

1. <sup>1</sup>Show that state  $|\psi\rangle_{RA}$  as defined in (5.3), is a purification of the density operator  $\rho_A$ , with a spectral decomposition as given in (5.2)

*proof.* We proceed to compute  $Tr_R\{|\psi\rangle\langle\psi|_{RA}\}$  as follows:

$$Tr_R\{|\psi\rangle\langle\psi|_{RA}\} = Tr_R\left\{\left(\sum_x \sqrt{p_X(x)} |x\rangle_R |x\rangle_A\right) \left(\sum_x \sqrt{p_X(x)} |x\rangle_R |x\rangle_A\right)^\dagger\right\} \quad (1)$$

Evaluating the 'dagged' term,

$$Tr_R\{|\psi\rangle\langle\psi|_{RA}\} = Tr_R\left\{\left(\sum_x p_X(x) |xx\rangle_{RA} \langle xx|_{RA}\right)\right\} \quad (2)$$

Now, before we take the partial trace of the bracketed expression, refer to the alternate notion of the partial trace operator as described in exercise (4.3.9) in Wilde. Re-expressing the tensor product<sup>a</sup>,

$$Tr_R\left\{\sum_x p_X(x) |x\rangle_R \langle x|_R \otimes |x\rangle_A \langle x|_A\right\} = \sum_x p_X(x) \langle x|x\rangle_R |x\rangle \langle x|_A \quad (3)$$

which leaves us with the desired expression:

$$\sum_x p_X(x) \langle x|x\rangle_R |x\rangle \langle x|_A = \rho_A \quad (4)$$

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<sup>a</sup>this overall procedure described in 4.3.9 saves a lot of time and effort – add it to the toolbox.

2. **(Canonical Purification)** Let  $\rho_A$  be a density operator and let  $\sqrt{\rho_A}$  be its unique positive semidefinite square root. We define the canonical purification of  $\rho_A$  as follows:

$$(I_R \otimes \sqrt{\rho_A}) |\Gamma\rangle_{RA} \quad (5)$$

where  $|\Gamma\rangle_{RA}$  is the maximally unnormalized vector from (3.233).<sup>2</sup> Show that (5.4) is a purification of  $\rho_A$

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<sup>1</sup><https://arxiv.org/pdf/1106.1445> (the copy of Wilde we are using)

<sup>2</sup>Recall exercise (3.7.12) in Wilde, *the transpose trick or ricochet* – didn't use it here but helpful to know

*proof.* Spectrally decomposing  $\rho_A = \sum_x p_X(x) |x\rangle \langle x|_A$ , where  $\{|x\rangle_A\}$  forms an O.N eigenbasis over  $\mathcal{H}_A$ . Then,  $\sqrt{\rho_A} = \sum_x \sqrt{p_X(x)} |x\rangle \langle x|_A$ . Take the orthonormal basis over  $\mathcal{H}_R$ , denoting it by  $\{|x\rangle_R\}$ . Then the unnormalized maximally entangled vector <sup>a</sup>  $|\Gamma\rangle_{RA} = \sum_x |x\rangle_R |x\rangle_A$ . <sup>b</sup> Now, all we have left to do is compute  $(I_R \otimes \sqrt{\rho_A})(|\Gamma\rangle_{RA})(I_R \otimes \sqrt{\rho_A})^\dagger |\Gamma\rangle_{RA}$  and trace over  $\mathcal{H}_R$ .

$$(I_R \otimes \sqrt{\rho_A}) |\Gamma\rangle_{RA} = (I_R \otimes \sum_x \sqrt{p_X(x)} |x\rangle_A \langle x|_A) \sum_{x'} |x'\rangle_R |x'\rangle_A \quad (6)$$

$$(I_R \otimes \sum_x \sqrt{p_X(x)} |x\rangle_A \langle x|_A) \sum_{x'} |x'\rangle_R |x'\rangle_A = \sum_x |x\rangle_R \otimes \sqrt{p_X(x)} |x\rangle_A \quad (7)$$

since the only non-vanishing terms are when  $x' = x$ .  $(I_R \otimes \sqrt{\rho_A})(|\Gamma\rangle_{RA})(I_R \otimes \sqrt{\rho_A})^\dagger |\Gamma\rangle_{RA}$  and trace over  $\mathcal{H}_R$  is then

$$\sum_x |x\rangle \langle x|_R \otimes p_X(x) |x\rangle \langle x|_A \quad (8)$$

Tracing out  $R$  gives us the following:

$$\sum_x \langle x|x\rangle_R p_X(x) |x\rangle \langle x|_A = \sum_x p_X(x) |x\rangle \langle x|_A = \rho_A, \quad (9)$$

as desired.

<sup>a</sup><https://physics.stackexchange.com/questions/267293/does-the-entanglement-depend-on-the-basis>

<sup>b</sup><https://quantumcomputing.stackexchange.com/questions/9242/maximally-entangled-state-definition-and-orthonormal-basis-of-maximally-entangled>

[Link to the template used](#)