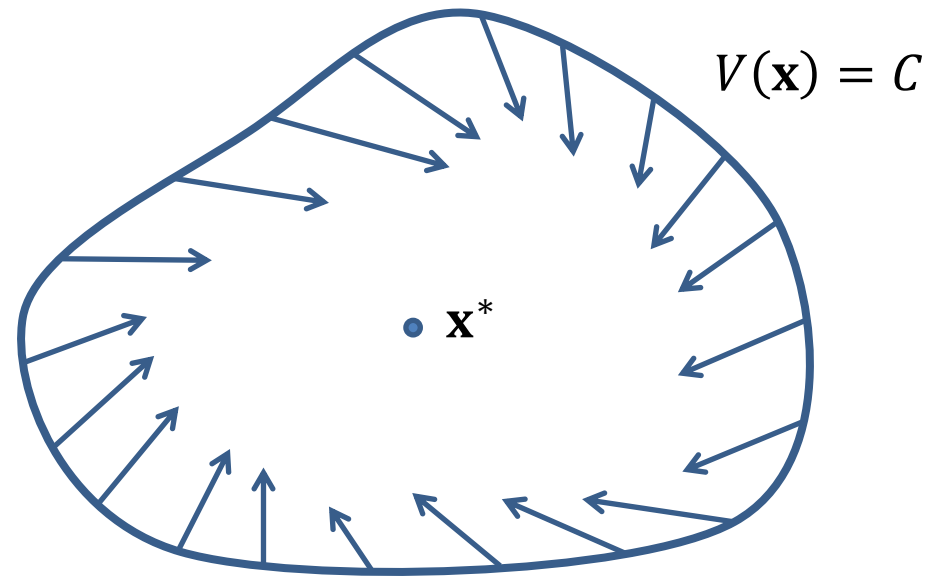


# Lecture 4: Lyapunov Functions

- Suppose there exists a connected orientable region (meaning there is an inside and outside) defined by  $\{\mathbf{x}: V(\mathbf{x}) \leq C\}$  surrounding an equilibrium point  $\mathbf{x}^*$  so that all flows crossing the boundary point remain inside the region.
- Once inside the region, the flow cannot escape



# Nested boundaries

Consider a nested sequence of surfaces defined by a reducing set of constants. The increasing normal to each surface is given by  $\underline{\nabla}V$

We require the flow to point inwards, i.e.

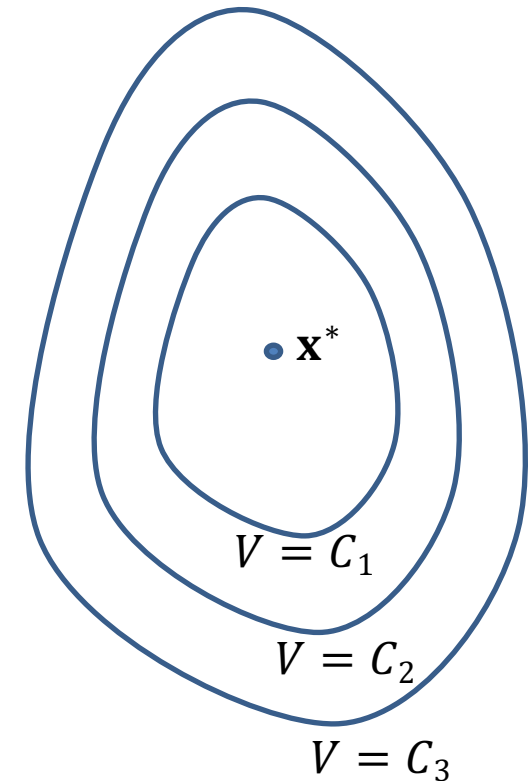
$\underline{\nabla}V \cdot \dot{\mathbf{x}} \leq 0$  But  $\dot{\mathbf{x}} = f(\mathbf{x})$ . We thus require

$$\underline{\nabla}V \cdot f(\mathbf{x}) \leq 0$$

But

$$\frac{dV}{dt} = \sum \frac{\delta V}{\delta x_i} \frac{dx_i}{dt} = \underline{\nabla}V \cdot \dot{\mathbf{x}} \leq 0$$

so flows go downhill and end up at the bottom



# Lyapunov's Theorem

- Let  $\mathbf{x}^*$  be an equilibrium point of  $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x})$ , i.e.  $\mathbf{f}(\mathbf{x}^*) = 0$ . Let  $D$  be an open set surrounding  $\mathbf{x}^*$  and let  $V(\mathbf{x}): D \rightarrow \mathbb{R}$  be a continuously differentiable function on  $D$  such that

1.  $V(\mathbf{x}^*) = 0$  and  $V(\mathbf{x}) > 0$  for all  $\mathbf{x} \neq \mathbf{x}^*$
2.  $\dot{V}(\mathbf{x}) = \underline{\nabla} V \cdot \mathbf{f}(\mathbf{x}) \leq 0$

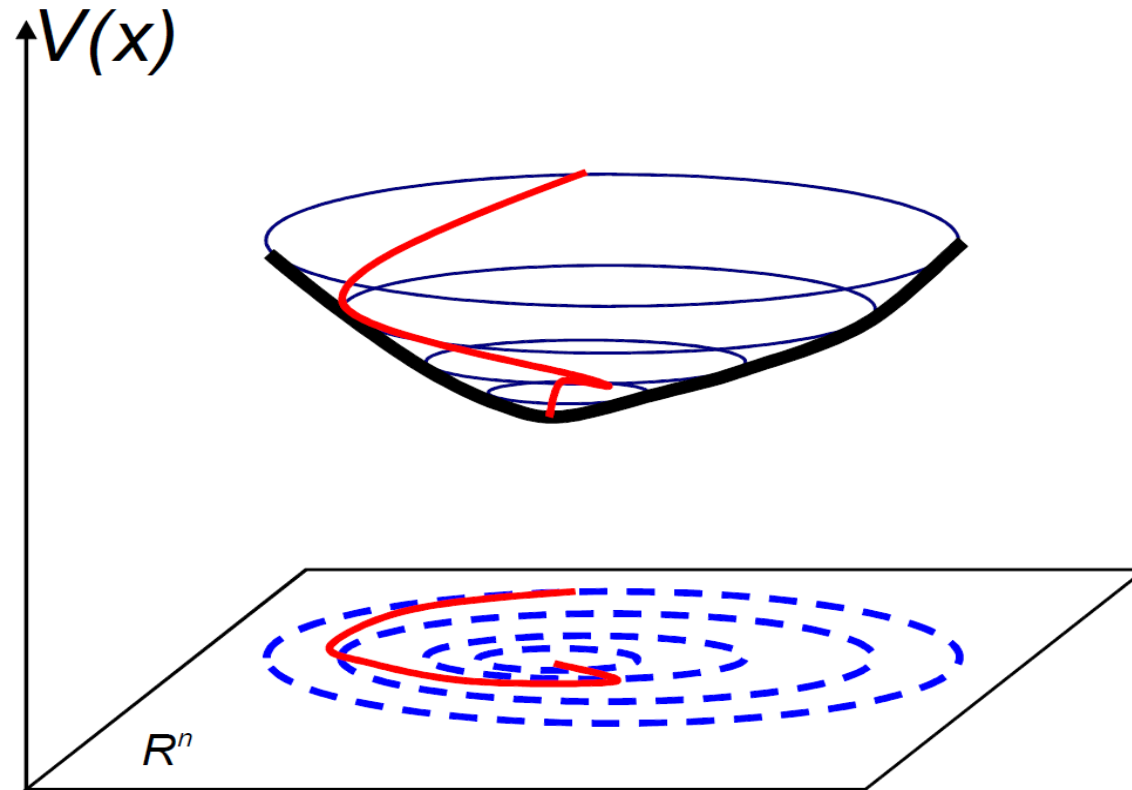
then  $\mathbf{x}^*$  is **stable**. If, in addition

3.  $\dot{V}(\mathbf{x}) < 0$  for all  $\mathbf{x} \neq \mathbf{x}^*$

then  $\mathbf{x}^*$  is **asymptotically stable**

- $V(\mathbf{x})$  is called a Lyapunov function
- If  $\lim_{\|\mathbf{x}\| \rightarrow \infty} V(\mathbf{x}) = \infty$  and  $D = \mathbb{R}^n$  then  $\mathbf{x}^*$  is **globally asymptotically stable**

# Illustration



$V(\mathbf{x})$  decreases along solution trajectories

# Example 1

Consider the dynamical system

$$\dot{x} = y$$

$$\dot{y} = -x + \epsilon x^2 y$$

Equilibrium:  $(0,0)$  has Jacobian:  $\begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$  with eigenvalues:  $\lambda = \pm j$

so Hartman-Grobman doesn't apply

Let  $V(x, y) = \frac{1}{2}(x^2 + y^2)$ ,

$$\frac{dV}{dt} = \underline{\nabla} V \cdot \dot{\mathbf{x}} = x\dot{x} + y\dot{y} = xy - xy + \epsilon x^2 y^2 = \epsilon x^2 y^2$$

If  $\epsilon < 0$  then  $(0,0)$  is stable

## Example 2

$$\dot{x}_1 = -2x_2 + x_2x_3$$

$$\dot{x}_2 = x_1 - x_1x_3$$

$$\dot{x}_3 = x_1x_2$$

- Equilibrium point:  $(0, 0, 0)$  is a linear centre. Let

$$V(\mathbf{x}) = \frac{1}{2}(c_1x_1^2 + c_2x_2^2 + c_3x_3^2)$$

Then

$$\begin{aligned}\dot{V} &= \underline{\nabla} V \cdot \dot{\mathbf{x}} = c_1x_1(-2x_2 + x_2x_3) + c_2x_2(x_1 - x_1x_3) + c_3x_3x_1x_2 \\ &= (c_1 - c_2 + c_3)x_1x_2x_3 + (-2c_1 + c_2)x_1x_2\end{aligned}$$

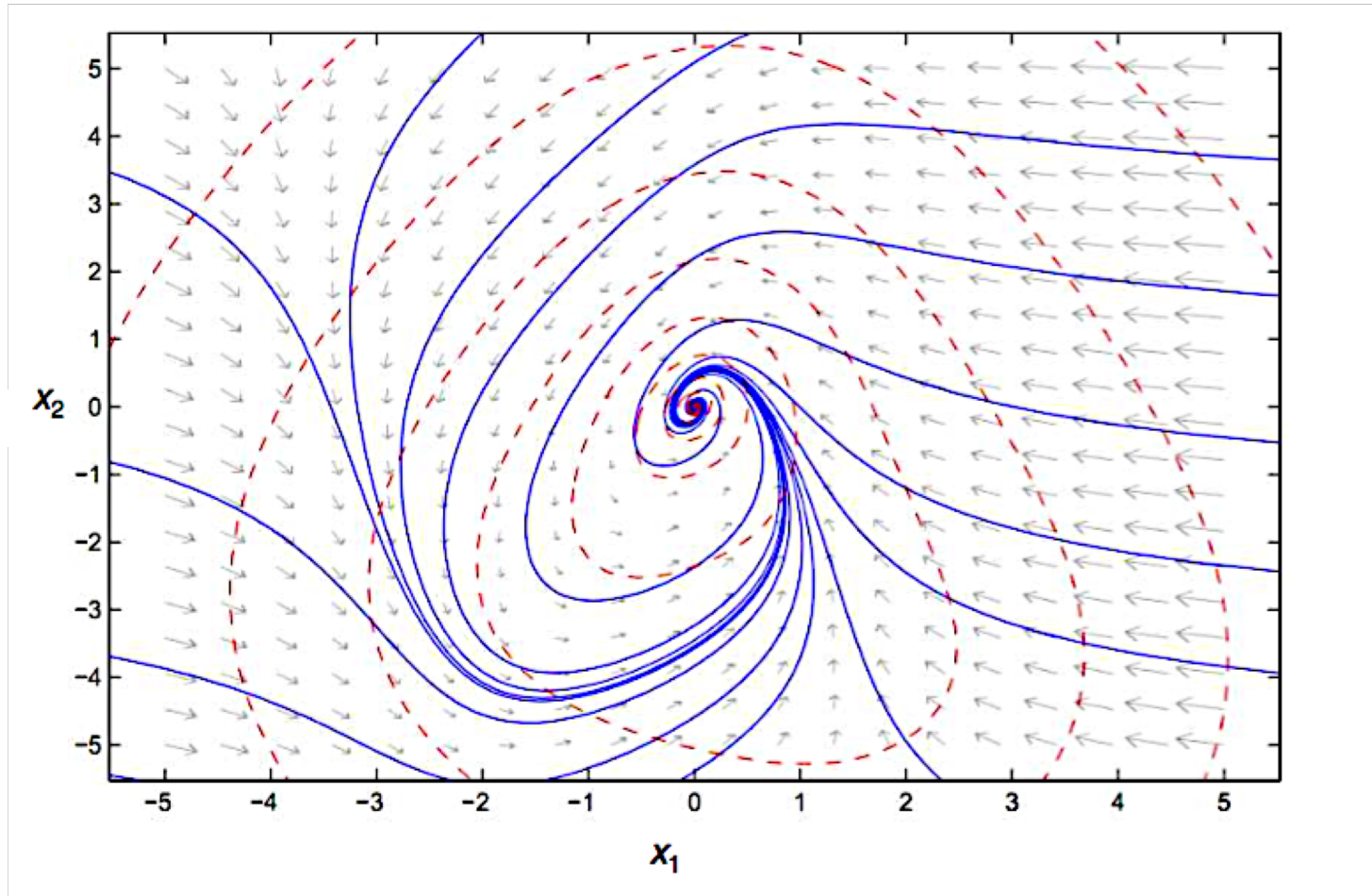
- Choose  $c_2 = 2c_1 > 0$  and  $c_3 = c_1$ , then  $\dot{V} = 0$ , so equilibrium is stable
- $\dot{V} = 0$  on  $V(\mathbf{x}) = \frac{1}{2}(x_1^2 + 2x_2^2 + x_3^2)$  so  $\mathbf{x}(t)$  lies on  $V(\mathbf{x}) = \text{const}$

# Jet Engine Example

$$\begin{aligned}\dot{x}_1 &= -x_2 + 1.5x_1^2 - 0.5x_1^3 \\ \dot{x}_2 &= 3x_1 - x_2\end{aligned}$$

- Equilibrium point: (0,0) has Jacobian  $\begin{bmatrix} 0 & -1 \\ 3 & -1 \end{bmatrix}$  with  $\lambda = \frac{-1 \pm j\sqrt{11}}{2}$  so is a linear stable focus.
- Hartman-Grobman theorem states that the non-linear system is stable (but only close to the origin)
- Lyapunov functions can extend this result globally using specially constructed functions – see the lecture notes

# Jet Engine Example



Level curves of the Lyapunov function showing global stability of the Jet engine model



# Vector fields possessing an integral

- Consider the flow associated with the solution of  $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x})$  as a vector field

- This is said to have an integral  $I(\mathbf{x})$  (a scalar function) if

$$\frac{dI(\mathbf{x})}{dt} = \frac{\partial I(\mathbf{x})}{\partial \mathbf{x}} \dot{\mathbf{x}} = \frac{\partial I(\mathbf{x})}{\partial \mathbf{x}} \mathbf{f}(\mathbf{x}) = 0$$

- $\frac{\partial I(\mathbf{x})}{\partial \mathbf{x}}$  is the gradient vector of  $I(\mathbf{x})$
- $I(\mathbf{x})$  defines level sets which contain the flow

# Pendulum example

$$\begin{aligned}\dot{q} &= p \\ \dot{p} &= -\frac{g}{l} \sin q\end{aligned}$$

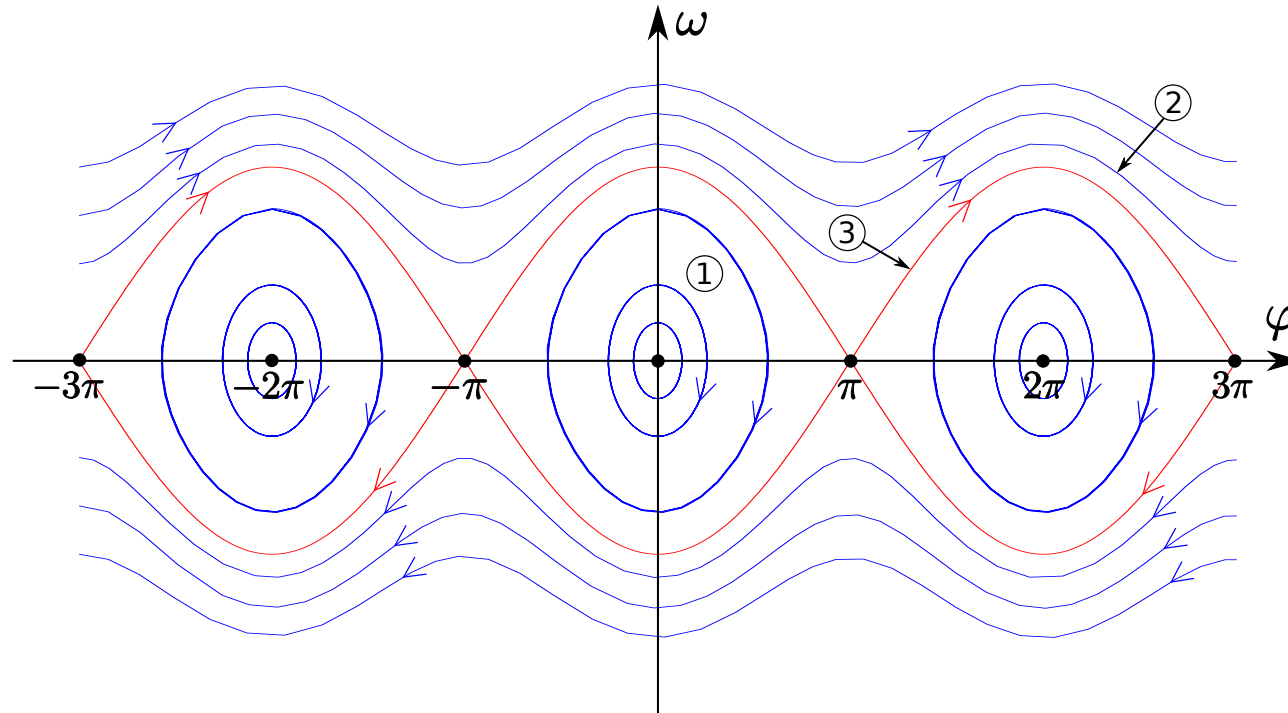
The total stored energy  $E$  is conserved

$$E = \frac{1}{2}p^2 - \frac{g}{l} \cos q$$

i.e.

$$\frac{dE}{dt} = p\dot{p} + \dot{q} \frac{g}{l} \sin q = 0$$

# Pendulum example



Phase plane of pendulum and level sets of constant energy

# Duffing Oscillator for $\delta=0$

$$\begin{aligned}\dot{x} &= y \\ \dot{y} &= x - x^3\end{aligned}$$

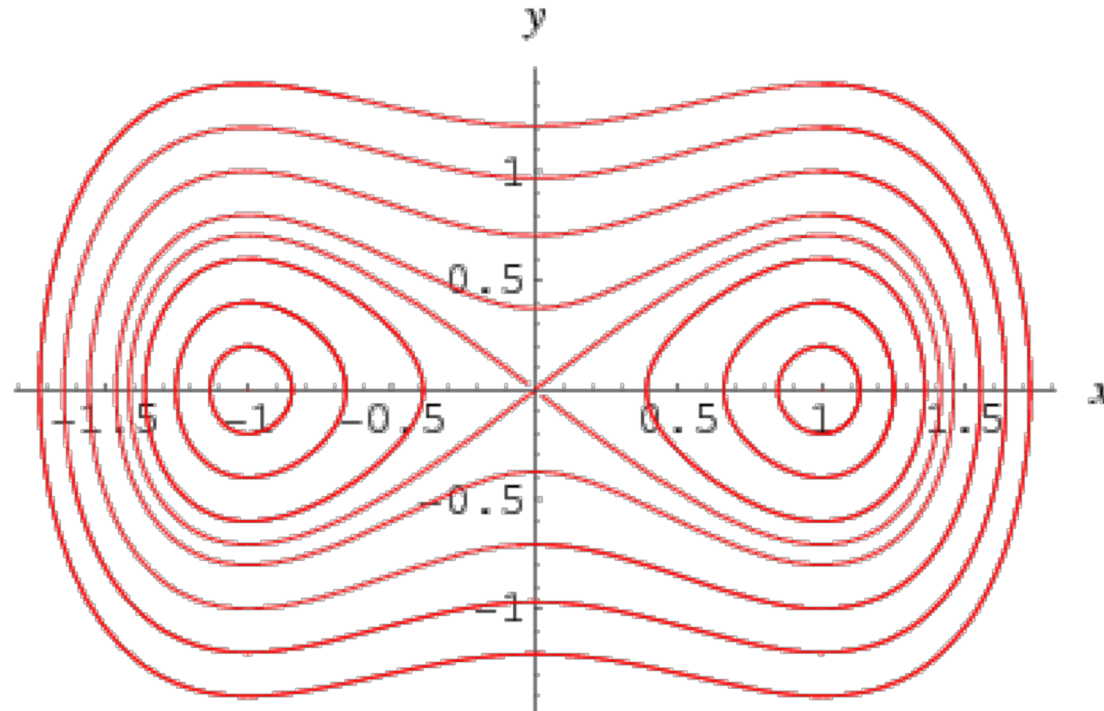
We require

$$\begin{aligned}\frac{dI}{dt} &= \frac{\partial I}{\partial x} \dot{x} + \frac{\partial I}{\partial y} \dot{y} = 0 \\ \frac{\partial I}{\partial x} y + \frac{\partial I}{\partial y} (x - x^3) &= 0\end{aligned}$$

So, for example

$$I = \frac{y^2}{2} - \frac{x^2}{2} + \frac{x^4}{4}$$

# Duffing Oscillator for $\delta=0$



Level sets of  $I(\mathbf{x})$  in the phase plane of the duffing oscillator

# Hamiltonian systems

Hamiltonian systems have vector fields that possess an integral

**Definition:** Systems of the form

$$\begin{aligned}\dot{\mathbf{p}} &= \mathbf{f}(\mathbf{p}, \mathbf{q}) \\ \dot{\mathbf{q}} &= \mathbf{g}(\mathbf{p}, \mathbf{q})\end{aligned}$$

such that

$$\mathbf{f}(\mathbf{p}, \mathbf{q}) = \partial H(\mathbf{p}, \mathbf{q}) / \partial \mathbf{q}, \quad \mathbf{g}(\mathbf{p}, \mathbf{q}) = -\partial H(\mathbf{p}, \mathbf{q}) / \partial \mathbf{p}$$

are called **Hamiltonian Systems**.

- $\mathbf{p}$  and  $\mathbf{q}$  are real vectors with  $n$  elements
- $H$  is a twice differentiable function called the Hamiltonian
- $\mathbf{q}$  is the vector of generalised positions,  $\mathbf{p}$  the vector of generalised momenta
- All Hamiltonian systems are conservative by construction

# More on Hamiltonian systems

- If  $(\mathbf{p}^*, \mathbf{q}^*)$  is an equilibrium and  $H(\mathbf{p}, \mathbf{q}) > 0$  in a region surrounding the equilibrium, then the equilibrium is stable
- A Newtonian system of the form  $\ddot{x} = f(x)$  can be written as a Hamiltonian system by summing the potential energy and kinetic energy

$$\begin{aligned}\dot{x} &= v \\ \dot{v} &= f(x) \\ H(x, v) &= \frac{v^2}{2} - \int_{x_0}^x f(s) ds\end{aligned}$$

# Gradient Systems

**Definition:** Let  $V(\mathbf{x})$  be a twice differentiable function in a region  $D \subseteq \mathbb{R}^n$ . The system

$$\dot{x}_i = -\frac{\partial V}{\partial x_i}$$

is called a **gradient** system.

- Equilibrium points are the critical points of  $V$ . Away from critical points the trajectories are orthogonal to the level sets of  $V$ .
- If  $\mathbf{x}^*$  is a strict local minimum of  $V$  then  $V(\mathbf{x}) - V(\mathbf{x}^*)$  is a Lyapunov function for  $\mathbf{x}^*$ , showing that  $\mathbf{x}^*$  is asymptotically stable. If  $\mathbf{x}^*$  is a strict local maximum, then the equilibrium is unstable.



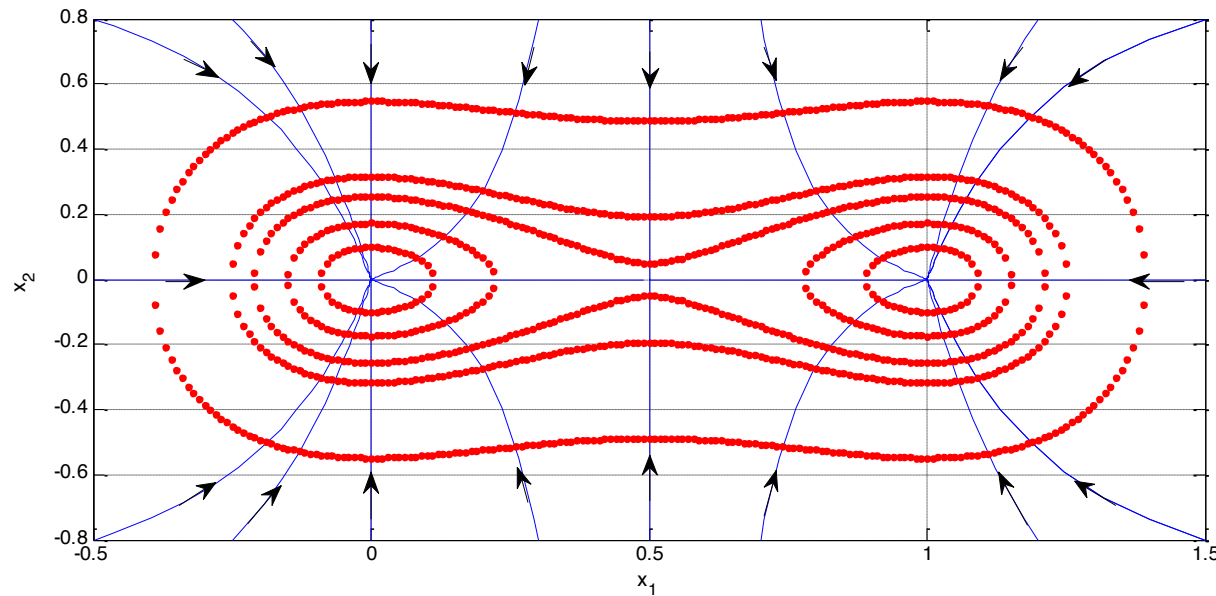
# Example Gradient System

$$\dot{x} = -4x(x - 1)(x - 0.5)$$

$$\dot{y} = -2y$$

Has

$$V(x, y) = x^2(x - 1)^2 + y^2$$



# Relationship between Gradient and Hamiltonian Systems

- The system

$$\begin{aligned}\dot{x} &= f(x, y) = \frac{\partial H}{\partial y} \\ \dot{y} &= g(x, y) = -\frac{\partial H}{\partial x}\end{aligned}$$

is orthogonal to

$$\begin{aligned}\dot{x} &= g(x, y) \\ \dot{y} &= -f(x, y)\end{aligned}$$

- They have the same equilibria, centres map to nodes, saddles to saddles and foci to foci.