

# TP4 Homework 4

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## Fourier Transform and Gaussian Wave Functions

Fourier transform and Gaussian wave functions.

### 1 Fourier Transform

Determine the Fourier transform of

$$\Psi(x) = (\pi\omega_0^2)^{-1/4} e^{ip_0 x/\hbar - (x-x_0)^2/(2\omega_0^2)}$$

where  $x_0, p_0, \omega_0$  are real parameters.

**Solution**

### 2 Normalization

(b) Show that  $\psi$  is normalized

### 3 Mean Values

### 4 Adjoints

### 5 Variance

For the wave function in (a), show that the variances of position and momentum operator are  $\hat{x} = x_0 + \frac{\hbar}{2m\omega_0}$ ,  $\hat{p} = p_0 - \frac{\hbar m\omega_0}{2}$ . (5)

### 6 Differentiable Wave Function

f) For any differentiable wave function having sufficiently rapid decay as  $|x| \rightarrow \infty$  (as well as its derivative), show that  $\langle x \rangle(k) = i \frac{d}{dk} \langle k \rangle$ ,  $\langle k \rangle(x) = -i \frac{d}{dx} \langle x \rangle$ , (6) so Fourier transform exchanges position and momentum operators. Show that if  $\psi$  is an  $L^2$ -wave function, then  $\langle x \rangle = i \hbar \frac{d}{dk} \langle k \rangle$  has Fourier transform  $\langle k \rangle = -i \hbar \frac{d}{dx} \langle x \rangle$ . How are these statements related?

## 7 Hamiltonian for Schrodinger Equation

We now consider the Hamilton operator  $\hat{H} = \frac{\hat{p}^2}{2m} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$  for a free particle on the real line. We wish to solve the time-dependent Schrodinger equation  $i\hbar \frac{\partial \psi(x,t)}{\partial t} = \hat{H} \psi(x,t)$ ,  $\psi(x,0) = \psi_0(x)$  with initial wave function  $\psi_0(x)$  as in (a). Using (e), show that such a solution satisfies  $\frac{\partial \psi(k,t)}{\partial t} = -\frac{\hbar k^2}{2m} \psi(k,t)$ . Using this and the result of (a), find  $\psi(k,t)$ . Applying the inverse Fourier transform, find  $\psi(x,t)$ .

## 8 General Solution Schrodinger Equation

Using the same type of method as in (g), show that the general solution to the time-dependent Schrodinger equation for our Hamiltonian may be written as  $\psi(x,t) = \int_{-\infty}^{\infty} K(x,y,t) \psi_0(y) dy$ ,  $K(x,y,t) = \left(\frac{m}{2\pi i \hbar t}\right)^{1/2} e^{im(x-y)^2/2\hbar t}$ . Hint: You may use the convolution theorem: If  $f(x)$  and  $g(x)$  are decaying sufficiently rapidly, then  $(f \cdot g)(k) = \int_{-\infty}^{\infty} f(k-p)g(p)dp$ , together with a similar theorem for the inverse Fourier transform.