

Exponential suppression of the topological gap in self-consistent intrinsic Majorana nanowires



Francisco Lobo, Elsa Prada, Pablo San-Jose
Consejo Superior de Investigaciones Científicas (CSIC - Spain)



arXiv:2412.15174
To appear in PRB

Predictions of topological p-wave superconductivity and Majorana zero modes (MZMs) in hybrid superconductor-semiconductor nanowires have been difficult to realize experimentally. Consequently, researchers are actively exploring alternative platforms for MZMs. Here, we theoretically study depleted nanowires with intrinsic superconductivity (as opposed to proximity-induced). Using a **self-consistent Hartree-Fock-Bogoliubov mean field theory**, we compute the topological phase diagram versus Zeeman field and filling for intrinsic wires with attractive interactions. We find that, although intrinsic wires could be less vulnerable than hybrids to topology-adverse effects, such as disorder and metallization, they are hindered by a fundamental limitation of their own. Although a **topological p-wave gap** is indeed possible, it is far less robust than in hybrid Majorana nanowires. Instead of remaining stable beyond the topological transition, it is found to **decay exponentially with Zeeman field**, greatly reducing the parameter region with an appreciable topological gap.

Abstract

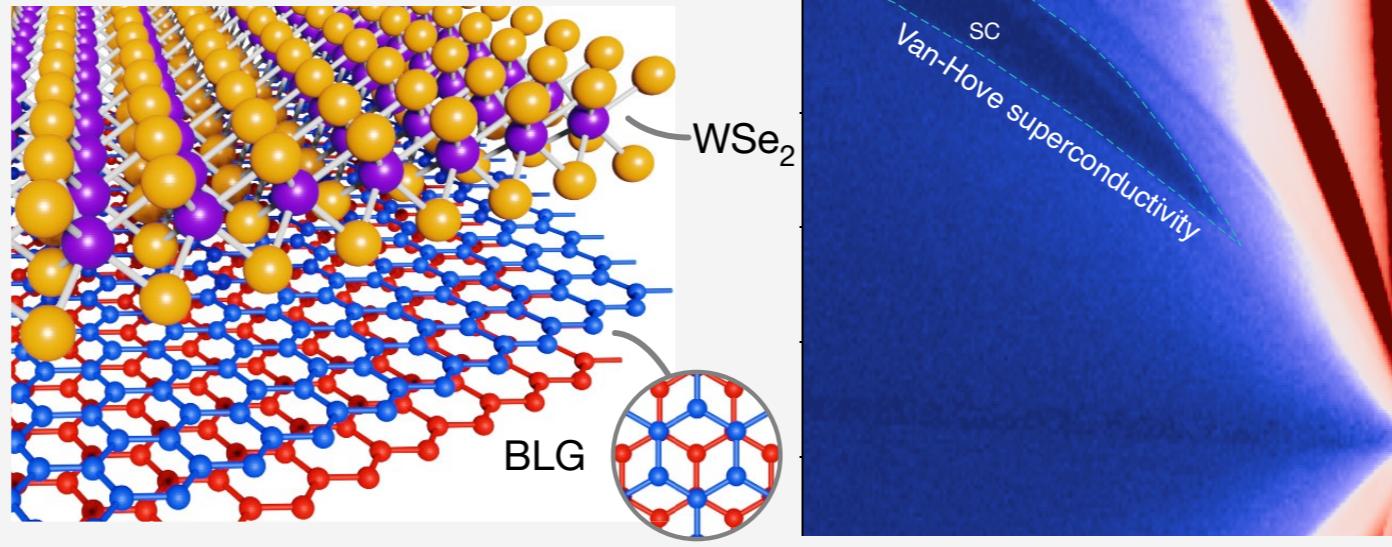
Statement of the problem

- The Oreg-Lutchyn recipe [1,2] for Majorana zero modes (MZMs) is based on **hybrid superconductor-semiconductor (SM/SC)** nanowires
- Key idea: **topological p-wave SC** and MZMs emerges for Zeeman field exceeding a **critical value**
- Hybrid systems are complex and prone to multiple issues, such as **disorder**, **metallization** and **interface effects**. MZMs not convincingly demonstrated yet
- Could **nanowires with intrinsic SC** develop p-wave without these shortcomings?

Note on low-density SCs

The Oreg-Lutchyn recipe requires **low density systems**, since the critical Zeeman grows as the Fermi energy. **Can these develop intrinsic SC?**

While traditional low-density semiconductors have strong SOC, they do not develop superconductivity at any reasonable temperature. However, some new materials (notably **2D heterostructures**) do [3].



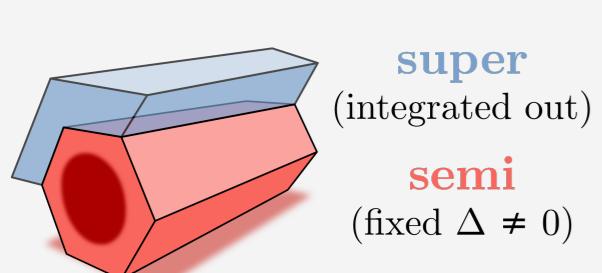
(left) Bilayer graphene/WSe₂ heterostructure. (right)
Phase diagram vs carrier density and displacement field

Majorana nanowire models

We consider three Majorana nanowire tight-binding models of increasing complexity. These are implemented on a quasi-1D nanowire, either infinite or of finite length L

Oreg-Lutchyn (OL) model

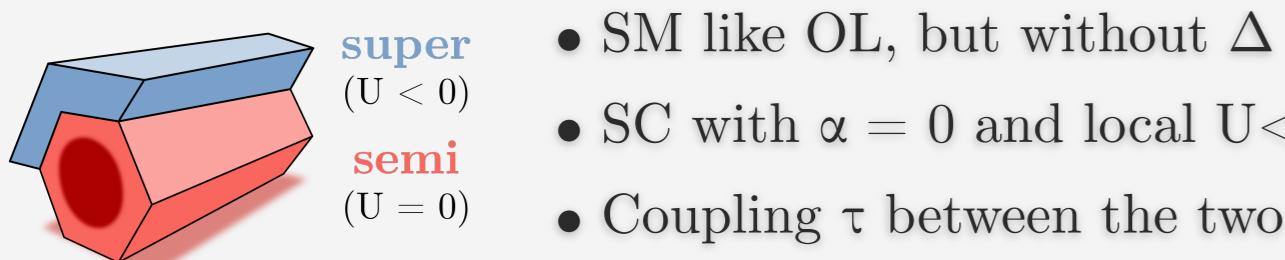
Hybrid SM/SC nanowire with SC integrated out (self-energy from the SC approximated at $\omega = 0$)



$$H^{\text{OL}} = H_0 + H_{\text{SOC}} + H_Z + H_{\Delta^{ii}}$$

Self-consistent hybrid nanowire model

Hybrid SM/SC nanowire with the SC included self-consistently via an attractive **on-site interaction U**



$$H^{\text{hyb}} = H_0^{\text{SM}} + H_Z^{\text{SM}} + H_{\text{SOC}}^{\text{SM}} + H_0^{\text{SC}} + H_Z^{\text{SC}} + \Sigma_{\text{HFB}}^{\text{SC}} + H^{\text{SC-SM}}$$

Self-consistent intrinsic nanowires

Low-density nanowire with SOC and **intrinsic SC** via an attractive on-site interaction U



$$H^{\text{int}} = H_0 + H_{\text{SOC}} + H_Z + \Sigma_{\text{HFB}}$$

Interactions

Electron-electron interactions are modelled by v^{ij}

$$H^{\text{int}} = \sum_{ij} n_i v^{ij} n_j = \sum_{ij\sigma\sigma'} c_{i\sigma}^\dagger c_{i\sigma} v^{ij} c_{j\sigma'}^\dagger c_{j\sigma'}$$

We approximate this by a self-consistent mean field $H^{\text{int}} \rightarrow \Sigma_{\text{HFB}}$ (including anomalous terms)

$$\Sigma_{\text{HFB}} = \Sigma_H + \Sigma_F + \Sigma_\Delta = \frac{1}{2} \sum_{ij\sigma\sigma'} (c_{i\sigma}^\dagger, c_{i\sigma}) \tilde{\Sigma}_{\sigma\sigma'}^{ij} \begin{pmatrix} c_{j\sigma'} \\ c_{j\sigma'}^\dagger \end{pmatrix}$$

Hartree Fock Bogoliubov (anomalous)

The Nambu self energy matrix depends on the Nambu density matrix as

$$\tilde{\Sigma}^{ij} = \frac{1}{2} \delta_{ij} \tau_z \sum_l v^{il} \text{Tr}(\tau_z \tilde{\rho}^{il}) - v^{ij} \tau_z \tilde{\rho}^{ij} \tau_z \quad \rho_{ee}^{ij} = \langle c_j^\dagger \otimes c_i \rangle$$

$$\tilde{\rho}^{ij} = \begin{pmatrix} \rho_{ee}^{ij} & \rho_{eh}^{ij} \\ \rho_{he}^{ij} & -(\rho_{ee}^{ii})^T \end{pmatrix} \quad \rho_{eh}^{ij} = \langle c_j^\dagger \otimes c_i \rangle \quad \rho_{he}^{ij} = \langle c_j^\dagger \otimes c_i^\dagger \rangle$$

Self-consistency equations

The mean field depends on the density matrix, and the density matrix depends on the mean field. Self-consistency is obtained by Anderson-accelerated fixed-point iteration.

The self-consistency equations for a SC without SOC ($\alpha = 0$) can be written out explicitly. This generalizes the BCS gap equation to:

$$\Delta = -\frac{U}{4N} \sum_k \frac{\Delta}{\tilde{E}_k} \left[\tanh\left(\frac{\tilde{E}_k + \tilde{V}_Z}{2k_B T}\right) + \tanh\left(\frac{\tilde{E}_k - \tilde{V}_Z}{2k_B T}\right) \right],$$

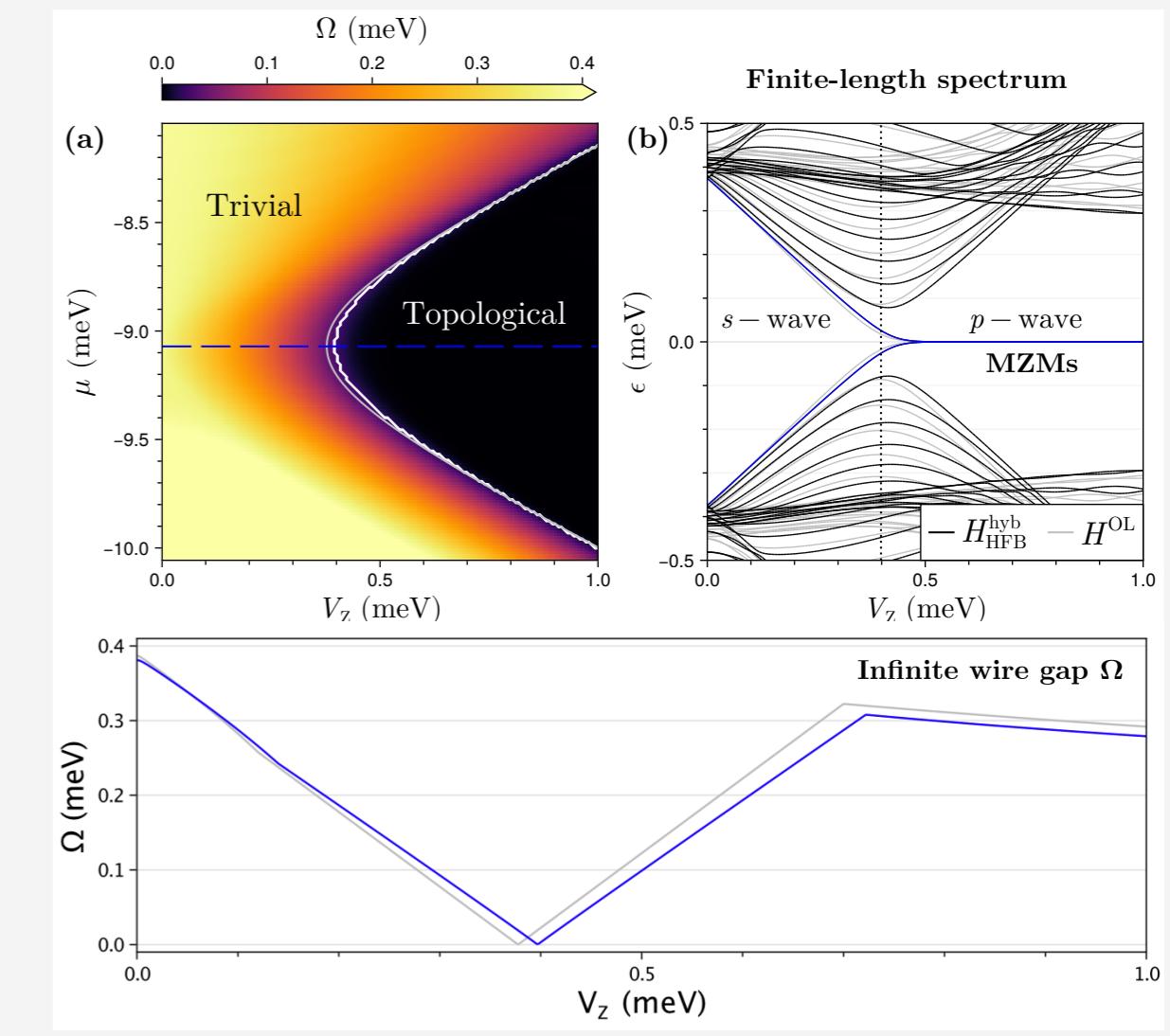
$$\delta\mu = -\frac{U}{2} + \frac{U}{4N} \sum_k \frac{\tilde{\epsilon}_k}{\tilde{E}_k} \left[\tanh\left(\frac{\tilde{E}_k + \tilde{V}_Z}{2k_B T}\right) + \tanh\left(\frac{\tilde{E}_k - \tilde{V}_Z}{2k_B T}\right) \right]$$

$$\delta V_Z = \frac{U}{2N} \sum_k \frac{\sinh(\tilde{V}_Z/k_B T)}{\cosh(\tilde{V}_Z/k_B T) + \cosh(\tilde{E}_k/k_B T)}.$$

$$\tilde{E}_k = \sqrt{\Delta^2 + \tilde{\epsilon}_k^2} \quad \tilde{\epsilon}_k = \epsilon_k - \mu - \delta\mu \quad \tilde{V}_Z = V_Z + \delta V_Z$$

OL vs Hybrid

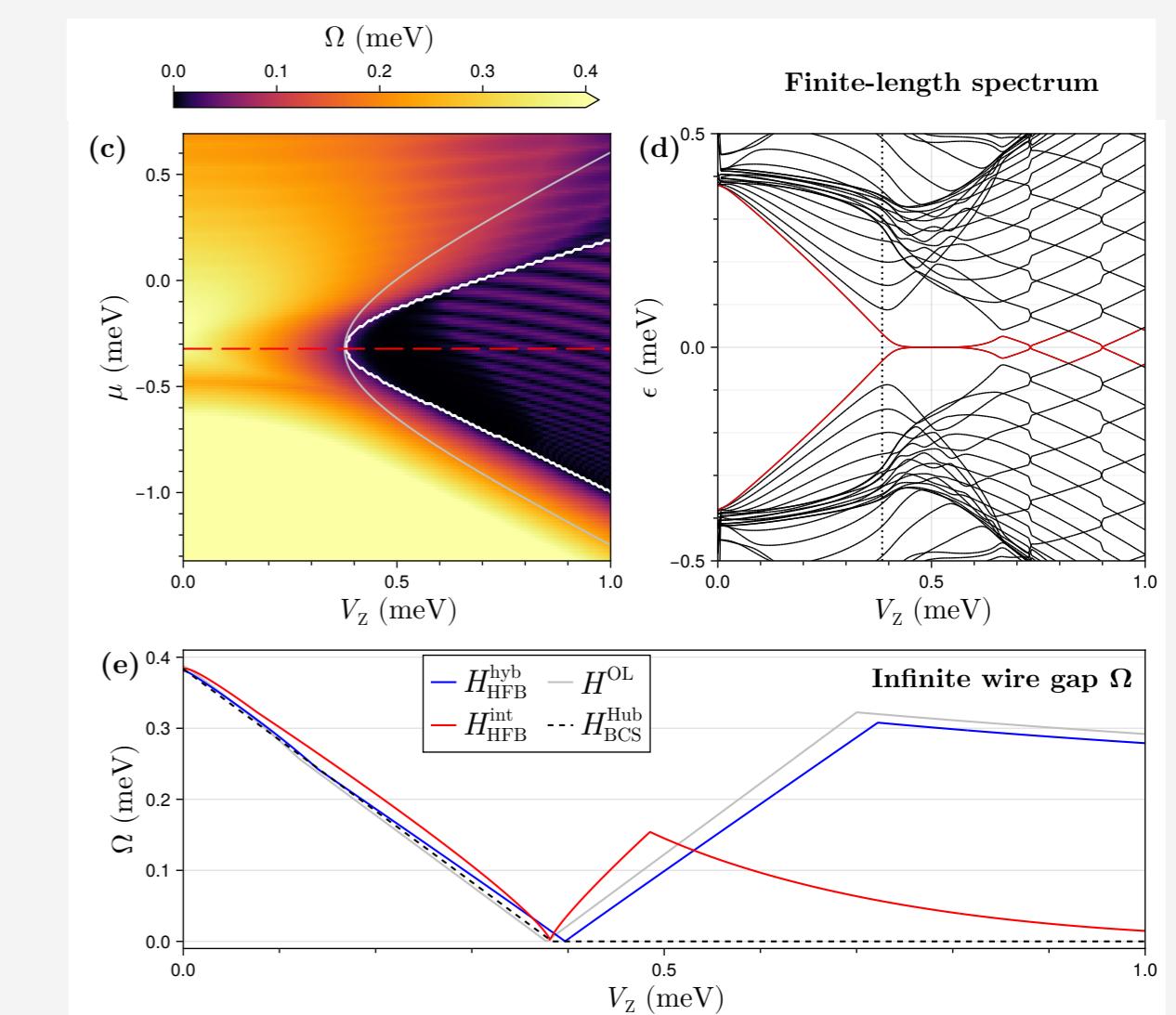
The two models are essentially identical



Intrinsic U<0 nanowire

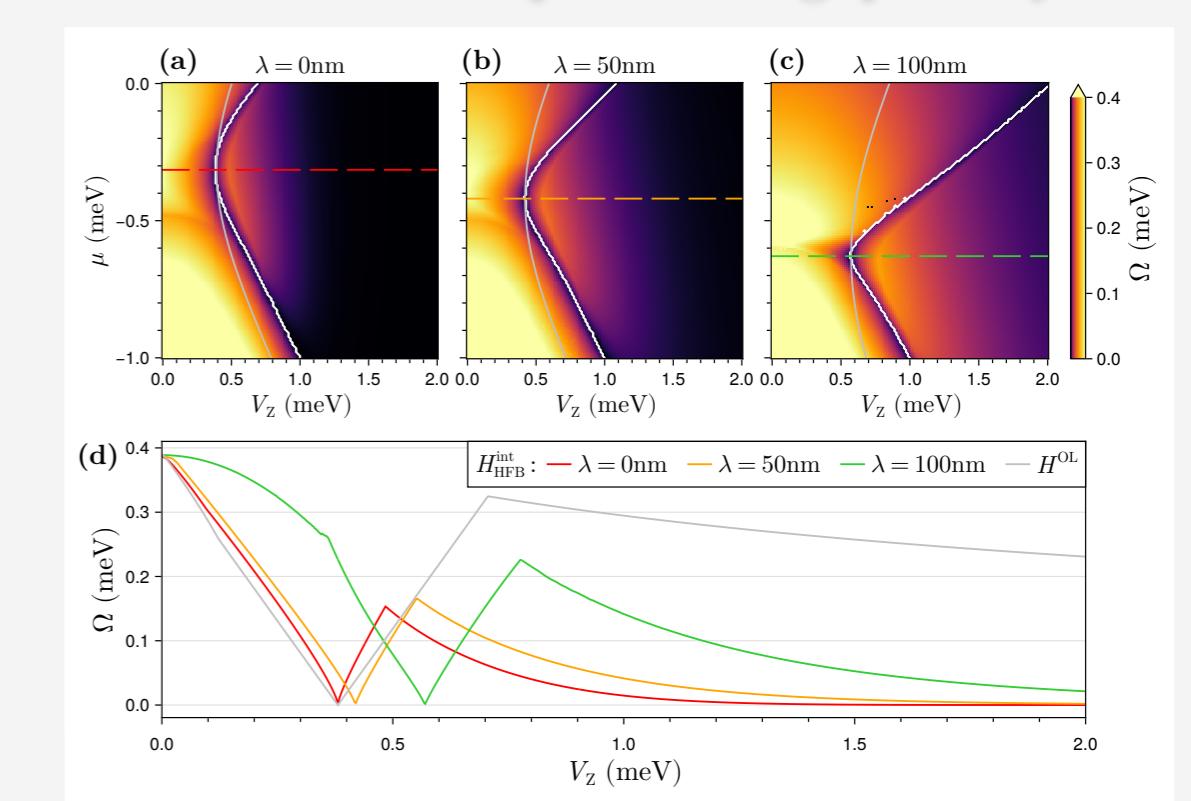
Intrinsic nanowires are qualitatively different

- Topological p-wave region is reduced
- The gap Ω decays exponentially with V_Z
- MZMs are destabilized
- Conclusion: hybrid is better than intrinsic!



Long-range interactions

Increasing the interaction range λ helps, but does not solve the exponential gap decay



[1] Oreg, Refael, von Oppen. *PRL* **105**, 177002 (2010)

[2] Lutchyn, Sau, Das Sarma. *PRL* **105**, 1707001 (2010)

[3] Zhang, Polski et al. *Nature* **613**, 268–273 (2023)