipse perturbed by a corner of angle } \phi \in (0, \pi), \mathcal{C}^1 \text{ junction.}$



Euler coordinates ($z = \ln(r), \theta$).

$$H_e := \{u \in H^1(\Omega \setminus \bar{D}) \mid u|_{\partial\Omega} = 0\}$$

$$H_c := \{\check{u} \in H^1(S) \mid \check{u}(\cdot, \pi) = \check{u}(\cdot, -\pi)\}$$

⚠ Mesh symmetry at $\partial\Omega_m$ to avoid spurious plasmons.

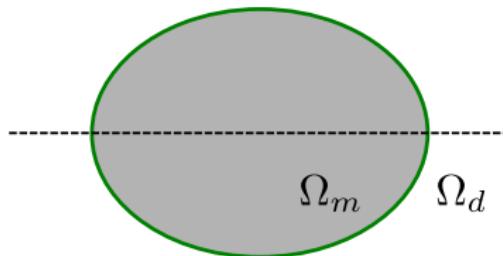
Proof for polygonal interfaces: (Bonnet-Ben Dhia, Carvalho, and Ciarlet 2018).

Methodology to deal with curvilinear $\partial\Omega_m$:

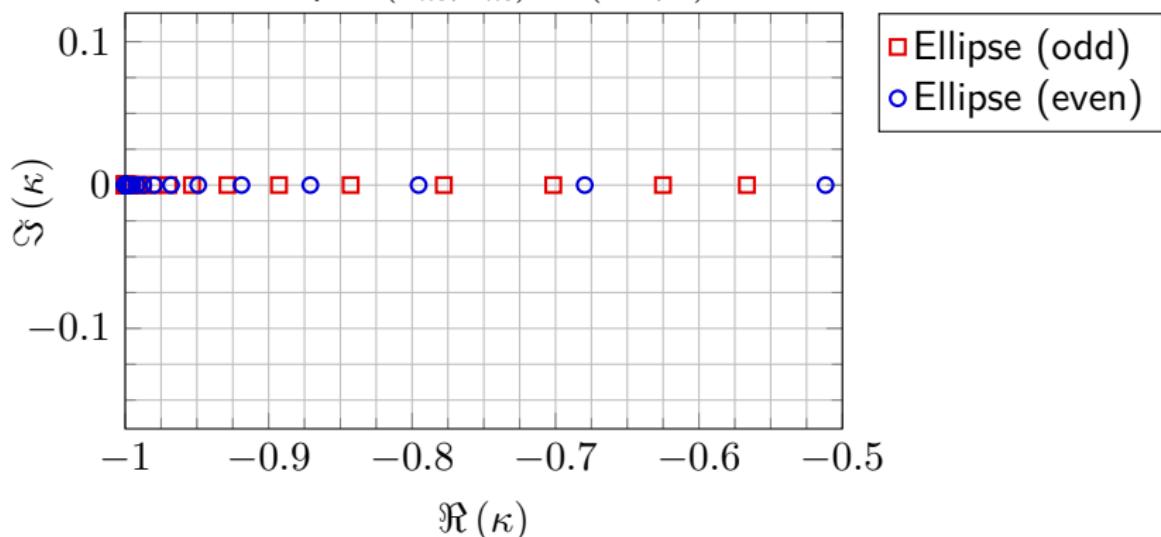
- ▶ One-cell thick **structured layer**.
- ▶ Symmetry w.r.t. elliptic coordinates (μ, θ) using isoparametric Q^2 .

Implementations COMSOL 5.4 and gmsh/dolfinx/PETSc/SLEPc.

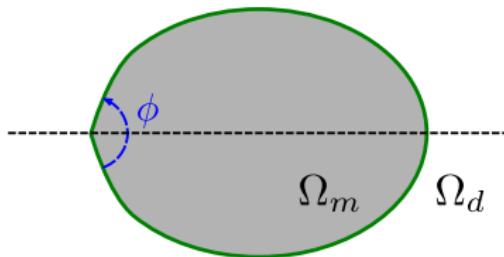
Results: corner perturbation along major axis (1)



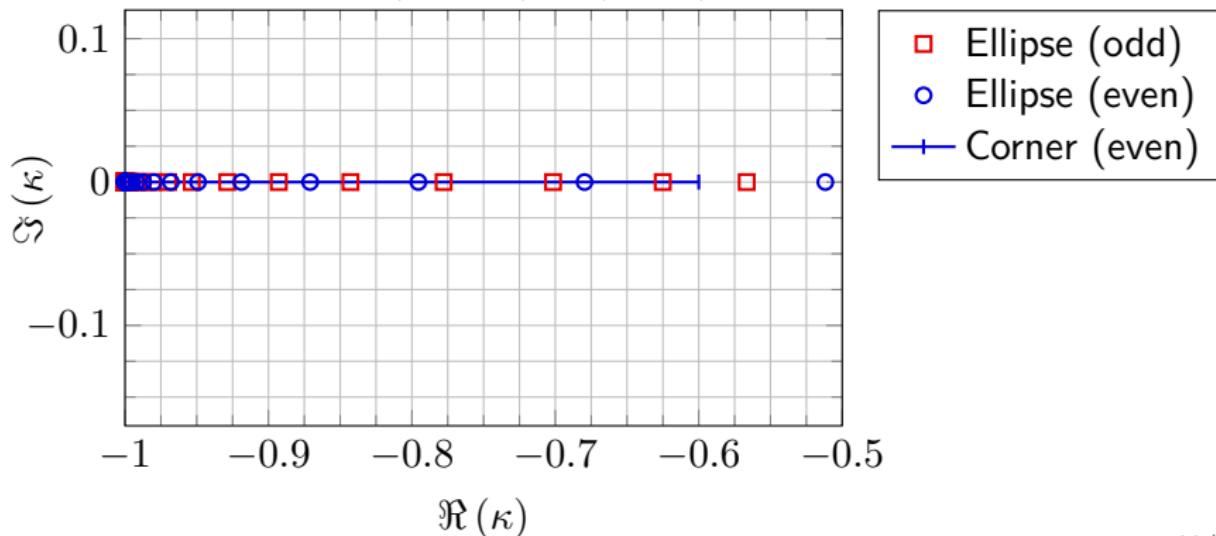
Ellipse $(a_m, b_m) = (2.5, 1)$



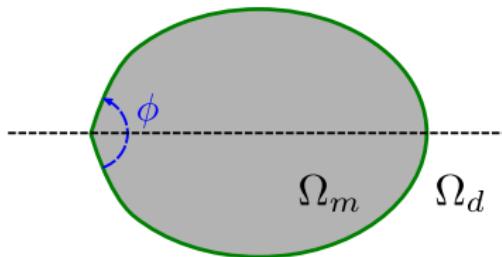
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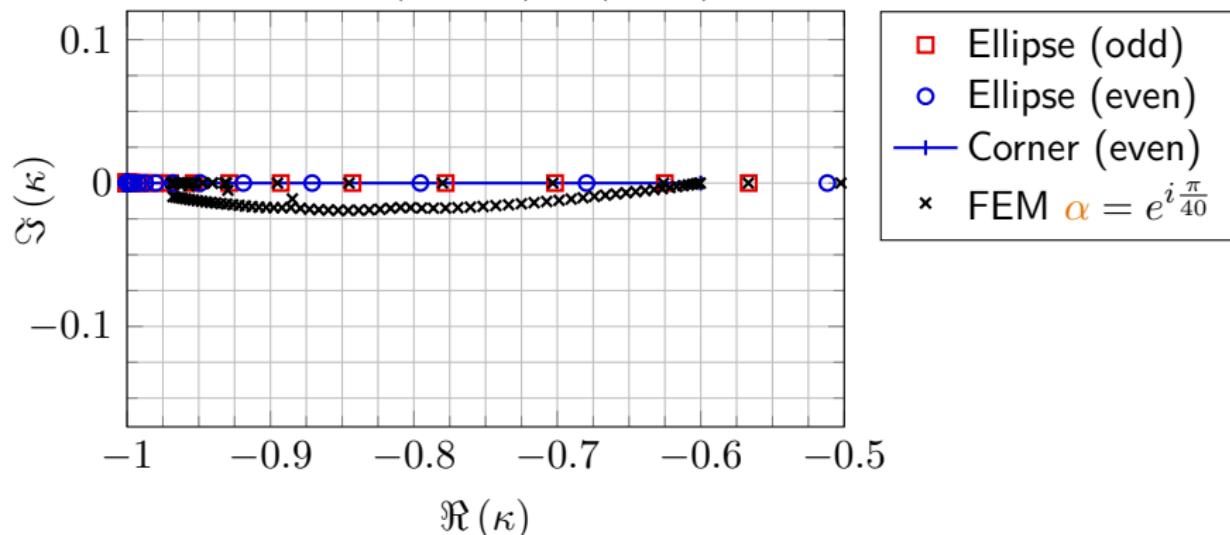
Perturbed ellipse: $(a_m, b_m) = (2.5, 1)$, $\phi = 0.75\pi$



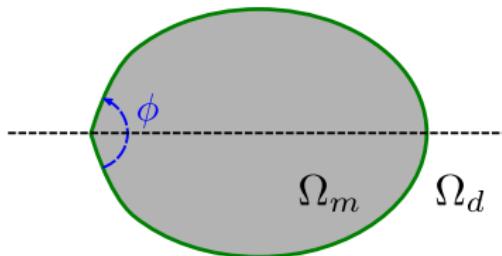
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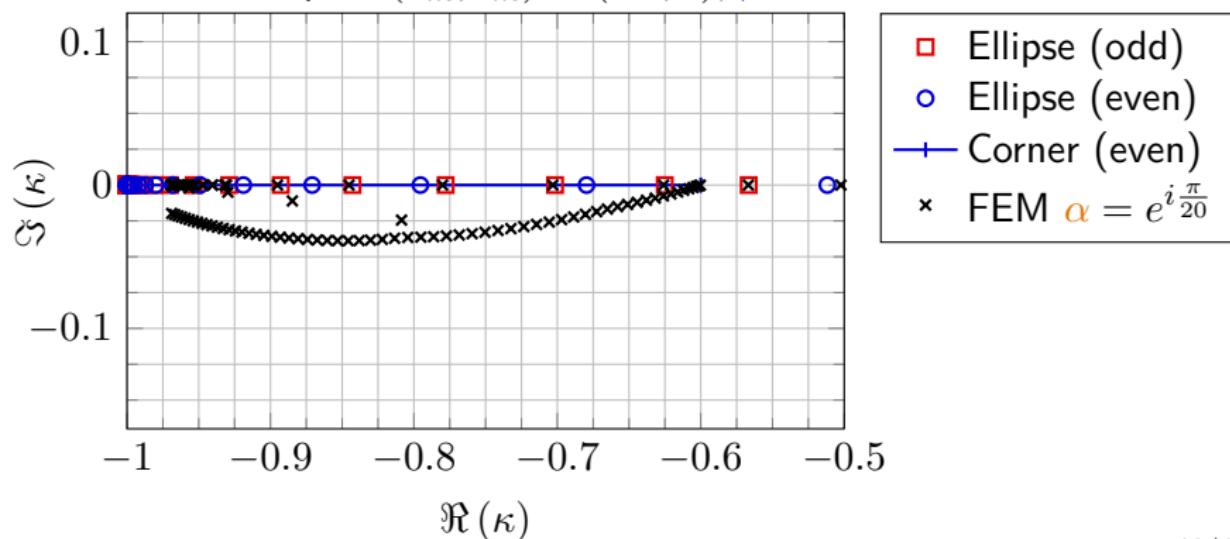
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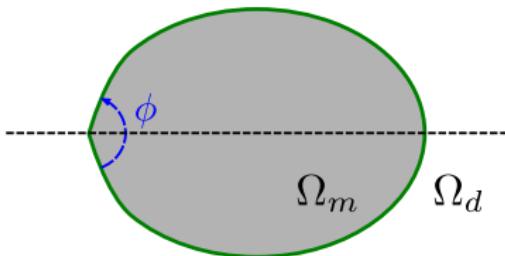
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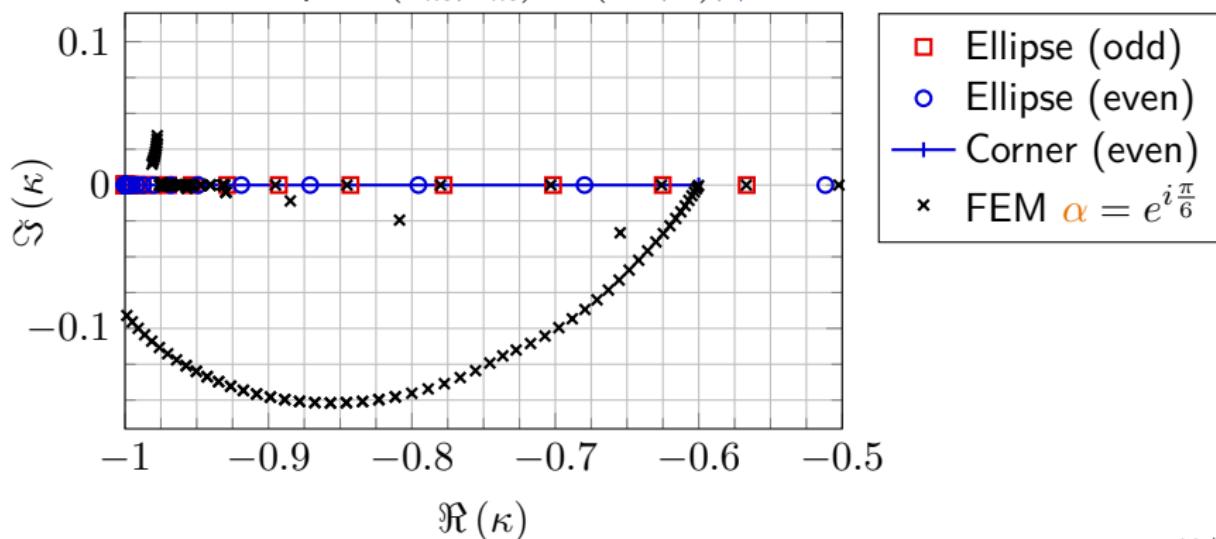
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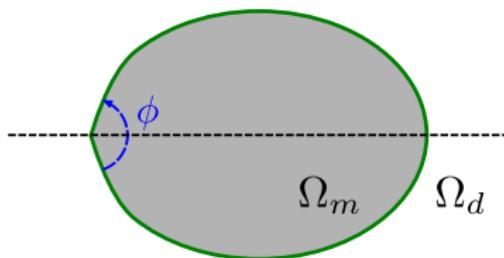
Results: corner perturbation along major axis (1)



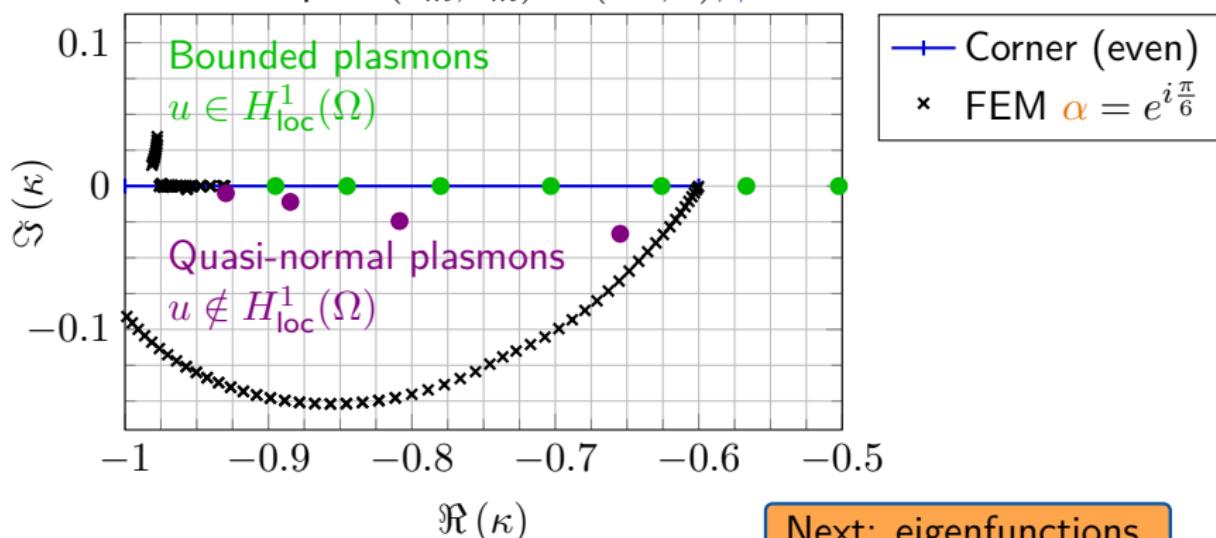
Perturbed ellipse: $(a_m, b_m) = (2.5, 1)$, $\phi = 0.75\pi$



Results: corner perturbation along major axis (1)



Perturbed ellipse: $(a_m, b_m) = (2.5, 1)$, $\phi = 0.75\pi$



Results: corner perturbations along major axis (2)

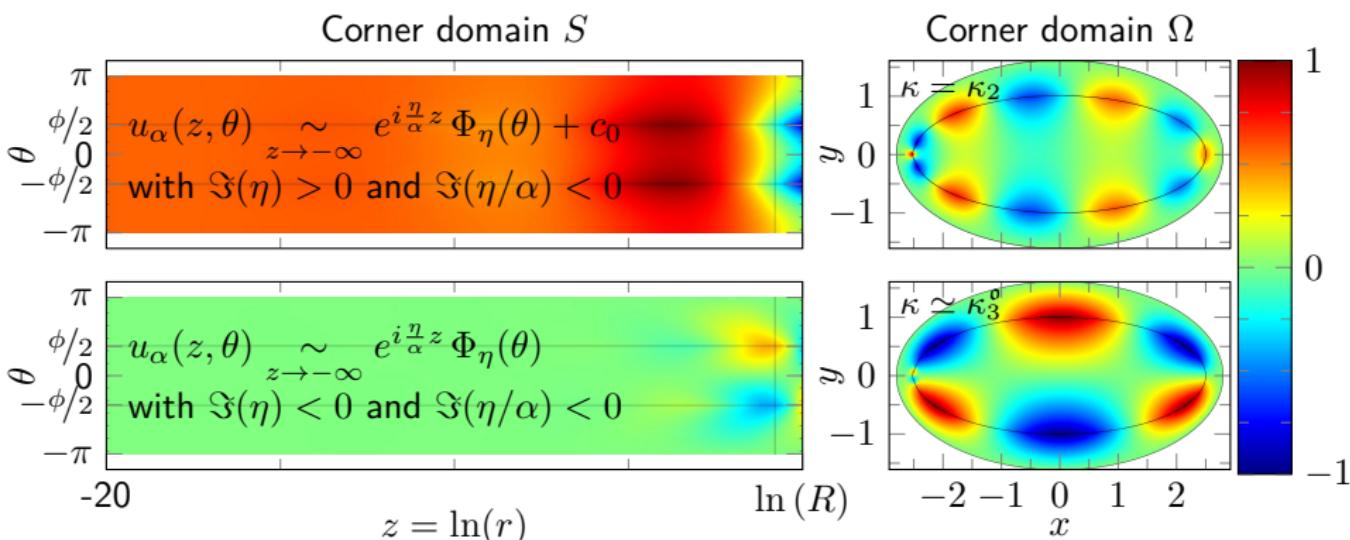
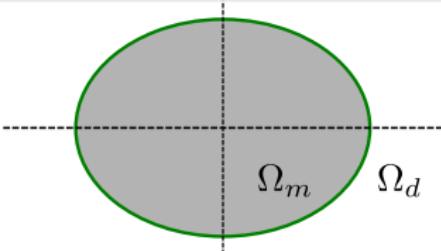


Fig. Eigenfunctions $\Re(u_\alpha)/\|u_\alpha\|_\infty$ of PEP- α with $\alpha = e^{i\frac{\pi}{6}}$.

(Top row) $\kappa = \kappa_2 \simeq 0.8086 - 0.02445i$, complex plasmonic resonance,

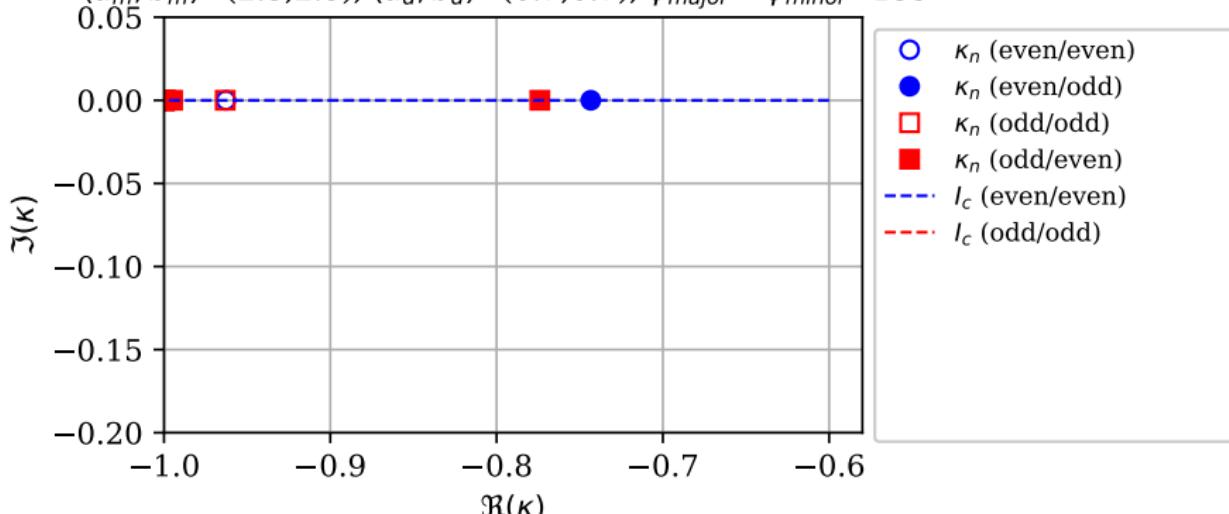
(Bottom row) $\kappa \simeq 0.70313 - 8.0357 \cdot 10^{-8}i \simeq \kappa_3^0$, embedded eigenvalue.

Results: corner perturbations along both major/minor axes

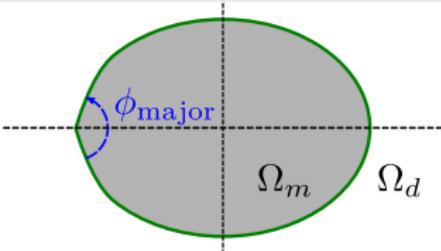


Elliptical particle perturbed by two corners

$(a_m, b_m) = (2.5, 2.5)$, $(a_d, b_d) = (6.7, 6.7)$, $\phi_{\text{major}} = \phi_{\text{minor}} = 135^\circ$

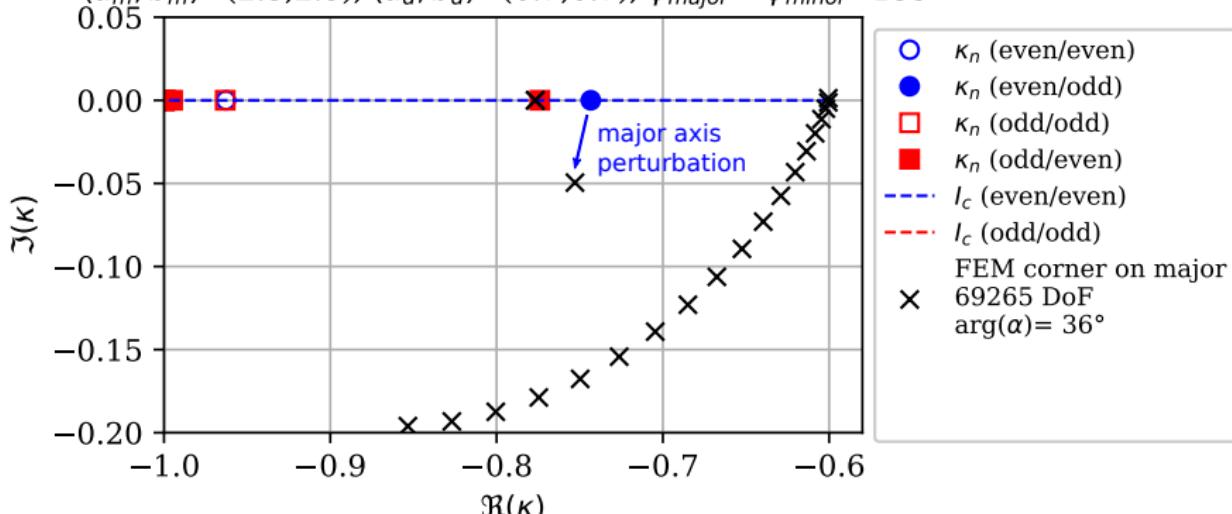


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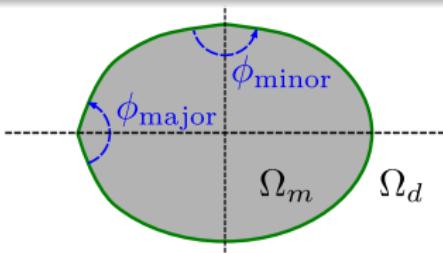


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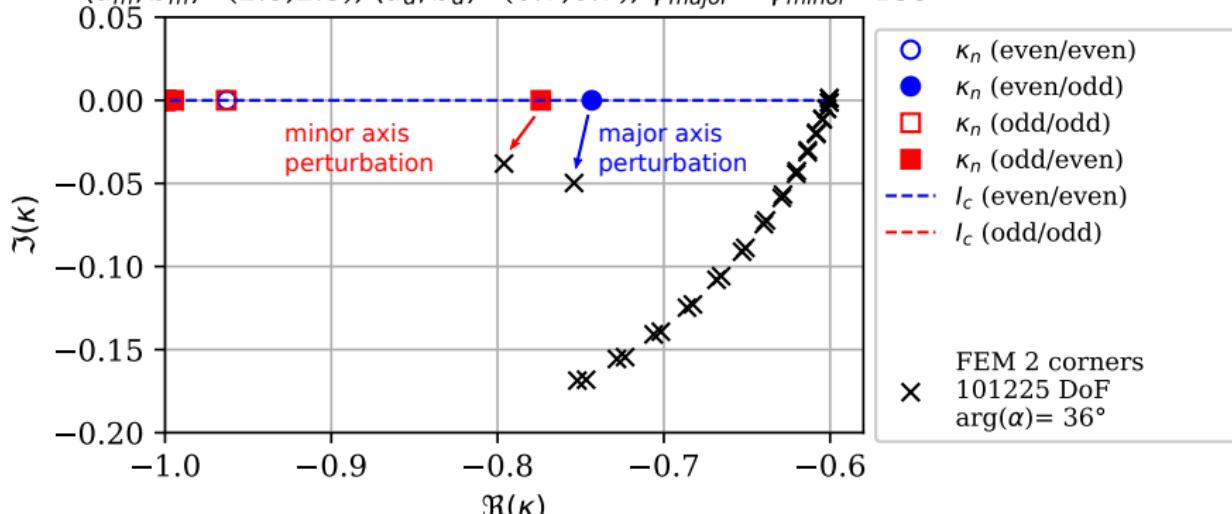


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 - Conclusion and outlook

Conclusions & outlook

◀ Appendix TOC

Takeaways

- ▶ Quasi-normal surface plasmons $u \notin H_{\text{loc}}^1(\mathbb{R}^2)$:
 - trap energy at corners,
 - are associated with complex resonances $\kappa = \frac{\epsilon}{\epsilon_0} \notin \mathbb{R}$,
 - are analogous to quasi-normal modes (“infinity \Leftrightarrow corner”)
- ▶ FE with corner complex scaling \Rightarrow linear eigenvalue problem
(Bonnet-Ben Dhia, Carvalho, Chesnel, and Ciarlet 2016)
- ▶ Agreement with (Li and Shipman 2019)

Outlook

- ▶ Interest of working with $\alpha(\kappa)$. (Nannen and Wess 2018)
- ▶ Properties and application of QNSP expansions. (Truong et al. 2020)
- ▶ Extension to e.g. $\Omega_m \subset \mathbb{R}^3$, Maxwell.
(Helsing and Perfekt 2018) (Li, Perfekt, and Shipman 2020)
(Bonnet-Ben Dhia, Chesnel, and Rihani 2022)

Conclusions & outlook

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