

# QCQI Chapter 4 Exercises

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4.2: Let  $x$  be a real number and  $A$  a matrix such that  $A^2 = I$ . Show that

$$\exp(iAx) = \cos(x)I + i \sin(x)A$$

Answer:

$$\exp(iAx) = \sum_{n=1}^{\infty} \frac{1}{n!} (iAx)^n \quad (1a)$$

$$= \sum_{n=1}^{\infty} \frac{1}{(2n)!} (iAx)^{2n} + \frac{1}{(2n+1)!} (iAx)^{2n+1} \quad (1b)$$

$$= I \sum_{n=1}^{\infty} \frac{(-1)^n}{(2n)!} x^n + \frac{1}{(2n+1)!} (iAx)^{2n+1} \quad (1c)$$

$$= I \sum_{n=1}^{\infty} \frac{(-1)^n}{(2n)!} x^n + iA \frac{(-1)^n}{(2n+1)!} x^n \quad (1d)$$

$$= \cos(x)I + i \sin(x)A \quad (1e)$$

4.3: Show that, up to a global phase, the  $\pi/8$  gate satisfies  $T = R_z(\pi/4)$

Answer:

$$\pi/8 \text{ gate} = \exp(i\pi/8) \begin{bmatrix} e^{-i\pi/8} & 0 \\ 0 & e^{-i\pi/8} \end{bmatrix}$$

and

$$R_z(\pi/4) = \begin{bmatrix} e^{-i\pi/8} & 0 \\ 0 & e^{-i\pi/8} \end{bmatrix}$$

$\exp(i\pi/8)$  is just a global phase factor dude.

4.4: Express the Hadamard gate  $H$  as a product of  $R_x$  and  $R_z$  rotations and  $e^{i\varphi}$

Answer:

$$\text{Hadamard} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad (2a)$$

$$R_z(\theta) = \begin{bmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{bmatrix} \quad (2b)$$

$$R_x(\theta) = \begin{bmatrix} \cos \frac{\theta}{2} & -i \sin \frac{\theta}{2} \\ -i \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{bmatrix} \quad (2c)$$

$$R_z(-2\pi) = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad (2d)$$

$$R_x(\pi) = \begin{bmatrix} 0 & -i \\ -i & 0 \end{bmatrix} \quad (2e)$$

$$e^{i \cdot \frac{5\pi}{2}} \cdot R_x(\pi/2) \cdot R_z(\pi/2) \cdot R_x(\pi/2) = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \quad (2f)$$

The insight here is that this series of rotations is equivalent (within a global phase factor) to the Hadamard's more obvious set of rotations: 90 degrees around the y axis followed by 180 degrees around the x axis.

4.5 Prove that  $(\hat{n} \cdot \sigma)^2 = I$  and use this to verify this equation:

$$R_{\hat{n}}(\theta) \equiv \exp(-i\theta \hat{n} \cdot \sigma/2) = \cos\left(\frac{\theta}{2}\right)I - i \sin\left(\frac{\theta}{2}\right)(n_x X + n_y Y + n_z Z)$$

Okay, proof by equational reasoning:

$\hat{n}$  is a unit vector so  $n_z^2 + n_x^2 + n_y^2 = 1$ .

$$(\hat{n} \cdot \sigma)^2 = \left( \begin{bmatrix} 0 & n_x \\ n_x & 0 \end{bmatrix} + \begin{bmatrix} 0 & -in_y \\ -in_y & 0 \end{bmatrix} + \begin{bmatrix} n_z & 0 \\ 0 & -n_z \end{bmatrix} \right)^2 \quad (3a)$$

$$= \left( \begin{bmatrix} n_z & n_x - in_y \\ n_x + in_y & n_z \end{bmatrix} \right)^2 = \begin{bmatrix} n_z^2 + n_x^2 + n_y^2 & 2n_z(n_x - in_y) \\ 2n_z(n_x + in_y) & n_z^2 + n_x^2 + n_y^2 \end{bmatrix} \quad (3b)$$

$$= \begin{bmatrix} 1 & 2n_z(n_x - in_y) \\ 2n_z(n_x + in_y) & 1 \end{bmatrix} \quad (3c)$$

Hermitian operators must be equal to their adjoint and so

$$2n_z(n_x + in_y) = 2n_z(n_x - in_y)$$

We'll call this term "scooby-doo".

The determinant of a rotation matrix must be 1, and so we have

$$1 - \text{scooby-doo}^2 = 1 \quad (4a)$$

$$-\text{scooby-doo}^2 = 0 \quad (4b)$$

$$\text{scooby-doo} = 0 \quad (4c)$$

and therefore

$$\begin{bmatrix} 1 & \text{scooby-doo} \\ \text{scooby-doo} & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = I$$

Having done this convoluted and scooby-dooful proof, I now realize that since all quantum operators must be Hermitian and unitary, then  $R_{\hat{n}}^2$  must trivially be equal to  $I$  by:

$$R_{\hat{n}}^\dagger R_{\hat{n}} = I \text{ by being unitary} \quad (5a)$$

$$R_{\hat{n}}^\dagger = R_{\hat{n}} \text{ by being Hermitian} \quad (5b)$$

$$R_{\hat{n}}^2 = I \quad (5c)$$

4.6: I'm getting kinda tired so I'm not gonna write the full question for this one. It's so long dude. They want me to prove that  $R_{\hat{n}}(\theta)$  rotates a Bloch vector around the  $\hat{n}$  axis.

For a state vector  $a|0\rangle + b|1\rangle$

$$R_n(\theta) = \begin{bmatrix} \cos \frac{\theta}{2} - i \sin \frac{\theta}{2} n_z & i \sin \frac{\theta}{2} (n_x - i n_y) \\ i \sin \frac{\theta}{2} (n_x + i n_y) & \cos \frac{\theta}{2} + i \sin \frac{\theta}{2} n_z \end{bmatrix} \quad (6a)$$

$$R_n \theta (a|0\rangle + b|1\rangle) = e^{-i \frac{\theta}{2} n_z} a|0\rangle + i \sin \frac{\theta}{2} (n_x - i n_y) b|1\rangle + i \sin \frac{\theta}{2} (n_x + i n_y) a|0\rangle + e^{i \frac{\theta}{2} n_z} b|1\rangle \quad (6b)$$

$$\frac{\partial R_n \theta (a|0\rangle + b|1\rangle)}{\partial a|0\rangle} = e^{-i \frac{\theta}{2} n_z} + i \sin \frac{\theta}{2} (n_x + i n_y) \quad (6c)$$

$$\frac{\partial R_n \theta (a|0\rangle + b|1\rangle)}{\partial b|0\rangle} = e^{i \frac{\theta}{2} n_z} + i \sin \frac{\theta}{2} (n_x - i n_y) \quad (6d)$$

$$\frac{\partial R_n \theta (a|0\rangle + b|1\rangle)}{\partial a|0\rangle} - \frac{\partial R_n \theta (a|0\rangle + b|1\rangle)}{\partial b|0\rangle} = 2i \sin \frac{\theta}{2} n_z - 2 \sin \frac{\theta}{2} n_y \quad (6e)$$

which shows up in the polar coordinates basis  $(\theta, \varphi)$  as

$$a = \cos \frac{\theta}{2} \text{ by definition of the Bloch vector} \quad (7a)$$

$$b = e^{i\varphi} \sin \frac{\theta}{2} \text{ by definition of the Bloch vector} \quad (7b)$$

$$\theta = 2 \arccos a \quad (7c)$$

$$\varphi = -i \ln \frac{b}{\sin \frac{\theta}{2}} \quad (7d)$$