From Cellular Automata to the Lattice Boltzmann Method

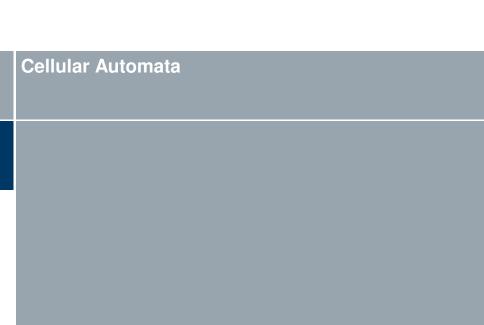
Theresa Pollinger, Marcial Gaißert Ferienakademie Sarntal 2017 September 9, 2017

Outline

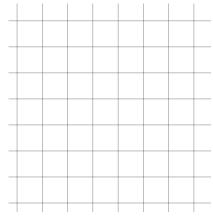
Cellular Automata

Lattice Gas Automata

The Lattice Boltzmann Method

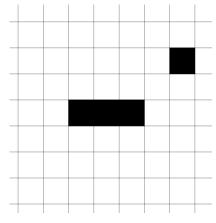


- · regular arrangement of cells



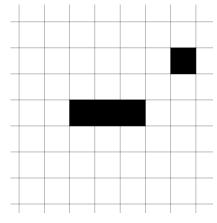
a grid of cells

- · regular arrangement of cells
- · each with discrete state



a grid of cells with state

- · regular arrangement of cells
- · each with discrete state
- updated simultaneously at each step



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- update rules:
 - same for each cell
 - only dependent on

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	2,3	
	> 3	
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Rules of Conway's Game of Life

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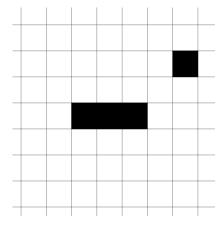
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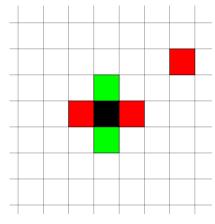
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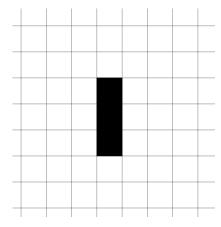
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changes in first step

- regular arrangement of cells
- · each with discrete state
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after first step

advantages

disadvantages

- simple structure
- easy to implemen
- easy to parallelize

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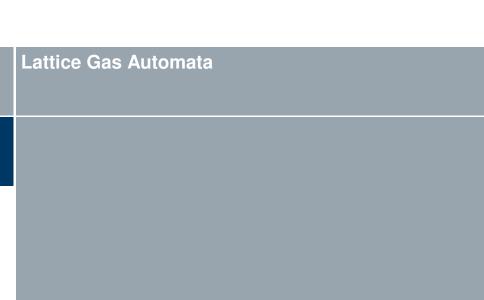
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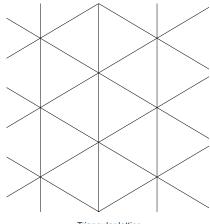
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- Regular arrangement of nodes
 - HPP: square lattice
 - FHP: triangular lattice
- 6 (7) cells per node
- for each step, do:

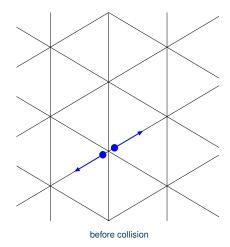
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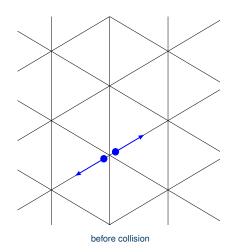


Triangular lattice

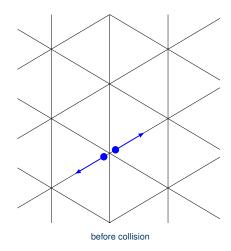
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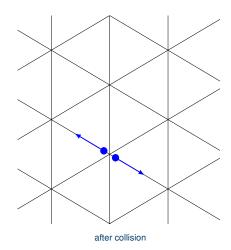
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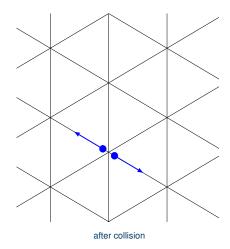
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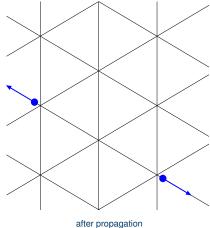
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advantages

- boolean state no rounding
- easy to parallelize
- FHP yields Navier-Stokes

- statistical noise

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- FHP yields Navier-Stokes

- statistical noise
- lack of Galilean invariance

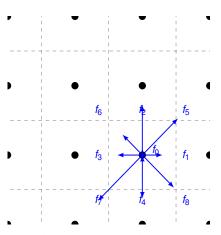


Collide Step

Particle distribution function f_i : discrete in space and time, continuous in probability values

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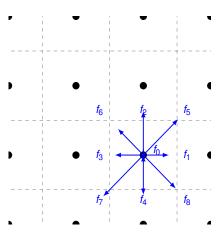
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One collide step for a D2Q9-stencil

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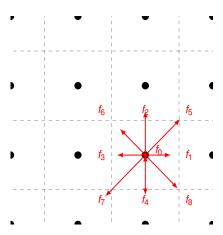
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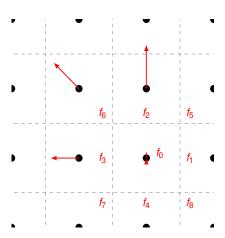
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Particle distribution function f_i : discrete in space and time, continuous in probability values



One stream step for a D2Q9-stencil

...is what we get when we apply both steps:

$$f_i(\mathbf{x} + \mathbf{c}_i \Delta t, t + \Delta t) = f_i(\mathbf{x}, t) + \Omega_i(\mathbf{x}, t)$$

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With collision operator $\Omega_i(\mathbf{x},t)$, e.g.:

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and the equilibrium particle distribution function $f_i^{\text{eq}} = w_i \rho \left(1 + \frac{\vec{e}_i \vec{u}}{c_s^2} + \frac{(\vec{e}_i \vec{u})^2}{2c_s^4} - \frac{\vec{u}^2}{2c_s^2} \right)$

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Cf. Boltzmann Equation

$$\frac{\partial f}{\partial t} + \frac{\mathbf{p}}{m} \cdot \nabla f = \left(\frac{\partial f}{\partial t}\right)_{coll}$$

Moments - the Interesting Outputs in LBM

Density and velocity are given locally

$$\rho(\mathbf{x},t) = \sum_{i} f_{i}(\mathbf{x},t)$$

$$\rho \mathbf{u}(\mathbf{x},t) = \sum_{i} \mathbf{c}_{i} f_{i}(\mathbf{x},t)$$

= Normalization in the three physical dimensions time, space, mass

$$\Delta x^* = 1 = C_x \cdot \Delta x$$
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$$v^* = (\tau - 0.5)c_s^2 = (\frac{1}{\omega} - 0.5)c_s^2, \ c_s^2 = \frac{1}{3}$$

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 \Rightarrow 0.5 < τ ... and not much larger than 1!

advantages

- (moderately) easy to implement
- easy to parallelize
- yields Navier-Stokes (by Chapman-Enskog analysis
- Galilean invariance

- needs parameter fine-tuning
- explicit scheme → may need fine time resolution for stability, compared to other CFD methods

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Thanks for listening.

Any questions?



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