

# Simulation Description

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## 1 Simulations

We used a large N-body simulation dubbed **Multidark** to extract statistics of halo parameters on cosmological scales. The data we use for this paper is publicly available through a database interface first presented by Riebe et al. [2011]. Here we summarize the main characteristics of the **Multidark** volume. More details can be found in Prada et al. [2012].

The simulation was run using an Adaptive-Mesh Refinement (AMR) code called ART. Details about the technical aspects and comparisons with other N-body codes are given in Klypin et al. [2009]. The simulation follows the non-linear evolution of a dark matter density field sampled with  $2048^3$  particles in a volume of  $1000h^{-1}\text{Mpc}^3$ . The physical resolution of the simulation is almost constant in time  $\sim 7h^{-1}\text{kpc}$  between redshifts  $z = 0 - 8$ . The cosmological parameters in the simulation are  $\Omega_m = 0.27$ ,  $\Omega_\Lambda = 0.73$ ,  $n_s = 0.95$ ,  $h = 0.70$  and  $\sigma_8 = 0.82$  for the matter density, dark energy density, slope of the matter fluctuations, the Hubble constant at  $z = 0$  in units of  $100\text{km s}^{-1} \text{Mpc}^{-1}$  and the normalization of the power spectrum, respectively. These cosmological parameters are consistent with the results from WMAP5 and WMAP7 [Komatsu et al., 2009, Jarosik et al., 2011]. With these characteristics the mass per simulation particle is  $m_p = 8.63 \times 10^9 h^{-1} \text{M}_\odot$ , which means that group-like halos of masses  $\sim 10^{13} h^{-1} \text{M}_\odot$  are sampled with at least 1100 particles.

Dark matter halos are identified using a Bound-Density-Maxima algorithm (BDM). The code starts by finding the density maxima at the particles' positions in the simulation volume. For each maxima it finds the radius  $R_{200}$  of a sphere containing a mass overdensity given by

$$M_{200} = \frac{4\pi}{3} \Delta \rho_{\text{cr}}(z) R_{200}^3 \quad (1)$$

where  $\rho_{\text{cr}}$  is the critical density of the Universe and  $\Delta = 200$  is the desired overdensity. This procedure allows for the detection of both halos and sub-halos. In our analysis we kept only the halos.

## 1.1 Concentration Estimates

The estimation for the concentration values is done using an analytical property of the NFW profile (see Sect. ??) that relates the circular velocity at the virial radius:

$$V_{200} = \left( \frac{GM_{200}}{R_{200}} \right)^{1/2}, \quad (2)$$

with the maximum circular velocity

$$V_{\max}^2 = \max \left[ \frac{GM(< r)}{r} \right]. \quad (3)$$

The  $V_{\max}/V_{200}$  velocity ratio is used to determine the halo concentration,  $c$  (the ratio between  $R_{200}$  and the scale radii of the NFW profile), using the following relation [Klypin et al., 2011]:

$$\frac{V_{\max}}{V_{200}} = \left( \frac{0.216c}{c} \right)^{1/2} \quad (4)$$

where  $f(c)$  is

$$f(x) = \ln(1 + c) - \frac{c}{(1 + c)}. \quad (5)$$

For each BDM overdensity the  $V_{\max}/V_{200}$  ratio in order to find the concentration  $c$  by solving numerically the previous two equations. This method provides a robust estimate of the concentration compared to a radial fitting to the NFW profile, which is strongly dependent on the radial range used for the fit [Klypin et al., 2011, Meneghetti and Rasia, 2013]. Comparison of these two methods to using halos where the NFW functional fit yield a small systematic offset of (5 – 15)%, with the concentration derived by the velocity ratio method being higher [Prada et al., 2012].

## References

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