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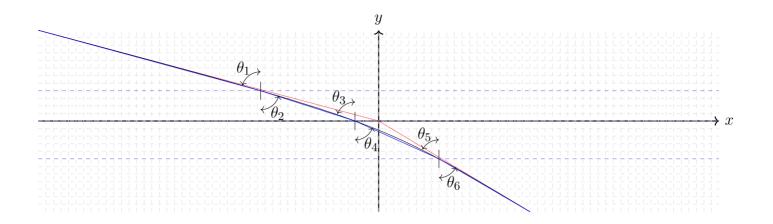
# Joint Research of Optics and Fluid Interface

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## 1 Differential Analytic

$$\forall i \in \mathbb{N}: \frac{\sin \theta_{2i-1}}{\sin \theta_{2i}} = \frac{v_{2i-1}}{v_{2i}} \tag{1}$$

Let the ray inbound from upper toward lower and refracs at origin. Let  $\theta(y)$  be the acute angle between the velocity vector at that instant and the y-axis.

Let y = P(x) be the path light actually taken, now our goal is to solve P(x) and also the function of speed of light within specific y v(y) = v(P(x)).

$$\lim_{\Delta x \to 0} \frac{\sin \theta(P(x))}{\sin(P(x + \Delta x))} = \frac{v(P(x))}{v(P(x + \Delta x))}$$
 (2)

$$\lim_{\Delta y \to 0} \frac{\sin \theta(y)}{\sin(y + \Delta y)} = \frac{v(y)}{v(y + \Delta y)} \tag{3}$$

$$\lim_{r \to 0} \frac{\sin \theta(y)}{\sin (\theta(y) + r\theta'(y))} = \frac{v(y)}{v(y) + rv'(y)} \tag{4}$$

$$\lim_{r \to 0} \sin \theta(y) v(y) + rv'(y) = v(y) \sin \left(\theta(y) + r\theta'(y)\right)$$
(5)

#### 2 Setup

Let  $\Delta$  be the total width the light travels with path not being straight, hence:

$$\forall x \notin \left[ -\frac{\Delta}{2}, \frac{\Delta}{2} \right] : \frac{d^2}{dx^2} y = 0 \tag{6}$$

Now define the light ray incoming from left-hand side towards right-hand side and (suppose transition layers' width approaches zero) refracs at origin. And let h be the thickness of the transition layers.

Now let's assume the incoming transition layer has the same thickness as outbounding transition layer, hence we get:

$$\Delta = \frac{h}{2} (\tan \theta_{in} + \tan \theta_{out}) \tag{7}$$

Now we can first confirm our path function is a function of  $\theta_{in}$ ,  $\theta_{out}$ , h

$$P(\theta_{in}, \theta_{out}, h) = \tag{8}$$

#### 3 Playing with Bezier Curve

First setup the environment

```
h, t=var("h t")
hh=h/2
theta=var("theta")
tprime=var("tprime", latex_name=r"\theta^\prime")
assume(h, "constant")
assume(h>0)
assume(0<theta, theta<pi/2)
assume(0<tprime, tprime<pi/2)</pre>
```

And now we could try to solve functions like y(x) from parametric bezier curve:

```
# x(t)=(t-1)^2*tan(theta)*hh*-1+t^2*tan(tprime)*hh y=(t-1)^2*hh-t^2*hh
```

First we try to find t(x)

```
result=[]
```

[((abs(res.rhs()(x=tan(theta)\*hh\*-1)(h=2, theta=0.2, tprime=0.4))<1e-12 and abs(res.rhs()(x=tan(tprime)\*hh)(h=2, theta=0.2, tprime=0.4)-1)<1e-12) and

```
result.append(res))
for res in solve((t-1)^2*tan(theta)*hh*-1+t^2*tan(tprime)*hh==x, t)]
t=result[0].rhs() if len(result)==1 else None
```

Where we get two solutions and only one will satisfied our assumptions (the first print is 0 and the second yields  $1^1$ ), which is:

$$t(x) = \frac{h \tan(\theta) - \sqrt{h^2 \tan(\theta) \tan(\theta') - 2(h \tan(\theta) - h \tan(\theta'))x}}{h \tan(\theta) - h \tan(\theta')}$$

And substituting back to y(t) we could get

## 4 Analyzing

Given a full path from  $v_1$  to  $v_n$ , whose path will be  $y(v_1, v_n, h)(x)$ , we could separate it into different sections.

$$\Delta_{i \to i+1} = \frac{h}{n-1} \frac{1}{2} \left[ -\tan \theta_i, \tan \theta_{i+1} \right]$$

$$\sum_{i=1}^{n-1} \Delta_{i \to i+1} = \frac{h}{2} (\tan \theta_1 + \tan \theta_n)$$
(9)

$$\forall x \in \underset{i \to i+1}{\Delta} : \tag{10}$$

$$(x, y(v_1, v_2, h_{12})(x)) + (\sum_{2 \to n} h_{2 \to n})$$

<sup>&</sup>lt;sup>1</sup> Here we use EPS = 1e - 12 as we're testing it with arbritarily chosen values which makes them no longer symbolic.

## 5 Finding path

#### 5.1 Constructing Layers

Let function  $v(\theta_{in}, \theta_{out}, \Delta, \delta)$  be the velocity of layer  $\delta$  (from the entry interface), whose entry velocity is  $\theta_{in}$ , exit velocity is be  $\theta_{out}$ , and the whole layer thickness is  $\Delta$ .<sup>2</sup>

$$a = v(x, y, \beta - \alpha, \alpha) \tag{11a}$$

$$b = v(x, y, \beta - \alpha, \beta) \tag{11b}$$

$$\alpha \le \beta$$
 (11c)

Let 
$$a = v(x, y, \alpha)$$
 and  $b = v(x, y, \beta)$   $(\alpha \le \beta), \forall p \in [\alpha, \beta] : v(a, b, p - \alpha) = v(x, y, p)$ 

#### 5.2 Fractal Structure

Let 
$$P(\theta_{in}, \theta_{out}, h)$$
 be the path taken (12)

$$P(\theta_{in\to\alpha}, \alpha) = P(\theta_{in}, \theta_{\alpha}) \tag{13}$$

<sup>&</sup>lt;sup>2</sup> Here we suppose every layer paraelles

## 6 Regarding Special Relativity

So first let the real velocity v, whom maintains the linear properties of traditional non-relativistic velocity.

Linear to non-linear (relativistic) velocity:

$$v_S = \frac{v_{ij} \, \mathcal{I}_S \, \mathcal{I}}{\mathcal{I}_S S} = v_{ij} \, \mathcal{I}_S \, \mathcal{I} \tag{14}$$

#### 6.1 Constant force/acceleration

Let a point located at origin, with initial velocity v(0) = 0.

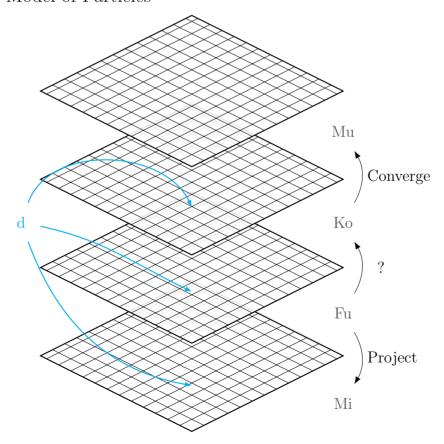
$$a_{\mathcal{L}}(v, F, m_{rest}) = \frac{F}{m_{rest}} \sqrt{1 - \frac{v^2}{c^2}}^3$$
 (15a)

$$a_{11}(t) = k \tag{15b}$$

$$v_{\succeq}(t) = v_{\mathcal{Y}} \, \mathcal{I}_S \, \mathcal{Y} = kt \, \mathcal{I}_S \, \mathcal{Y} \tag{15c}$$

$$v_{\geq}(t) = \int \frac{F}{m_{rest}} \sqrt{1 - \frac{v_{\geq}(t)^2}{c^2}}^3$$
 (15d)

## 7 Model of Particles



## 8 Conclusion

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## Acknowledgment

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#### References

[1] https://www.mail-archive.com/dou-geometry@googlegroups.com/msg00004/\_\_\_\_\_\_\_

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