

## Syllabus

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- Background required
  - Vector calculus
  - Maxwell's equations for electro and magnetostatics; we will revise the latter
  - Scalar (and vector) potential
  - Perfect conductors in vacuum
  - Electric current
  - Definition of capacitance, resistance, and inductance (for steady currents)
- Topics
  - Electro and magnetostatics in depth
    - \* Laplace's elementary law and vector potential
    - \* Lorentz force and relativity
    - \* Weak form of Poisson's problem and Maxwell's inductance
    - \* Mechanical effects due to the electro and magnetostatic field
    - \* Multipole expansion for the scalar (in depth) and vector potential
  - The electrostatic field in presence of dielectric materials (in depth)
    - \* Polarization and polarization vector
    - \* Electric displacement and boundary conditions
    - \* Electric susceptibility
    - \* Macroscopic electric field inside a dielectric
    - \* Microscopic polarization phenomena
  - The magnetostatic field in presence of magnetic materials
    - \* Magnetization and magnetization vector
    - \* Magnetic field intensity and boundary conditions
    - \* Diamagnetism, paramagnetism, and ferromagnetism
  - The electromagnetic field
    - \* Electromagnetic induction
    - \* Displacement currents
    - \* Maxwell's equations for the electromagnetic field
    - \* Electromagnetic potentials
    - \* Energy of the electromagnetic field

- Grades

- Attendance: 7%
- Home assignments: 18%
- Midterm exam: 27%
- Final exam: 48%
- No changes to the above scheme will be allowed.

- Home assignments

There will be four (4) sets of problems, one for each topic.

- Set 1: Handed out on 2015-09-21; handed in on 2015-10-19.
- Set 2: Handed out on 2015-11-02; handed in on 2015-11-12.
- Set 3: Handed out on 2015-11-16; handed in on 2015-11-26.
- Set 4: Handed out on 2015-11-23; handed in on 2015-12-04.

- Tutorials

There will be six (6) tutorials. The first will be a review of magnetostatics (basic phenomena from PHYS 242). The next four will be a detailed review of the home assignments. The last one will be a review of the course before the final exam. The schedule for the tutorials will be decided in class, the first day of lectures. The tutorials will likely be scheduled a few days before the midterm and final exams, in order to prepare for these events.

- Office hours

Contact me (or the TAs) at [matteo.mariantoni@uwaterloo.ca](mailto:matteo.mariantoni@uwaterloo.ca). I prefer to meet with small group of students rather than single students, so that many people can benefit from the same discussion.

- Literature

- R.P.Feynman, R.B.Leighton, and M.Sands, The Feynman Lectures on Physics - Volume 2, any edition
- E.M.Purcell and D.J.Morin, Electricity and Magnetism, Cambridge University Press; 3rd edition, 2013

- My notes on Learn

## PHYS 342 - EM 02

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### Lecture 1

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#### 1.1 Review of Electrostatics

Consider a volume charge distribution with density  $\rho \in C^0(\Omega)$  in a domain  $\Omega$  of the 3D Euclidean space.

According to Gauss' theorem

$$\oiint_{\Sigma} \vec{E} \cdot \vec{n} dA = \frac{1}{\epsilon_0} \iiint_{\Omega} \rho dV \quad (1.1)$$

Where  $\Sigma$  is a surface enclosing  $\Omega$ . This equation can be written in differential form as (if  $\rho$  is continuous and limited)

$$\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho \quad (1.2)$$

which is also valid when  $\rho = 0$ .

The irrotational property of  $\vec{E}$  can be written in integral form as

$$\oint \vec{E} \cdot \vec{t} dl = 0 \quad (1.3)$$

and, in differential form, as

$$\vec{\nabla} \times \vec{E} = \vec{0} \quad (1.4)$$

In the case of a surface charge distribution with density  $\rho$  in  $\Omega$ , the local form of Gauss' theorem and the irrotational property is

$$(\vec{E}_1 - \vec{E}_2) \cdot \vec{n} = \frac{1}{\epsilon_0} \sigma \quad (1.5)$$

and

$$(\vec{E}_1 - \vec{E}_2) \cdot \vec{t} = 0 \quad (1.6)$$

respectively (cf. conventions in the next figure).

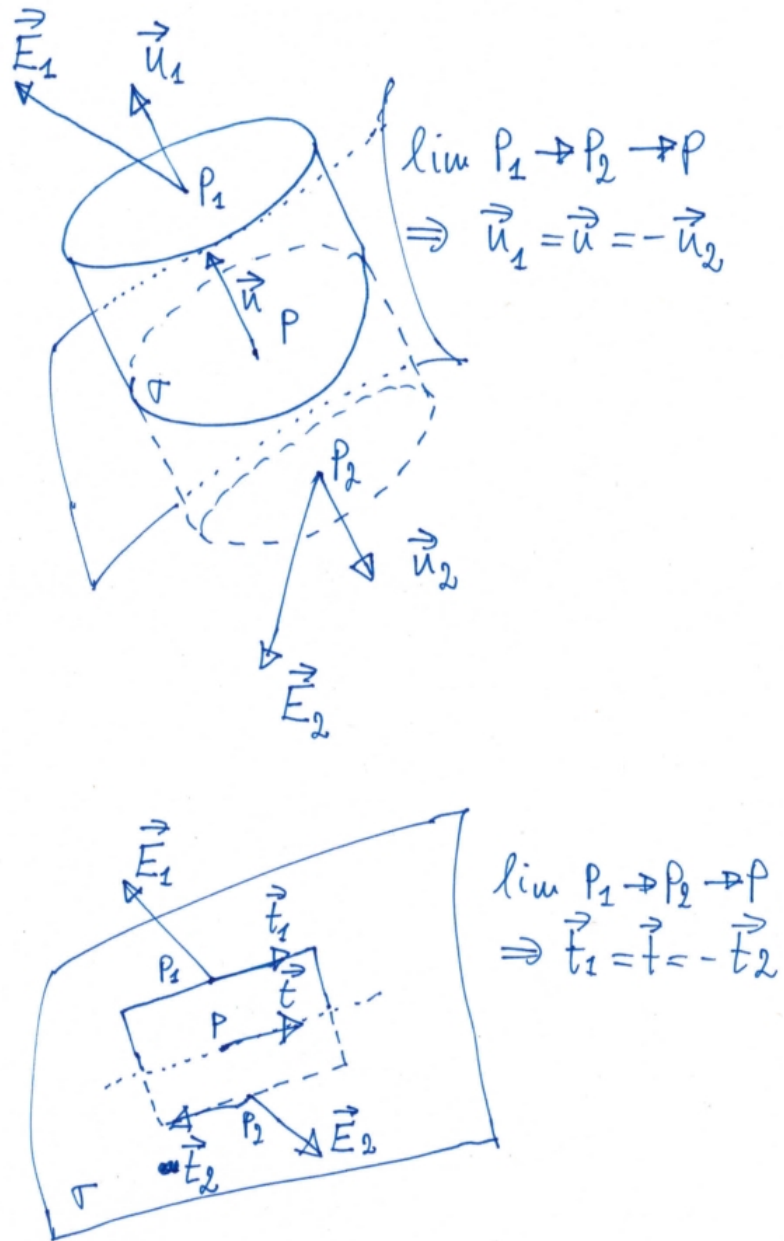


Figure 1.1

In a simply connected (or star) domain, the electrostatic field  $\vec{E}$  is not only irrotational, but also conservative. In this case, a scalar potential  $\phi$  can be defined, such that

$$\vec{E} = -\vec{\nabla}\phi \quad (1.7)$$

From this definition and the differential form of Maxwell's equations for  $\vec{E}$ , it can be shown that

$$\nabla^2\phi = 0 \quad (1.8)$$

when  $\rho = 0$  (Laplace equation) and

$$\nabla^2\phi = -\frac{1}{\epsilon_0}\rho \quad (1.9)$$

when  $\rho \neq 0$  (Poisson equation).

## 1.2 Perfect Conductors in Vacuum

Consider a perfect and homogeneous conductor in vacuum.

The conductor is said to be in electrostatic equilibrium when no macroscopic motion of charges exists in the conductor. Under these conditions

$$\vec{E} = \vec{0} \quad (1.10)$$

at each point inside the conductor. Following the notation in the figure below,

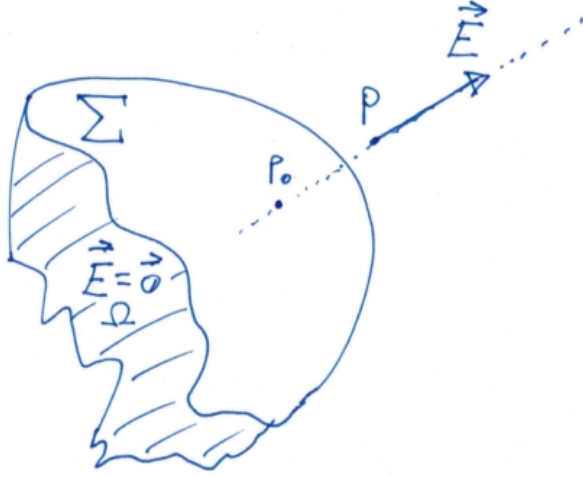


Figure 1.2

when the conductor is in an external field  $\vec{E}_{ex}$ , a reaction field  $\vec{E}_r$  is generated until

$$\vec{E} = \vec{E}_{ex} + \vec{E}_r = \vec{0} \quad (1.11)$$

at each point inside the conductor. Under these conditions, a surface charge density  $\sigma$  is generated on the conductor surface.  $\sigma$  depends on the point  $P_0$  on the surface. According to Coulomb's theorem

$$\lim_{P \rightarrow P_0^+} E_n(P) = \frac{\sigma(P_0)}{\epsilon_0} \quad (1.12)$$

where  $P_0^+$  indicates that the point is located on the outer skin of the conductor.

From the definition of  $\phi$ , it follows

$$\phi = \text{const} \Big|_{\Omega \cup \Sigma} \quad (1.13)$$

and

$$E_n = - \frac{\partial}{\partial n} \phi \Big|_{\Sigma^+} \quad (1.14)$$

where  $\Sigma^+$  indicates the outer skin of the conductor.

If the conductor is charged with  $Q$ ,

$$Q = \epsilon_0 \oint_{\Sigma' \supset \Sigma} E_n dA = \oint_{\Sigma} \sigma dA \quad (1.15)$$

In general,

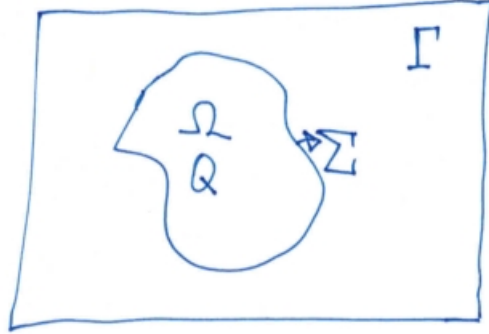


Figure 1.3

- i)  $\phi \in C^0(\Gamma \cup \Sigma)$
- ii)  $\vec{\nabla}^2 \phi = 0 \quad \forall P \in \Gamma$
- iii)  $\phi = \text{const} \quad \forall P \in \Omega \cup \Sigma$
- iv)  $-\iint_{\Sigma^+} \frac{\partial}{\partial n} \phi dA = \frac{1}{\epsilon_0} Q$
- v)  $\lim_{P \rightarrow \infty} \phi(P) = 0$

- Definition of capacitance

Given two conductors,  $A$  and  $B$ , with charges  $+|Q|$  and  $-|Q|$ , respectively, the capacitance of the system is defined as

$$C \equiv \frac{|Q|}{\Delta\phi} \quad (1.16)$$

where  $\Delta\phi = \phi_A - \phi_B$ .