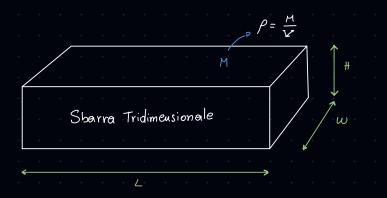
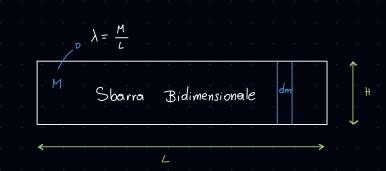
$$I = mR^2$$
 -0 $I_1 = m_1 R_1^2 = 0$ $I_{TOT} = \mathcal{E}_1 m_1 R_1^2$





$$= 0 \quad \text{I}_{\tau \circ \tau} = \sum_{i} m_{i} R_{i}^{2} \quad - 0 \quad \text{I}_{\tau \circ \tau} = \int R^{2} dm$$

• Moltiplico e divido per
$$dV: I = \int R^2 \frac{dm}{dv} dV$$

$$I = \int_{V} R^{2} \frac{dm}{dv} dV$$

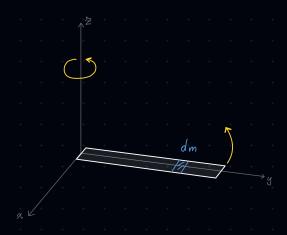
pongo
$$f = \frac{M}{V}$$

$$- > I = \int_{V} R^{2} \int_{V} dV$$

$$\beta - 0 \lambda = \frac{M}{L}$$
 $R - 0 x$

$$- \Rightarrow I = \int_{A} \lambda x^{2} dx$$

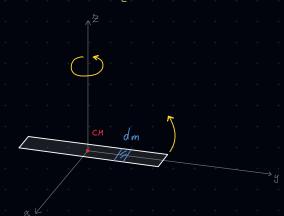
lungo un estremo Rotazione



$$T = \int_{\Lambda}^{B} \chi^{2} dy \longrightarrow T = \int_{\Lambda}^{L} y^{2} dy$$

$$-0 \quad T = \lambda \int_{0}^{y^{2}} dy = \lambda \left[\frac{y^{3}}{3} \right]_{0}^{L} = \lambda \left[\frac{L^{3}}{3} \right] = \frac{1}{3} \lambda L^{3}$$

$$ma \quad \lambda = \frac{M}{L} = 0 \quad T = \frac{1}{3} \frac{M}{L} L^{3^{2}} - \rho \left(T = \frac{1}{3} M L^{2} \right)$$



$$T = \lambda \int y^2 dy = \frac{M}{L} \left[\frac{y^3}{3} \right]^{\frac{L}{2}} = \frac{1}{3} \frac{M}{L} \left[\frac{L^3}{8} + \frac{L^3}{8} \right]$$

$$= \frac{1}{3} \frac{M}{L} \cdot \frac{L^3}{4} = \frac{1}{12} \frac{ML^2}{12}$$

Conclusioni: E' più semplice mettere in rotazione una sbarra attorno al suo ceutro di massa che rispetto ad un suo estremo