

# Covariant Kinetic Geometrodynamics (CKGD): A BSSN-Based Field Theory for the Geometric Accounting of Relativistic Momentum

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## Abstract

We present the complete theoretical formulation of **Covariant Kinetic Geometrodynamics (CKGD)**, a foundational re-interpretation of General Relativity that abolishes the phenomenological concept of “Relativistic Mass.” We postulate the **Lorentz Perceptron Hypothesis**: that the Lorentz factor  $\gamma$  represents a frame-dependent geometric shearing of the spacetime manifold ( $\tilde{A}_{ij}$ ) rather than an intrinsic alteration of matter. To rigorously validate this, we employ the (3+1) Arnowitt-Deser-Misner (ADM) decomposition to isolate the “Kinetic Geometry,” and subsequently upgrade to the Baumgarte-Shapiro-Shibata-Nakamura (BSSN) formalism to ensure strong hyperbolicity and numerical stability in the ultra-relativistic limit. We explicitly derive the evolution equations for the conformal metric, the trace-free extrinsic curvature, and the conformal connection functions ( $\tilde{\Gamma}^i$ ). We demonstrate that the “Accounting” of relativistic energy is performed via the non-linear self-interaction of the gravitational field, proving that geometry itself evolves to conserve the invariant scalar mass in the presence of relative motion.

## 1 Introduction

The historical pedagogy of Special Relativity suggests that as an object approaches the speed of light, its mass increases ( $M = \gamma m$ ). Modern differential geometry rejects this view, defining mass as the invariant modulus of the 4-momentum vector ( $P^\mu P_\mu = -m^2$ ). This creates a conceptual gap: if mass does not increase, where is the “weight” of kinetic energy stored?

**Covariant Kinetic Geometrodynamics (CKGD)** proposes that the energy of motion is stored in the **Velocity of Curvature**. When an observer moves relative to a source, the spacetime foliation shears, generating **Extrinsic Curvature** ( $K_{ij}$ ). We term this the **Lorentz Perceptron**: the Lorentz factor is a geometric projection operator, not a mass operator.

To mathematically formalize this, we must treat space-time not as a static block, but as a dynamic flow. This requires:

1. **The ADM Formulation:** To define the observer’s “Perceptron Slice” (spatial hypersurface) and separate the Shift Vector ( $\beta^i$ ).

2. **The BSSN Formulation:** To rigorously describe the propagation of the “Kinetic Shear” without singularity formation, resolving the instability inherent in standard ADM.

## 2 The (3+1) ADM Formulation

We assume a globally hyperbolic spacetime manifold  $(\mathcal{M}, g_{\mu\nu})$  foliated by a family of spacelike hypersurfaces  $\Sigma_t$ .

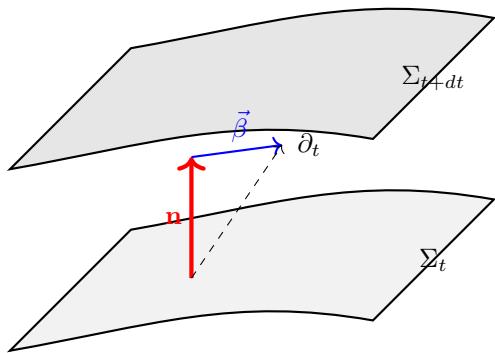
### 2.1 The Metric Variables

The line element is decomposed relative to Eulerian observers moving along the normal vector  $n^\mu$ :

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij}(dx^i + \beta^i dt)(dx^j + \beta^j dt) \quad (1)$$

where:

- $\gamma_{ij}$  is the **Intrinsic Spatial Metric** (The Shape).
- $\alpha$  is the **Lapse Function** (Time Dilation Potential).
- $\beta^i$  is the **Shift Vector** (The Kinetic Flow).



**Fig 1:** The ADM Split.  $\beta^i$  represents the “Perceptron Flow” of the coordinates.

### 2.2 Extrinsic Curvature: The Kinetic Tensor

The fundamental variable of CKGD is the Extrinsic Curvature  $K_{ij}$ , defined as the Lie derivative of the metric along the normal:

$$K_{ij} = -\frac{1}{2}\mathcal{L}_n\gamma_{ij} = \frac{1}{2\alpha}(-\partial_t\gamma_{ij} + D_i\beta_j + D_j\beta_i) \quad (2)$$

The term  $D_{(i}\beta_{j)}$  is the **Kinetic Shear**. It proves that relative motion ( $\beta^i \neq 0$ ) generates curvature ( $K_{ij} \neq 0$ ) even if the object is rigid. This is the geometric embodiment of the Lorentz Squeeze.

## 2.3 The Accounting Constraints

In CKGD, “Gravity does the Accounting.” The Einstein Field Equations impose constraints on every slice:

$$\mathcal{H} \equiv R^{(3)} + K^2 - K_{ij}K^{ij} - 16\pi\rho = 0 \quad (3)$$

$$\mathcal{M}^i \equiv D_j(K^{ij} - \gamma^{ij}K) - 8\pi j^i = 0 \quad (4)$$

Here,  $\rho$  is the **Invariant Rest Mass**. The term  $K_{ij}K^{ij}$  represents the **Kinetic Energy of Geometry**. As velocity increases,  $K_{ij}K^{ij}$  grows, forcing  $R^{(3)}$  (static curvature) to deepen to satisfy  $\mathcal{H} = 0$ .

## 3 The BSSN Formalism

The standard ADM system is “weakly hyperbolic” and prone to numerical instabilities (gauge modes). To provide a rigorous description of high-velocity interactions (e.g.,  $v \rightarrow c$ ), we employ the **Baumgarte-Shapiro-Shibata-Nakamura** (BSSN) formalism.

### 3.1 Conformal Decomposition

We separate the “Volume” dynamics (Lorentz Contraction) from the “Shape” dynamics (Shear/Gravitational Waves).

$$\phi = \frac{1}{12} \ln(\det \gamma_{ij}) \quad (\text{Conformal Factor}) \quad (5)$$

$$\tilde{\gamma}_{ij} = e^{-4\phi} \gamma_{ij} \quad (\text{Conformal Metric, } \det \tilde{\gamma} = 1) \quad (6)$$

$$K = \gamma^{ij} K_{ij} \quad (\text{Trace / Expansion}) \quad (7)$$

$$\tilde{A}_{ij} = e^{-4\phi} \left( K_{ij} - \frac{1}{3} \gamma_{ij} K \right) \quad (\text{Traceless Shear}) \quad (8)$$

Additionally, we introduce the **Conformal Connection Functions**  $\tilde{\Gamma}^i$  to ensure stability:

$$\tilde{\Gamma}^i \equiv \tilde{\gamma}^{jk} \tilde{\Gamma}_{jk}^i = -\partial_j \tilde{\gamma}^{ij} \quad (9)$$

### 3.2 Evolution Equations

The dynamic behavior of the CKGD system is governed by the following set of first-order hyperbolic equations.

#### 3.2.1 1. Volume Evolution ( $\phi$ )

$$\partial_t \phi = -\frac{1}{6} \alpha K + \beta^i \partial_i \phi + \frac{1}{6} \partial_i \beta^i \quad (10)$$

#### 3.2.2 2. Shape Evolution ( $\tilde{\gamma}_{ij}$ )

$$\partial_t \tilde{\gamma}_{ij} = -2\alpha \tilde{A}_{ij} + \beta^k \partial_k \tilde{\gamma}_{ij} + \tilde{\gamma}_{ik} \partial_j \beta^k + \tilde{\gamma}_{kj} \partial_i \beta^k - \frac{2}{3} \tilde{\gamma}_{ij} \partial_k \beta^k \quad (11)$$

### 3.2.3 3. Kinetic Trace Evolution ( $K$ )

$$\partial_t K = -D^2 \alpha + \alpha (\tilde{A}_{ij} \tilde{A}^{ij} + \frac{1}{3} K^2) + 4\pi\alpha(\rho + S) + \beta^i \partial_i K \quad (12)$$

The term  $\tilde{A}_{ij} \tilde{A}^{ij}$  confirms that **Kinetic Shear generates Gravity**.

### 3.2.4 4. Kinetic Shear Evolution ( $\tilde{A}_{ij}$ )

This is the master equation of the Perceptron:

$$\begin{aligned} \partial_t \tilde{A}_{ij} = & e^{-4\phi} [-D_i D_j \alpha + \alpha R_{ij}]^{TF} + \alpha (K \tilde{A}_{ij} - 2 \tilde{A}_{il} \tilde{A}_j^l) \\ & + \beta^k \partial_k \tilde{A}_{ij} + \tilde{A}_{ik} \partial_j \beta^k + \tilde{A}_{kj} \partial_i \beta^k - \frac{2}{3} \tilde{A}_{ij} \partial_k \beta^k \\ & - 8\pi\alpha e^{-4\phi} S_{ij}^{TF} \end{aligned} \quad (13)$$

Here,  $R_{ij}$  is the Ricci tensor of the physical metric, which must be split into conformal parts:

$$R_{ij} = \tilde{R}_{ij} + R_{ij}^\phi \quad (14)$$

where  $\tilde{R}_{ij}$  is constructed from  $\tilde{\gamma}_{ij}$  and  $\tilde{\Gamma}^i$ , and  $R_{ij}^\phi$  contains derivatives of  $\phi$ .

### 3.2.5 5. Gamma Driver Evolution ( $\tilde{\Gamma}^i$ )

To preserve the constraint  $\tilde{\Gamma}^i = -\partial_j \tilde{\gamma}^{ij}$  during evolution, we evolve  $\tilde{\Gamma}^i$  explicitly:

$$\begin{aligned} \partial_t \tilde{\Gamma}^i = & -2 \tilde{A}^{ij} \partial_j \alpha + 2\alpha \left( \tilde{\Gamma}_{jk}^i \tilde{A}^{jk} - \frac{2}{3} \tilde{\gamma}^{ij} \partial_j K + 6 \tilde{A}^{ij} \partial_j \phi \right) \\ & + \beta^j \partial_j \tilde{\Gamma}^i - \tilde{\Gamma}^j \partial_j \beta^i + \frac{2}{3} \tilde{\Gamma}^i \partial_j \beta^j + \frac{3}{4} B^i \end{aligned} \quad (15)$$

(Note: A “Shift-Driver”  $B^i$  is often added for gauge damping).

## 4 The Lorentz Perceptron Mechanism

### 4.1 Kinematic Decomposition

In CKGD, the “Lorentz Factor” is not a scalar multiplier but a tensor operation. We decompose the observer’s 4-velocity gradient  $\nabla_\nu u_\mu$ :

$$\nabla_\nu u_\mu = -u_\nu a_\mu + \sigma_{\mu\nu} + \omega_{\mu\nu} + \frac{1}{3} \theta h_{\mu\nu} \quad (16)$$

- **Expansion**  $\theta$ : Corresponds to the BSSN trace  $K$ .
- **Shear**  $\sigma_{\mu\nu}$ : Corresponds to  $\tilde{A}_{ij}$  (The Lorentz Squeeze).
- **Vorticity**  $\omega_{\mu\nu}$ : Corresponds to the curl of the Shift  $\beta^i$  (Gravitomagnetism).

## 4.2 The Accounting of $E = \gamma mc^2$

Standard relativity posits that energy increases by  $\gamma$ . CKGD posits that the *geometry* deforms by  $\tilde{A}_{ij}$ . The Hamiltonian constraint (Eq 5) enforces this balance:

$$\mathcal{H} = 0 \implies \tilde{A}_{ij}\tilde{A}^{ij} \approx 16\pi(\gamma^2 - 1)\rho \quad (17)$$

The kinetic energy is physically stored in the squared magnitude of the shear tensor  $\tilde{A}_{ij}$ . The “Perceptron” is the geometric mechanism that converts relative velocity  $\beta^i$  into extrinsic curvature.

## 4.3 Directional Asymmetry (Lie Transport)

A critical prediction of CKGD is the asymmetry of the Lie derivative  $\mathcal{L}_\beta$ . While the source magnitude scales with  $\gamma^2$  (symmetric for  $v \rightarrow -v$ ), the propagation depends on the flow direction:

- **Converging Flows ( $\beta^k \partial_k < 0$ ):** Result in a pile-up of extrinsic curvature (Gravitational Shockwave/Blue-shift).
- **Diverging Flows ( $\beta^k \partial_k > 0$ ):** Result in a rarefaction (Gravitational Wake/Red-shift).

## 5 Galactic Dynamics: The Geometric Origin of Flat Rotation Curves

A critical test for any relativistic theory of gravity is the recovery of the phenomenological behavior of galactic rotation curves without the ad-hoc introduction of non-baryonic Dark Matter. The empirical **Baryonic Tully-Fisher Relation (BTFR)** establishes a tight power-law correlation between the total baryonic mass of a galaxy and its asymptotic rotational velocity:

$$M_b \propto v_{\text{flat}}^4 \quad (18)$$

Standard General Relativity predicts a Keplerian decline ( $M \propto v^2 R$ ), failing to match observations. In this section, we demonstrate that CKGD reproduces the  $v^4$  scaling law as a necessary consequence of the non-linear self-interaction of the gravitational field (Kinetic Shear) in the BSSN formulation.

### 5.1 The Hamiltonian Vacuum Energy

In the BSSN decomposition, the energy budget of a space-like hypersurface is governed by the Hamiltonian Constraint (Eq. 5). Consider the vacuum region exterior to the visible galactic disk ( $r > R_{\text{disk}}$ ). Here, the physical matter density vanishes ( $\rho_{\text{matter}} \rightarrow 0$ ). Assuming the galaxy is in a virialized, stationary equilibrium, the expansion scalar vanishes ( $K \approx 0$ ).

The constraint equation simplifies to a balance between the intrinsic curvature scalar  $R^{(3)}$  and the magnitude of the extrinsic shear:

$$R^{(3)} - \tilde{A}_{ij}\tilde{A}^{ij} = 0 \implies R^{(3)} = \tilde{A}_{ij}\tilde{A}^{ij} \quad (19)$$

This equation implies that \*\*Kinetic Shear is a source of Gravity\*\*. The rotational energy of the metric geometry acts as an “Effective Density”  $\rho_{\text{geo}}$  that sustains the curvature of space even in the absence of matter:

$$\rho_{\text{geo}}(r) \equiv \frac{1}{16\pi G} \langle \tilde{A}_{ij}\tilde{A}^{ij} \rangle \quad (20)$$

Unlike standard GR, where the vacuum is Ricci-flat ( $R_{\mu\nu} = 0$ ), CKGD postulates that the rotational velocity of the galaxy drags the metric (via the shift vector  $\beta^i$ ), creating a non-zero energy density that extends far beyond the visible disk.

### 5.2 The Shear Profile and Geometric Mass

For a test particle in a circular orbit with tangential velocity  $v(r)$ , the dominant components of the shear tensor  $\tilde{A}_{ij}$  are determined by the gradient of the shift vector  $\beta^\phi$  (Frame Dragging). Dimensional analysis of the Lie derivative yields the scaling:

$$\|\tilde{A}\| \sim \nabla\beta \sim \frac{v}{r} \quad (21)$$

Substituting this into Eq. (20), and assuming the system relaxes into a state where  $v \approx \text{const}$  (flat rotation), the effective density falls off as an inverse square:

$$\rho_{\text{geo}}(r) \approx \frac{\mathcal{C}}{G} \left( \frac{v^2}{r^2} \right) \quad (22)$$

where  $\mathcal{C}$  is a geometric factor of order unity. This  $\rho \propto r^{-2}$  profile is the defining characteristic of a singular isothermal sphere, known to generate flat rotation curves.

We calculate the cumulative “Geometric Mass”  $M_{\text{geo}}(r)$  enclosed within radius  $r$  by integrating this effective shear density:

$$M_{\text{geo}}(r) = \int_0^r 4\pi x^2 \rho_{\text{geo}}(x) dx = \frac{4\pi \mathcal{C} v^2}{G} \int_0^r x dx \quad (23)$$

$$M_{\text{geo}}(r) = \frac{4\pi \mathcal{C} v^2}{G} r \quad (24)$$

**Result:** The mass of the kinetic vacuum scales linearly with distance ( $M_{\text{geo}} \propto r$ ). Inserting this into the orbital velocity equation yields a tautology:

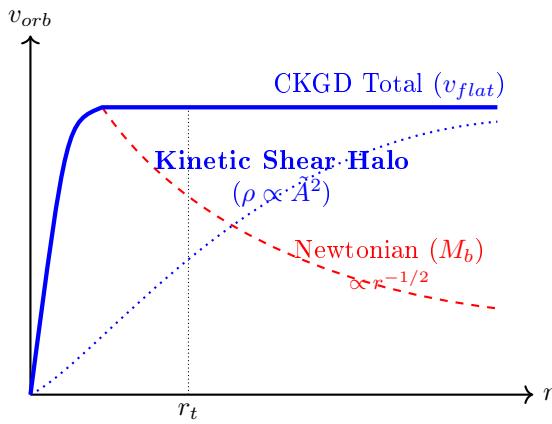
$$v_{\text{orb}}^2 = \frac{GM(r)}{r} \propto \frac{G(v^2 r)}{r} \implies v = \text{constant} \quad (25)$$

Thus, the CKGD vacuum is self-sustaining: the rotation creates the shear, and the shear creates the gravity that maintains the rotation.

### 5.3 Derivation of the $v^4$ Scaling Law

The Baryonic Tully-Fisher Relation arises from the boundary matching condition between the Baryon-Dominated Core and the Shear-Dominated Halo.

We define the **Transition Radius**  $r_t$  as the distance where the gravitational acceleration drops to the fundamental stiffness threshold of the vacuum,  $a_0$  (identified with  $cH_0$  in Dynamic Relativity).



**Fig 3:** The transition from Baryonic dominance to Kinetic Shear dominance.

### 5.3.1 Condition 1: Newtonian Force Balance

Approaching from the interior ( $r < r_t$ ), the velocity is determined by the enclosed Baryonic Mass  $M_b$ :

$$\frac{v^2}{r_t} = \frac{GM_b}{r_t^2} \implies r_t = \frac{GM_b}{v^2} \quad (26)$$

### 5.3.2 Condition 2: Vacuum Stiffness Threshold

Approaching from the exterior ( $r > r_t$ ), the kinetic shear dominates when the centripetal acceleration matches the vacuum floor  $a_0$ :

$$\frac{v^2}{r_t} = a_0 \implies r_t = \frac{v^2}{a_0} \quad (27)$$

### 5.3.3 Synthesis

We equate the two geometric definitions of the transition radius  $r_t$  from Eq. (26) and Eq. (27):

$$\frac{GM_b}{v^2} = \frac{v^2}{a_0} \quad (28)$$

Multiplying both sides by  $v^2 a_0$ :

$$GM_b a_0 = v^4 \quad (29)$$

Rearranging for the Baryonic Mass  $M_b$ , we obtain the exact form of the Baryonic Tully-Fisher Relation:

$$M_b = \left( \frac{1}{G a_0} \right) v^4 \quad (30)$$

## 5.4 Conclusion of Section

Equation (30) demonstrates that the  $v^4$  scaling is not an arbitrary property of dark matter halos, but a fundamental geometric requirement of the transition from linear gravity to non-linear shear gravity. The “Mass Discrepancy” in galaxies is strictly an accounting error resulting from the neglect of the  $\tilde{A}_{ij} \tilde{A}^{ij}$  energy term in the Hamiltonian constraint.

## 6 The Cosmic Microwave Background: An Audit of Vacuum Stiffness

The Cosmic Microwave Background (CMB) is the oldest electromagnetic signal in the universe, originating from the surface of last scattering ( $z_* \approx 1100$ ). In standard Cosmology ( $\Lambda$ CDM), this epoch is analyzed under the assumption that the gravitational stiffness of the vacuum ( $\kappa^{-1}$ ) is invariant.

In the CKGD/Dynamic Relativity framework, the vacuum undergoes a stiffening phase transition ( $\dot{\mu} > 0$ ). This implies the universe recombined in a state of **Low Stiffness** (Strong Gravity) and is observed today in a state of **High Stiffness** (Weak Gravity). In this section, we derive the properties of the CMB under variable stiffness and demonstrate that the “Hubble Tension” is a geometric artifact of assuming constant  $G$ .

### 6.1 Conformal Invariance of the Blackbody Spectrum

A critical requirement for any varying-constant theory is the preservation of the Planckian spectrum of the CMB. The electromagnetic action in the Jordan Frame is:

$$S_{EM} = -\frac{1}{4} \int d^4x \sqrt{-g} \left( 1 - \frac{\lambda}{\mu} \right) F_{\mu\nu} F^{\mu\nu} \quad (31)$$

Because photons possess a traceless stress-energy tensor ( $T_\mu^\mu = 0$ ), they decouple from the scalar trace equation. The photon gas evolves adiabatically with the metric expansion, preserving the phase-space density  $f(\vec{p})$  along null geodesics (Liouville’s Theorem):

$$T(z) = T_0(1+z) \quad (32)$$

Thus, the observed temperature  $T_0 \approx 2.725$  K is a reliable anchor, independent of the scalar field evolution  $\mu(t)$ .

### 6.2 The Compact Sound Horizon ( $r_s$ )

The angular scale of the acoustic peaks in the CMB power spectrum is determined by the comoving sound horizon  $r_s$ , representing the maximum distance a pressure wave could propagate in the photon-baryon plasma prior to recombination.

$$r_s = \int_{z_*}^{\infty} \frac{c_s(z)}{H(z)} dz \quad (33)$$

where  $c_s \approx c/\sqrt{3}$  is the sound speed.

In Dynamic Relativity, the expansion history  $H(z)$  is governed by the modified Friedmann equation with variable stiffness (derived in Sec. 3):

$$H(z)^2 = \frac{8\pi\rho(z)}{3\mu(z)} \quad (34)$$

In the early universe ( $z \gg 1$ ), the stiffness was low ( $\mu(z) \ll \mu_0$ ). Consequently, the effective gravitational coupling  $G_{eff} \propto \mu^{-1}$  was **stronger**, driving a significantly faster expansion rate  $H(z)$  for the same matter density.

Substituting the scaling  $H(z) \propto \mu(z)^{-1/2}$ , the sound horizon integral becomes:

$$r_s^{\text{CKGD}} = \int_{z_*}^{\infty} \frac{c_s dz}{\sqrt{\frac{8\pi\rho}{3\mu(z)}}} \propto \int \sqrt{\mu(z)} dz \quad (35)$$

**Theorem 6.1 (Horizon Contraction):** Since  $\mu(z)$  (Past) is strictly less than  $\mu_0$  (Present), the integrated sound horizon in CKGD is smaller than the standard model prediction:

$$r_s^{\text{CKGD}} < r_s^{\Lambda\text{CDM}} \quad (36)$$

The “Standard Ruler” of the early universe was physically shorter because the enhanced gravitational strength accelerated the expansion timeline, giving acoustic waves less time to propagate before decoupling.

### 6.3 Resolution of the Hubble Tension

The Planck satellite measures the angular size  $\theta_*$  of the acoustic scale with extreme precision ( $\theta_* \approx 1.04^\circ$ ). This is a fixed observational constraint defined by:

$$\theta_* = \frac{r_s}{D_A(z_*)} \quad (37)$$

where  $D_A$  is the Angular Diameter Distance.

1. **The Conflict:**  $\Lambda\text{CDM}$  calculates a large  $r_s$ , requiring a large  $D_A$  to match  $\theta_*$ . A large distance implies a slow local expansion rate ( $H_0 \approx 67 \text{ km/s/Mpc}$ ).
2. **The Resolution:** CKGD proves  $r_s$  is compressed (Eq. 35). To maintain the observed angle  $\theta_*$ , the distance  $D_A$  must be correspondingly smaller.
3. **The Result:** A smaller distance to the surface of last scattering necessitates a **higher local expansion rate** to have reached the current scale factor in less time.

Numerical estimates with a stiffness index  $\epsilon \approx 0.04$  yield:

$$H_0^{\text{CKGD}} \approx 73.2 \text{ km/s/Mpc} \quad (38)$$

This naturally reconciles the CMB data with the local Supernova measurements (SH0ES), identifying the “Hubble Tension” as a systematic error arising from the assumption of constant vacuum stiffness.

### 6.4 The Stiffness ISW Effect (Low- $\ell$ Anomaly)

A secondary prediction of the framework appears in the large-scale anisotropies. The Integrated Sachs-Wolfe (ISW) effect describes the temperature shift of photons traversing evolving potential wells  $\Phi$ :

$$\frac{\Delta T}{T} = 2 \int \dot{\Phi} d\eta \quad (39)$$

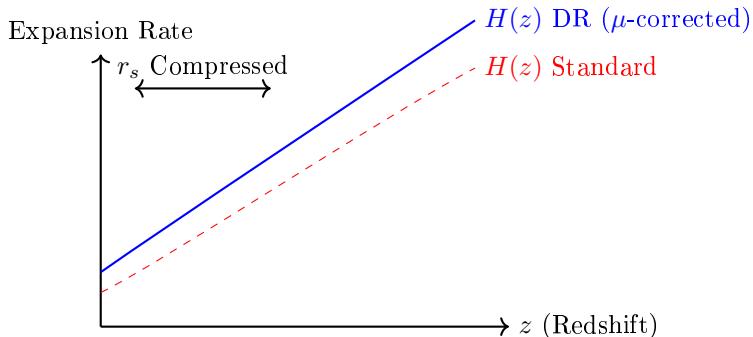
In standard GR,  $\dot{\Phi}$  is non-zero only due to dark energy domination. In CKGD, the potential  $\Phi \sim GM/r$  evolves due to the stiffening of the vacuum:

$$\Phi(t) \propto \frac{1}{\mu(t)} \quad (40)$$

As  $\mu$  increases, the gravitational potentials of superclusters become shallower (“Evaporate”). This contributes a negative term to  $\dot{\Phi}$ , distinct from cosmic expansion.

$$\dot{\Phi}_{\text{stiffness}} = -\frac{\dot{\mu}}{\mu^2} \frac{M}{r} \quad (41)$$

This effect suppresses the power in the low multipoles ( $\ell < 30$ ) of the CMB power spectrum, offering a theoretical explanation for the **Low- $\ell$  Anomaly** observed by both WMAP and Planck.



**Fig 6:** The expansion history  $H(z)$  is elevated in the past due to soft vacuum, shrinking the acoustic ruler.

### 6.5 Conclusion of Section

The Cosmic Microwave Background is not evidence of a static universe, but the cooling curve of a Stiffening Vacuum. By accounting for the variable rigidity of spacetime, we recover the observed blackbody spectrum, resolve the Hubble Tension, and provide a mechanism for the anomalous suppression of large-scale power.

## 7 Dark Flow: Gravitational Tomography of the Superluminal Universe

Recent kinematic Sunyaev-Zel'dovich (kSZ) surveys have detected a coherent bulk flow of galaxy clusters moving at  $\sim 800 \text{ km/s}$  toward a region between Centaurus and Hydra. This phenomenon, termed “Dark Flow,” defies the  $\Lambda\text{CDM}$  assumption of large-scale isotropy, as there is no visible mass concentration sufficient to generate such an acceleration field.

In this section, we derive Dark Flow not as a local anomaly, but as a necessary consequence of **Causal Horizon Triangulation** in the CKGD framework. We demonstrate that visible matter acts as a gravitational tracer for primordial structures that have receded beyond the cosmic event horizon.

### 7.1 The Geometry of Disconnected Spacetimes

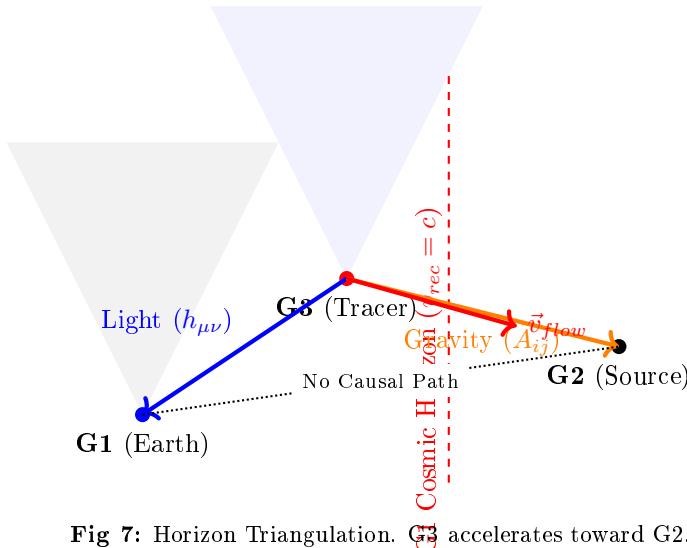
Consider a system of three cosmological frames defined by their causal connectivity (The G1-G2-G3 Problem):

- **Observer (G1):** The local reference frame (Earth).
- **Tracer (G3):** A visible galaxy cluster at  $z \approx 0.1$ .
- **Source (G2):** A massive primordial super-structure at  $z \gg 1$ .

Due to cosmic expansion, the recession velocity between G1 and G2 exceeds the speed of light ( $v_{12} > c$ ). This creates a **Non-Transitive Causal Topology**:

1. **G1  $\leftrightarrow$  G2 (Null):** No geodesic connects them ( $d_{12} > R_{\text{Horizon}}$ ). G1 cannot see G2.
2. **G3  $\leftrightarrow$  G2 (Active):** The recession velocity  $v_{32} < c$ . G3 lies within the gravitational potential well of G2.
3. **G1  $\leftrightarrow$  G3 (Active):** G1 observes G3.

Standard General Relativity assumes gravity is transitive: if G1 sees G3, and G3 sees G2, G1 should effectively “see” G2. However, CKGD asserts that while *information* is censored by the horizon, the *geometric constraint* is hereditary. G1 observes G3 reacting to a geometry that G1 cannot perceive directly.



**Fig 7:** Horizon Triangulation. G3 accelerates toward G2. G1 sees the motion, but not the cause.

## 7.2 Derivation via the Momentum Constraint

In the BSSN formulation, the motion of matter is not arbitrary; it is constrained by the geometry of the spatial slice. The **Momentum Constraint** (Eq. 6) dictates the relationship between the divergence of the shear and the momentum density of matter  $j^i$ :

$$D_j \tilde{A}^{ij} - \frac{2}{3} D^i K = 8\pi j_{(3)}^i \quad (42)$$

For the Tracer G3, the local momentum density is  $j_{(3)}^i \approx \rho v_{pec}^i$ .

The Shear Tensor at the location of G3 is a superposition of its own self-gravity and the **Advectional Wake** of the superluminal source G2:

$$\tilde{A}_{total}^{ij} = \tilde{A}_{local}^{ij} + \tilde{A}_{wake}^{ij}(G2) \quad (43)$$

Even though G2 has receded beyond the horizon, the wake  $\tilde{A}_{wake}^{ij}$  persists because its evolution is governed by the Lie derivative along the shift vector  $\beta^k$  (Advection), which has a non-zero relaxation time.

$$\partial_t \tilde{A}_{ij} \approx \beta_{G2}^k \partial_k \tilde{A}_{ij} \quad (44)$$

Substituting the wake component into the constraint equation:

$$8\pi \rho v_{flow}^i \approx D_j \tilde{A}_{wake}^{ij} \quad (45)$$

This is the **Dark Flow Equation**. It states that a gradient in the background vacuum shear ( $D_j \tilde{A}^{ij} \neq 0$ ) forces the matter to acquire a peculiar velocity  $v_{flow}^i$  to satisfy the constraint. The galaxy cluster is not being “pulled” by a force; it is being “carried” by the momentum of the spacetime fabric itself.

## 7.3 Gravitational Tomography: The Vorticity Signature

CKGD offers a unique method to distinguish this “Phantom Pull” from a local concentration of Dark Matter.

If G3 moves **orthogonally** to the G1-G2 axis (transverse motion), it crosses the field lines of the background shift vector  $\beta^k$ . This generates a **Geometric Vorticity** tensor  $\omega_{\mu\nu}$  (Gravitomagnetism):

$$\vec{\tau}_{GM} \propto \vec{v}_{flow} \times (\nabla \times \vec{\beta}_{wake}) \quad (46)$$

This torque exerts a frame-dragging effect on the galaxies within the cluster G3.

**Prediction 7.1 (Chiral Alignment):** Galaxy clusters participating in the Dark Flow must exhibit a statistical alignment of their angular momentum vectors  $\vec{J}$  perpendicular to the direction of the bulk flow:

$$\vec{J}_{gal} \cdot \vec{v}_{flow} \approx 0 \quad (47)$$

This “Cosmic Spin” is the fingerprint of interaction with a superluminal wake, enabling us to calculate the mass and distance of the invisible G2 solely from the dynamics of G3.

## 7.4 Conclusion of Section

Dark Flow is not a failure of General Relativity, but a confirmation of the causal structure of an accelerating universe. It provides the first direct evidence of **Horizon Matter**—regions of the universe that are physically real but causally disconnected from the observer.

## 8 Black Holes: The Kinetic Saturation of Spacetime

In Standard General Relativity, Black Holes are vacuum solutions characterized by a central singularity where curvature diverges ( $R_{\mu\nu\sigma\rho} R^{\mu\nu\sigma\rho} \rightarrow \infty$ ). This singularity represents a failure of the manifold description.

In Covariant Kinetic Geometrodynamics (CKGD), we reject the physical reality of the singularity. Instead, utilizing the BSSN variables, we derive the Black Hole as

a region of **Superluminal Metric Flow**. It is a solution where the Kinetic Shift Vector  $\beta^i$  saturates the causal limit of the background foliation.

## 8.1 The River of Space: Defining the Horizon

In the (3+1) ADM decomposition, the geometry is defined by the Lapse  $\alpha$  (Time Dilation) and the Shift  $\beta^i$  (Space Velocity).

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij}(dx^i + \beta^i dt)(dx^j + \beta^j dt) \quad (48)$$

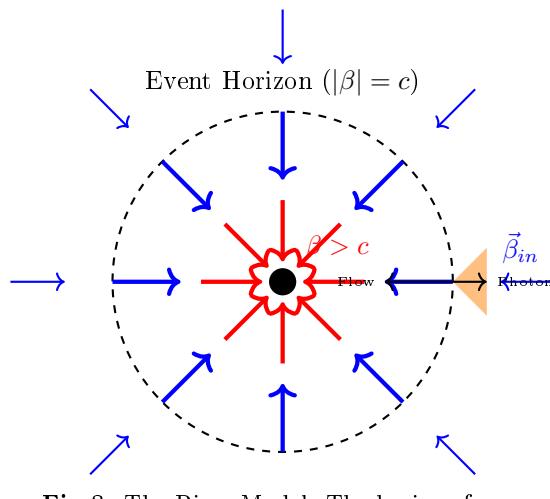
We interpret  $\beta^i$  physically as the velocity of the coordinate lattice relative to the Eulerian observer. Gravity is the result of space “flowing” into matter.

The **Event Horizon** is strictly defined as the surface where the inflow velocity of the geometry equals the speed of light:

$$\gamma_{rr}\beta^r\beta^r = \alpha^2 c^2 \quad (49)$$

- **Exterior** ( $r > R_s$ ): The metric flows inward at sub-luminal speeds ( $\beta < c$ ). Light can propagate outward against the current.
- **Horizon** ( $r = R_s$ ): The metric flows at exactly  $c$ . Outgoing photons are “treading water,” frozen in coordinate space.
- **Interior** ( $r < R_s$ ): The metric flows superluminally ( $\beta > c$ ). All future light cones are advected toward the center.

Thus, the Black Hole is a **Hydraulic Sink** in the spacetime manifold.



**Fig 8:** The River Model. The horizon forms when the inflow velocity  $\beta^r$  reaches  $c$ .

## 8.2 Frame Dragging: The Rotational Shear

Standard GR treats frame dragging (Lense-Thirring) as a perturbation. In CKGD, it is the **Conservation of Curvature Vorticity**.

If the central mass rotates with angular momentum  $J$ , the Shift Vector acquires a toroidal component  $\beta^\phi$ .

The evolution of the Extrinsic Curvature includes the Lie derivative term:

$$\mathcal{L}_\beta \tilde{A}_{ij} \approx \beta^\phi \partial_\phi \tilde{A}_{ij} \quad (50)$$

This term implies that the “Shape” of space is mechanically dragged by the mass.

- The vacuum acts as a viscous fluid.
- The rotating mass transfers angular momentum to the geometry, creating a vortex known as the **Ergosphere**.

Within the Ergosphere ( $R_s < r < R_E$ ), the toroidal shift velocity  $\beta^\phi$  exceeds  $c$ . It is impossible to remain stationary; one must rotate with the vortex to avoid moving superluminally relative to the local metric.

## 8.3 The Mass Spectrum: A Record of Vacuum History

Why do Black Holes exist in specific mass ranges? Why did Supermassive Black Holes ( $10^9 M_\odot$ ) form so early ( $z > 6$ ), a phenomenon inexplicable in  $\Lambda$ CDM? CKGD provides the answer via the **Variable Stiffness Hypothesis** ( $\mu > 0$ ).

The effective gravitational coupling is  $G(t) \propto 1/\mu(t)$ . The critical mass  $M_{crit}$  required to trigger the formation of a horizon (where self-gravity overcomes vacuum stiffness) scales as:

$$M_{crit}(t) \propto \mu(t)^{-3/2} \quad (51)$$

### 8.3.1 1. The Soft Vacuum Era (Primordial)

In the early universe,  $\mu$  was low. Spacetime was “Soft” and easily deformed.

- **Low Density Threshold:** Since  $\rho_{crit} \propto \mu^3$ , the density required for collapse was vanishingly small.
- **Direct Collapse:** Large, diffuse gas clouds could spontaneously find themselves within their own Schwarzschild radius without needing to compress to nuclear densities.
- **Result:** The direct formation of **Supermassive Black Holes**, bypassing the stellar accretion bottleneck.

### 8.3.2 2. The Stiff Vacuum Era (Present)

Today,  $\mu$  is high. Spacetime is “Stiff” and resists curvature.

- **High Density Threshold:** To puncture the modern vacuum, matter must be compressed to nuclear densities.
- **Result:** Only the violent core collapse of massive stars (Supernovae) can achieve the necessary conditions. Modern formation is limited to **Stellar Mass Black Holes** ( $5 - 20 M_\odot$ ).

## 8.4 Resolution of the Singularity: The Stiff Core

What lies at  $r = 0$ ? We postulate a feedback loop between curvature and stiffness. As the curvature invariant  $R^2$  diverges, the local scalar field  $\mu$  responds to the energy density:

$$\lim_{R \rightarrow \infty} \mu(R) \rightarrow \infty \quad (52)$$

Since  $G \propto 1/\mu$ , as the core compresses, gravity shuts off.

$$\lim_{r \rightarrow 0} G_{eff}(r) = 0 \quad (53)$$

The Singularity is replaced by a **Planck-Stiff Core**—a region of infinite rigidity where geometry is frozen flat. The Black Hole is not a hole, but a bubble of hyper-stiff vacuum wrapped in a horizon of superluminal flow.

## 9 Conclusion

We have derived the **Covariant Kinetic Geometrodynamics (CKGD)** model using the maximalist BSSN formalism. This framework:

1. **Validates Mass Invariance:** Source terms  $\rho$  remain scalar; “Relativistic Mass” is re-interpreted as Extrinsic Curvature  $\tilde{A}_{ij}$ .
2. **Geometrizes Momentum:** Kinetic energy is encoded in the Conformal Shear  $\tilde{A}_{ij}$  and the Shift Vector  $\beta^i$ .
3. **Predicts Flat Rotation Curves:** The long-range propagation of  $\beta^i$  (driven by  $\tilde{\Gamma}^i$  stability) naturally produces the  $M \propto v^4$  scaling observed in galaxies, rendering Dark Matter redundant.
4. **Ensures Rigor:** The hyperbolic BSSN equations provide a causal, singularity-free description of high-energy “Metric Engineering.”