

Measurement of Central Exclusive Production of $h\bar{h}$ pairs $(h=\pi,K,p)$ with Roman Pot detectors in diffractive proton-proton interactions at $\sqrt{s}=200~{\rm GeV}$

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Abstract

In this note we present analysis of the Central Exclusive Production process using data from proton-proton collisions collected in 2015. This data was collected using the Roman Pot detectors which ensured efficient triggering and measuring diffractively scattered protons. We describe all intermediate stages of analysis involving extraction of the acceptance and efficiency corrections, comparison of data with Monte Carlo simulations of detector response, and study of systematic uncertainties. Finally, we show the physics outcome of the analysis.

DRAFT

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1. Introduction

1.1 Central Exclusive Production

The Central Exclusive Production (CEP) takes place when interacting particles form in the mid-rapidity region a state ("central production") whose all constituents/decay products are measured in the detector ("exclusive"). The initial state particles can either dissociate, excite or stay intact. The latter case of CEP in proton-proton collisions can be written as

$$p + p \rightarrow p + X + p \tag{1.1}$$

and depicted as in Fig. 1.1. Mass and rapidity of state X is given by

$$M_X = \sqrt{s(\xi_1 \xi_2 \sin^2(\alpha/2) - (1 - \xi_1 - \xi_2) \cos^2(\alpha/2))} \stackrel{\alpha = \pi}{=} \sqrt{s\xi_1 \xi_2}, \quad (1.2) \qquad y_X = \frac{1}{2} \ln \frac{\xi_1}{\xi_2}, \quad (1.3)$$

where α is angle between scattered protons and $\xi = (p_0 - p)/p_0$ is the fractional momentum loss of proton.

1.2 Double IPomeron Exchange

Reaction from Eq. (1.1) can exhibit purely electromagnetic $(\gamma-\gamma)$, mixed $(\gamma-\mathcal{O})$ or purely strong nature $(\mathcal{O}-\mathcal{O})$. The last type is dominant at RHIC energies. It is characterized by the lack of hard scale (if protons are scattered at small angles), therefore perturbative QCD cannot be applied and Regge theory [1] is used instead. An object \mathcal{O} does not have unequivocal QCD representation - in Regge formalism it is the so-called "trajectory" (Reggeon, \mathbb{R}). Reggeon with quantum numbers of vacuum is called "Pomeron" (\mathbb{P}) and \mathbb{P} - \mathbb{P} reaction (Fig. 1.2) is called "Double Pomeron Exchange".

Processes involing IPomeron exchange are referred as diffraction due to cross-section in scattering angle resembling similar shape to instesity pattern of diffracted light. For low values of Mandelstam t (small scattering angles) cross-section takes exponential form

$$\frac{d\sigma}{d|t|} \propto e^{-B|t|},\tag{1.4}$$

where the slope parameter B reflects the size of target at which \mathbb{P} omerons scatter.

Diffractive events have specific property of the "rapidity gap" which is an angular region free of hadrons. In DIPE two such gaps are present, marked in Fig. 1.1 as $\Delta \eta_1$ and $\Delta \eta_2$.

DIPE is a spin-parity filter - from the fact that scattered particles have all quantum numbers unchanged after the interaction, central states must satisfy

$$I^G J^{PC} = 0^+ \text{even}^{++}.$$
 (1.5)

The lowest order QCD picture of the IPomeron is a pair of oppositely colored gluons (colour singlet). This fact makes the DIPE recognized as the gluon-rich environment process which should enhance production of the bound states of gluons ("glueballs") or hybrid mesons.

For detailed introduction to the topic of diffraction see Refs. [2,3].

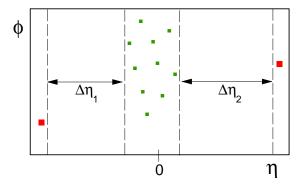


Figure 1.1: Central Exclusive Production in η - ϕ space.

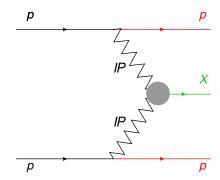


Figure 1.2: Diagram of DIPE process.

1.3 Physics motivation for the measurement

STAR collected in 2015 large dataset dedicated for the measurement of the Central Diffraction (DIPE in particular). Since that year the experiment was enriched with Roman Pot Phase II* subsystem and thus gained possibility of detection of forward protons. It enabled studies of properties of the central state with respect to observables related to exchanged IPomerons. No such measurement was performed before at that high c.m.s. energy ($\sqrt{s} = 200 \text{ GeV}$, contamination from IReggeon exchanges is small) which makes it particularly attractive. A list of physics issues that can be covered with the study described in this note is briefly introduced below.

1.3.1 DIPE differential cross-sections, mass spectrum

As stated in Sec. 1.2 DIPE is a soft process whose theoretical description is done mainly using phenomenological tools, thus measurement of differential cross-sections is needed to verify various production models.

The main focus is put on the simplest state (and most numerously) produced in DIPE, namely a pair of oppositely charged pions, $\pi^+\pi^-$. It can be formed either in a non-resonant or resonant mechanism. In the first case the $\pi^+\pi^-$ continuum is formed by the exchange of the off-shell pion between IPomerons. Currently there are two models of this reaction on the market [4,5], [6]. In the second case the IPomerons directly couple into resonance (e.g. $f_2(1270)$), which then decays to $\pi^+\pi^-$. Attempts to calculate cross-section for this production mechanism are presented in Ref. [5] and [7].

Understanding of the mass spectrum in $\pi^+\pi^-$ channel is important to learn about relative contribution from continuum and resonant production, as well as relative production of resonances. Recognition of resonant states may indicate candidates for low-mass glueballs of $J^{PC}=0^{++}$, however presence of underlaying scalar $q\bar{q}$ states makes this task challenging.

Other channels, like K^+K^- , are also of great interest. Comparison of the cross-sections for production of $\pi^+\pi^-$ and K^+K^- gives information about strength of the Pomeron coupling to different quark flavors. Also, structures in $d\sigma/dm$ can be easier attributed to resonances by measuring more than one channel and known branching ratios thereof.

Detection of intact protons scattered at very small angle with respect to the beamline enables determination of the reaction plane which makes the Partial Wave Analysis (PWA) possible. It also allows to look at the the cross-sections more differentially, especially with respect to properties of exchanged Pomerons, like carried squared four-momentum t, azimuthal separation of Pomerons in the transverse plane $\Delta \varphi$ or relative momentum of Pomerons Δp_T . The last quantity was proposed to distinguish pure $q\bar{q}$ states from these with gluonic content [8].

1.3.2 Absorption effects

One can imagine in diagram in Fig. 1.2 additional soft lines e.g. between protons in the initial state or one of Pomerons and final state proton. These so-called rescattering effects (or absorption effects) lead to production of hadrons other than these belonging to central state X hence the diffractive signature of an event in form of rapidity gap is no longer present. Measurement of the probability that the state X will remain exclusive and forward protons will remain intact, in other words the rapidity gap survival probability S^2 , would be valuable ingredient for development of absorption models.

1.3.3 Size of interaction region

From the measurement of protons in Roman Pots one is able to reconstruct squared four-momenta transferred in proton-Pomeron vertices and determine the differential cross-section $d\sigma/d|t|$. Fit of exponent allows to extract the slope parameter B, which may depend on the Pomeron-Pomeron c.m.s. energy, or in other words on the mass of diffractive system X. Knowledge on the slope parameter gives insight to the volume and distribution of Pomerons inside proton.

2. Data set

2.1 Trigger

The main trigger designed for studies of Central Diffraction in run 15 was RP_CPT2. It was formed of the following conditions combined by logical AND (&&):

- 1. (ET && !IT) || (!ET && IT) = signal in at least one RP on each side of the STAR central detector to ensure presence of two forward-scattered protons; a veto was imposed on simultaneous signal in RPs above and below the beamline, which might have originated either from proton dissociation, or pile-up event, or beam halo proton etc.,
- 2. !BBCE && !BBCW && !ZDCE && !ZDCW = veto on any signal in small BBC tiles or ZDCs on any side of STAR central detector such requirement is in accordance with the double-gap topology of CEP events, it mostly filtered out CEP events with parallel pile-up event(s),
- 3. $TOF \ge 2$ = at least 2 hits in TOF aim of this condition was to ensure activity in the mid-rapidity; since the lowest multiplicity allowed in CEP is 2, that was the lower threshold of L0 TOF multiplicity.

This trigger was running with an average prescale of 5 and average DAQ rate of 250 Hz, which allowed to collect in total about 560 M events corresponding to 16.5 pb⁻¹ of integrated luminosity. More information about number of events per run, rates etc. can be found under link provided in Ref. [9], which contains selected data from STAR run log [10]. Luminosity data used in this analysis comes from Ref. [11].

All RP triggers which were intended for usage in diffractive physics analyses or efficiency studies are listed in Tab. 2.1. Components used in definitions of these triggers are outlined in Fig. 2.1. Detailed explanation of all trigger bits can be found in Refs. [12,13]. Explanation of naming convention in Roman Pot system can be found in Ref. [14].

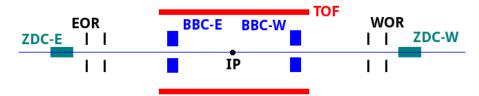


Figure 2.1: Sketch of the trigger components used in definitions of diffractive triggers in run 15.

Trigger name	Definition	Events [M] (available)	Comment
RP_CP	EOR && WOR	73.3 (46.0)	Loose trigger (mostly elastic events) designed for monitoring/trigger efficiency study
RP_CPT	EOR && WOR && !BBCE && !BBCW && !ZDCE && !ZDCW && TOF≥1	38.9 (0.0)	Intended to be main CEP trigger (later switched to RP_CPT2 due to large prescale)
RP_CPT2	(ET && !IT) (!ET && IT) && !BBCE && !BBCW && !ZDCE && !ZDCW && TOF≥2	556.5 (493.8)	Main CEP trigger Note: On Apr 14 added upper TOF limit (10)
RP_CPX	IT && !BBCE && !BBCW && !ZDCE && !ZDCW && TOF≥2	40.1 (0.0)	The same as RP_CPT2 but only IT configuration
RP_CPEI	ET && IT && !BBCE && !BBCW && !ZDCE && !ZDCW && TOF>2	15.6 (11.8)	Control trigger for CPT2 to estimate effect of !(ET && IT) veto

Table 2.1: Central Diffraction physics triggers and control triggers involving Roman Pot detectors in run 15.

2.2 Reconstruction software

Raw data was processed with STAR libraries in versions SL15k (trigger datasets: production_pp200trans_2015, production_pp200long2_2015, production_pp200long3_2015) and SL16d (production_pp200long_2015) [15].

2.3 Data format

The analyzed data was stored in ROOT files in the picoDST format which was in large part a skimmed MuDST (standard STAR format). The picoDST format was introduced in Ref. [16]. PicoDST description files (C++ headers etc.) can be found in the analysis code repository [17].

2.4 Bad runs

Analysis of CEP was performed on the data from runs with completion status "Successful" in the STAR run log. However, based on additional requirements explained below, some runs were omitted from analysis.

2.4.1 RP distance from the beamline

 $16065025\ 16065026\ 16065027\ 16065028\ 16072057\ 16072058\ 16077055\ 16083006\ 16083007\ 16106031$

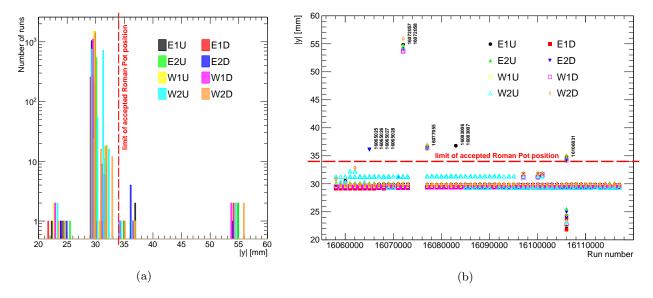


Figure 2.2: Histogram of beam-detector distance |y| (2.2a) and graph showing run-dependence of |y| (2.2b) for all Roman Pots.

3. Event selection

Complete list of analysis cuts used for signal extraction is presented in Sec. 3.1. Detailed description of each cut can be found in Sec. 3.2.

List of cuts¹ 3.1

- C1. Exactly 1 primary vertex with TPC track(s) matched with hits in TOF.
- C2. TPC vertex from C1 is placed within $|z_{vx}| < 100$ cm.
- C3. Exactly 2 opposite-sign primary TPC tracks (C3.2) of good quality and within high TPC acceptance region (C3.3) matched with hits in TOF (C3.1).
 - C3.1. Exactly 2 TOF-matched primary tracks,
 - C3.2. Tracks are of opposite signs,
 - C3.3. Both tracks satisfy criteria below:

$$|\eta|<1, \quad p_T>0.2~{\rm GeV}/c, \quad {\rm DCA}<2~{\rm cm}, \quad N_{\rm hits}^{\rm fit}\geq 25, \quad N_{\rm hits}^{\rm dE/dx}\geq 15, \quad N_{\rm hits}^{\rm fit}/N_{\rm hits}^{\rm poss}\geq 0.52,$$

- C4. Exactly 1 RP track on each side of STAR central detector (C4.3) of good quality (C4.1), with local angles consistent with the IP being the track origin (C4.2), lying within fiducial region of high geometrical acceptance (C4.4).
 - C4.1. RP tracks contain only track-points with at least 3 (out of 4) planes used in reconstrucion,
 - C4.2. Local angles [mrad] lie within $-1.2 < \theta_x^{\text{RP}} < 5.0$ and $1.5 < |\theta_y^{\text{RP}}| < 4.5$,
 - C4.3. Exactly 1 track passing cuts C4.1-C4.2 per side,
 - **C4.4.** Tracks passing cut C4.3 lie within the (p_x, p_y) region defined as: $|p_y| > 0.2 \text{ GeV}/c$, $(p_x/(\text{GeV}/c) + 0.3)^2 + (p_y/(\text{GeV}/c))^2 < 0.46^2, \qquad (p_x/(\text{GeV}/c) - 0.15)^2 + (p_y/(\text{GeV}/c))^2 < 0.42^2.$
- C5. Vertex z-positions measured in TPC and reconstructed from the difference of proton detection time in west and east RPs are consistent with each other within the resolution $|z_{\text{vx}}^{\text{TPC}} - z_{\text{vx}}^{\text{RP}}| < 36 \text{ cm}$.
- C6. No signal in any tile of BBC-large (east or west) with ADC > 40 and TDC > 100.
- C7. Maximally 2 reconstructed TOF clusters $N_{\rm clstrs}^{\rm TOF} \leq 2$.
- C8. Missing (total) momentum of TPC tracks and RP tracks $p_T^{\text{miss}} < 75 \text{ MeV}/c$.
- C9. Particle (pair) identification:

```
else event rejected.
```

3.2 Description of cuts

3.2.1(C1,C2) Primary vertex

Bla bla

¹Some cuts (e.g. C3) are decomposed to constituent sub-cuts. Cut is formed by the logical AND of all its sub-cuts.

- (C3) TPC tracks 3.2.2
- 3.2.3 (C4) RP tracks
- (C5) TPC-RP z-vertex matching 3.2.4
- 3.2.5 (C6) BBC-large signal veto
- 3.2.6 (C7) TOF clusters limit
- (C8) Exclusivity cut (missing p_T cut) 3.2.7
- 3.2.8 (C9) Particle identification

In addition to information from the TPC we use time of hit detection in the barrel TOF subsystem. From the simple algebra describing relation between track lengths, momenta and times of hit detection one can derive formula for the squared mass of two particles, assuming that their masses are equal (particles are of the same

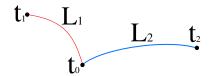


Figure 3.1: Scheme of two central tracks of lengths L_1 and L_2 , produced Figure 3.1: Scheme of two central tracks of lengths \mathcal{L}_1 and \mathcal{L}_2 , in common vertex in moment t_0 , hitting cells in TOF detector in moments t_1 and t_2 .

$$\begin{cases}
t_1 - t_0 = L_1 \sqrt{1 + \frac{m_1^2}{p_1^2}}, \\
t_2 - t_0 = L_2 \sqrt{1 + \frac{m_2^2}{p_2^2}},
\end{cases}$$
(3.1)

$$\Delta t = t_1 - t_2 = L_1 \sqrt{1 + \frac{m_1^2}{p_1^2}} - L_2 \sqrt{1 + \frac{m_2^2}{p_2^2}}.$$
 (3.2)

Assuming $m_1 = m_2 = m \rightarrow m^2$ from quadratic eq.

Parameters of the quadratic equation whose solution is suqared mass and the final formula for m_{TOF}^2 are given below:

$$\mathcal{A} = -2\frac{L_1^2 L_2^2}{p_1^2 p_2^2} + \frac{L_1^4}{p_1^4} + \frac{L_2^4}{p_2^4},\tag{3.3}$$

$$\mathcal{B} = -2L_1^2 L_2^2 \left(\frac{1}{p_1^2} + \frac{1}{p_2^2} \right) + \frac{2L_1^4}{p_1^2} + \frac{2L_2^4}{p_2^2} - 2\left(\Delta t\right)^2 \left(\frac{L_1^2}{p_1^2} + \frac{L_2^2}{p_2^2} \right),\tag{3.4}$$

$$C = (\Delta t)^4 - 2(\Delta t)^2 (L_1^2 + L_2^2) + L_1^4 + L_2^4 - 2L_1^2 L_2^2,$$
(3.5)

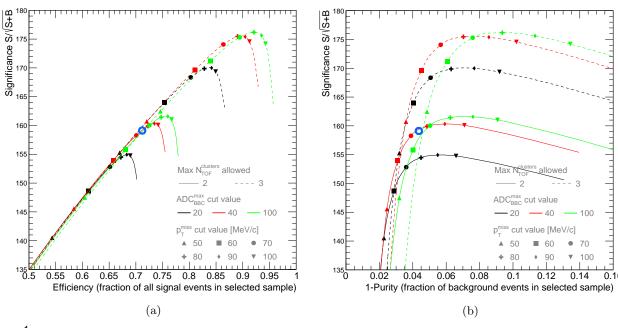
$$m_{\text{\tiny TOF}}^2 = \frac{-\mathcal{B} + \sqrt{\mathcal{B}^2 - 4\mathcal{A}\mathcal{C}}}{2\mathcal{A}}.$$
 (3.6)

Working point for cuts C6, C7 and C8 3.3

$$\text{Significance} = \frac{N_{\text{signal}}^{\text{cut}}}{\sqrt{N_{\text{signal}}^{\text{cut}} + N_{\text{bkgd}}^{\text{cut}}}}, \quad (3.7) \quad \text{Efficiency} = \frac{N_{\text{signal}}^{\text{cut}}}{N_{\text{signal}}^{\text{no cut}}}, \quad (3.8) \quad \text{Purity} = \frac{N_{\text{signal}}^{\text{cut}}}{N_{\text{signal}}^{\text{cut}} + N_{\text{bkgd}}^{\text{cut}}}, \quad (3.9)$$

Signal per integrated luminosity 3.4

3.5 Cut flow



2 ADC_{BBC} cut value 20 40 s cut value [MeV/c] 50 **6**0 70 . 80 100 0.5 0.06 0.08 0.12 0.14 0.04 0.1 1-Purity (fraction of background events in selected sample) (c)

Figure 3.2: Relation between $\pi^+\pi^-$ signal significance and efficiency (3.2a), significance and purity (3.2b), and efficiency and purity (3.2c) as a function of cut thresholds in BBC-large veto (C6), TOF cluster limit (C7) and exclusivity cut (C8). Lines show forementioned relations with changing $p_T^{\rm miss}$ cut whose some specific values are indicated with different markers. Color denotes ADC threshold in BBC-large veto (black, red or green). Style of line (solid or dashed) denotes $N_{\rm clstrs}^{\rm TOF}$ limit. Working point considered optimal is marked with opened blue circle.

4. Corrections

4.1 Method of corrections application

$$\frac{d\sigma}{dq} = \frac{1}{\Delta q} \times \frac{1}{\varepsilon} \times \frac{N^w - N_{\text{bkgd}}^w}{L_{\text{int}}^{\text{eff}}}$$
(4.1)

$$L_{\rm int}^{\rm eff} = \sum_{\rm run} L_{\rm int}^{\rm run} \times \epsilon_{\rm veto}(L^{\rm run})$$
(4.2)

$$\varepsilon = \epsilon_{\text{ET/IT}} \times \epsilon_{\text{vrtx}}(q) \times \epsilon_{C2} \times \epsilon_{C5} \times \epsilon_{C8} \times \epsilon_{\text{PID}}(q)$$
(4.3)

$$N^w = \sum_{\text{event}} w_{\text{event}} \tag{4.4}$$

$$w = \left[\prod_{\text{sign}} \epsilon_{\text{TOF}}(\text{sign}, \text{PID}, p_T, z_{vx}, \eta) \times \prod_{\text{sign}} \epsilon_{\text{TPC}}(\text{sign}, \text{PID}, p_T, z_{vx}, \eta) \times \prod_{\text{side}} \epsilon_{\text{RP}}^{\text{side}}(p_x, p_y) \right]^{-1},$$

$$\text{sign} = \{+, -\}, \text{ side} = \{E, W\}$$

$$(4.5)$$

4.2 Efficiencies and acceptances

4.2.1 Trigger efficiency

4.2.1.1 Online veto (BBC-small and ZDC veto)

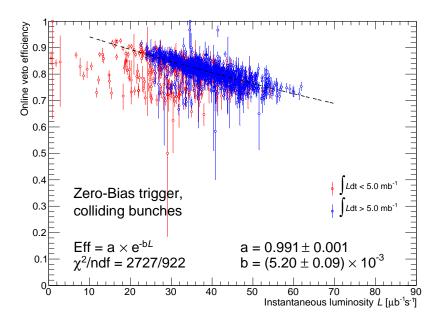


Figure 4.1: Overall efficiency of the online BBC-small and ZDC veto as a function of instantaneous luminosity.

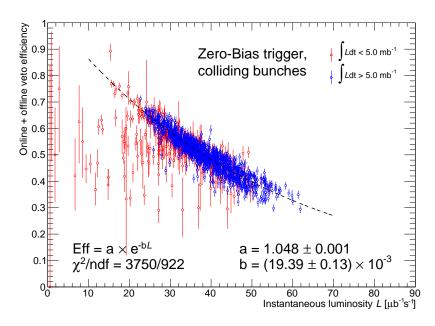


Figure 4.2: Overall efficiency of the online BBC-small and ZDC veto, primary vertices limit (C1), BBC-large veto (C6) and TOF clusters limit (C7) as a function of instantaneous luminosity.

- 4.2.1.2 RP triggering efficiency
- 4.2.1.3 Up and Down RP combination veto
- 4.2.2 Cuts efficiency
- 4.2.2.1 TPC z-vertex cut (C2)
- 4.2.2.2 TPC-RP z-vertex matching (C5)
- 4.2.2.3 Primary vertices limit (C1), BBC-large veto (C6) and TOF clusters limit (C7)
- 4.2.2.4 Missing p_T cut (C8)
- 4.2.2.5 Particle identification (C9)

It is possible to transform dE/dx in MC to make it follow the shape of dE/dx in the data. We know that nSigmaX (where X=Pion, Kaon, Proton, ...) variable follows a gaussian distribution (for particle X)

$$nSigmaX = \Big(\ln\frac{dE/dx}{\langle dE/dx\rangle_X}\Big)/\sigma_{dE/dx}, \qquad f(nSigmaX) = \mathcal{N}(nSigmaX; \mu = 0, \sigma = 1)$$

therefore dE/dx itself follows log-normal distribution:

$$f(dE/dx) = \mathcal{L}og\mathcal{N}(dE/dx; \mu = \langle dE/dx \rangle, \sigma = \sigma_{dE/dx}) = \frac{1}{\sqrt{2\pi} \cdot \sigma \cdot dE/dx} e^{-\frac{\ln^2 \frac{dE/dx}{\langle dE/dx \rangle}}{2\sigma^2}}$$

The transformation we want to apply should preserve the shape of dE/dx (so that it is still described by $\mathcal{L}og\mathcal{N}$), however it should change μ and σ so that these values are euqal to those seen in the data. The transformation that satisfies above postulate is

$$dE/dx' = c \cdot (dE/dx)^a$$

Parameters of the distribution $\mathcal{L}og\mathcal{N}(dE/dx')$ would be then

$$\mu' = c \cdot \mu^a, \quad \sigma' = a \cdot \sigma$$

From above we get formulae for parameters of the transformation:

$$a = \sigma'/\sigma, \quad c = \frac{\mu'}{\mu^a}$$

AlternativeToCrystallBall [18] Eq. (4.6)

$$f(dE/dx) = \begin{cases} \frac{A}{\sqrt{2\pi} \cdot \sigma \cdot dE/dx} \exp\left(-\frac{1}{2} \left(\frac{\ln \frac{dE/dx}{\langle dE/dx}\rangle}{\sigma}\right)^{2}\right) & \text{for } \frac{\ln \frac{dE/dx}{\langle dE/dx}\rangle}{\sigma} \leq k \\ \frac{A}{\sqrt{2\pi} \cdot \sigma \cdot dE/dx} \exp\left(-k \cdot \frac{\ln \frac{dE/dx}{\langle dE/dx}\rangle}{\sigma} + \frac{1}{2}k^{2} - k^{-1} \left(\frac{\ln \frac{dE/dx}{\langle dE/dx}\rangle}{\sigma} - 1\right)^{k}\right) & \text{for } \frac{\ln \frac{dE/dx}{\langle dE/dx}\rangle}{\sigma} > k \end{cases}$$
(4.6)

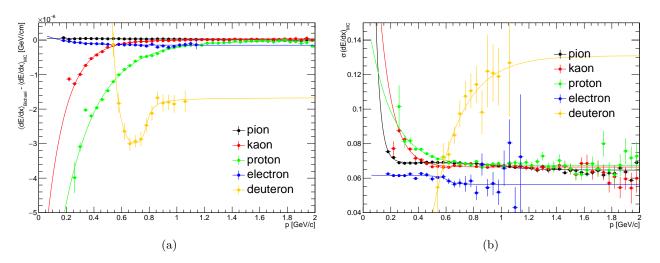


Figure 4.3: Difference between MPV of dE/dx predicted by Bichsel parametrization and obtained from the fit of Eq. (4.6) to dE/dx distribution from MC sample (4.3a) and dE/dx width parameter (4.3b) as a function of reconstructed particle momentum for a few particle species. Solid lines represent fits to points of corresponding color.

$$g(p) = P_1 + P_2 \cdot \exp(-P_3 \cdot p) + P_4 \cdot \arctan(P_5 \cdot (p - P_6))$$
(4.7)

PID	$\langle dE/dx angle_{ m Bichsel} - \langle dE/dx angle_{ m MC}$				$\sigma (dE/dx)_{ m MC}$							
1110	P_1	P_2	P_3	P_4	P_5	P_6	P_1	P_2	P_3	P_4	P_5	P_6
π^{\pm}	3.293e-8	2.237e-8	4.399				0.0668	1.186	28.683	-1.5e-3	7.887	1.028
K^{\pm}	3.068e-9	-8.190e-6	7.089				0.0664	0.323	-10.438			
$\overline{p,ar{p}}$	-3.061e-8	-1.061e-5	4.08				0.0674	0.103	-5.952			
e^{\pm}	-1.483e-7	3.710e-7	5.639				0.0590	-0.009	-194.421	-1.81e-3	26.392	0.607
$\overline{d}, ar{d}$	-2.474e-6	0.385	21.719	5.165e-7	29.642	0.781	0.131	-0.999	4.749			

Table 4.1: Parameters of functions from Fig. 4.3 describing track dE/dx as a function of reconstructed momentum for a few particle species (STARsim MC). Units of parameters P_i are such that if one provides momentum in Eq. (4.7) in GeV/c the resultant offset of dE/dx MPV with respect to Bichsel parametrization is in GeV/cm, and the resultant σ parameter is unitless.

- 4.2.3 RP track acceptance and reconstruction efficiency
- 4.2.4 TPC track acceptance and reconstruction efficiency
- 4.2.5 TOF acceptance, reconstruction and track-matching efficiency
- 4.2.6 TPC vertex reconstruction efficiency
- 4.3 Particle energy loss
- 4.4 Background subtraction
- 4.5 Unfolding

5. Backgrounds

5.1 Non-exclusive background

5.2 Exclusive background (particle misidentification)

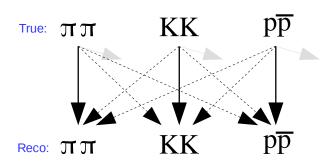


Figure 5.1: Graph illustrating the misidentification problem - the origin of exclusive background in selected samples. Gray arrows represent event rejection due to failed PID selection (C9). Solid black arrows represent successful identification, whereas dashed black lines show misidentification paths.

$$N_{R}^{\pi\pi} = \underbrace{\epsilon^{\pi\pi} \cdot N_{T}^{\pi\pi}}_{\text{true pion pairs}} + \underbrace{\lambda^{KK \to \pi\pi} \cdot N_{T}^{KK}}_{\text{kaon pairs reconstructed as pion pairs}}_{\text{kaon pairs reconstructed as pion pairs}} + \underbrace{\lambda^{p\bar{p} \to \pi\pi} \cdot N_{T}^{p\bar{p}}}_{\text{proton pairs reconstructed as pion pairs}} + \underbrace{N_{bkgd}^{\pi\pi}}_{\text{bkgd}}$$
(5.1a)
$$N_{R}^{KK} = \underbrace{\lambda^{\pi\pi \to KK} \cdot N_{T}^{\pi\pi}}_{\text{pion pairs reconstructed as kaon pairs}} + \underbrace{\epsilon^{KK} \cdot N_{T}^{KK}}_{\text{true kaon pairs}} + \underbrace{\lambda^{p\bar{p} \to KK} \cdot N_{T}^{p\bar{p}}}_{\text{proton pairs reconstructed as kaon pairs}} + \underbrace{N_{bkgd}^{KK}}_{\text{bkgd}}$$
(5.1b)
$$N_{R}^{p\bar{p}} = \underbrace{\lambda^{\pi\pi \to p\bar{p}} \cdot N_{T}^{\pi\pi}}_{\text{pion pairs reconstructed as proton pairs}} + \underbrace{\lambda^{KK \to p\bar{p}} \cdot N_{T}^{KK}}_{\text{kaon pairs reconstructed as proton pairs}} + \underbrace{\lambda^{KK \to p\bar{p}} \cdot N_{T}^{KK}}_{\text{true proton pairs}} + \underbrace{\lambda^{p\bar{p}} \cdot N_{D}^{p\bar{p}}}_{\text{true proton pairs}} + \underbrace{\lambda^{p\bar{p}} \cdot N_{D}^{p\bar{p}}}_{\text{bkgd}}}_{\text{true proton pairs}}$$
(5.1c)

Eqs. (5.1) can be written in the matrix form, as shown in Eq. (5.2).

$$\begin{bmatrix} N_{R}^{\pi\pi} - N_{bkgd}^{\pi\pi} \\ N_{R}^{KK} - N_{bkgd}^{KK} \\ N_{R}^{p\bar{p}} - N_{bkgd}^{p\bar{p}} \end{bmatrix} = \begin{bmatrix} \epsilon^{\pi\pi} & \lambda^{KK \to \pi\pi} & \lambda^{p\bar{p} \to \pi\pi} \\ \lambda^{\pi\pi \to KK} & \epsilon^{KK} & \lambda^{p\bar{p} \to KK} \\ \lambda^{\pi\pi \to p\bar{p}} & \lambda^{KK \to p\bar{p}} & \epsilon^{p\bar{p}} \end{bmatrix} \begin{bmatrix} N_{T}^{\pi\pi} \\ N_{T}^{KK} \\ N_{T}^{p\bar{p}} \end{bmatrix}$$

$$(5.2)$$

$$\begin{bmatrix} N_T^{\pi\pi} \\ N_T^{KK} \\ N_T^{p\bar{p}} \end{bmatrix} = \Lambda^{-1} \begin{bmatrix} N_R^{\pi\pi} - N_{bkgd}^{\pi\pi} \\ N_R^{KK} - N_{bkgd}^{KK} \\ N_R^{p\bar{p}} - N_{bkgd}^{p\bar{p}} \end{bmatrix}$$

$$(5.3)$$

6. Systematic errors

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