

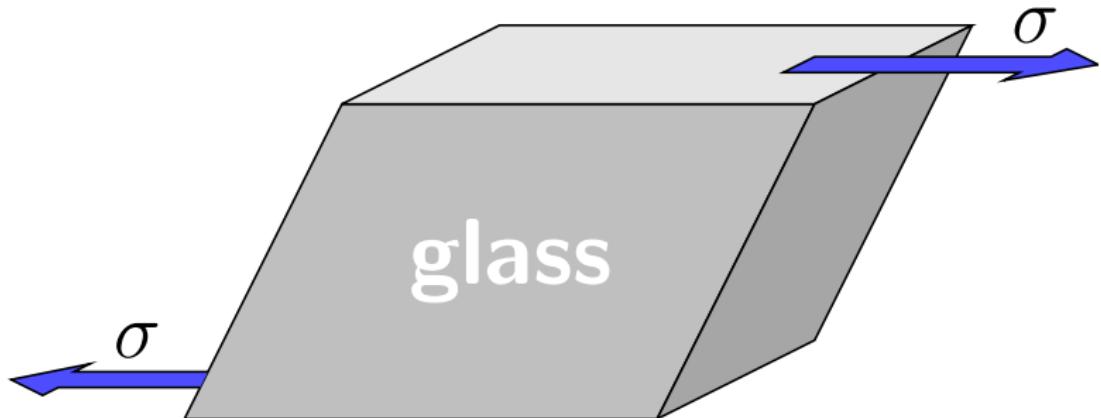
Nonlinear plastic modes – micromechanics and statistics

Edan Lerner

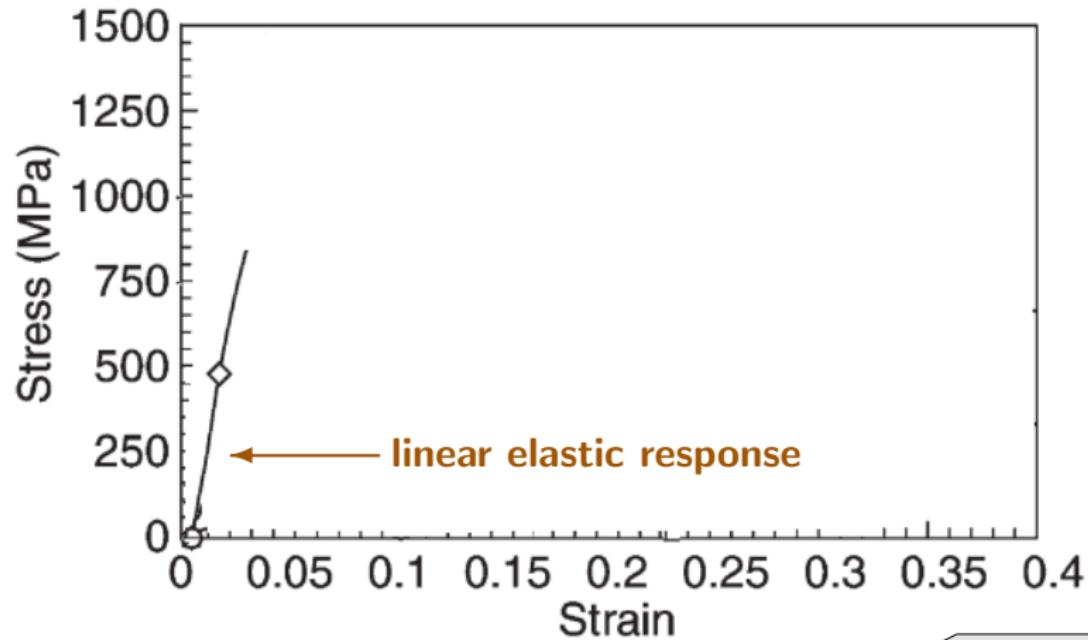
Institute for Theoretical Physics
University of Amsterdam

yielding of amorphous solids
ENS Paris
Oct 2017

what happens when a glass is deformed?

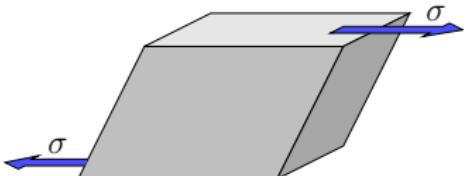


elasto-plasticity – macroscopic response & yielding

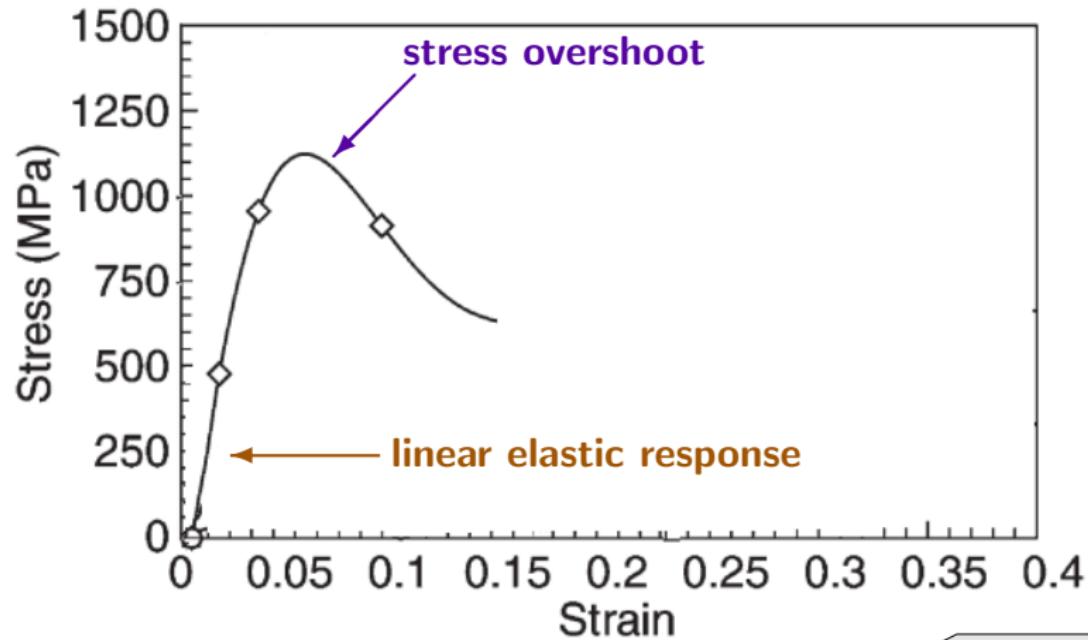


'Vitreloy 1' (metallic glass)

J. Lu, G. Ravichandran, W. Johnson, Acta Materialia 51 (2003)

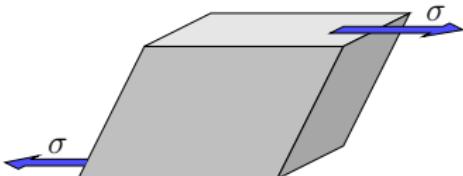


elasto-plasticity – macroscopic response & yielding

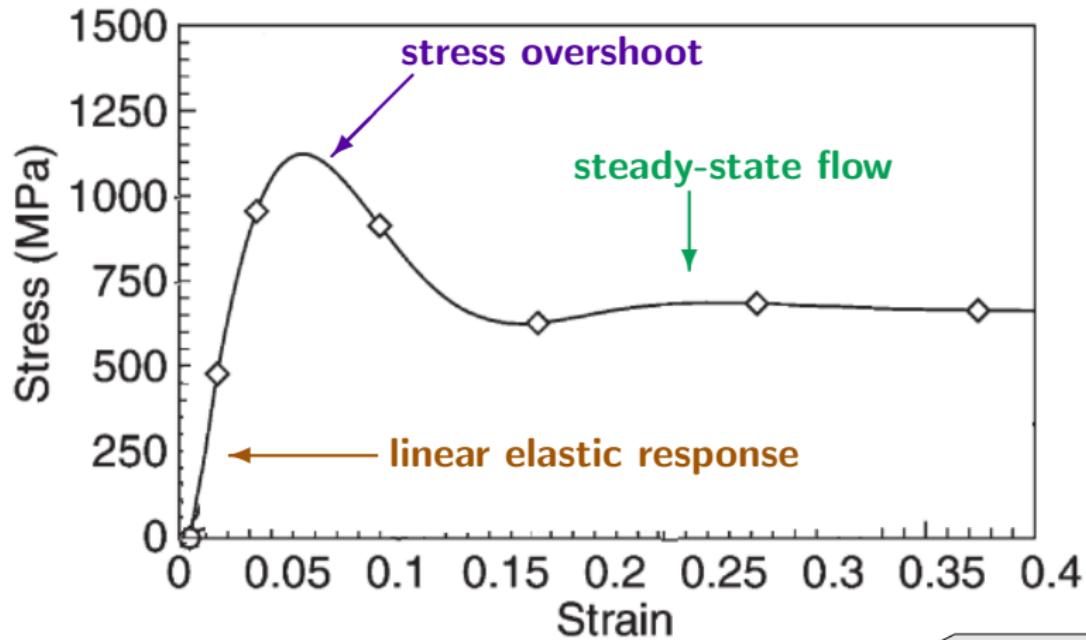


'Vitreloy 1' (metallic glass)

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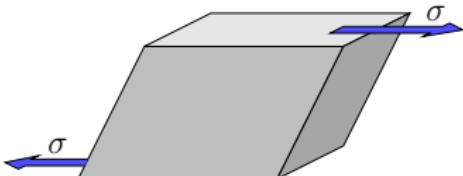


elasto-plasticity – macroscopic response & yielding



'Vitreloy 1' (metallic glass)

J. Lu, G. Ravichandran, W. Johnson, Acta Materialia 51 (2003)



what do we want to find?

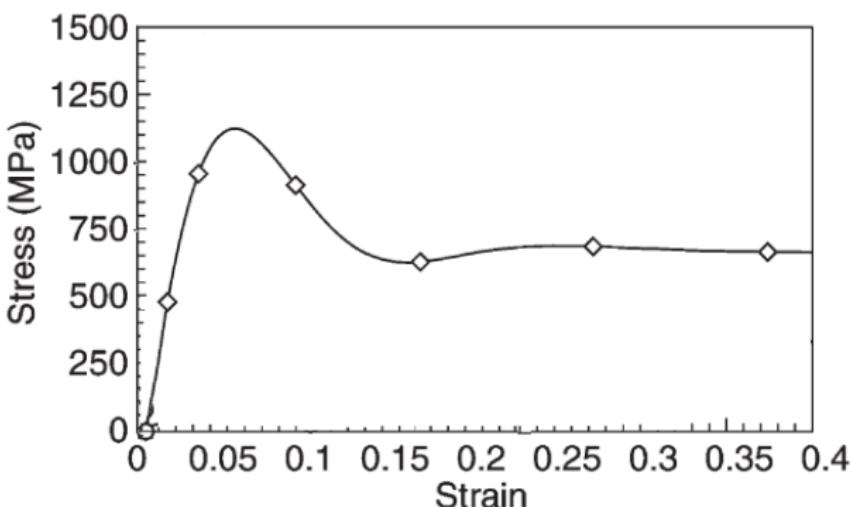
goal: theory for

$\sigma \equiv$ stress

$\dot{\gamma} \equiv$ deformation rate

$T \equiv$ Temperature

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, T, \dots)$$

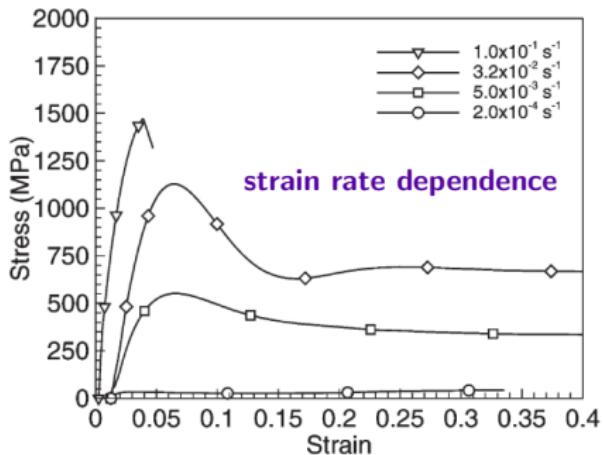


dependence on external parameters

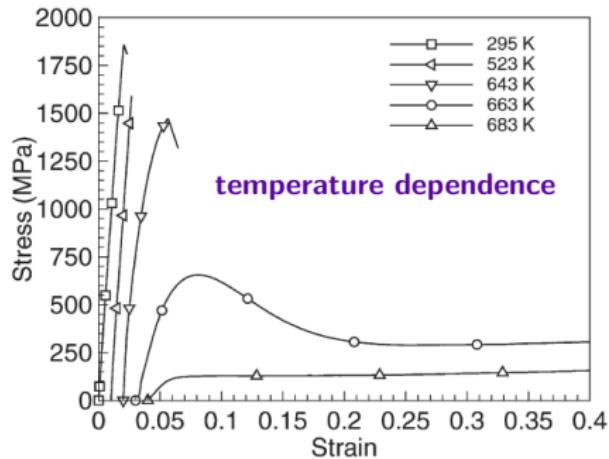
experiments, metallic glass

J. Lu, G. Ravichandran, W. Johnson, Acta Materialia 51 (2003)

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, T, \dots)$$



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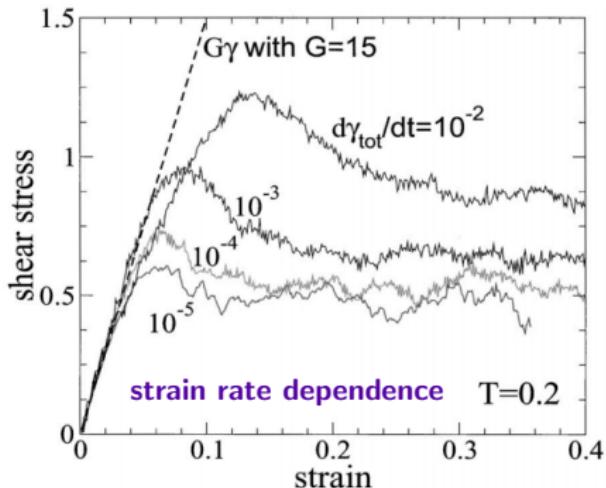


dependence on external parameters

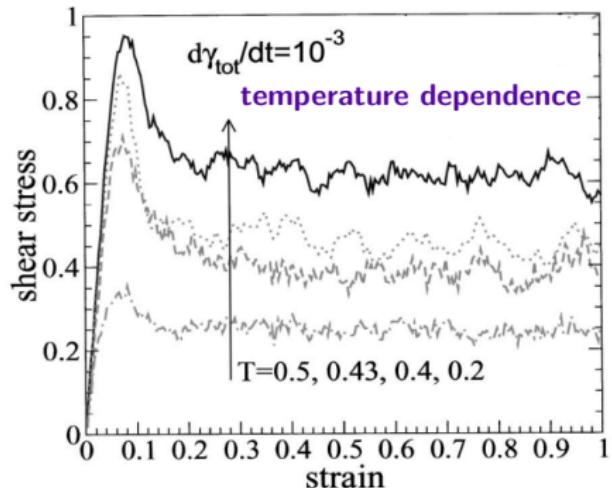
simulations of model glasses

F. Varnik, L. Bocquet, and J.-L. Barrat, J. Chem. Phys. 120, 2788 (2004)

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, T, \dots)$$



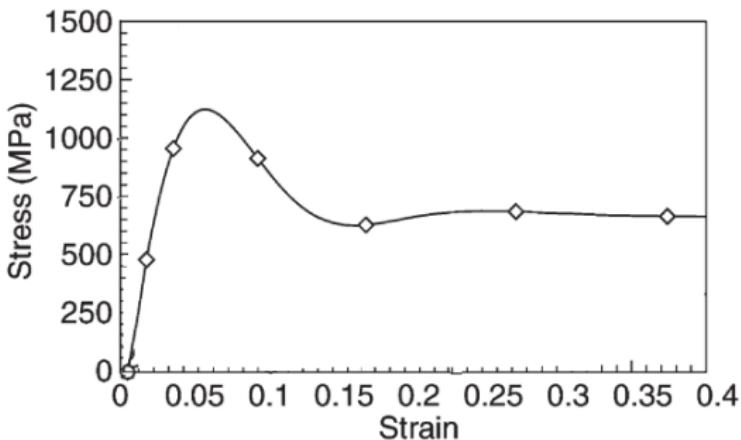
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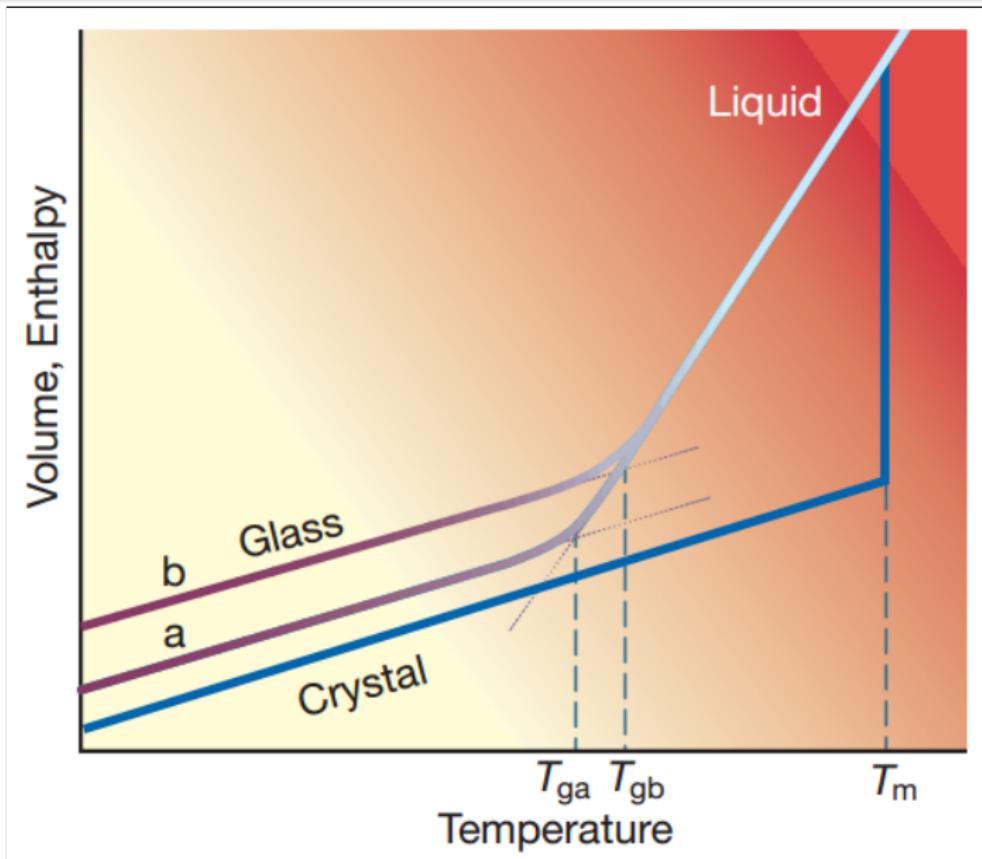
this talk: **structural** order parameters

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, T, ???)$$

what should go here?



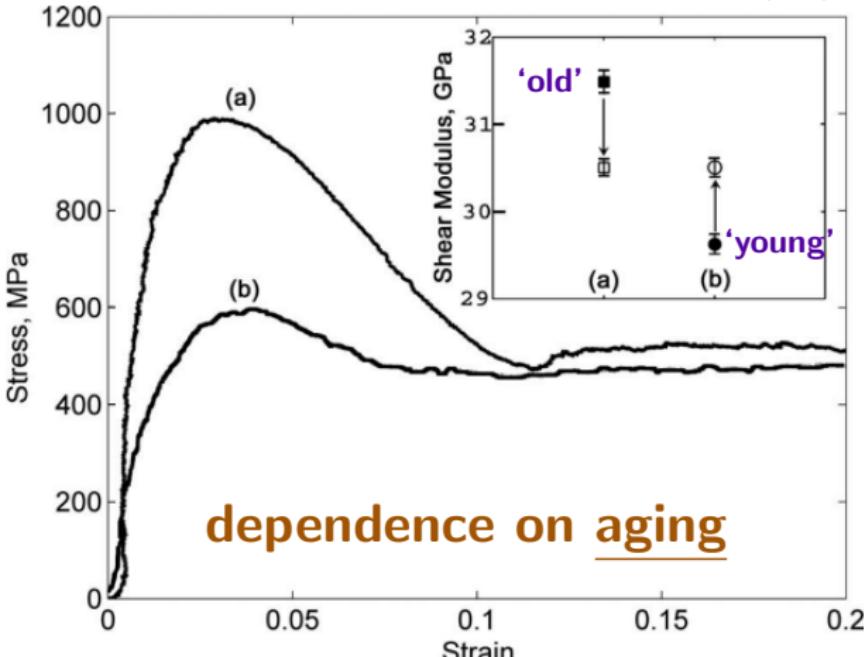
structural order parameters? aging effects



structural order parameters? aging effects

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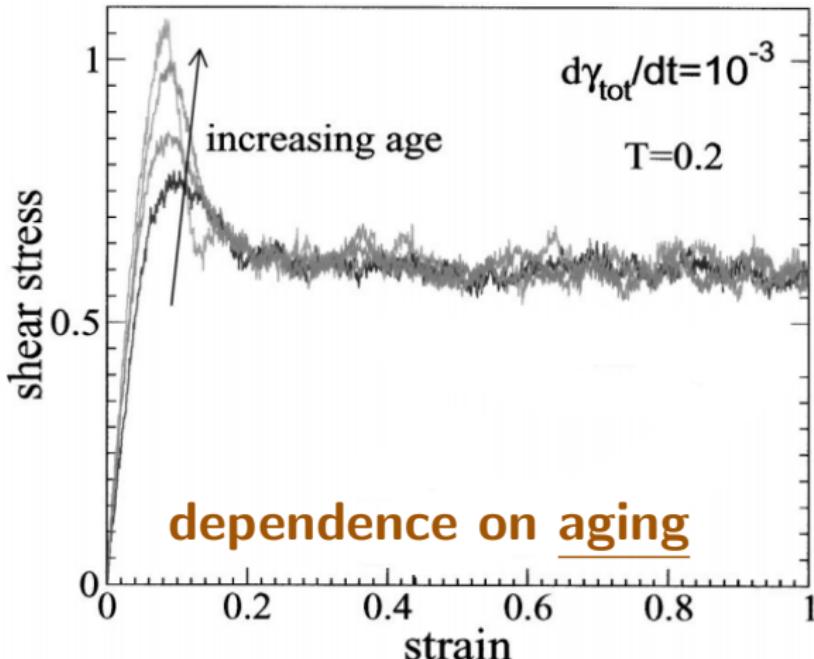
J. Lu, G. Ravichandran, W. Johnson, Acta Materialia 51 (2003)



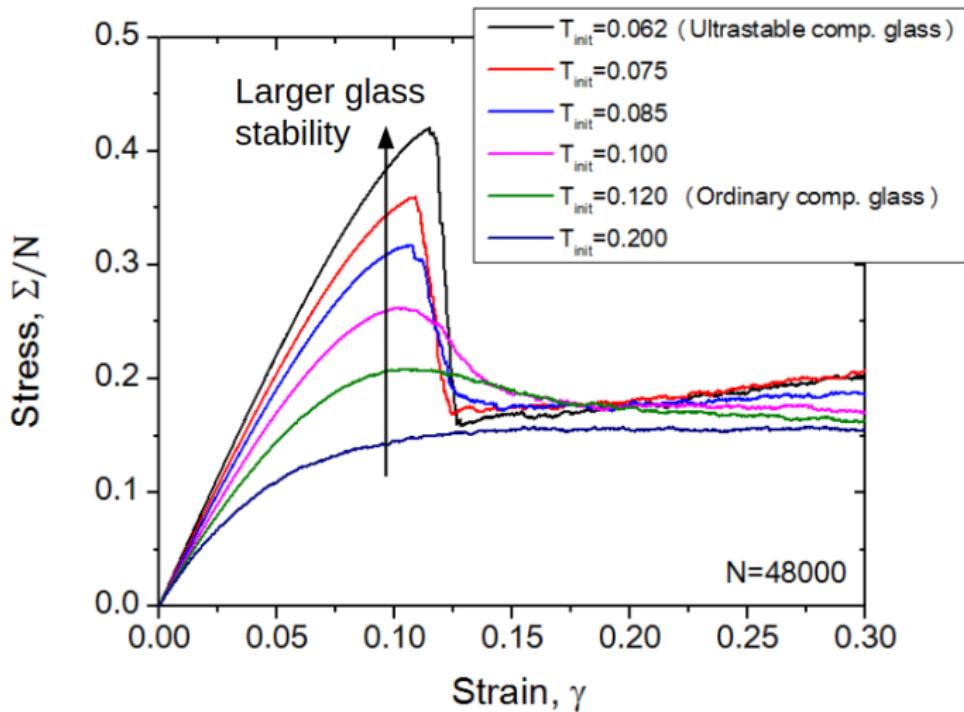
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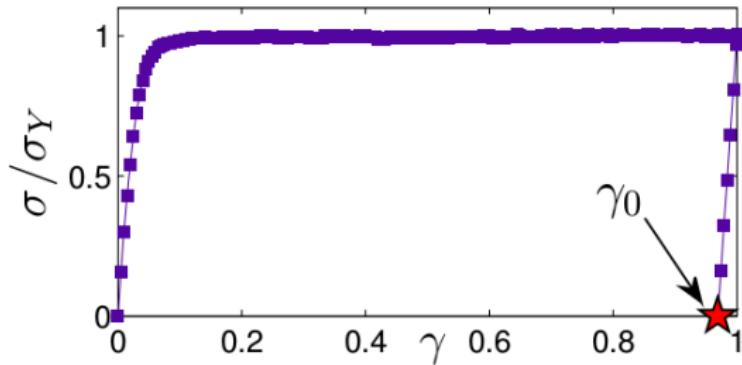
structural order parameters? aging effects



many thanks to Misaki Ozawa!

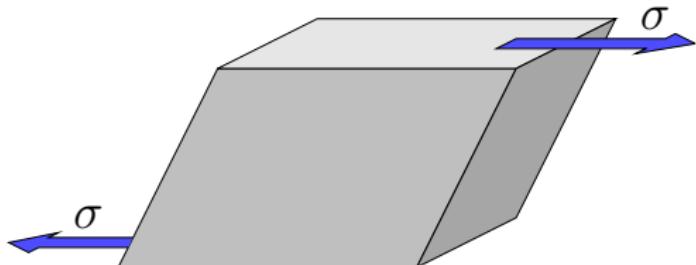
structural order parameters? anisotropy

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, T, ???)$$



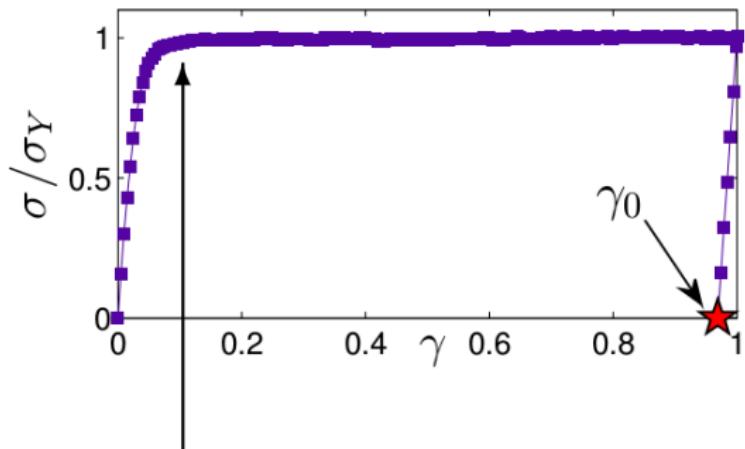
simulations @

$$T \rightarrow 0$$
$$\dot{\gamma} \rightarrow 0$$



structural order parameters? anisotropy

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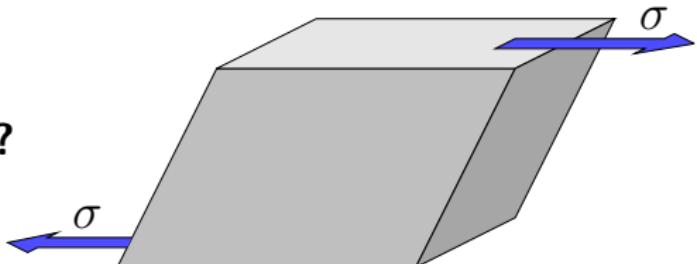


simulations @

$$T \rightarrow 0$$
$$\dot{\gamma} \rightarrow 0$$

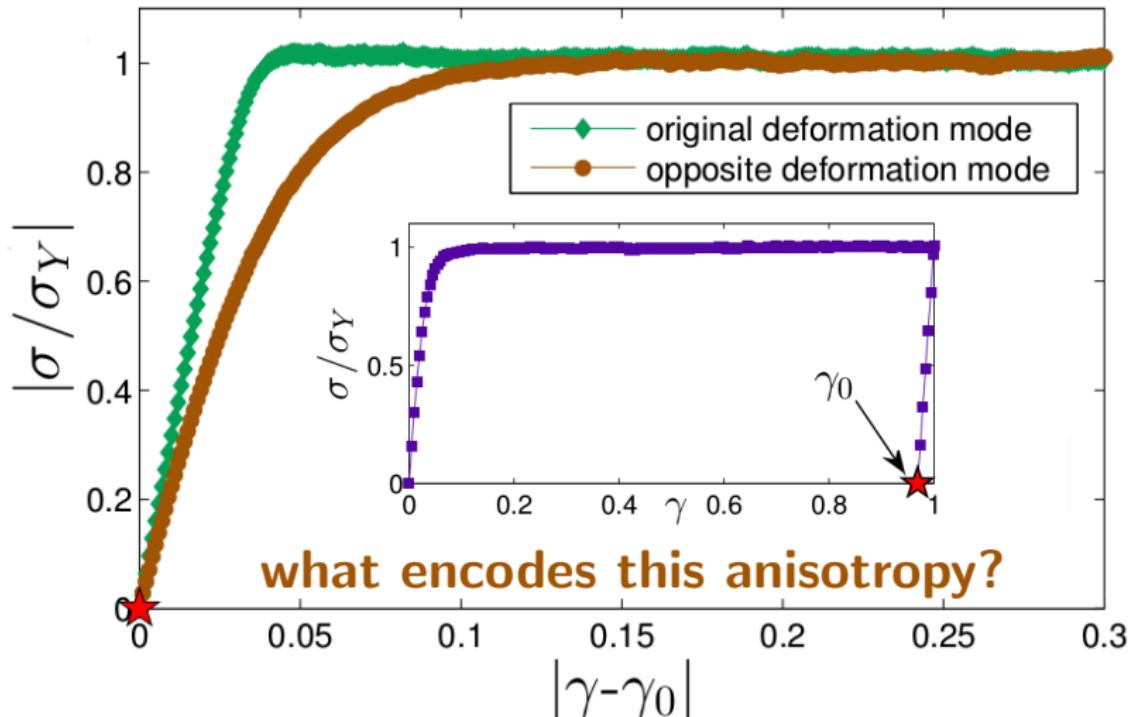
no stress overshoot

\Rightarrow does structure evolve?



structural order parameters? anisotropy

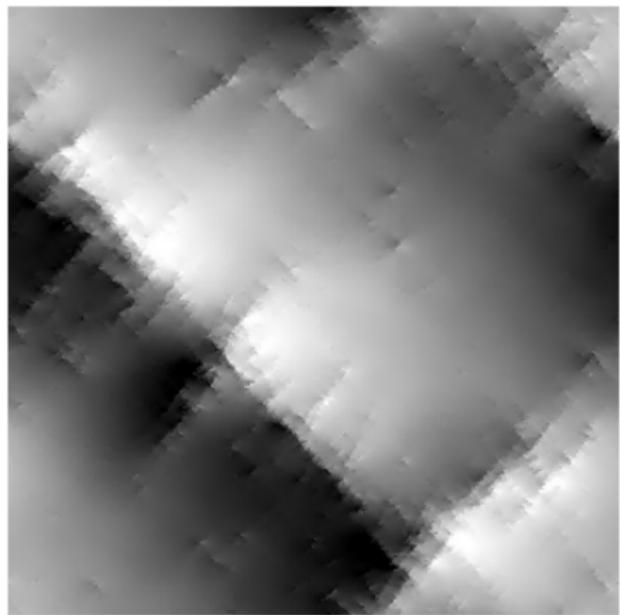
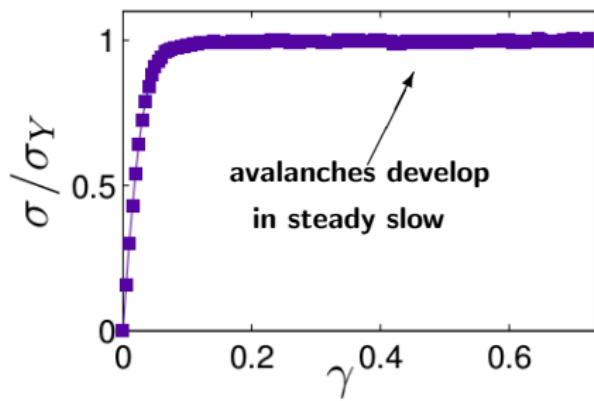
$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, \ddot{\gamma}, ???)$$



structural order parameters? development of avalanches

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\chi}, \ddot{\chi}, ???)$$

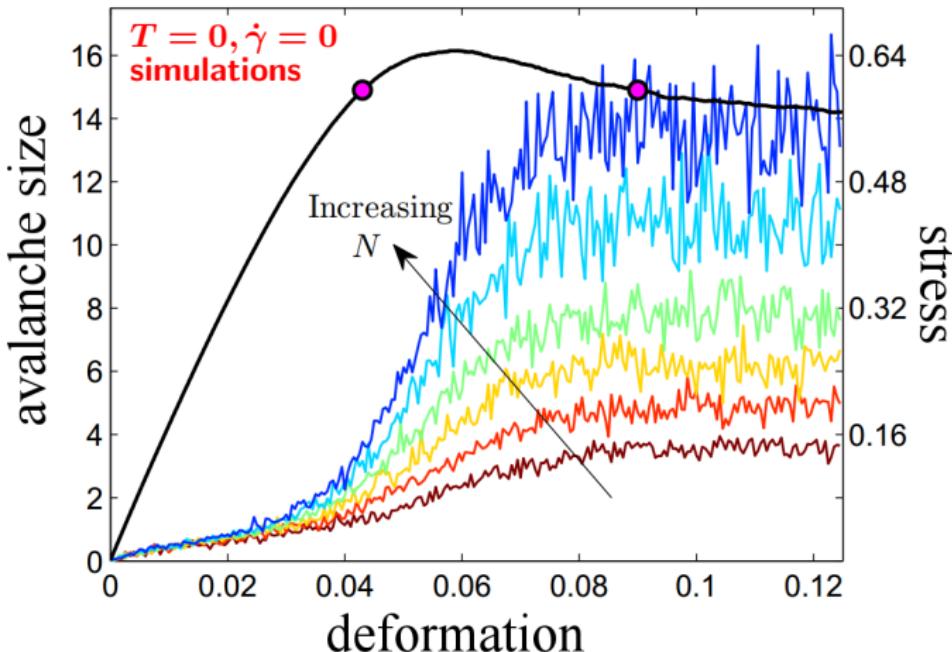
$T = 0, \dot{\gamma} = 0$
simulations



structural order parameters? development of avalanches

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\chi}, \chi, ???)$$

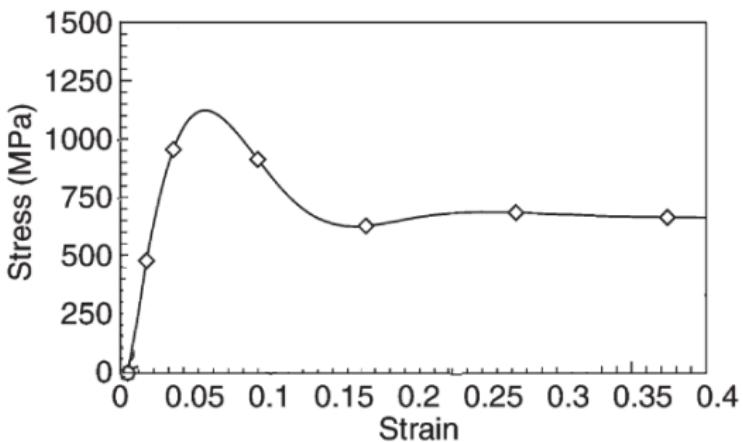
S. Karmakar, E.L., and I. Procaccia, Phys. Rev. E 82, 055103(R) (2010).



structural order parameters

$$\frac{d\sigma}{dt} = f(\sigma, \dot{\gamma}, T, ???)$$

what should go here?

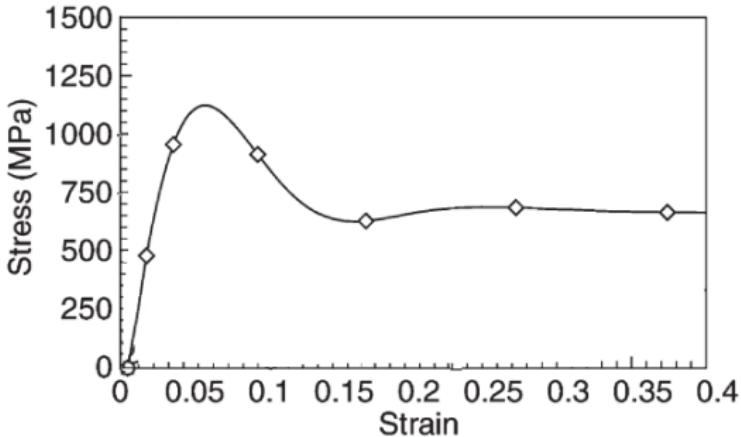


structural order parameters

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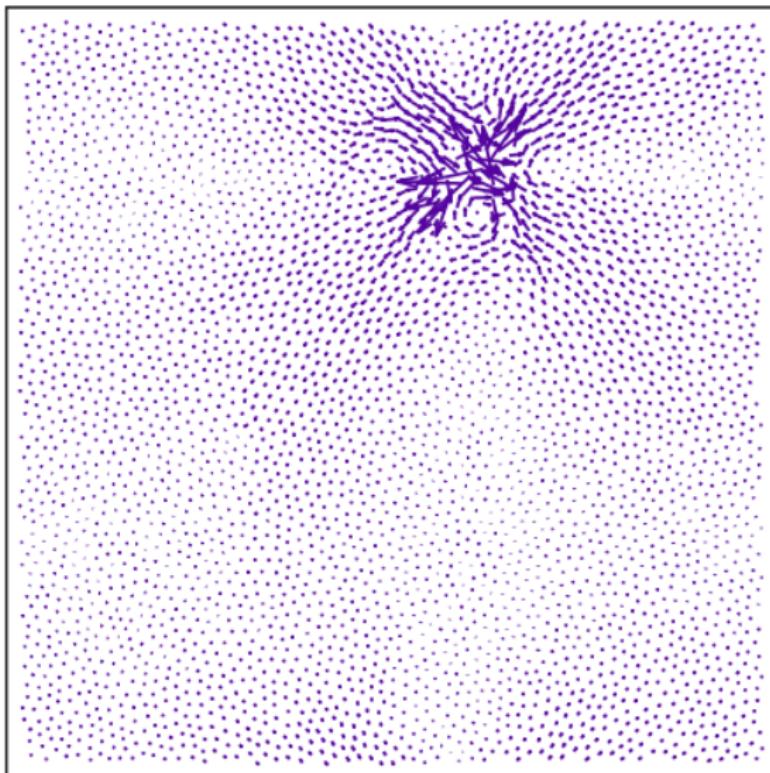
what should go here?

what actually is $p(x)$
from elasto-plastic models?



what is plasticity on the micro-scale?

what is plasticity on the micro-scale?



‘shear-
transformation’

or

‘shear-
transformation-
zone’

how are plastic instabilities triggered?

VOLUME 93, NUMBER 19

PHYSICAL REVIEW LETTERS

week ending
5 NOVEMBER 2004

Universal Breakdown of Elasticity at the Onset of Material Failure

Craig Maloney^{1,2} and Anaël Lemaître^{1,3}

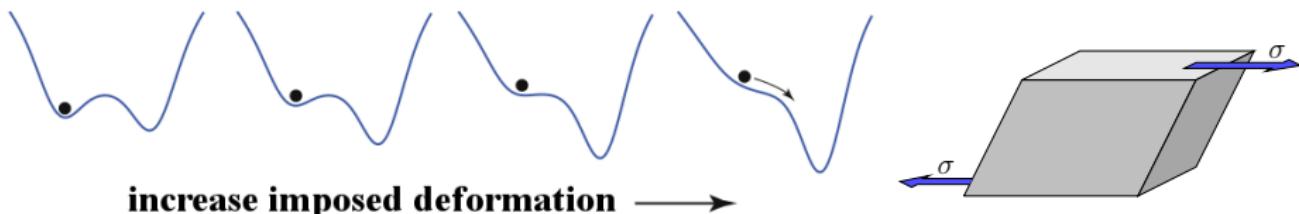
¹*Department of Physics, University of California, Santa Barbara, California 93106, USA*

²*Lawrence Livermore National Lab, CMS-MSTD, Livermore, California 94550, USA*

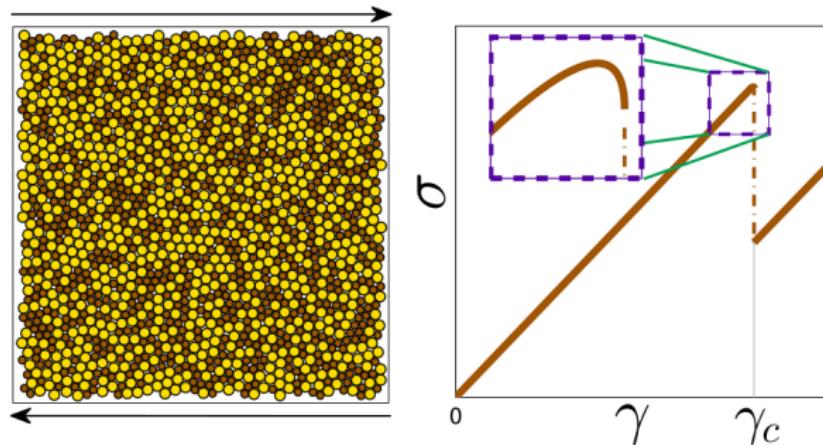
³*LMDH, Université Paris VI, UMR 7603, 4 place Jussieu, 75005 Paris, France*

(Received 6 May 2004; published 2 November 2004)

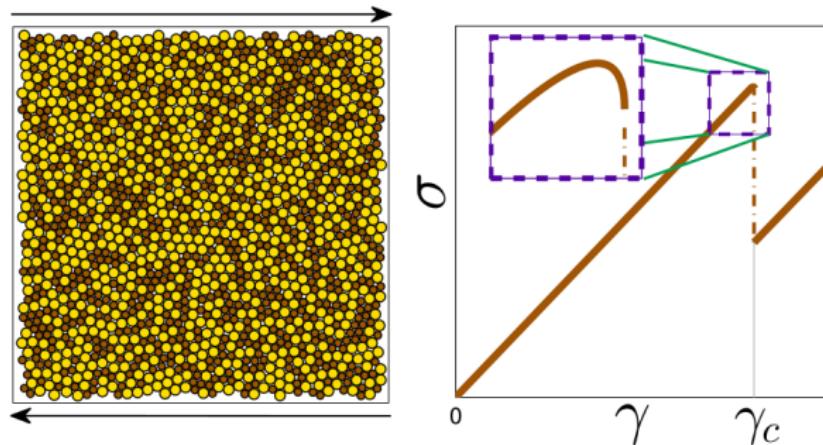
'energy landscape' picture:



micromechanics of plastic instabilities



micromechanics of plastic instabilities

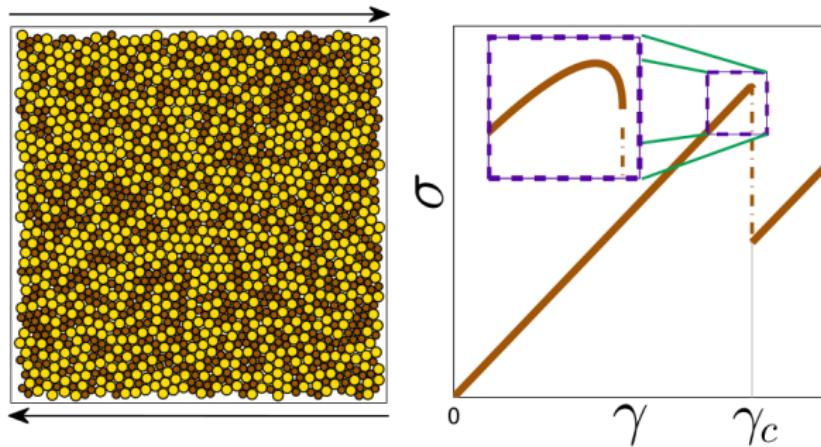


consider linear stability

$$\mathcal{M}_{jk} \equiv \frac{\partial^2 U}{\partial \vec{x}_j \partial \vec{x}_k}$$

dynamical matrix / hessian

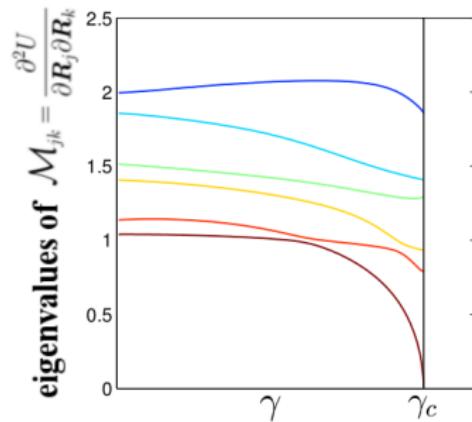
micromechanics of plastic instabilities



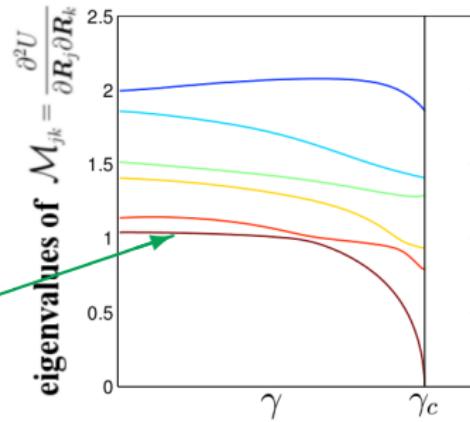
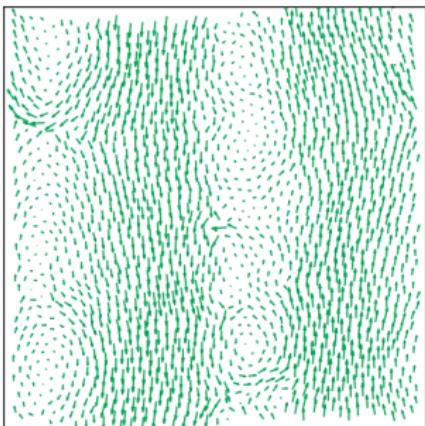
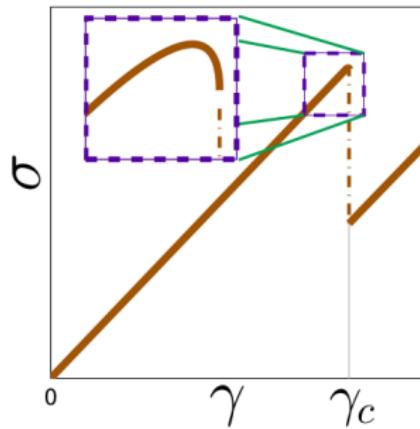
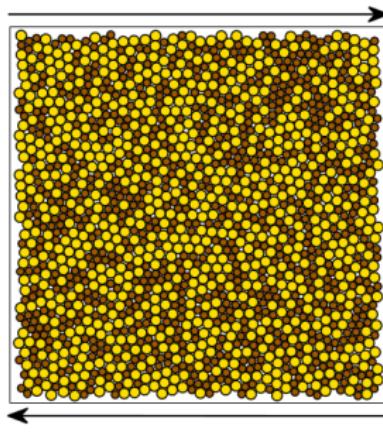
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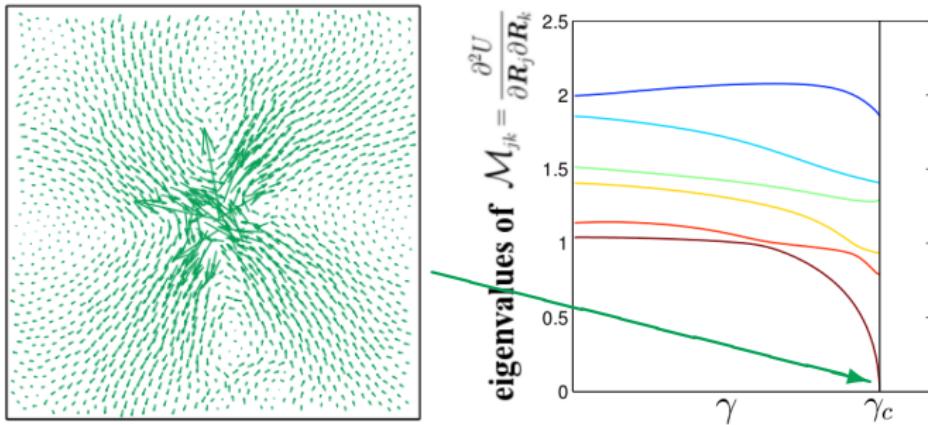
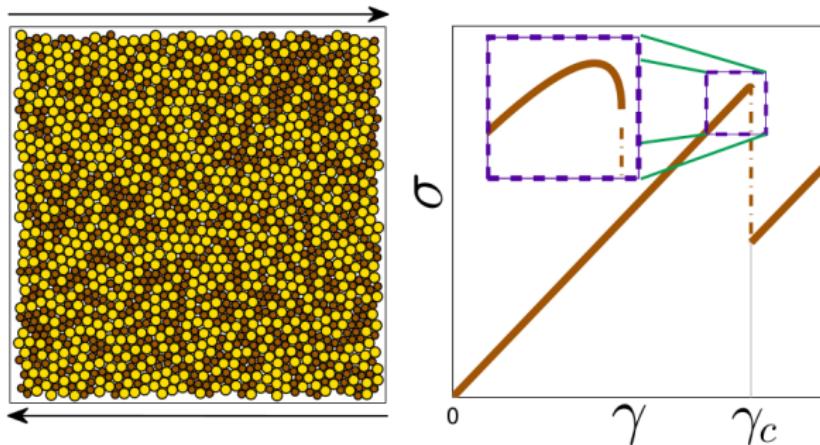
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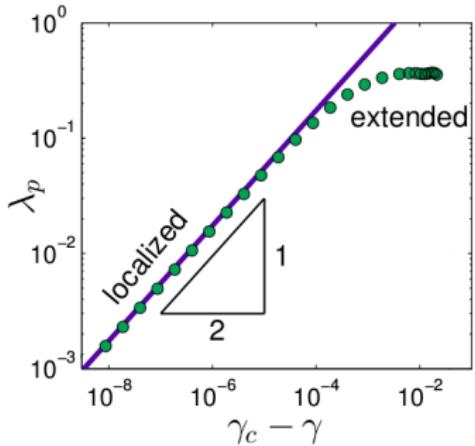
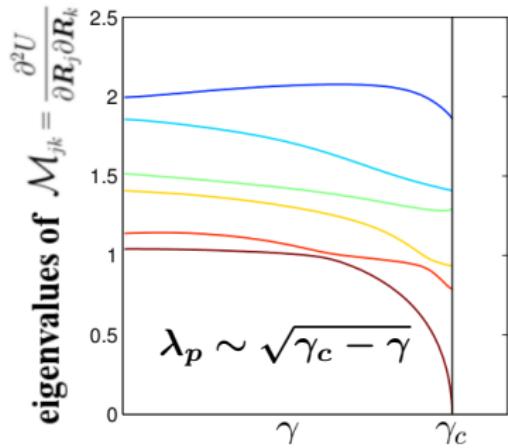
micromechanics of plastic instabilities



micromechanics of plastic instabilities



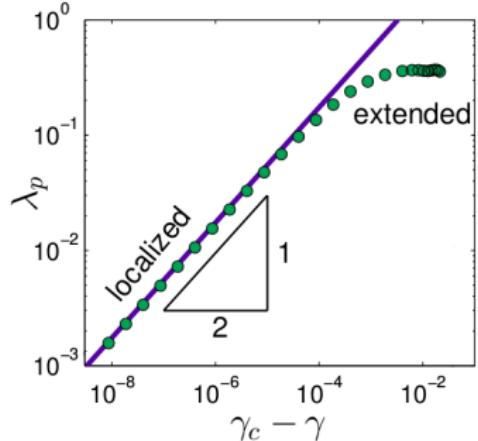
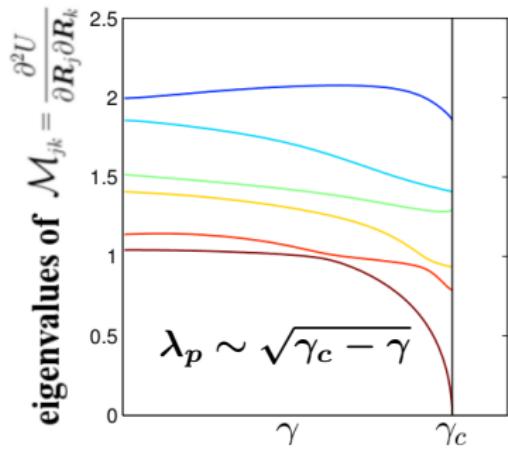
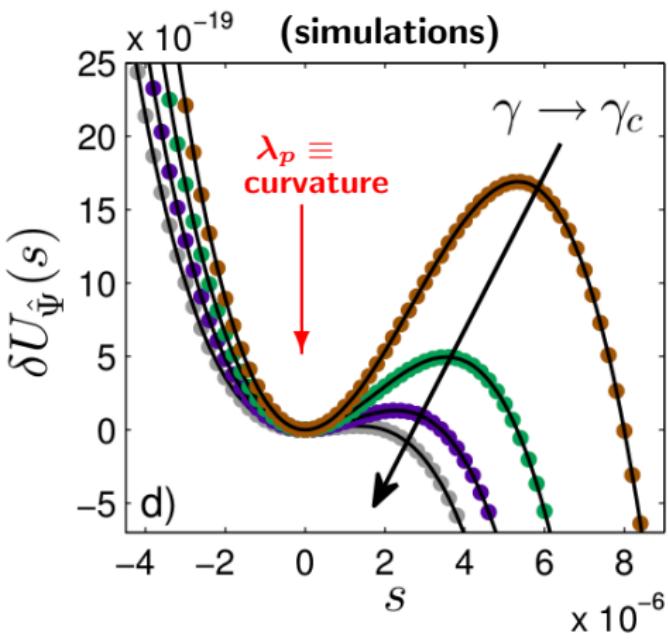
micromechanics of plastic instabilities



micromechanics of plastic instabilities



increase imposed deformation →



predicting plastic instabilities using normal modes

observation: ‘linear’ (normal) modes are indicative of plastic instabilities.

predicting plastic instabilities using normal modes

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proposition: use normal modes to detect ‘soft spots’

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 A LETTERS JOURNAL EXPLORING
THE FRONTIERS OF PHYSICS

EPL, 90 (2010) 16004 April 2010
doi: 10.1209/0295-5075/90/16004 www.epljournal.org

Vibrational modes as a predictor for plasticity in a model glass

A. TANGUY^(a), B. MANTISI and M. TSAMADOS

*Université de Lyon - F-69622, Lyon, France, EU and
CNRS, UMR5586, Laboratoire de Physique de la Matière Condensée et des Nanostructures, Université Lyon 1
F-69622, Villeurbanne Cedex, France, EU*

PRL 107, 108302 (2011) PHYSICAL REVIEW LETTERS week ending
2 SEPTEMBER 2011

Vibrational Modes Identify Soft Spots in a Sheared Disordered Packing

M. L. Manning*

*Princeton Center for Theoretical Science, Princeton, New Jersey 08544, USA
Department of Physics, Syracuse University, Syracuse, New York 13244, USA*

A. J. Liu

*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA
(Received 21 December 2010; published 31 August 2011)*

PHYSICAL REVIEW E 89, 042304 (2014)

Predicting plasticity with soft vibrational modes: From dislocations to glasses

Jörg Rottler,¹ Samuel S. Schoenholz,² and Andrea J. Liu²

¹*Department of Physics and Astronomy, University of British Columbia, 6224 Agricultural Road,
Vancouver, British Columbia, Canada V6T 1Z4*

²*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*

(Received 19 December 2013; revised manuscript received 26 February 2014; published 14 April 2014)

predicting plastic instabilities using normal modes

observation: ‘linear’ (normal) modes are indicative of plastic instabilities.

proposition: use normal modes to detect ‘soft spots’

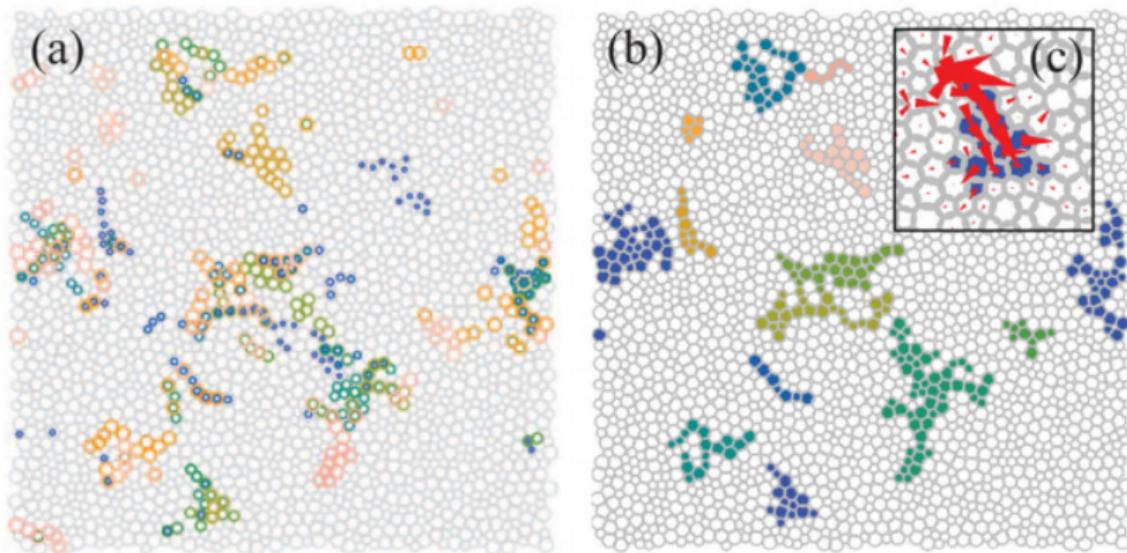
problems:

predicting plastic instabilities using normal modes

observation: ‘linear’ (normal) modes are indicative of plastic instabilities.

proposition: use normal modes to detect ‘soft spots’

problems: 1) no quantitative information



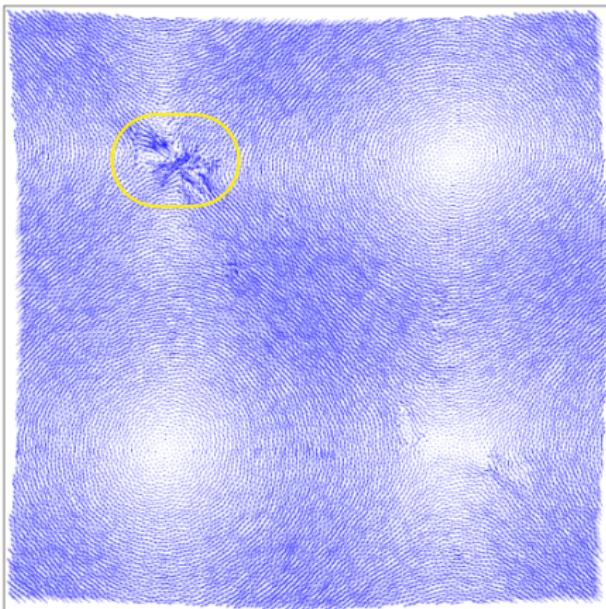
(Manning and Liu, PRL 2011)

predicting plastic instabilities using normal modes

observation: ‘linear’ (normal) modes are indicative of plastic instabilities.

proposition: use normal modes to detect ‘soft spots’

problems: 2) hybridizations with plane waves



predicting plastic instabilities using normal modes

observation: ‘linear’ (normal) modes are indicative of plastic instabilities.

proposition: use normal modes to detect ‘soft spots’

problems: 2) hybridizations with plane waves

compare lowest energy

$$\text{plane wave freq. } \omega^2 \sim 1/L^2$$

$$\text{with } \lambda_p \sim \sqrt{\gamma_c - \gamma}$$

$$\Rightarrow \gamma_c - \gamma \sim L^{-4}$$

dehybridization strain scale

predicting plastic instabilities using normal modes

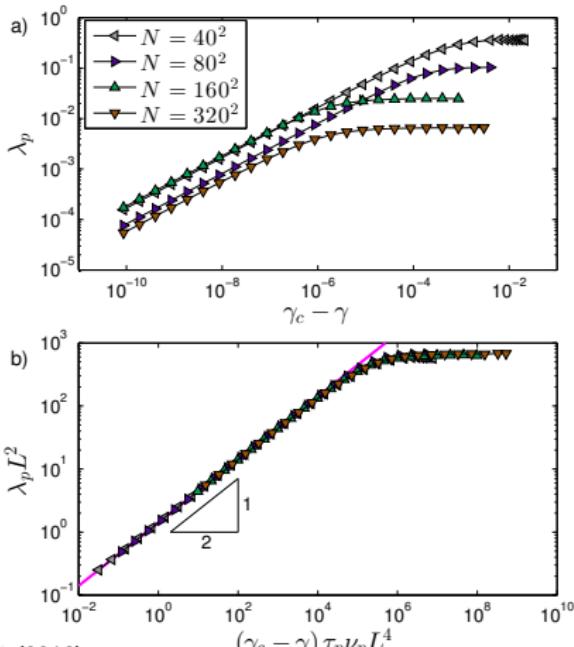
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predicting plastic instabilities using normal modes

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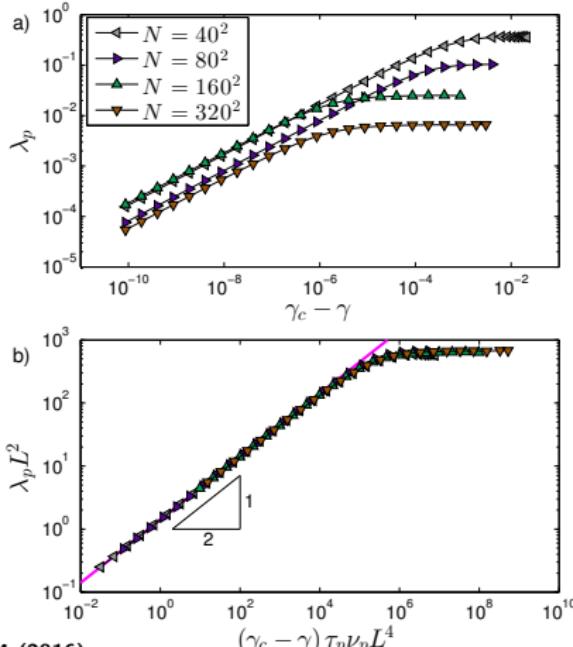
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with $\lambda_p \sim \sqrt{\gamma_c - \gamma}$

$\Rightarrow \gamma_c - \gamma \sim L^{-4}$
dehybridization strain scale

problematic as L increases...



is there a way to define and detect
plastic modes far from instability strains,
deep in the hybridized regime?

yes

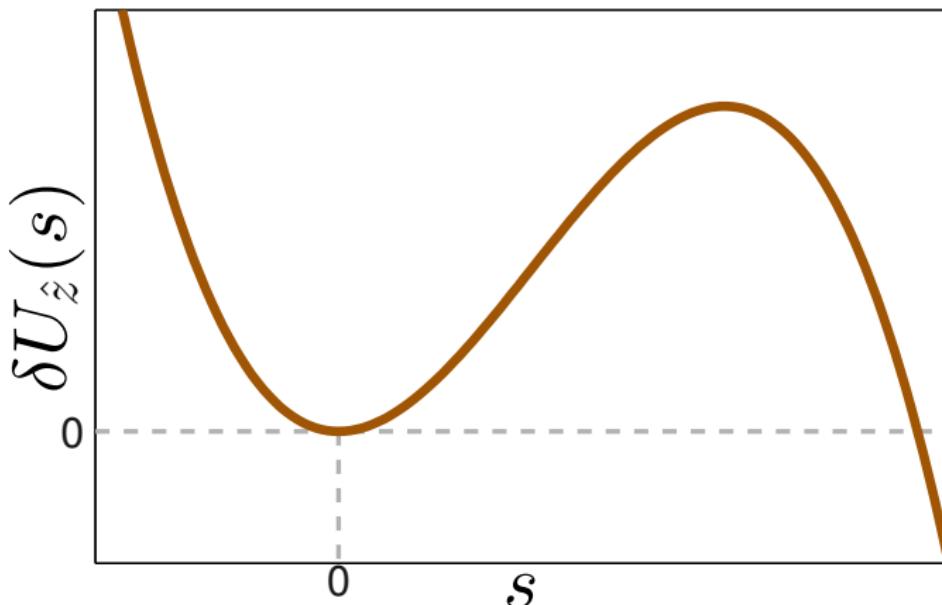
probably
(TBD...)

is there a way to define and detect
plastic modes far from instability strains,
deep in the hybridized regime?

the barrier function – definition

consider the energy variation upon displacing particles' coordinates \vec{x} according to $\delta\vec{x} = s\hat{z}$

$$\delta U_{\hat{z}}(s) \simeq \frac{1}{2}\kappa_{\hat{z}}s^2 + \frac{1}{6}\tau_{\hat{z}}s^3$$

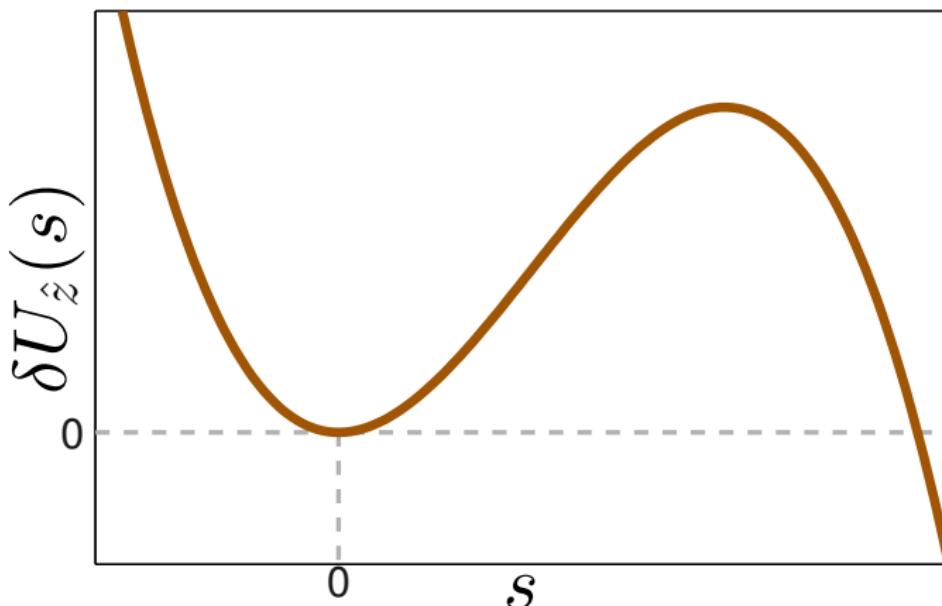


the barrier function – definition

consider the energy variation upon displacing particles' coordinates \vec{x} according to $\delta\vec{x} = s\hat{z}$

\hat{z} not necessarily a normal mode!

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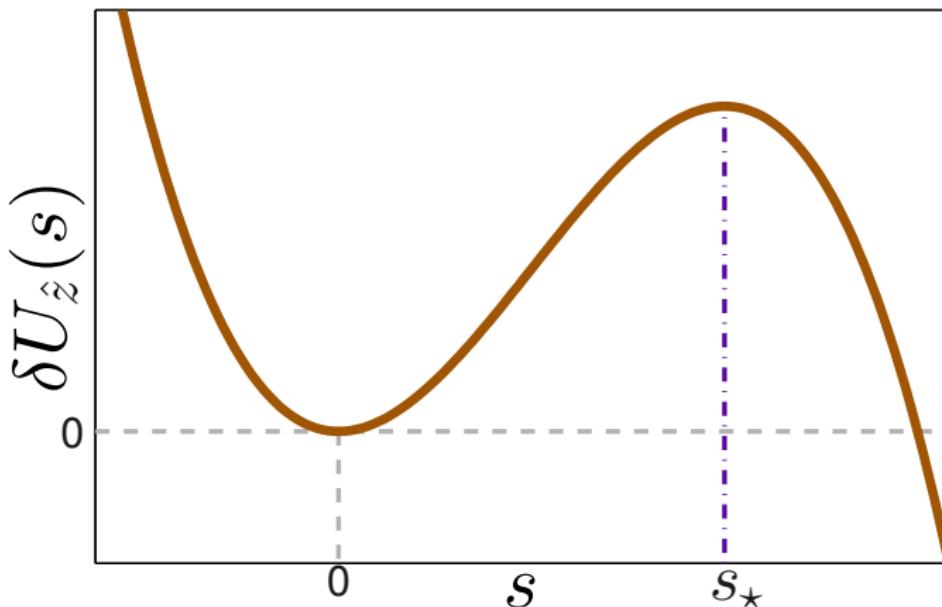


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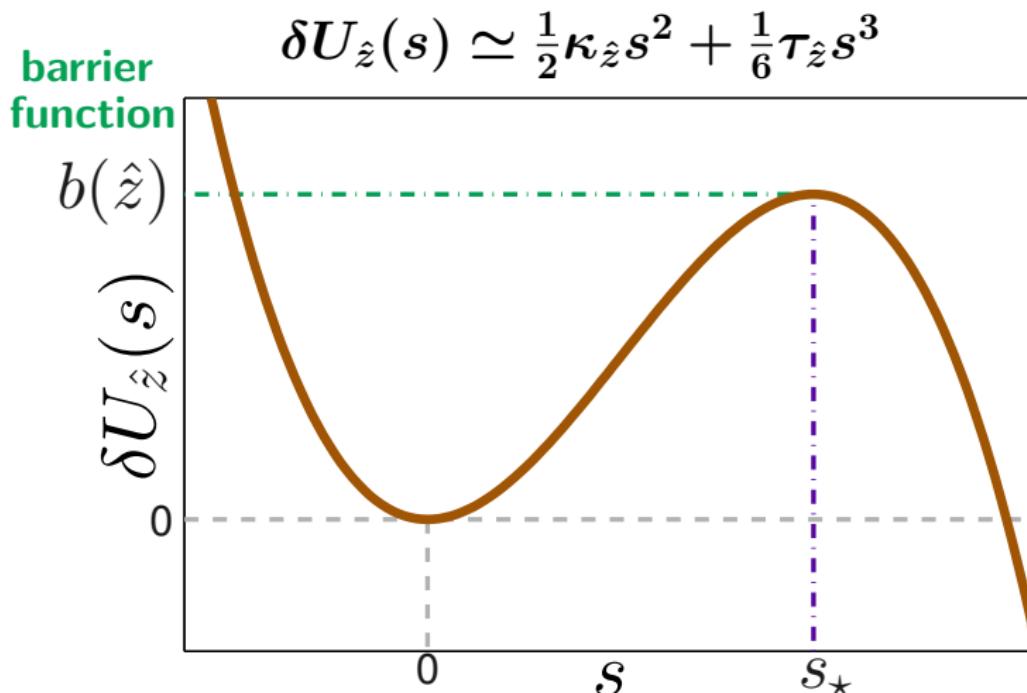
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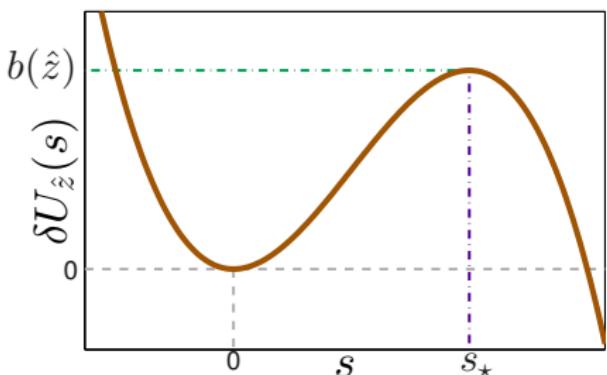
consider the energy variation upon displacing particles' coordinates \vec{x} according to $\delta\vec{x} = s\hat{z}$

\hat{z} not necessarily a normal mode!

barrier
function

$$\delta U_{\hat{z}}(s) \simeq \frac{1}{2}\kappa_{\hat{z}}s^2 + \frac{1}{6}\tau_{\hat{z}}s^3$$

$$b(\hat{z}) \equiv \delta U_{\hat{z}}(s_*) = \frac{2\kappa_{\hat{z}}^3}{3\tau_{\hat{z}}^2} = \frac{2}{3} \frac{(\mathcal{M} : \hat{z}\hat{z})^3}{(U''' : \hat{z}\hat{z}\hat{z})^2}$$



the barrier function – definition

consider the energy variation upon displacing particles' coordinates \vec{x} according to $\delta\vec{x} = s\hat{z}$

\hat{z} not necessarily a normal mode!

barrier function

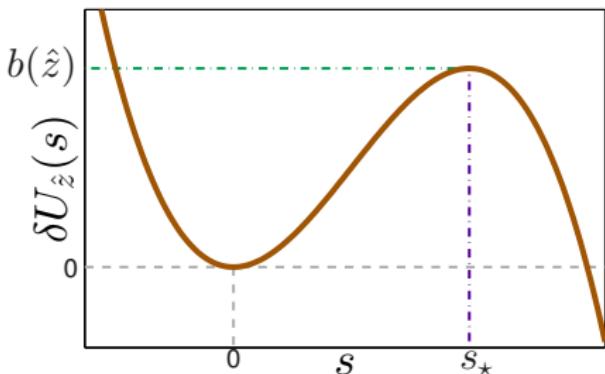
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shorthand notations:

$$\mathcal{M} \equiv \frac{\partial^2 U}{\partial \vec{x} \partial \vec{x}} \quad \text{'dynamical matrix'}$$

$$U''' \equiv \frac{\partial^3 U}{\partial \vec{x} \partial \vec{x} \partial \vec{x}} \quad \text{'cubic tensor'}$$



the barrier function – definition

consider the energy variation upon displacing

particles' coordinates \vec{x} according to $\delta\vec{x} = s\hat{z}$

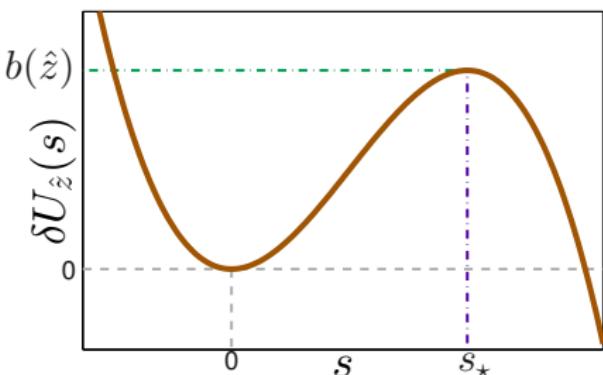
\hat{z} not necessarily
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barrier
function

$$\delta U_{\hat{z}}(s) \simeq \frac{1}{2}\kappa_{\hat{z}}s^2 + \frac{1}{6}\tau_{\hat{z}}s^3$$

$$b(\hat{z}) \equiv \delta U_{\hat{z}}(s_*) = \frac{2\kappa_{\hat{z}}^3}{3\tau_{\hat{z}}^2} = \frac{2}{3} \frac{(\mathcal{M} : \hat{z}\hat{z})^3}{(U''' : \hat{z}\hat{z}\hat{z})^2}$$

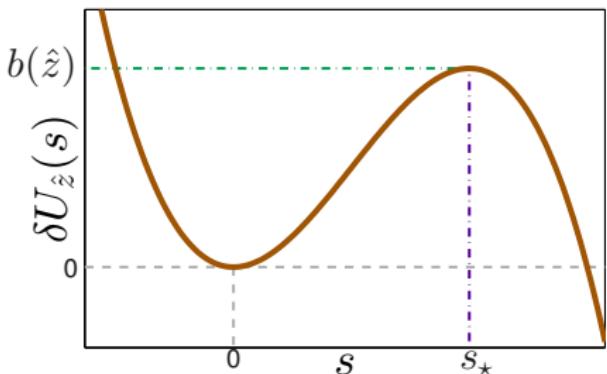
only a function of
inherent state information,
and the direction \hat{z}



the barrier function – definition

barrier function $b(\hat{z}) \equiv \delta U_{\hat{z}}(s_*) = \frac{2}{3} \frac{\kappa_{\hat{z}}^3}{\tau_{\hat{z}}^2} = \frac{2}{3} \frac{(\mathcal{M} : \hat{z}\hat{z})^3}{(U''' : \hat{z}\hat{z}\hat{z})^2}$

directions \hat{z} which take the system over
low saddle points will have small $b(\hat{z})$'s

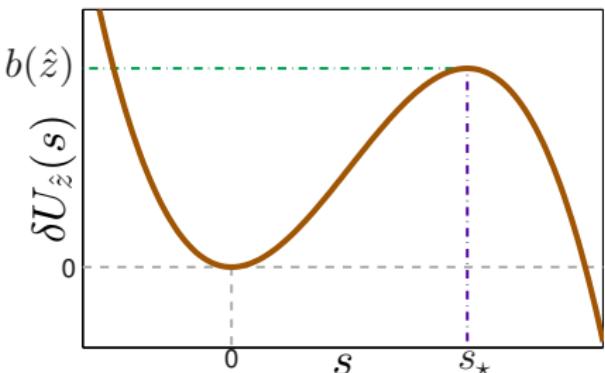


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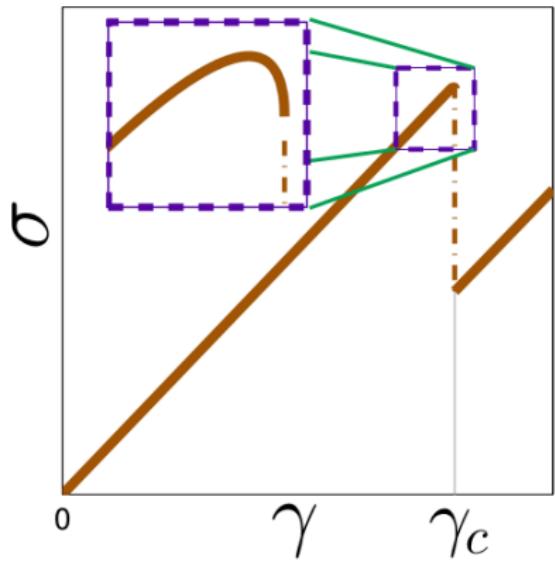
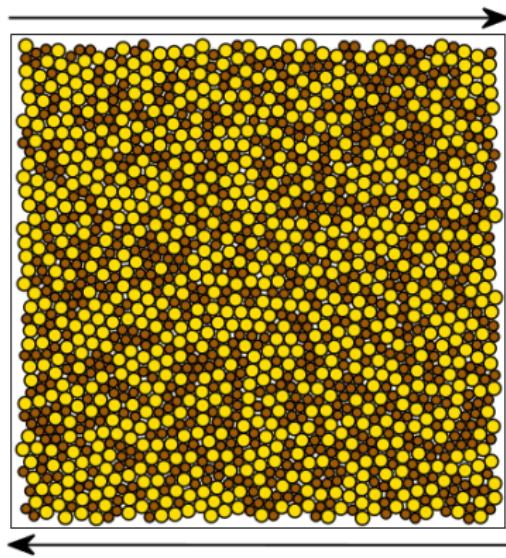
directions \hat{z} which take the system over
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⇒ find directions with
small $b(\hat{z})$ by minimizing
 $b(\hat{z})$ over directions \hat{z}



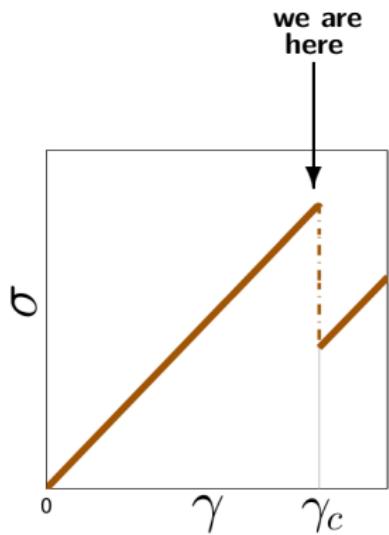
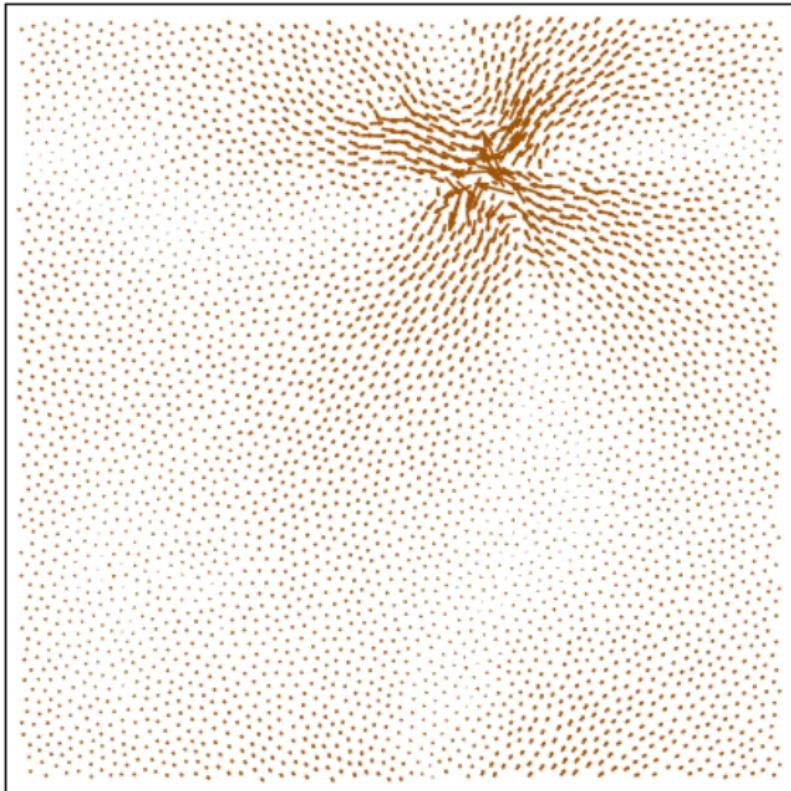
finding small $b(\hat{z})$'s

setup:



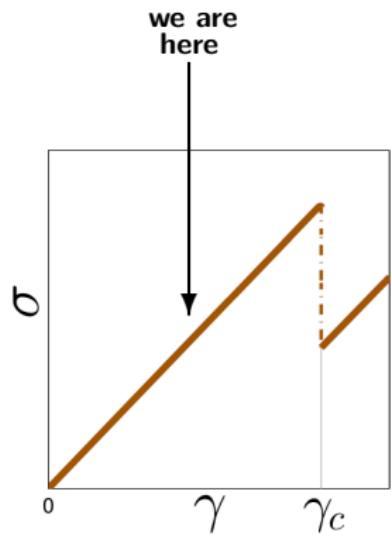
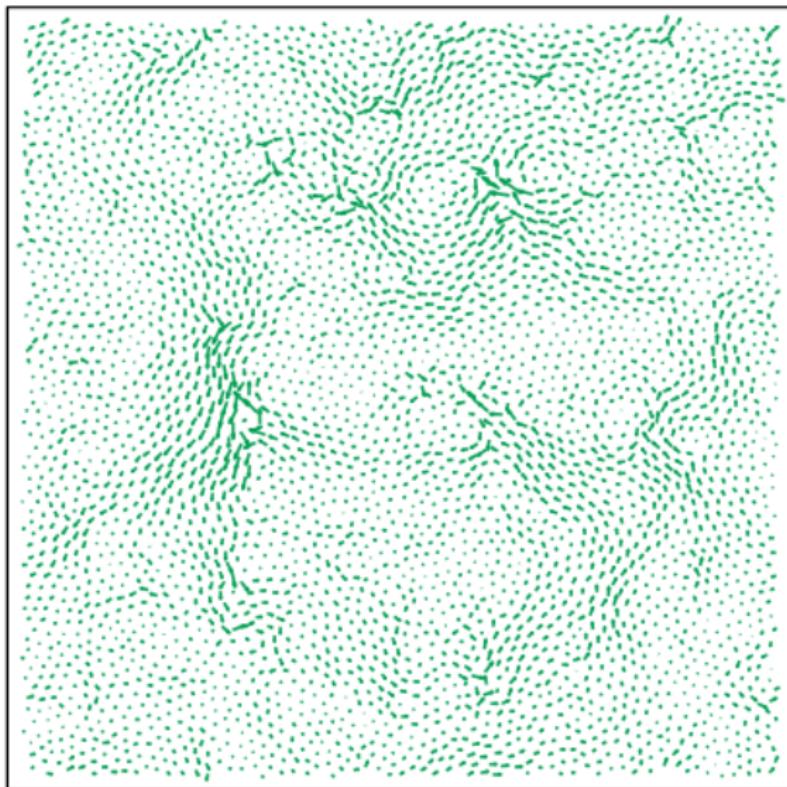
finding small $b(\hat{z})$'s

destabilizing **linear** mode $\hat{\Psi}_p$, $\gamma \rightarrow \gamma_c$



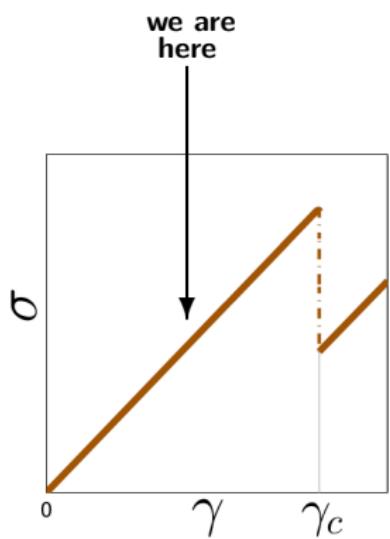
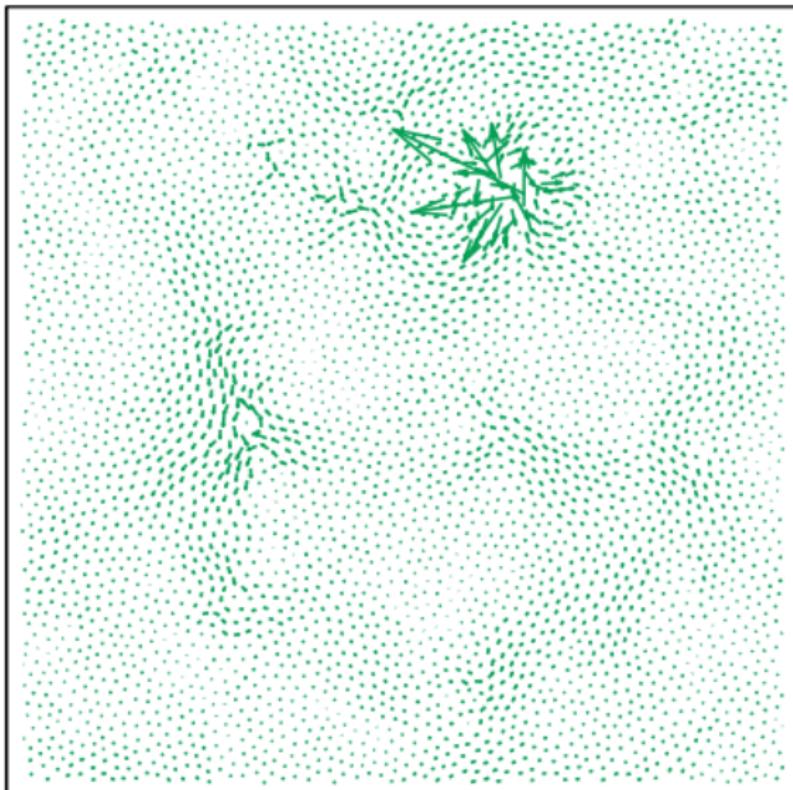
finding small $b(\hat{z})$'s

linear response to shear, $\gamma_c - \gamma \sim 10^{-2}$, use as initial guess \hat{z}_{ini}



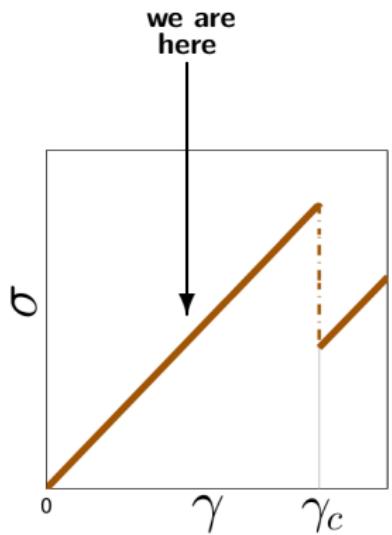
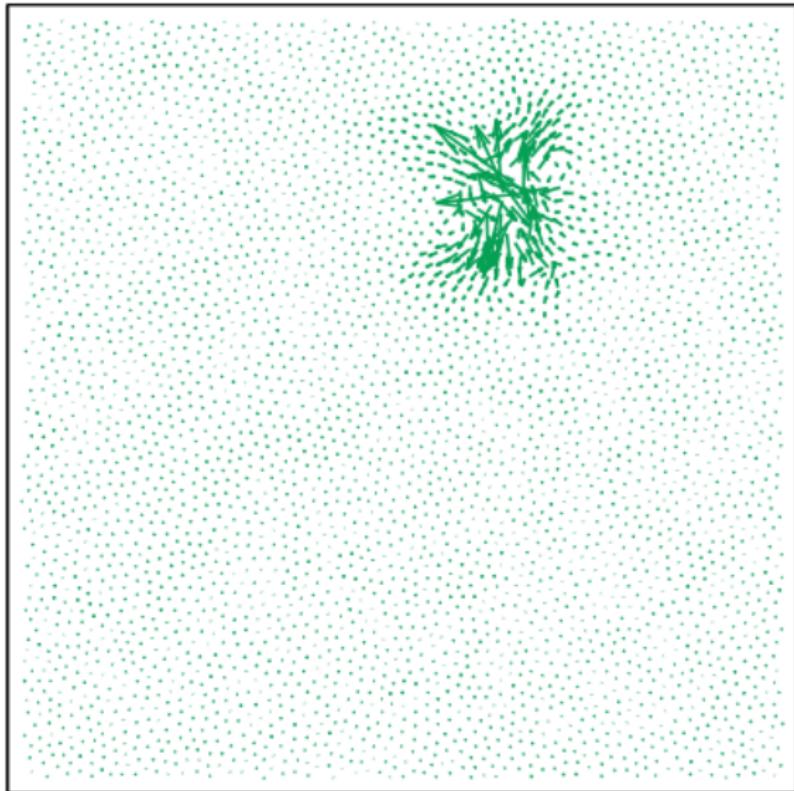
finding small $b(\hat{z})$'s

minimize $b(\hat{z})$, after 12 iterations, $\gamma_c - \gamma \sim 10^{-2}$



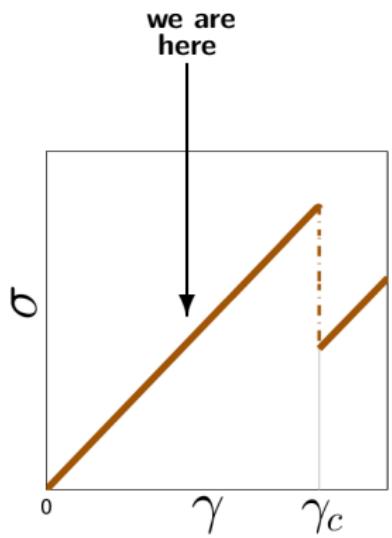
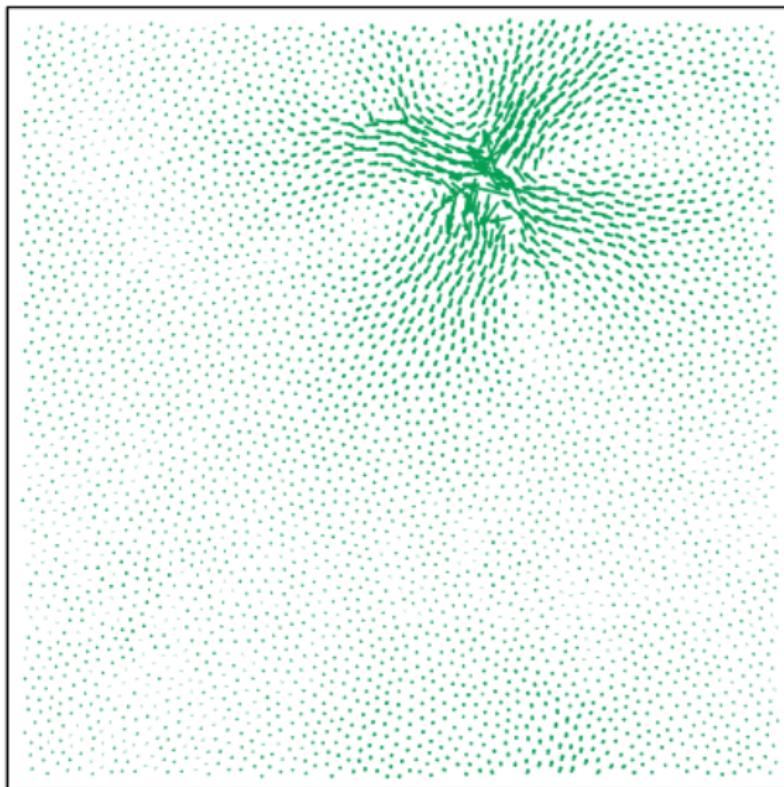
finding small $b(\hat{z})$'s

minimize $b(\hat{z})$, after 24 iterations, $\gamma_c - \gamma \sim 10^{-2}$



finding small $b(\hat{z})$'s

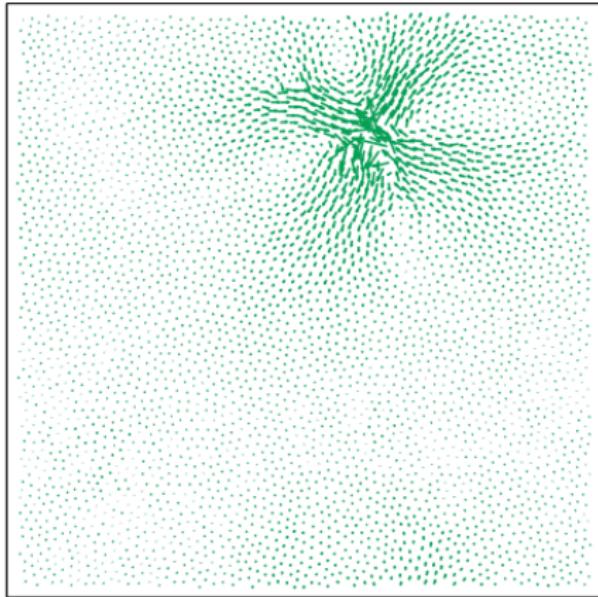
minimization converged, $\gamma_c - \gamma \sim 10^{-2}$



finding small $b(\hat{z})$'s

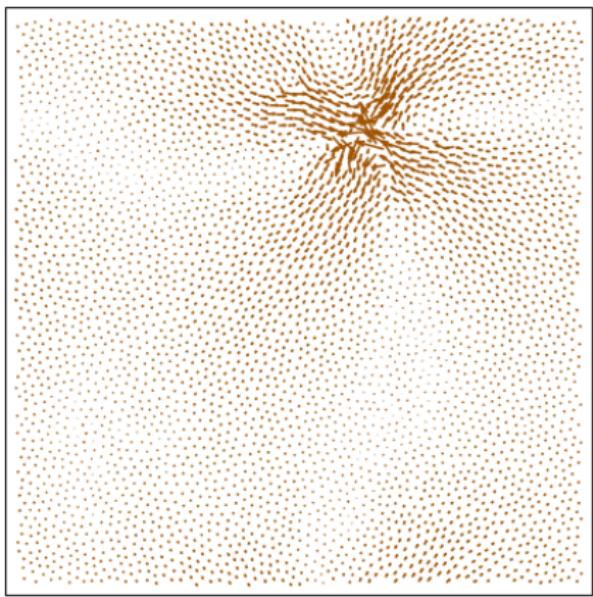
direction \hat{z} that **minimizes** $b(\hat{z})$

$$\gamma_c - \gamma \sim 10^{-2}$$



destabilizing **linear mode** $\hat{\Psi}_p$

$$\gamma_c - \gamma \sim 10^{-7}$$

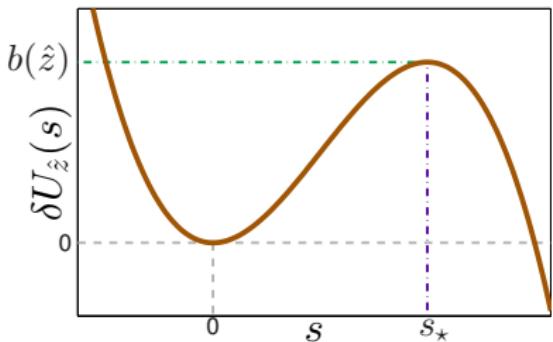


nonlinear plastic modes – definition

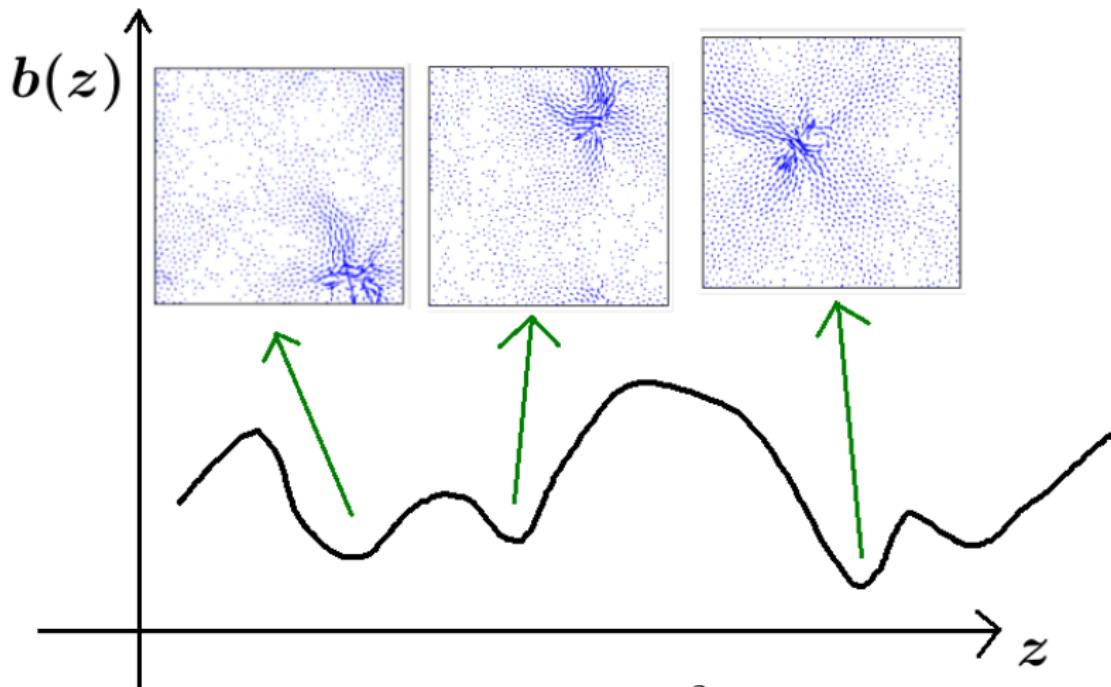
nonlinear plastic modes are collective displacement directions $\hat{\pi}$ for which the barrier function $b(\hat{z})$ displays a **local minimum**

$$\frac{\partial b}{\partial \vec{z}} \Big|_{\hat{\pi}} = 0, \quad \frac{\partial^2 b}{\partial \vec{z} \partial \vec{z}} \Big|_{\hat{\pi}} > 0$$

$$b(\hat{z}) \equiv \frac{2}{3} \frac{(\mathcal{M} : \hat{z}\hat{z})^3}{(U''' : \hat{z}\hat{z}\hat{z})^2}$$

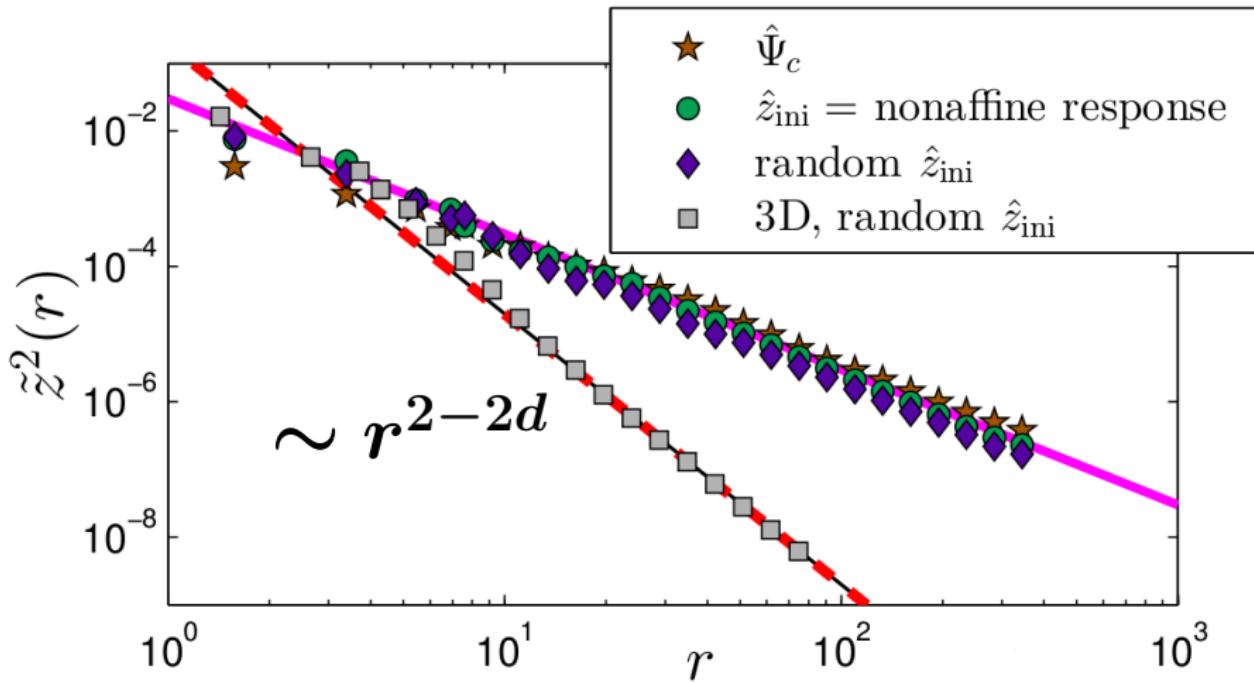


nonlinear plastic modes – illustration



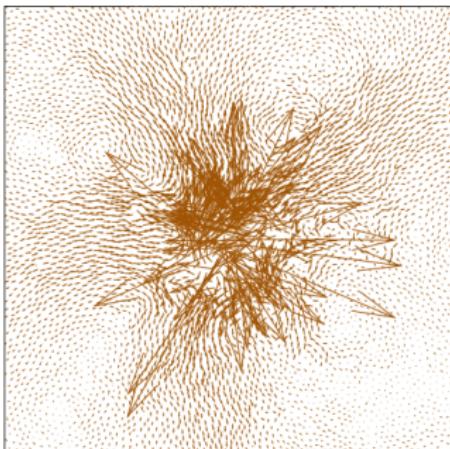
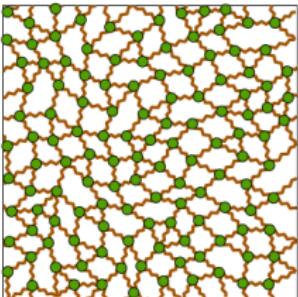
$$b(\hat{z}) \equiv \frac{2}{3} \frac{(\mathcal{M} : \hat{z}\hat{z})^3}{(U''' : \hat{z}\hat{z}\hat{z})^2}$$

nonlinear plastic modes decay like the response to local perturbation $|\hat{\pi}| \sim r^{1-d}$

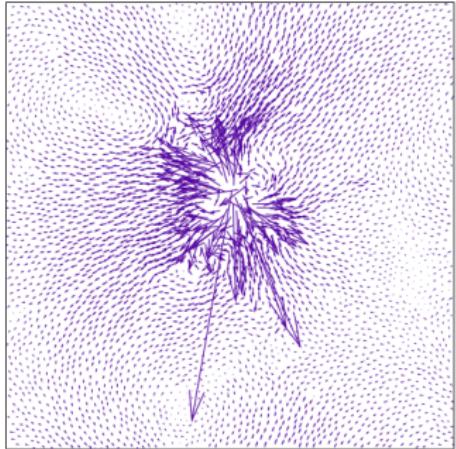
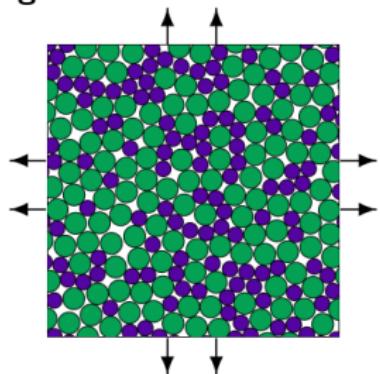


nonlinear plastic modes – spatial structure

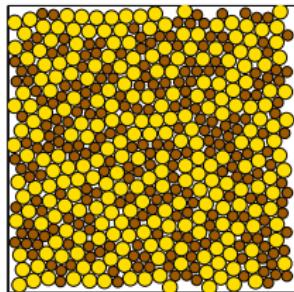
no internal stresses



glass under tension

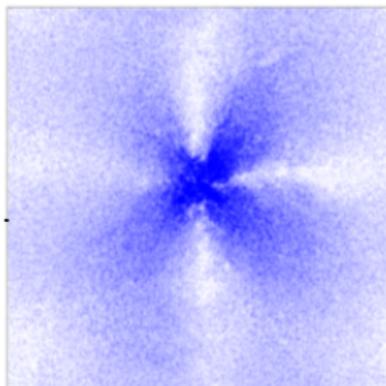


nonlinear plastic modes – core size

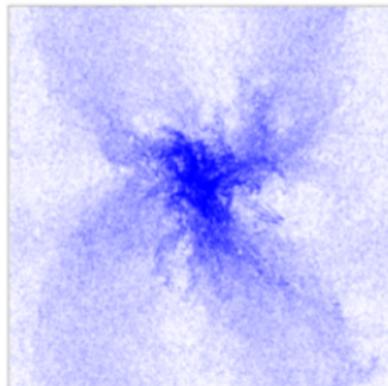


packings of harmonic discs

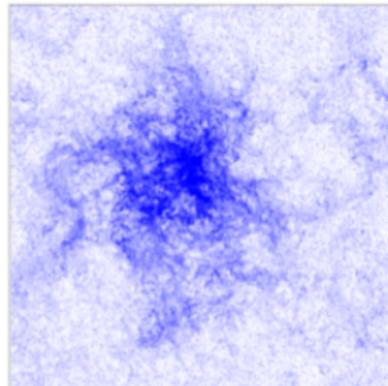
'unjamping'



$$p = 10^{-1}$$

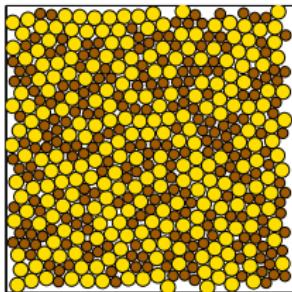


$$p = 10^{-3}$$



$$p = 10^{-5}$$

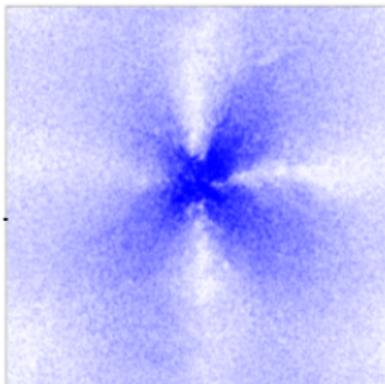
nonlinear plastic modes – core size



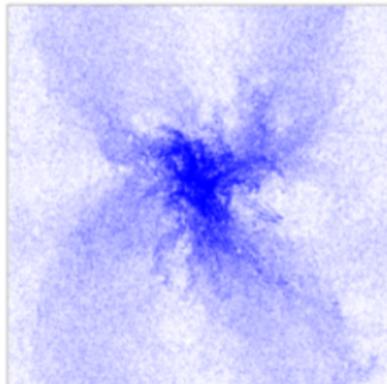
packings of harmonic discs

$$\ell_c \sim \frac{1}{\sqrt{z - z_c}} \sim p^{-1/4}?$$

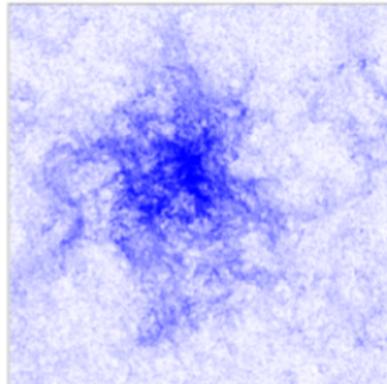
'unjamming'



$p = 10^{-1}$



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$p = 10^{-5}$

usefulness of nonlinear plastic modes

nonlinear plastic modes are collective displacement directions $\hat{\pi}$ for which the barrier function $b(\hat{z})$ displays a **local minimum**

$$\frac{\partial b}{\partial \vec{z}} \Big|_{\hat{\pi}} = 0, \quad \frac{\partial^2 b}{\partial \vec{z} \partial \vec{z}} \Big|_{\hat{\pi}} > 0$$

why are nonlinear plastic modes the
natural micromechanical objects
to consider in plasticity studies?

we defined $\hat{\pi}$ via $\frac{\partial b}{\partial \vec{z}}|_{\hat{\pi}} = 0$

nonlinear plastic modes – deformation dynamics

we defined $\hat{\pi}$ via $\frac{\partial b}{\partial \vec{z}}|_{\hat{\pi}} = 0$

\Rightarrow modes $\hat{\pi}$ solve the nonlinear equation:

$$(*) \quad \mathcal{M} \cdot \hat{\pi} = \frac{\kappa_{\hat{\pi}}}{\tau_{\hat{\pi}}} \mathbf{U}''' : \hat{\pi} \hat{\pi}$$

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how do the stiffnesses $\kappa_{\hat{\pi}}$ depend on deformation?

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how do the stiffnesses $\kappa_{\hat{\pi}}$ depend on deformation?

$$\frac{d\kappa_{\hat{\pi}}}{d\gamma} \simeq \frac{d\mathcal{M}}{d\gamma} : \hat{\pi} \hat{\pi} = -\mathbf{U}''' : \hat{\pi} \hat{\pi} \left(\mathcal{M}^{-1} \cdot \frac{\partial^2 \mathbf{U}}{\partial \vec{x} \partial \gamma} \right)$$

following $(*)$:

$$= -\frac{\tau_{\hat{\pi}} \hat{\pi} \cdot \mathcal{M} \cdot \mathcal{M}^{-1} \cdot \frac{\partial^2 \mathbf{U}}{\partial \vec{x} \partial \gamma}}{\kappa_{\hat{\pi}}}$$
$$= -\frac{\tau_{\hat{\pi}} \nu_{\hat{\pi}}}{\kappa_{\hat{\pi}}}, \quad \tau_{\hat{\pi}} \equiv \frac{\partial^3 U}{\partial \vec{x} \partial \vec{x} \partial \vec{x}} : \hat{\pi} \hat{\pi} \hat{\pi} \text{ asymmetry}$$
$$\nu_{\hat{\pi}} \equiv \hat{\pi} \cdot \frac{\partial^2 U}{\partial \vec{x} \partial \gamma} \text{ shear coupling}$$

nonlinear plastic modes – deformation dynamics

we found a simple form for

$$\frac{d\kappa_{\hat{\pi}}}{d\gamma} \simeq -\frac{\tau_{\hat{\pi}}\nu_{\hat{\pi}}}{\kappa_{\hat{\pi}}}$$

nonlinear plastic modes – deformation dynamics

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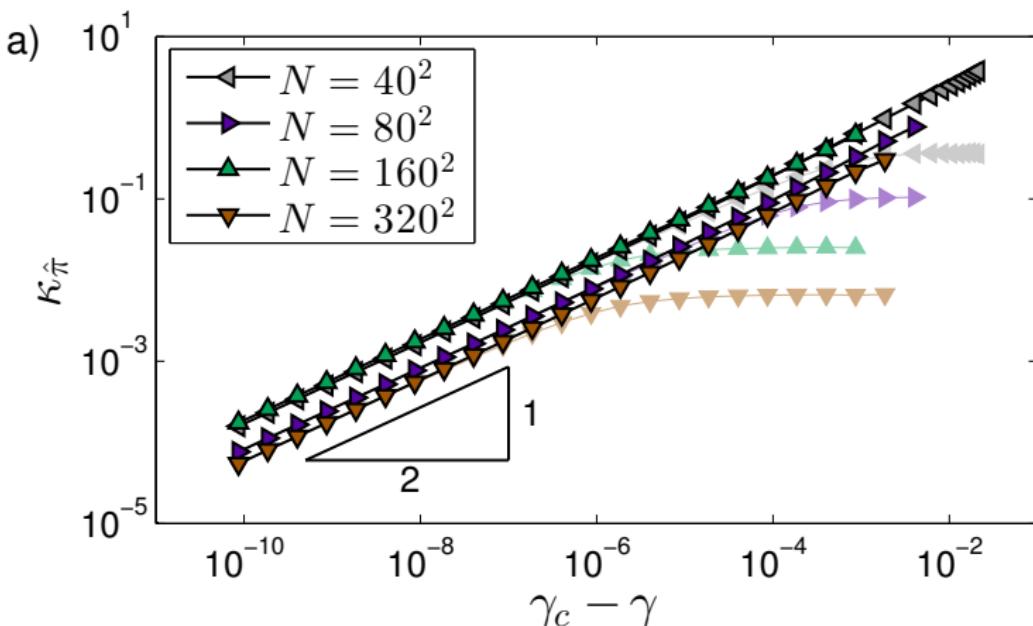
trivially solved as $\kappa_{\hat{\pi}} \simeq \sqrt{2\nu_{\hat{\pi}}\tau_{\hat{\pi}}}\sqrt{\gamma_c - \gamma}$

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nonlinear plastic modes – deformation dynamics

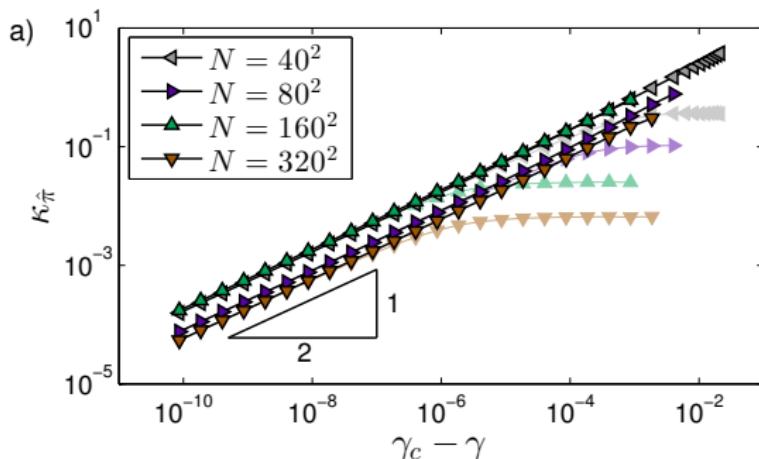
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important points:

- 1) deformation dynamics only weakly coupled to other modes
- 2) N -independent range of validity, in stark contrast with linear modes

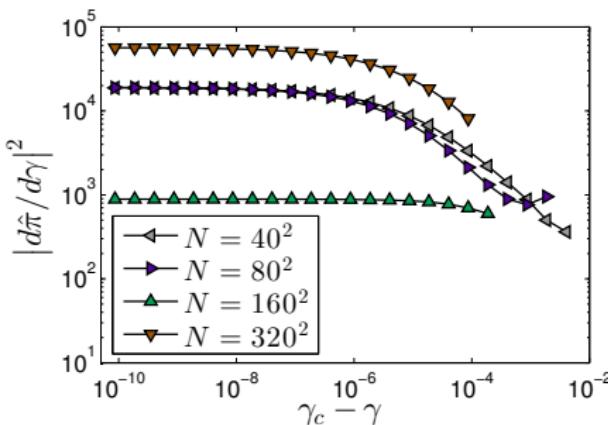


nonlinear plastic modes – deformation dynamics

linear modes' variations are **singular**, plastic modes' are **regular**:

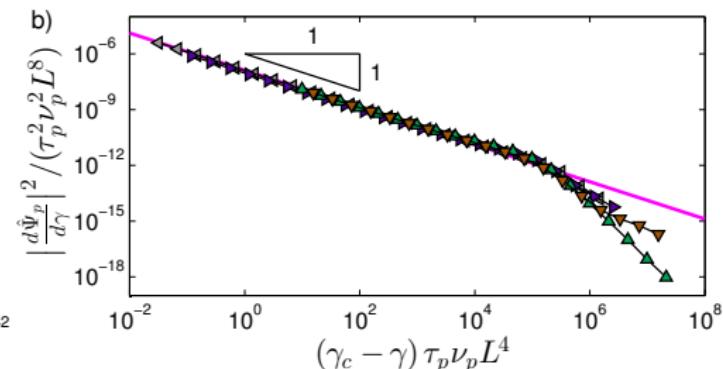
nonlinear plastic modes

$$\left| \frac{d\hat{\pi}}{d\gamma} \right| \sim \text{const.}$$



linear destabilizing mode

$$\left| \frac{d\hat{\Psi}_p}{d\gamma} \right| \sim \frac{L^2}{\sqrt{\gamma_c - \gamma}}$$



nonlinear plastic modes – deformation dynamics

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nonlinear plastic modes

$$\left| \frac{d\hat{\pi}}{d\gamma} \right| \sim \text{const.}$$

linear destabilizing mode

$$\left| \frac{d\hat{\Psi}_p}{d\gamma} \right| \sim \frac{L^2}{\sqrt{\gamma_c - \gamma}}$$

this is odd since both stiffnesses follow same EOM

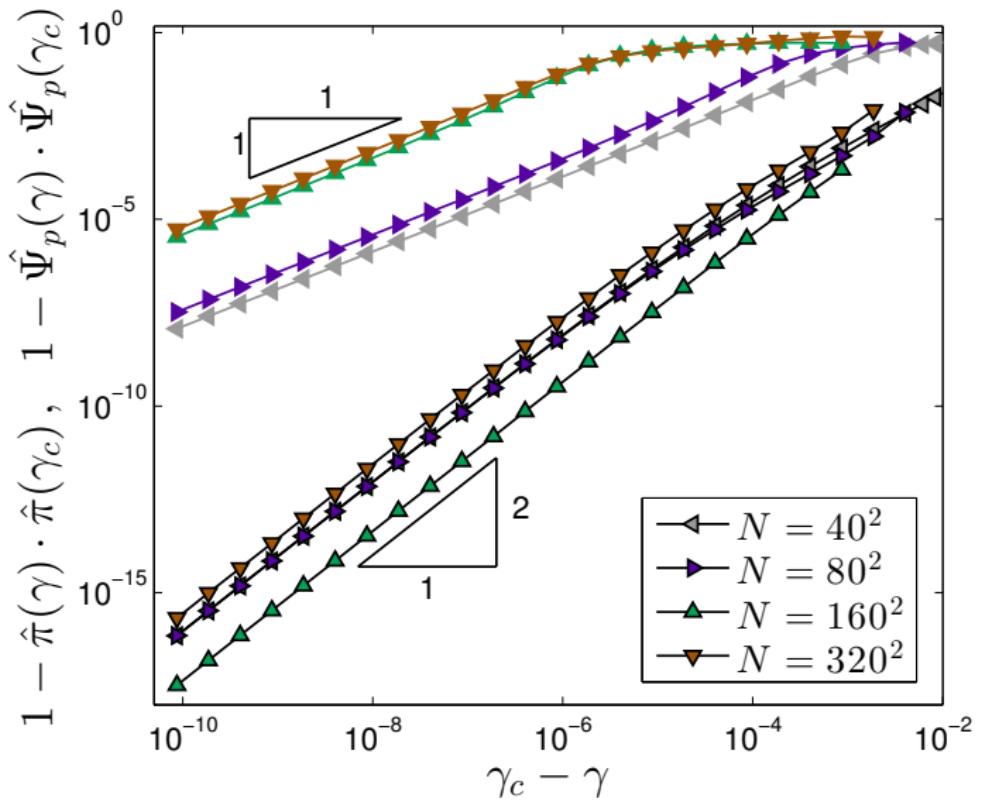
$$\frac{d\kappa}{d\gamma} \sim \frac{1}{\kappa}$$

$$(\kappa \equiv \mathcal{M} : \hat{\pi}\hat{\pi})$$

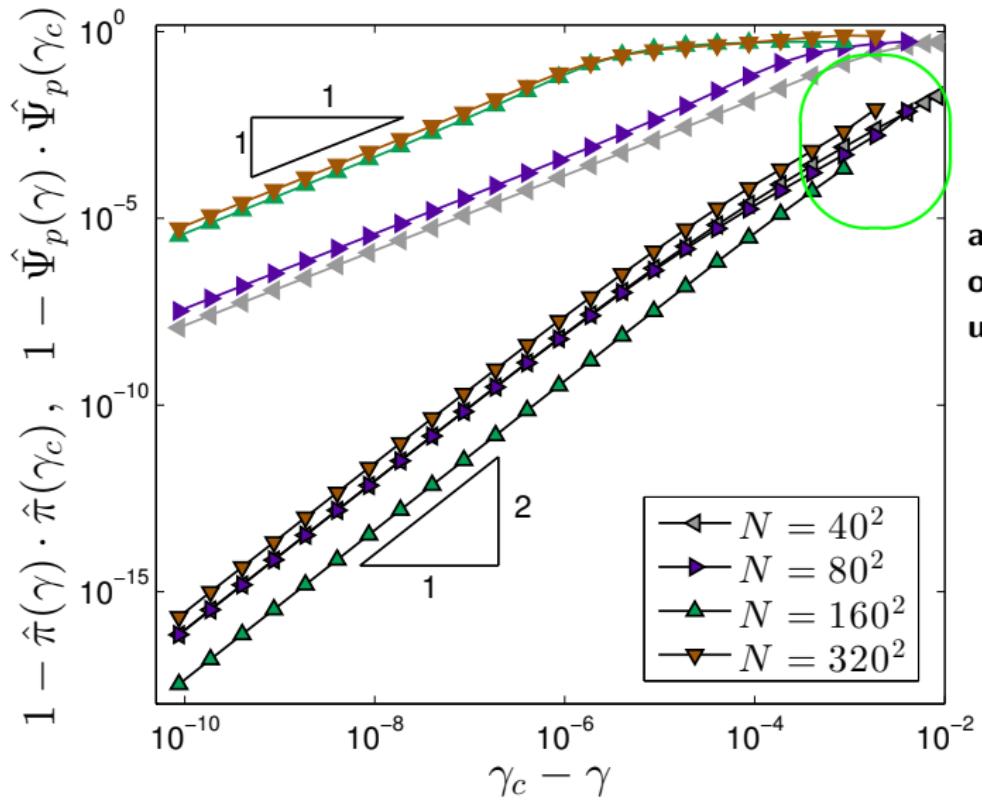
$$\frac{d\lambda_p}{d\gamma} \sim \frac{1}{\lambda_p}$$

$$(\lambda_p \equiv \mathcal{M} : \hat{\Psi}_p \hat{\Psi}_p)$$

predictiveness of nonlinear plastic modes

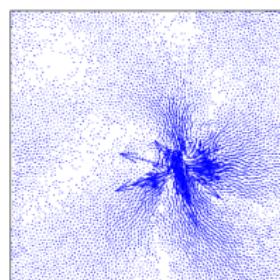
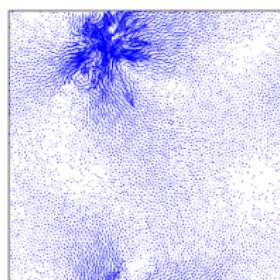
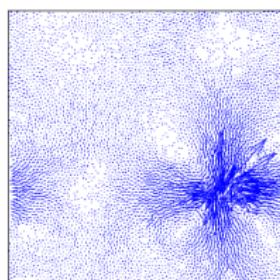
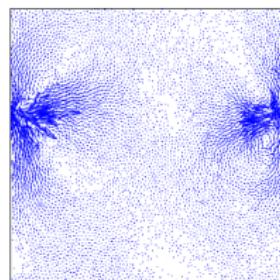
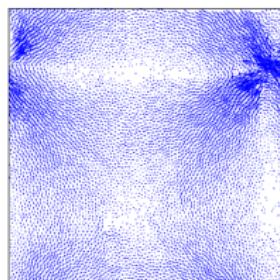
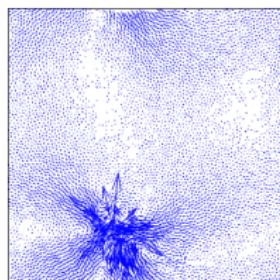
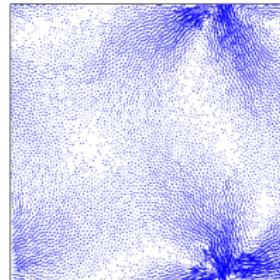
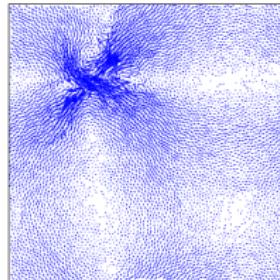
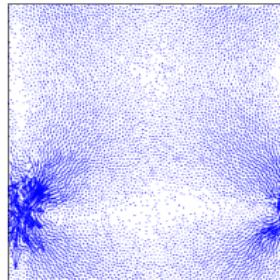


predictiveness of nonlinear plastic modes



as soon as detected,
overlap with instability
up to more than 99%!

TBD: detecting the **field** of nonlinear plastic modes



**modes detected
in a single sample**

in progress...

statistics of nonlinear plastic modes

what attributes of NPM's should we care about?

statistics of nonlinear plastic modes

what attributes of NPM's should we care about?

recall: NPM's are characterized by:

- their stiffnesses $\kappa = \frac{\partial^2 U}{\partial \vec{x} \partial \vec{x}} : \hat{\pi} \hat{\pi}$
- their asymmetries $\tau = \frac{\partial^3 U}{\partial \vec{x} \partial \vec{x} \partial \vec{x}} : \hat{\pi} \hat{\pi} \hat{\pi}$
- their deformation coupling $\nu = \frac{\partial^2 U}{\partial \gamma \partial \vec{x}} \cdot \hat{\pi}$

statistics of nonlinear plastic modes

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we can construct a **field of local destabilization strains** $\delta \gamma_c(\hat{\pi})$:

$$\delta \gamma_c(\hat{\pi}) = \gamma_c(\hat{\pi}) - \gamma = \frac{\kappa}{2 \frac{d\kappa}{d\gamma}} = \frac{\kappa^2}{2\nu\tau}$$

(recall that $\kappa = \sqrt{2\tau\nu} \sqrt{\gamma_c - \gamma}$, and $\frac{d\kappa}{d\gamma} = -\frac{\tau\nu}{\kappa}$)

statistics of nonlinear plastic modes

what attributes of NPM's should we care about?

recall: NPM's are characterized by:

- their stiffnesses $\kappa = \frac{\partial^2 U}{\partial \vec{x} \partial \vec{x}} : \hat{\pi} \hat{\pi}$ assume τ and ν have non-interesting distributions,
- their asymmetries $\tau = \frac{\partial^3 U}{\partial \vec{x} \partial \vec{x} \partial \vec{x}} : \hat{\pi} \hat{\pi} \hat{\pi}$ focus on stiffnesses κ
- their deformation coupling $\nu = \frac{\partial^2 U}{\partial \gamma \partial \vec{x}} \cdot \hat{\pi}$

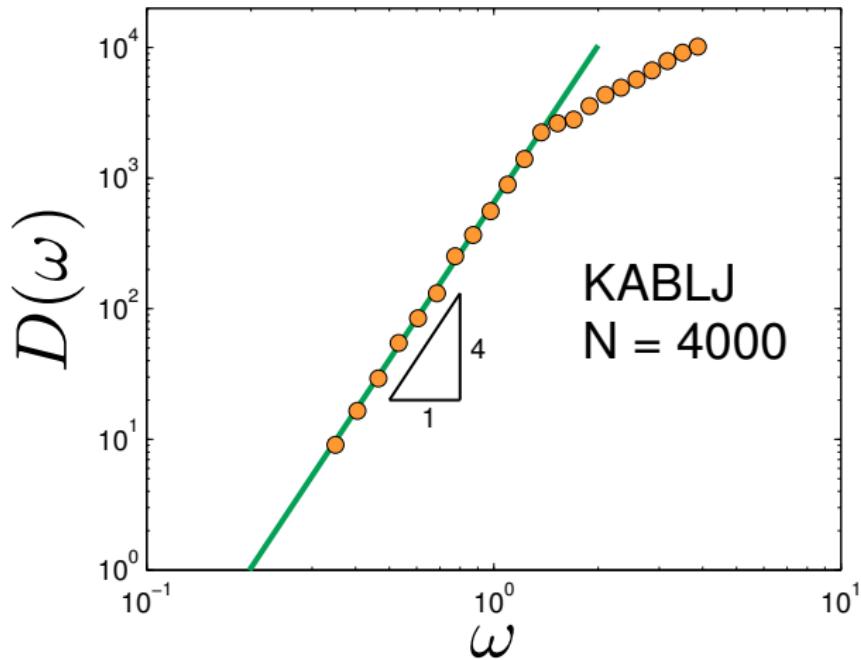
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how are NPM stiffnesses κ distributed?

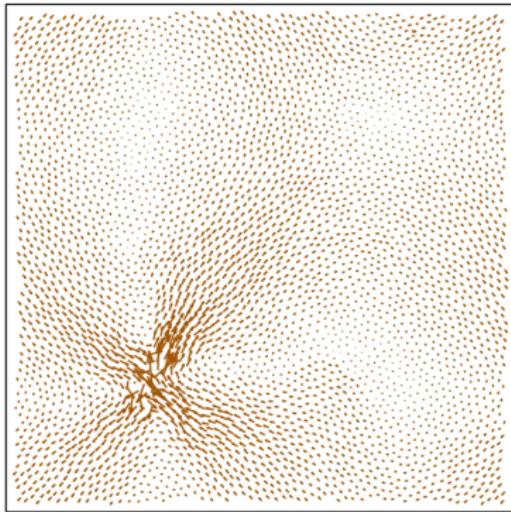
it was recently observed that a **universal** distribution $D(\omega) \sim \omega^4$ of quasi-localized **glassy modes** appears at low frequencies



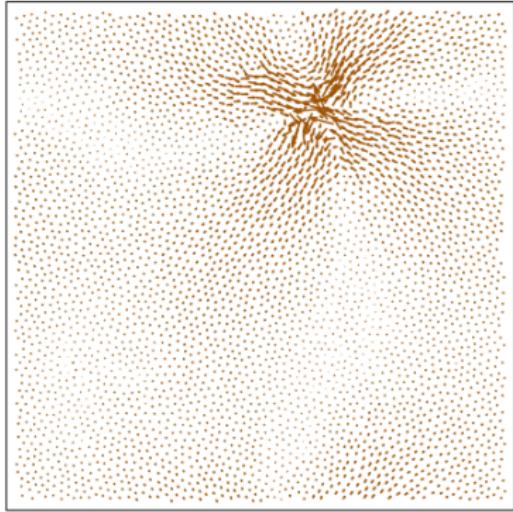
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harmonic glassy mode
in **undeformed** sample



plastic instability
upon imposing **shear**



how are NPM stiffnesses κ distributed?

recall we assume that strain couplings ν and asymmetries τ have **characteristic** (κ independent) values, then we expect

$$p(\kappa) \sim \kappa^{3/2} \quad \Rightarrow \quad p(\delta\gamma_c) \sim \delta\gamma_c^{1/4}$$

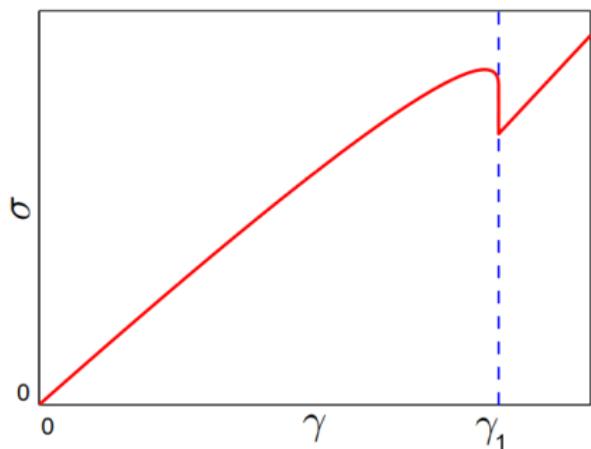
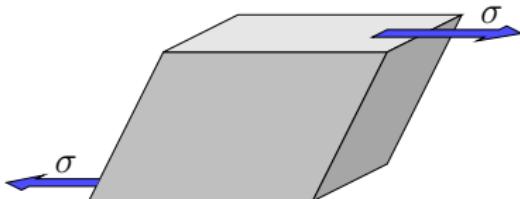
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assume now different NPMs are independent, then we expect

$$\gamma_1(N) \sim N^{-\frac{4}{5}}$$



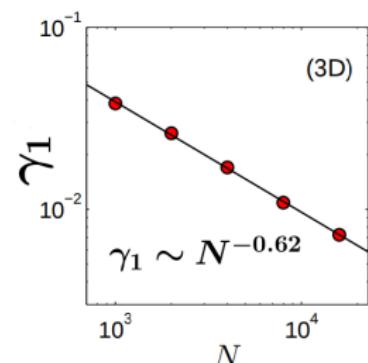
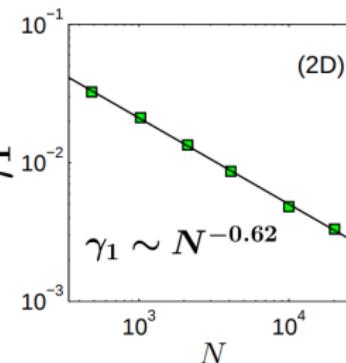
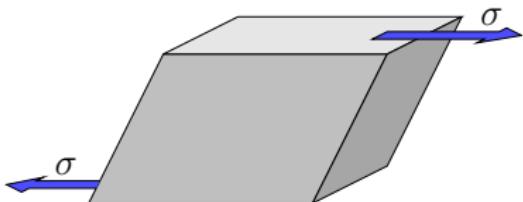
extent of first elastic branch

recall we assume that strain couplings ν and asymmetries τ have **characteristic** (κ independent) values, then we expect

$$p(\kappa) \sim \kappa^{3/2} \Rightarrow p(\delta\gamma_c) \sim \delta\gamma_c^{1/4}$$

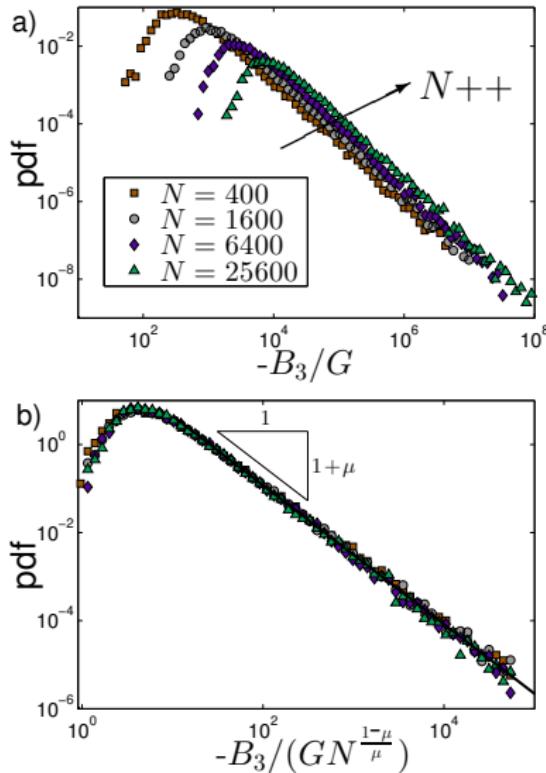
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finite-size scaling of nonlinear elasticity

a similar discrepancy appears for nonlinear elastic moduli

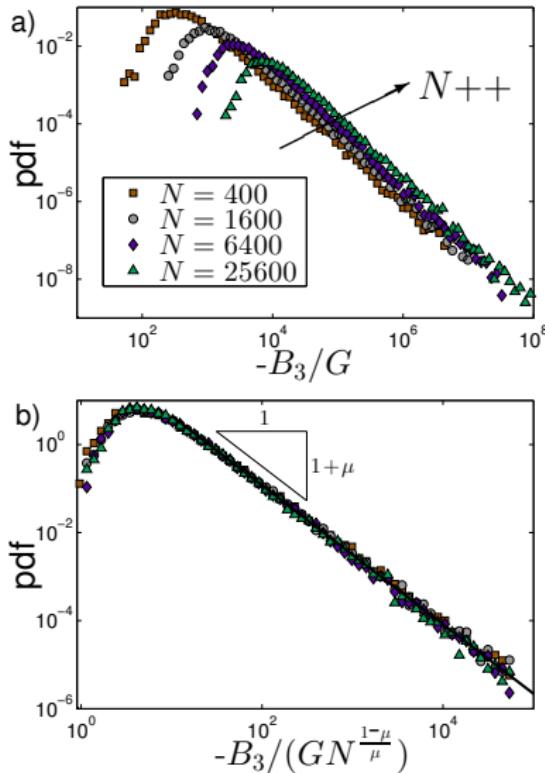


$$B_3 \equiv \frac{d^3\sigma}{d\gamma^3}$$

we find $\mu \approx 0.57$,
whereas $D(\omega) \sim \omega^4$
implies $\mu = 1/2$

finite-size scaling of nonlinear elasticity

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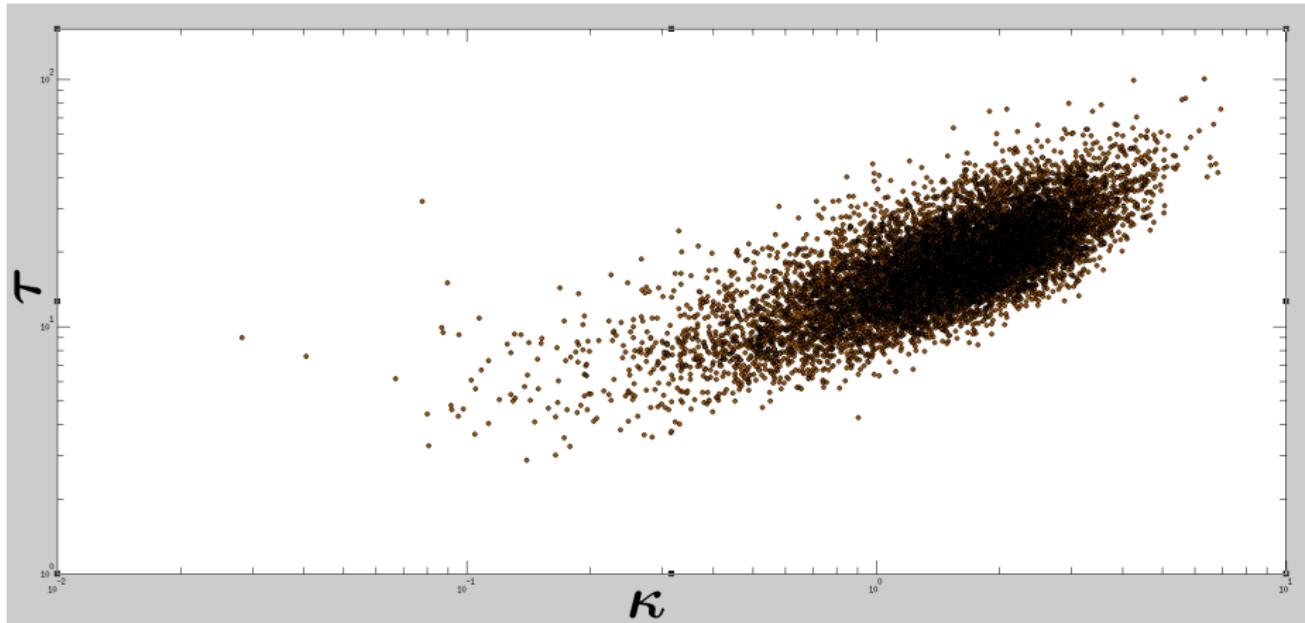
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asymmetries should depend
on stiffnesses

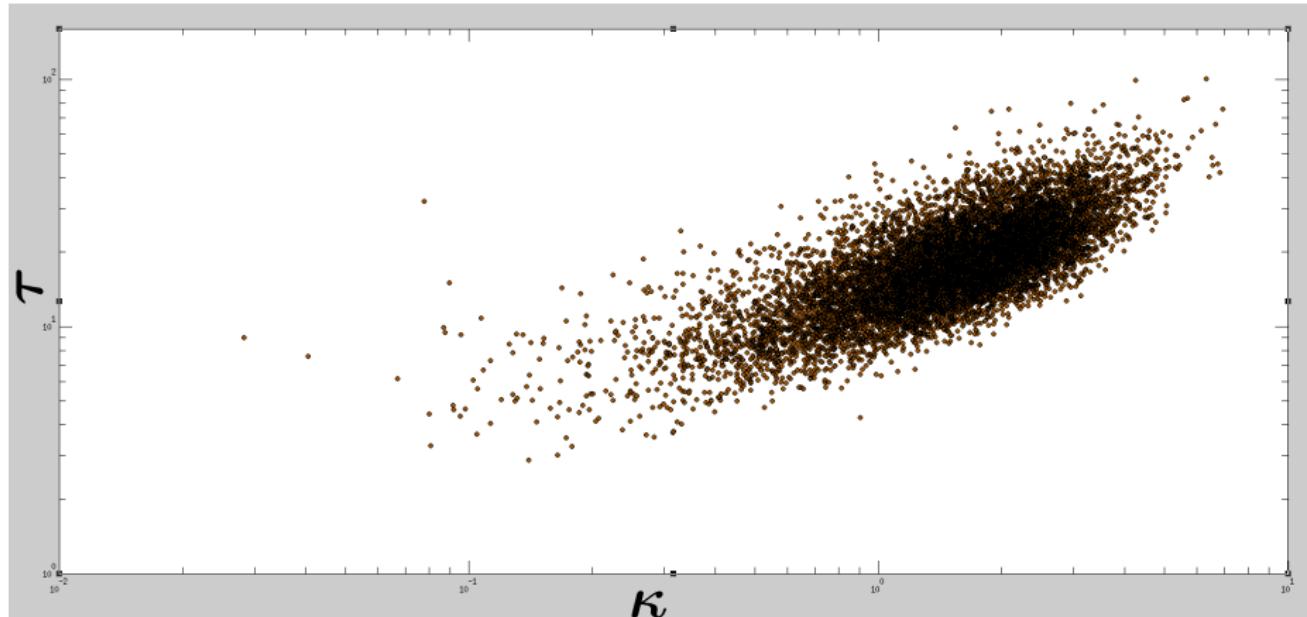
dependence of asymmetries on stiffnesses

data measured for low-energy NPMs in 3D with $N = 2000$



dependence of asymmetries on stiffnesses

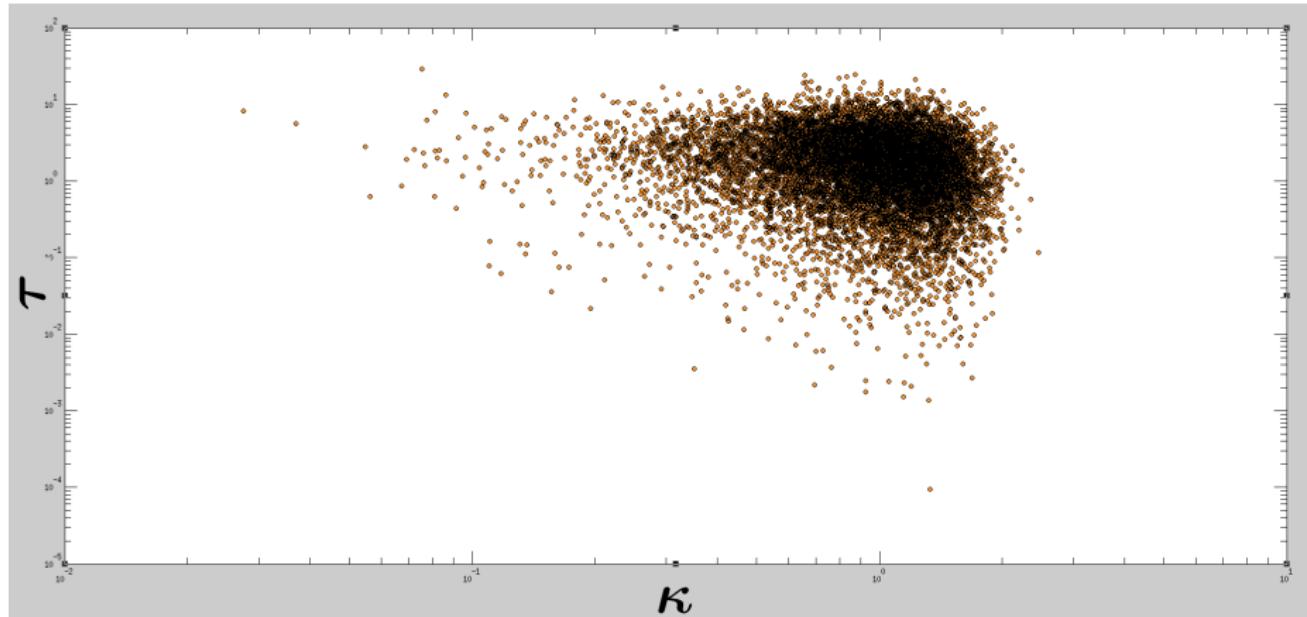
data measured for low-energy NPMs in 3D with $N = 2000$



does this trend persist to $\kappa \rightarrow 0$?

dependence of asymmetries on stiffnesses

data measured for low-energy **harmonic modes** in 3D with $N = 2000$

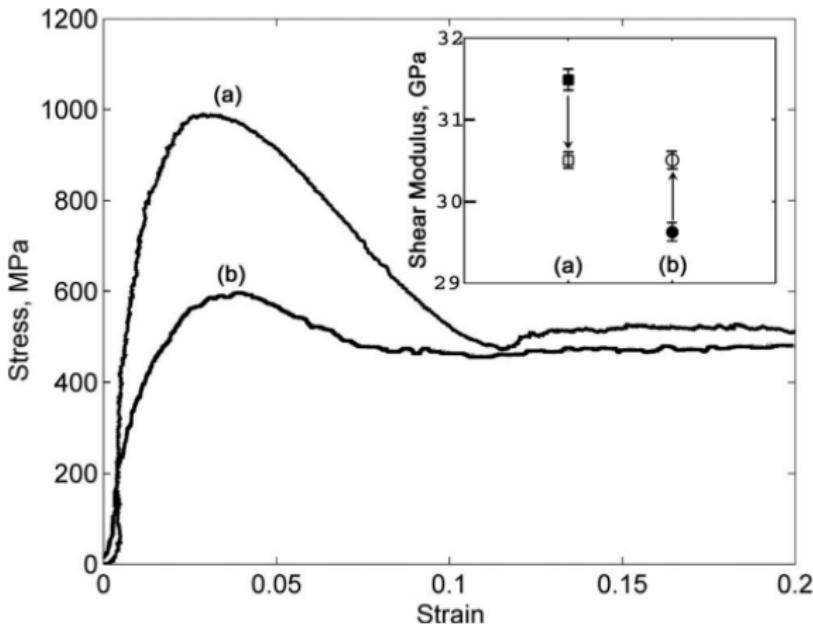


asymmetries appear to be **stiffness independent** for harmonic modes, but **not** for plastic modes

summary: nonlinear plastic modes

- understanding elasto-plasticity and yielding requires the proper identification of the relevant structural state variables

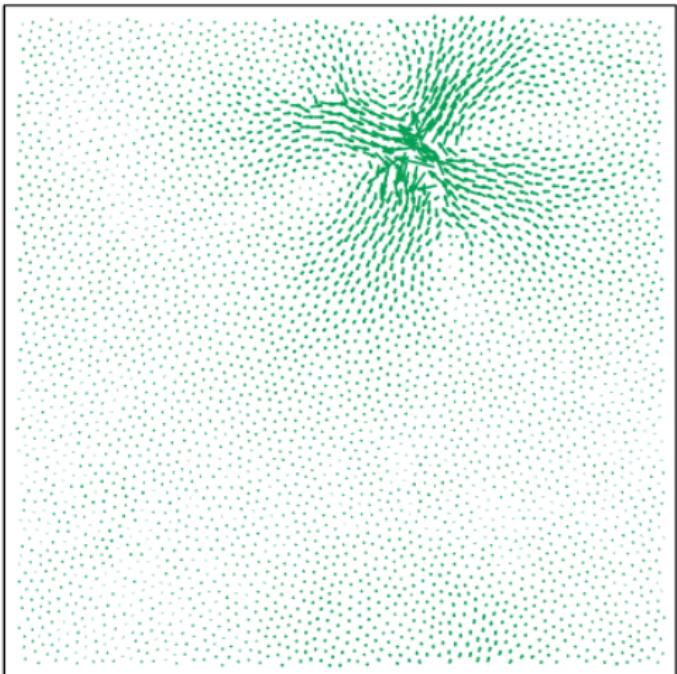
J. Lu, G. Ravichandran, W. Johnson, Acta Materialia 51 (2003)



summary: nonlinear plastic modes

- NPMs offer a robust **micromechanical** definition of plasticity carriers, based solely on **inherent state information**

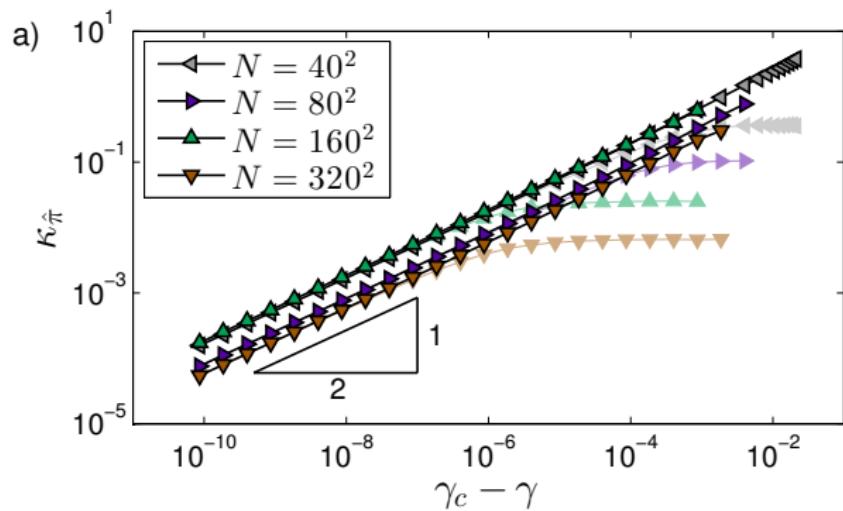
$$\frac{\partial b}{\partial \vec{z}} \Big|_{\hat{\pi}} = 0$$



summary: nonlinear plastic modes

- deformation dynamics of NPMs: N -independent,
no hybridizations

$$\frac{d\kappa}{d\gamma} \underset{\sim}{=} -\frac{\tau \nu}{\kappa}$$



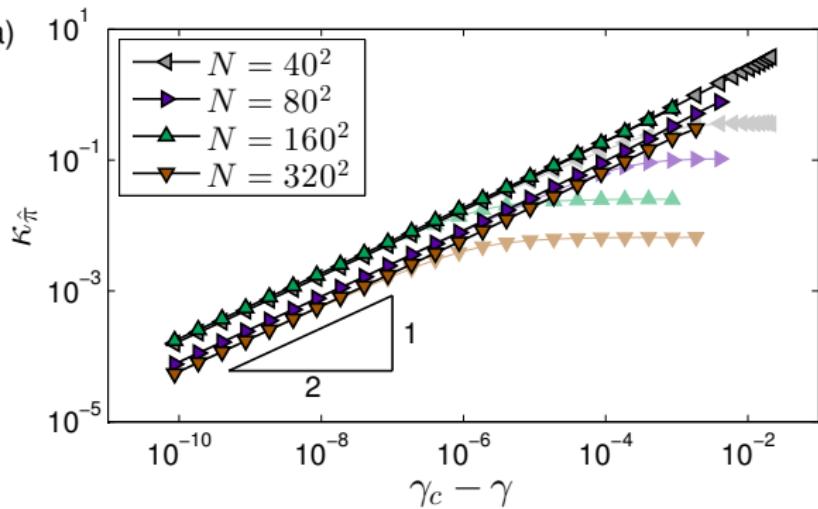
summary: nonlinear plastic modes

- deformation dynamics of NPMs: N -independent,
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asymmetry

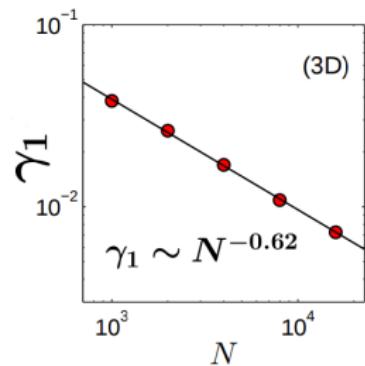
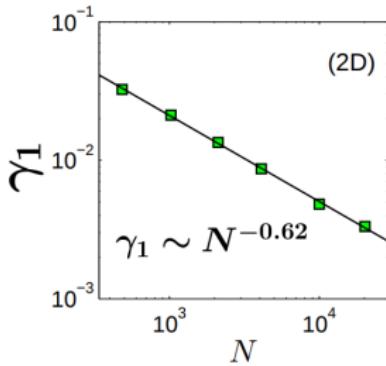
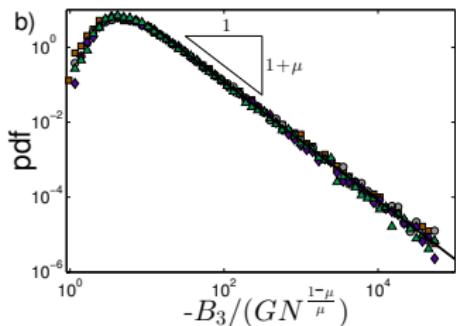
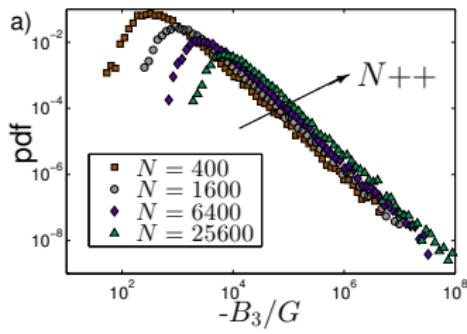
deformation coupling

$$\frac{d\kappa}{d\gamma} \underset{\sim}{=} -\frac{\tau \nu}{\kappa}$$



summary: nonlinear plastic modes

- still something left to understand regarding the statistics of NPMs, and the stiffness-dependence of asymmetries & deformation coupling



Ph.D. & postDoc positions available!

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thanks for your attention! questions?

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