An Introduction to String Theory

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Table of Contents

- Introduction
 - Notation
- The Relativistic String
 - The Relativistic Point Particle
 - Dynamics of a Relativistic String The Nambu-Goto Action
- 3 Covariant Quantization of the Solutions
- Quantization in the Lightcone Gauge

Introduction – Notation

- Metric:
 - \bullet $\eta_{\mu
 u}=\mathsf{diag}(-1,+1,\,\ldots,\,+1)$

The Relativistic Point Particle – The Action

Consider the action of a point particle (with fixed coordinates $X_{\mu}=(t,\vec{x})$ in a given frame):

$$S = -m \int dt \sqrt{1 - \dot{\vec{x}}\dot{\vec{x}}}$$

 \rightarrow not Lorentz-invariant, due to mixture of spacial and temporal coordinates under a Lorentz-transformation Λ .

lacktriangle Consider instead for a generalized coordinate au along the line element:

$$S = -m \int d\tau \sqrt{-\frac{dX^{\mu}}{d\tau}} \frac{dX^{\nu}}{d\tau} \eta_{\mu\nu}$$



Ref: [1]

4 / 11

Remark: S is proportional to the integral over the worldline of the particle

The Relativistic Point Particle – The Action

Description Consider the action of a point particle (with fixed coordinates $X_{\mu}=(t,\vec{x})$ in a given frame):

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 \rightarrow not Lorentz-invariant, due to mixture of spacial and temporal coordinates under a Lorentz-transformation Λ .

Consider instead for a generalized coordinate τ along the line element:

$$S = -m \int d\tau \sqrt{-\frac{dX^{\mu}}{d\tau} \frac{dX^{\nu}}{d\tau} \eta_{\mu\nu}}$$



Ref: [1]

Remark: S is proportional to the integral over the worldline of the particle

The Relativistic Point Particle – Action Symmetries

What are the symmetries of

$$S = -m \int d\tau \sqrt{-\frac{dX^{\mu}}{d\tau} \frac{dX^{\nu}}{d\tau} \eta_{\mu\nu}}$$

► Reparametrization invariance:

Let $\tilde{\tau} = \tilde{\tau}(\tau)$. Then

$$S' = -m \int d\tilde{\tau} \sqrt{-\frac{dX^{\mu}}{d\tilde{\tau}}} \frac{dX^{\nu}}{d\tilde{\tau}} \eta_{\mu\nu} = S$$

ightarrow gauge symmetry of the action ightarrow still D-1 dof!

▶ Poincaré invariance:

Let:

$$X'^{\mu} = \Lambda^{\mu}{}_{\nu}X^{\nu} + c^{\mu}$$

Then

$$S'=S, \qquad ext{as} \qquad \Lambda^{\mu}{}_{
ho}\,\eta_{\mu
u}\,\Lambda^{
u}{}_{\sigma}=\eta_{
ho\sigma}$$

The Relativistic Point Particle – Action Symmetries

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 \rightarrow gauge symmetry of the action \rightarrow still D-1 dof!

Poincaré invariance:

Let:

$$X'^{\mu} = \Lambda^{\mu}{}_{\nu}X^{\nu} + c^{\mu}$$

Then:

$$S' = S,$$
 as $\Lambda^{\mu}{}_{\rho} \eta_{\mu\nu} \Lambda^{\nu}{}_{\sigma} = \eta_{\rho\sigma}$

Dynamics of a Relativistic String – The Nambu-Goto Action

Action of the Relativistic String?

particle
$$\Leftrightarrow$$
 worldline \Leftrightarrow $S = -m \int \underbrace{d\tau \sqrt{-\frac{dX^{\mu}}{d\tau}\frac{dX^{\nu}}{d\tau}\eta_{\mu\nu}}}_{\text{line element }ds}$

closed string
$$\Leftrightarrow$$
 ? \Leftrightarrow

Boundary condition of closed strings living in D dimensions:

$$X^{\mu}(\tau,\sigma) = X^{\mu}(\tau,\sigma+2\pi)$$
 for $\mu = 0, 1, ..., D-1; \ \sigma \in [0;2\pi)$

Shorthand notation: $\sigma^{\alpha}=(\tau,\sigma)$ for $\alpha\in\{0;1\}$

image of a string

Dynamics of a Relativistic String – The Nambu-Goto Action

Action of the Relativistic String?

particle
$$\Leftrightarrow$$
 worldline \Leftrightarrow $S=-m\int d\tau\,\sqrt{-\frac{dX^{\mu}}{d\tau}\frac{dX^{\nu}}{d\tau}\eta_{\mu\nu}}$

closed string
$$\Leftrightarrow$$
 worldsheet \Leftrightarrow

Boundary condition of closed strings living in D dimensions

$$X^{\mu}(\tau,\sigma) = X^{\mu}(\tau,\sigma+2\pi)$$
 for $\mu = 0, 1, ..., D-1; \ \sigma \in [0;2\pi)$

Shorthand notation: $\sigma^{\alpha} = (\tau, \sigma)$ for $\alpha \in \{0, 1\}$

Dynamics of a Relativistic String – The Nambu-Goto Action

Action of the Relativistic String?

particle
$$\Leftrightarrow$$
 worldline \Leftrightarrow $S = -m \int d\tau \sqrt{-\frac{dX^{\mu}}{d\tau} \frac{dX^{\nu}}{d\tau} \eta_{\mu\nu}}$

closed string \Leftrightarrow worldsheet \Leftrightarrow Nambu-Goto/Dirac Action Boundary condition of closed strings living in D dimensions:

$$X^{\mu}(\tau,\sigma)=X^{\mu}(\tau,\sigma+2\pi)\quad\text{for }\mu=0,1,...,D-1;\,\sigma\in[0;2\pi)$$

Shorthand notation: $\sigma^{\alpha}=(\tau,\sigma)$ for $\alpha\in\{0;1\}$ image of a string

How to describe the surface area to construct the action?

$$S \propto \int_{\partial V} d^2x = \int_{\partial V} d^2\sigma \left| \det J \right| \tag{1}$$

For this, first remember that any metric is given by:

$$g_{\alpha\beta} = g(e_{\alpha}, e_{\beta})$$

So the metric on the worldsheet (the pullback metric) can be written as:

$$\gamma_{\alpha\beta} = \underbrace{\frac{\partial X^{\mu}}{\partial \sigma^{\alpha}} \frac{\partial X^{\nu}}{\partial \sigma^{\beta}} \eta_{\mu\nu}}_{(J^{T}\eta J)_{\alpha\beta}}$$

Thus,

$$\det \gamma = \det^2 J \det \eta = -\det^2 J$$
$$|\det J| = \sqrt{-\det \gamma} = \sqrt{-\gamma}$$

- ▶ first item
 - subitem
 - subsubitem
- second item
 - 1 item 1
 - 1.1 subitem 1
 - 1.2 subitem 2
 - 2. item 2
- ► third item

- ▶ first item
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- third item

Covariant Quantization of the Solutions of the Nambu-Goto Action

$$S = \frac{1}{2\pi\alpha'} \int \sqrt{-\det g} \,\partial_{\mu} \tag{2}$$

Quantization of X^{μ} in the Lightcone Gauge

References

[1] David Tong. "Lectures on String Theory". In: arXiv:0908.0333 [hep-th] (Feb. 2012). arXiv: 0908.0333. URL: http://arxiv.org/abs/0908.0333 (visited on 07/13/2020).