EE 160 SIST, Shanghai Tech

Quadratic Integral Forms

Introduction

Analyzing quadratic integral forms

Quadratic integral forms with quadratic end term

Reverse Lyapunov differential equations

Boris Houska 9-1

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We focus on homogeneous linear differential equations of the form

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 with $x(0) = x_0$.

We are interested in analyzing quadratic integrals,

$$q_0 = \int_0^T x(\tau)^{\mathsf{T}} Q(\tau) x(\tau) \, \mathrm{d}\tau \ .$$

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Simple examples

- For $Q(t) = \frac{1}{T}I$ the term $q_0 = \frac{1}{T}\int_0^T x(\tau)^\intercal x(\tau)\,\mathrm{d}\tau$ is the average of the square of the Euclidean norm of the trajectory x.
 - If Q(t) is a positive definite function, the term

$$q_0 = ||x||_{L_2[0,T]}^2 = \int_0^T x(\tau)^T Q(\tau) x(\tau) d\tau$$

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Assume x(t) = I(t) is the current in an RC-circuit

$$\dot{I}(t) = -\frac{1}{RC}I(t) \qquad \text{with} \qquad I(0) = \frac{V_0}{R} \; . \label{eq:interpolation}$$

The power that is consumed by the resistor at time t is given by $P(t)=RI(t)^2$. Quadratic performance measure

$$q_0 = \int_0^T x(\tau)^T Q(\tau) x(\tau) dt = \int_0^T RI(\tau)^2 d\tau$$

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In this example, we can work out q explicitly:

$$q_0 \; = \; \int_0^T RI(\tau)^2 \, \mathrm{d}t \; = \; \int_0^T \frac{V_0^2}{R} e^{-\frac{2t}{RC}} \, \mathrm{d}t \; = \; \frac{CV_0^2}{2} \left[1 - e^{-\frac{2T}{RC}}\right] \; .$$

Overall energy consumption for $T \to \infty$ is given by

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Linear integral forms can be written as

$$\int_0^T g(\tau)^{\mathsf{T}} x(\tau) \,\mathrm{d}\tau$$

with $g: \mathbb{R} \to \mathbb{R}^{n_x}$ being a vector-valued function.

Important: Linear integral forms can be re-written as quadratic integral forms!

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Important: Linear integral forms can be re-written as quadratic integral forms!

Main idea: consider the augmented state $y(t) = [x(t)^{\mathsf{T}}, 1]^{\mathsf{T}}$,

$$\dot{y}(t) = \left(\begin{array}{cc} A(t) & 0 \\ 0 & 0 \end{array} \right) \quad \text{with} \quad y(0) = \left(\begin{array}{c} x_0 \\ 1 \end{array} \right) \; .$$

With
$$Q(t)=\frac{1}{2}\begin{pmatrix}0&g(t)\\g(t)^{\mathsf{T}}&0\end{pmatrix}$$
 we have
$$\int_0^Tg(\tau)^{\mathsf{T}}x(\tau)\,\mathrm{d}\tau=\int_0^Ty(\tau)^{\mathsf{T}}Q(\tau)y(\tau)\,\mathrm{d}\tau\;.$$

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$$\int_0^T g(\tau)^\intercal x(\tau) \,\mathrm{d}\tau = \int_0^T y(\tau)^\intercal Q(\tau) y(\tau) \,\mathrm{d}\tau \;.$$

Integrals over bi-linear terms

Integrals over products of components of x, e.g.,

$$q_0 = \int_0^T x_1(\tau) x_2(\tau) \,\mathrm{d}\tau$$

are called bi-linear forms. They can be written as quadratic forms, e.g.,

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Short Quiz

If you see a term like

$$q_0 = \int_0^T \left\{ x_1(\tau) x_2(\tau) + x_4(\tau)^2 + x_3(\tau) + \sin(\tau) + x_2(\tau) \cos(\tau) \right\} d\tau ,$$

what can you say about it?

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Let us consider the auxiliary function

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- defined for all $t \leq T$ and satisfies $q(0) = q_0$
- If G denotes the fundamental solution,

$$q(t) = \int_{t}^{T} x(\tau)^{\mathsf{T}} Q(\tau) x(\tau) \, \mathrm{d}\tau$$
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- $ullet \ q(t)$ can be interpreted as a quadratic form in x(t).
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- If $Q(\tau)$ is symmetric [positive semi-definite] then P(0) is symmetric [positive semi-definite].

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Quadratic integral form with end term

Notation:

$$q(t) = \int_{t}^{T} x(\tau)^{\mathsf{T}} Q(\tau) x(\tau) d\tau + x(T)^{\mathsf{T}} Q_{T} x(T) ,$$

• $Q_T \in \mathbb{R}^{n_x \times n_x}$ is the weighting matrix of the end term.

Quadratic integral form with end term

The function q can be written in the form

$$q(t) = x(t)P(t)x(t)$$

with

$$P(t) = \int_t^T G(\tau, t)^{\mathsf{T}} Q(\tau) G(\tau, t) d\tau + G(T, t)^{\mathsf{T}} Q_T G(T, t) .$$

Proof (direct verification)

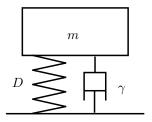
$$q(t) = \int_{t}^{T} x(\tau)^{\intercal} Q(\tau) x(\tau) d\tau + x(T)^{\intercal} Q_{T} x(T)$$

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$$= x(t)^{\intercal} P(t) x(t) .$$

Kinetic energy of a spring

Consider a spring-damper system (no damping, $\gamma=0$)

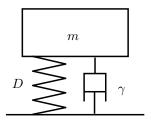


ullet States: elongation s(t) and velocity v(t)

Differential equation:
$$\begin{cases} \dot{s}(t) &= v(t) \\ \dot{v}(t) &= -\frac{D}{m}s(t) \end{cases}$$

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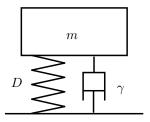
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Kinetic energy:

$$E_{\rm kin}(T) = \frac{1}{2} m \left[v(T) \right]^2$$

Can be written as

$$E_{\rm kin}(T) = x(T)^{\mathsf{T}} Q_T x(T)$$
 with $Q_T = \begin{pmatrix} 0 & 0 \\ & & \\ 0 & \frac{m}{2} \end{pmatrix}$

where $x(T) = (s(T), v(T))^{\mathsf{T}}$ is the state vector.

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where $x(T) = (s(T), v(T))^{\mathsf{T}}$ is the state vector.

We work out P(0) explicitly:

$$P(0) = G(T,0)^{\mathsf{T}} Q_T G(T,0)$$

$$= \frac{m}{2} \begin{pmatrix} \omega^2 \sin(\omega T)^2 & -\omega \sin(\omega T) \cos(\omega T) \\ -\omega \sin(\omega T) \cos(\omega T) & \omega^2 \cos(\omega T)^2 \end{pmatrix}$$

with
$$\omega = \sqrt{\frac{D}{m}}$$
.

• If we start at if we start $x(0) = (0, v_0)^{\mathsf{T}}$, the kinetic energy of the mass point is

$$E_{\rm kin}(T) = x(T)^{\intercal} Q_T x(T) = x(0)^{\intercal} P(0) x(0) = \frac{1}{2} m v_0^2 \cos(\omega T)^2 \; .$$

Homework 5: details of the derivation & potential energy.

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The matrix-valued function

$$P(t) = \int_t^T G(\tau, t)^{\mathsf{T}} Q(\tau) G(\tau, t) d\tau + G(T, t)^{\mathsf{T}} Q_T G(T, t) .$$

satisfies

$$\begin{split} \dot{P}(t) &= -Q(t) - \int_t^T A(t)^\intercal G(\tau,t)^\intercal Q(\tau) G(\tau,t) \,\mathrm{d}\tau \\ &- A(t)^\intercal G(T,t)^\intercal Q_T G(T,t) \\ &- \int_t^T G(\tau,t)^\intercal Q(\tau) G(\tau,t) A(t) \,\mathrm{d}\tau - G(T,t)^\intercal Q_T G(T,t) A(t) \\ &= -Q(t) - A(t)^\intercal P(t) - P(t) A(t) \;. \end{split}$$

The differential equation

$$-\dot{P}(t) = A(t)^{\mathsf{T}}P(t) + P(t)A(t) + Q(t)$$
 with $P(T) = Q_T$

is called a reverse (inhomogeneous) Lyapunov differential equation.

- ullet We "start" with the end value $P(T)=Q_T$ and
- simulate "backwards" to find P(0).
- We could also "mirror" the free variable $t \leftarrow T t$. (Exercise).

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Summary

The reverse (inhomogeneous) Lyapunov differential equation

$$-\dot{P}(t) = A(t)^{\mathsf{T}}P(t) + P(t)A(t) + Q(t) \quad \text{with} \quad P(T) = Q_T$$

allows us to write q(t) in the form

$$\int_{t}^{T} x(\tau)^{\mathsf{T}} Q(\tau) x(\tau) \, \mathrm{d}\tau + x(T)^{\mathsf{T}} Q_{T} x(T) = x(t)^{\mathsf{T}} P(t) x(t) \; .$$

 This result is a preparation for understanding optimal control concepts such as Dynamic Programming; we'll come back to this later.

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