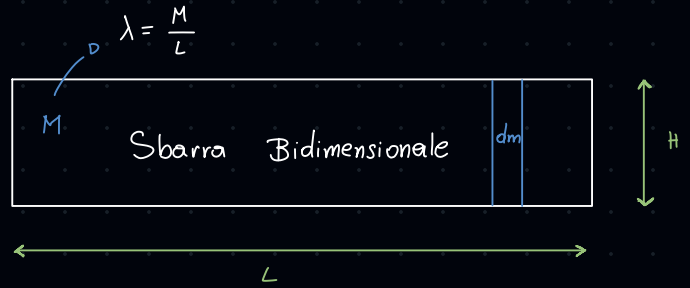
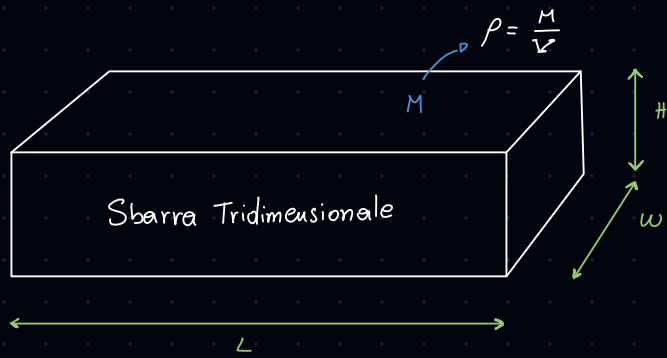


Momento di Inerzia di una sbarra

$$I = m R^2 \quad \rightarrow \quad I_i = m_i R_i^2 \quad \rightarrow \quad I_{\text{TOT}} = \sum_i m_i R_i^2$$



\Rightarrow La sbarra è un susseguirsi di punti CONTINUO $\Rightarrow \sum m_i \rightarrow \int dm$

$$\Rightarrow I_{\text{TOT}} = \sum_i m_i R_i^2 \quad \rightarrow \quad I_{\text{TOT}} = \int R^2 dm$$

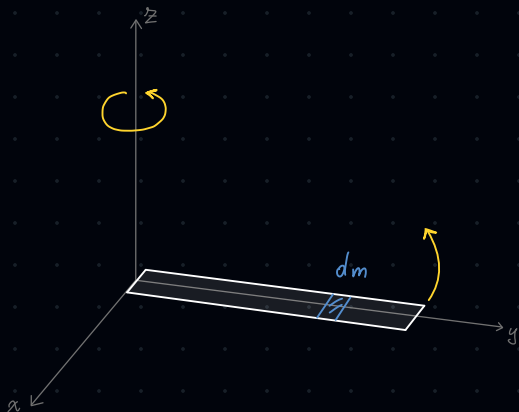
• Moltiplico e divido per dV :
$$I = \int_V R^2 \frac{dm}{dV} dV \quad \text{pongo } \rho = \frac{M}{V}$$

$$\rightarrow I = \int_V R^2 \rho dV$$

• Se la sbarra è a 2 dimensioni $\rho \rightarrow \lambda = \frac{M}{L}$ $R \rightarrow x$

$$\rightarrow I = \int_A^B \lambda x^2 dx$$

Rotazione lungo un estremo

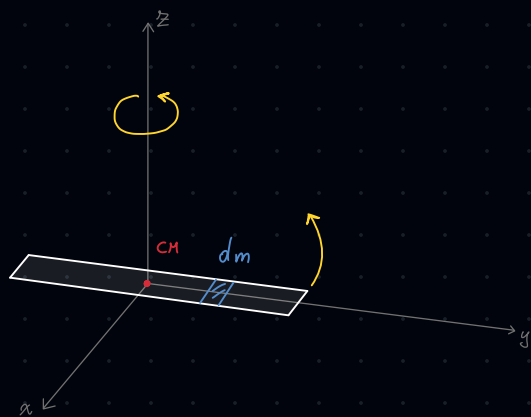


$$I = \int_A^B \lambda R^2 dV \quad \rightarrow \quad I = \int_0^L \lambda y^2 dy$$

$$\rightarrow I = \lambda \int_0^L y^2 dy = \lambda \left[\frac{y^3}{3} \right]_0^L = \lambda \left[\frac{L^3}{3} \right] = \frac{1}{3} \lambda L^3$$

ma $\lambda = \frac{M}{L} \Rightarrow I = \frac{1}{3} \frac{M}{L} L^3 \Rightarrow I = \frac{1}{3} M L^2$

Rotazione lungo il centro di massa

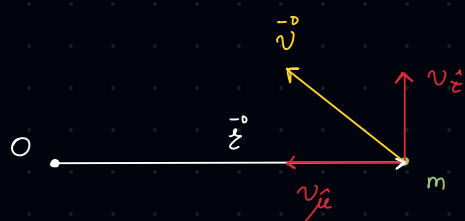


$$I = \lambda \int_{-\frac{L}{2}}^{\frac{L}{2}} y^2 dy = \frac{M}{L} \left[\frac{y^3}{3} \right]_{-\frac{L}{2}}^{\frac{L}{2}} = \frac{1}{3} \frac{M}{L} \left[\frac{L^3}{8} + \frac{L^3}{8} \right]$$

$$= \frac{1}{3} \frac{M}{L} \cdot \frac{L^3}{4} = \frac{1}{12} M L^2$$

Conclusioni: È più semplice mettere in rotazione una sbarra attorno al suo centro di massa che rispetto ad un suo estremo

Bonus: Possiamo trovare $I = m R^2$ con:



$$\vec{L}^0 = \vec{r}^0 \wedge \vec{p}^0 = \vec{r}^0 \wedge m \cdot \vec{v}^0 = \vec{r}^0 \wedge m (\vec{v}_\mu^0 + \vec{v}_\tau^0)$$

$$= \vec{r}^0 \wedge m \vec{v}_\mu^0 + \vec{r}^0 \wedge m \vec{v}_\tau^0$$

$\vec{r}^0 \parallel \vec{v}_\mu^0$

$\rightarrow |\vec{L}| = l m v_\tau \cdot \sin \alpha = l m v$ ma $v = \omega l \rightarrow |\vec{L}| = \overset{\text{Inerzia}}{(l^2 m)} \omega$