

Simulating Schwarzschild and Kerr Black Hole Geodesics

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MSc Physics
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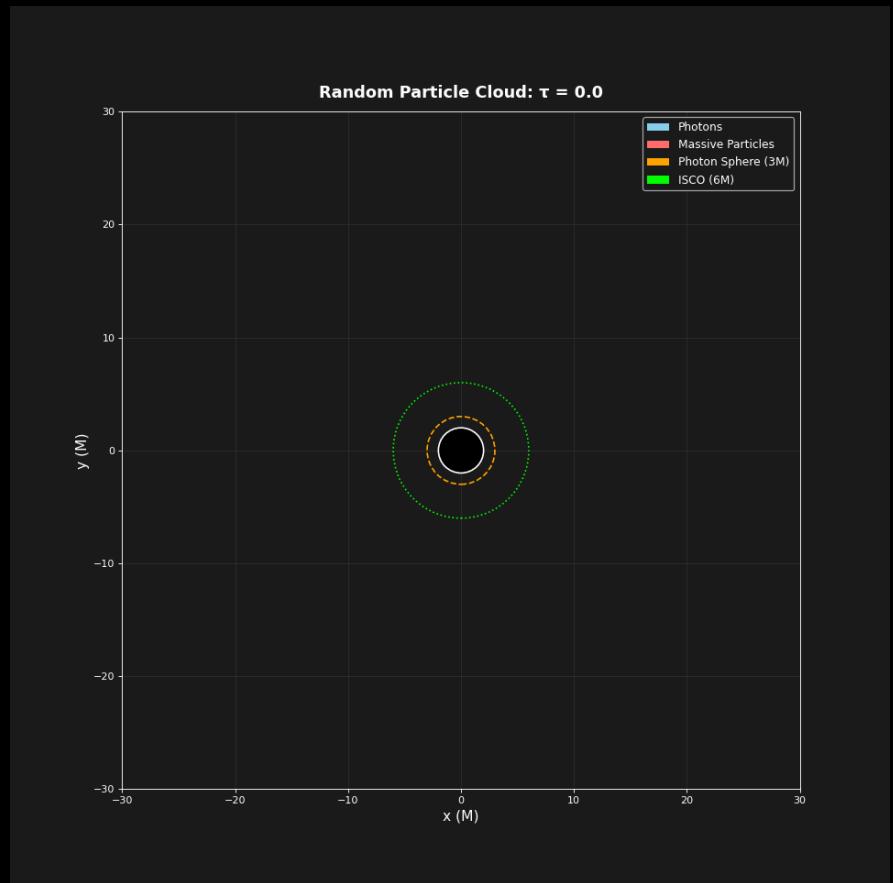


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Introduction

- **Problem – Solve geodesic equations in curved spacetime for timelike and null trajectories**

- **Geodesic Equation:**

- $$\frac{d^2x^\mu}{d\lambda^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\lambda} \frac{dx^\beta}{d\lambda} = 0$$

- **Christoffel Symbols:**

- $$\Gamma_{\alpha\beta}^\mu = \frac{1}{2} g^{\mu\nu} (\partial_\alpha g_{\nu\beta} + \partial_\beta g_{\nu\alpha} - \partial_\nu g_{\alpha\beta})$$

- **Metrics Considered:**

- Schwarzschild ($a = 0$): Spherical Symmetry
- Kerr ($a \neq 0$): Axial Symmetry

- **Numerical Method:**

- DOP853 (8th order RK)

Schwarzschild – Theory and Implementation

- **Schwarzschild Metric (Boyer-Lindquist, $\theta = \pi/2$):**
 - $ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 d\phi^2$, where $f(r) = 1 - \frac{r_s}{r} = 1 - \frac{2M}{r}$
- **Killing Vectors & Conserved Quantities:**
 - Time translation: $E = -p_t = f(r) u^t$
 - Axial rotation: $L = p_\phi = r^2 u^\phi$
- **Geodesic Equations from Christoffel Symbols (equatorial, $\theta = \pi/2$):**
 - $\frac{d^2 t}{d\tau^2} = (2M/r^2 f) u^t u^r$
 - $\frac{d^2 r}{d\tau^2} = -(Mf/r^2) u^t u^t + (M/r^2 f) u^r u^r + (r - 2M) u^\phi u^\phi$
 - $\frac{d^2 \phi}{d\tau^2} = -\frac{2}{r} u^r u^\phi$
- **Normalization:**
 - $g_{\mu\nu} u^\mu u^\nu = -1$ (**timelike**) | $g_{\mu\nu} u^\mu u^\nu = 0$ (**null**)

Schwarzschild – Theory and Implementation

Implementation Details:

- State Vector:
 - $\mathbf{y} = [t, r, \phi, \mathbf{u}^t, \mathbf{u}^r, \mathbf{u}^\phi]$
- Initial Conditions (from conserved quantities):
 - Given: r_0, φ_0, E, L ($orb = L/E$)
 - $u_t = E/f(r_0)$
 - $u^\varphi = L/r_0^2$
 - from normalization constraints:
 - Null: $(u^r)^2 = E^2 - f(r_0) \frac{L^2}{r_0^2}$
 - Timelike: $(u^r)^2 = E^2 - \left(1 + \frac{L^2}{r_0^2}\right) f(r_0)$

Schwarzschild – Results: Null Geodesics

Results:

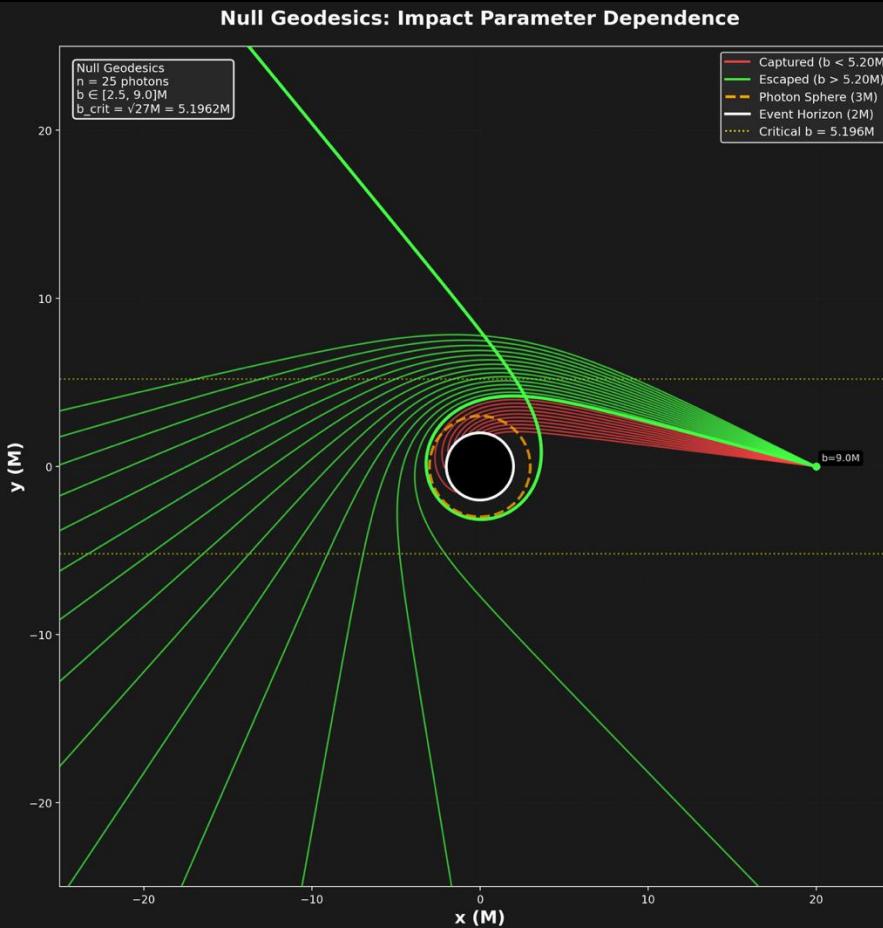
- $b < b_{crit} \rightarrow$ captured, eventually crosses horizon
- $b > b_{crit} \rightarrow$ scattered, returns to $r \rightarrow \infty$
- $b \approx b_{crit} \rightarrow$ unstable circular orbit, sensitive to initial cond.
- Clear bifurcation at critical impact parameter

Constraint Violations (over full trajectory):

$$\text{Energy: } \left| \frac{\Delta E}{E} \right| < 10^{-15}$$

$$\text{Angular Momentum: } \left| \frac{\Delta L}{L} \right| < 10^{-15}$$

$$\text{Null constraint: } \left| g_{\mu\nu} u^\mu u^\nu \right| < 10^{-14}$$



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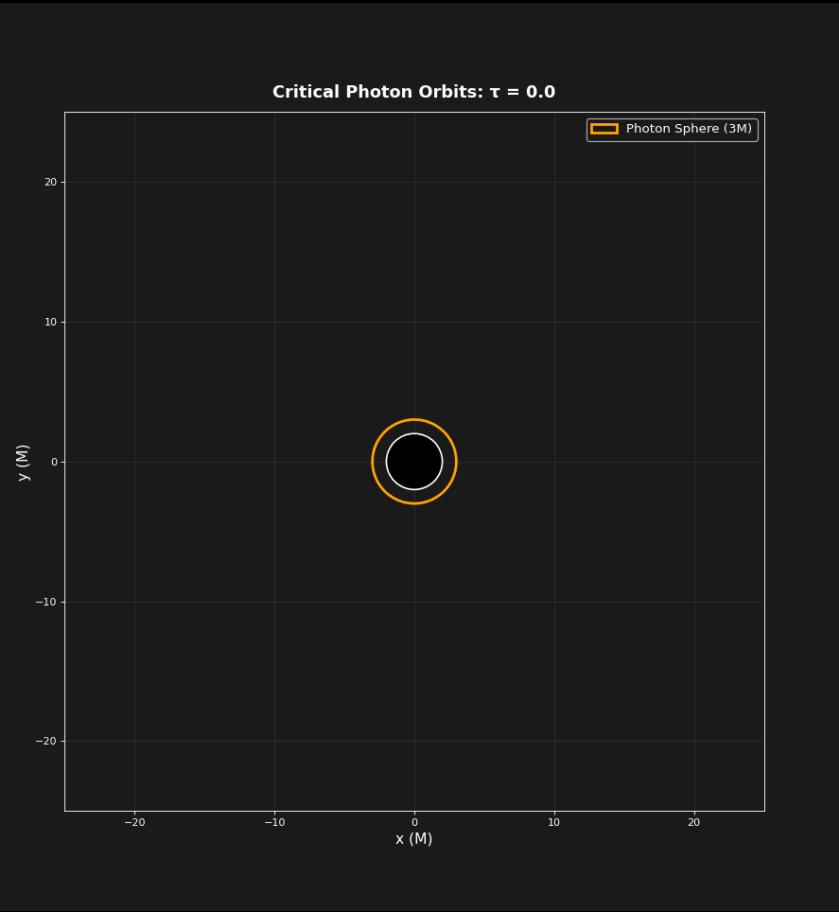
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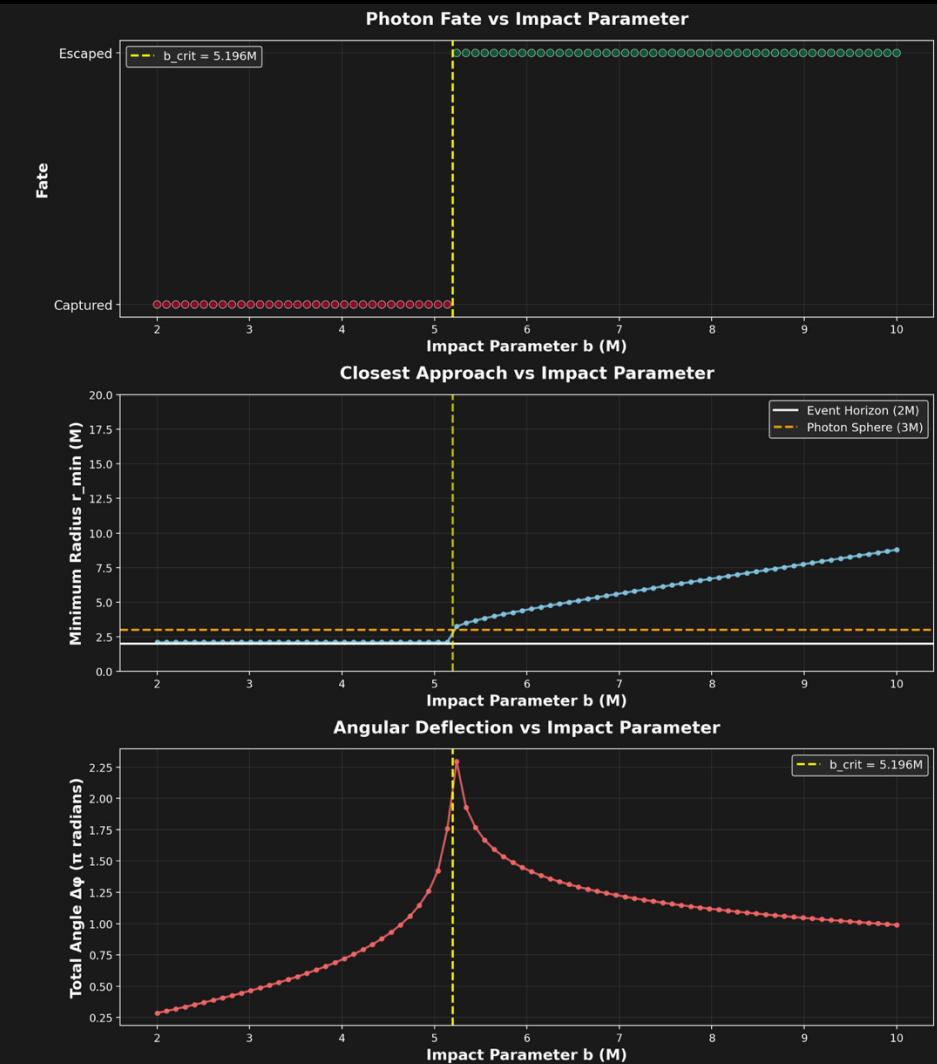
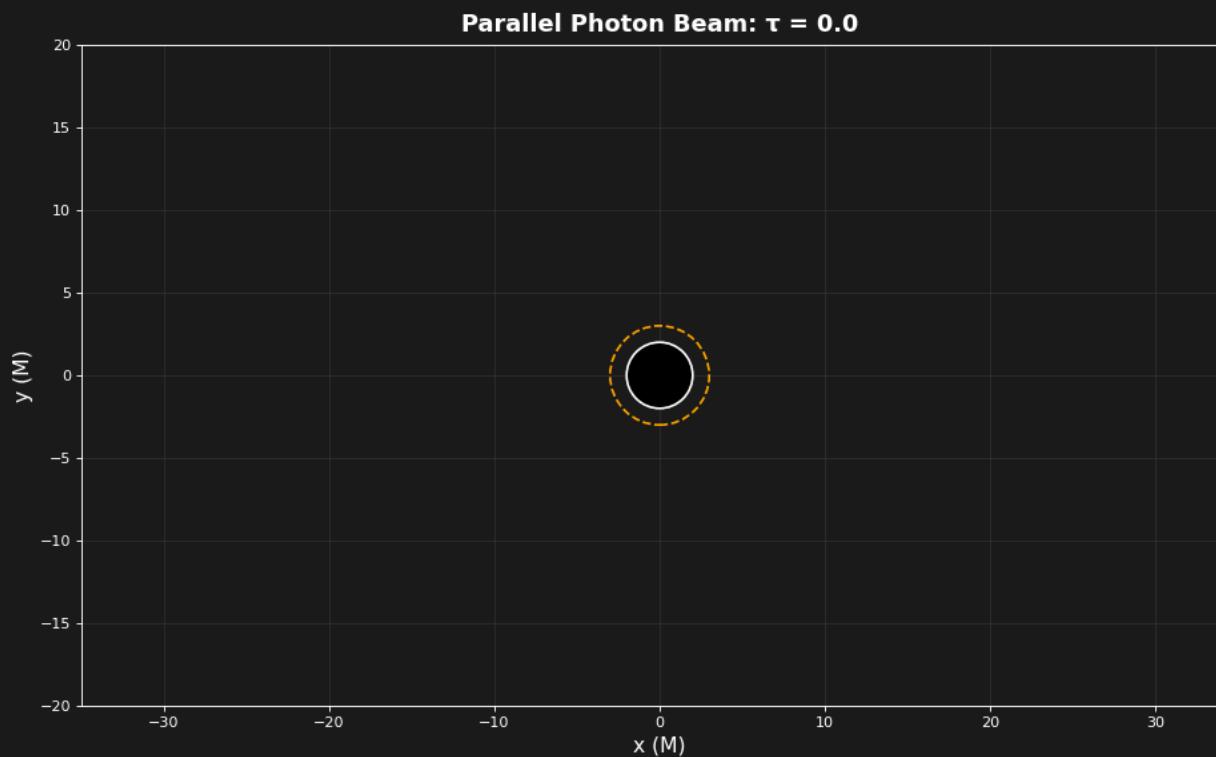
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Schwarzschild – Scattering Experiment



Schwarzschild – Results: Timelike Geodesics

Energy Constraint Analysis:

Effective Potential:

$$V_{\text{eff}}^{\text{massive}} = (1 + L^2/r^2)(1 - 2M/r)$$

Radial turning points satisfy:

$$E^2 = V_{\text{eff}}(r_{\text{crit}})$$

For circular orbits at r :

$$E_{\text{circ}}(r) = \sqrt{(r - 2M)/(r - 3M)} / \sqrt{r}$$

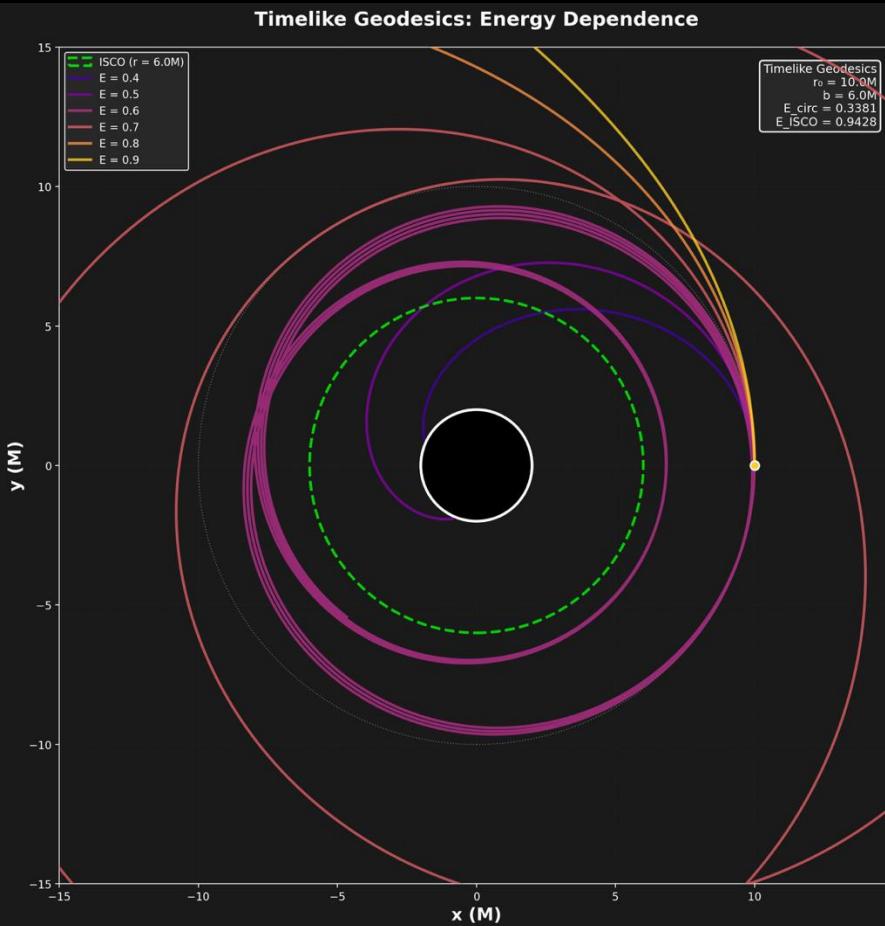
$$L_{\text{circ}}(r) = M\sqrt{r/(r - 3M)}$$

ISCO (innermost stable circular orbit) Analysis ($r=6M$):

$$E_{\text{ISCO}} \approx 0.943$$

$$L_{\text{ISCO}} \approx 3.464M$$

$E < E_{\text{circ}}(r) \rightarrow$ Bound orbits



Schwarzschild – Results: Timelike Geodesics

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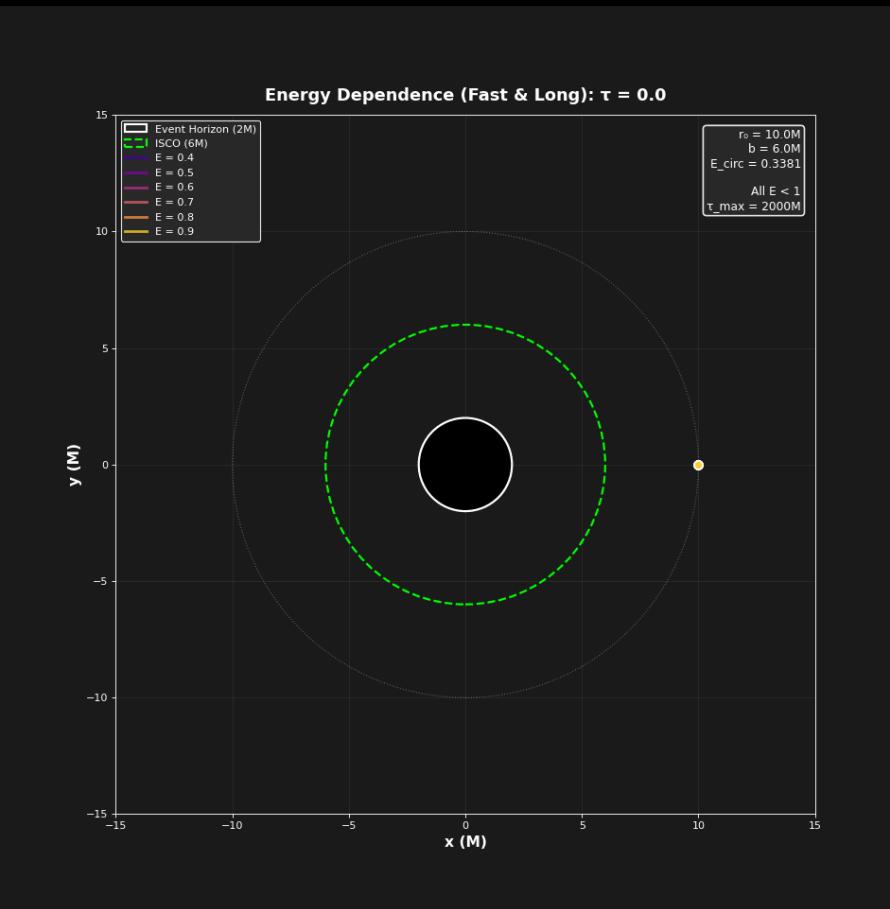
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Kerr – Theory and Implementation

- Kerr Metric (Boyer-Lindquist, $\theta = \pi/2$):

- $- ds^2 = -\left(1 - \frac{2Mr}{\Sigma}\right) dt^2 - \frac{4Mra}{\Sigma} dt d\phi + \frac{\Sigma}{\Delta} dr^2 + \frac{A}{\Sigma} d\phi^2$

- $-$ Where:

- $- \Sigma(r, \theta) = r^2 + a^2 \cos^2 \theta,$
- $- \Delta(r) = r^2 - 2Mr + a^2,$
- $- A(r, \theta) = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta$

- $-$ At equator ($\theta = \pi/2$):

- $- \Sigma = r^2$
- $- A = (r^2 + a^2)^2 - a^2 \Delta = r^4 + a^2 r^2 + 2Ma^2 r$

- Metric components (equatorial):

- $- g_{\phi\phi} = r^2 + a^2 + \frac{2Ma^2}{r}$
- $- g_{rr} = \frac{r^2}{\Delta}$
- $- g_{t\phi} = -\frac{2Ma}{r}$
- $- g_{tt} = -\left(1 - \frac{2M}{r}\right)$

- Event Horizons ($\Delta = 0$):

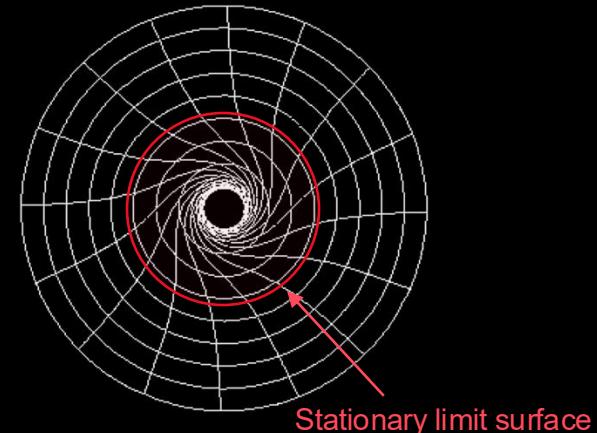
- $- r_{\pm} = M \pm \sqrt{M^2 - a^2}$
- $-$ Outer horizon: r_+
- $-$ Inner horizon: r_- ,

Extremal limit: $a = M \Rightarrow r_{\pm} = M$

- Ergosphere ($g_{tt} = 0$):

- $- r_{\text{ergo}} = 2M$
- $- r_+ < r < r_{\text{ergo}}$

Visualization of Frame Dragging



- Frame Dragging (angular velocity of ZAMO¹):

- $- \omega_H = \frac{a}{r_+^2 + a^2}, \text{ at horizon}$
- $- \omega = -\frac{g_{t\phi}}{g_{\phi\phi}} = \frac{2Mra}{A}$

Kerr – Theory and Implementation

Non-zero Christoffel Symbols (equatorial):

Limit $a \rightarrow 0$: Schwarzschild

$$\Gamma_{tr}^\phi = -\frac{Ma}{\Sigma^2 \Delta},$$

$$\Gamma_{r\phi}^\phi = \frac{1}{r},$$

$$\Gamma_{t\phi}^r = \frac{2aM^2 r}{\Sigma^3},$$

$$\Gamma_{\phi\phi}^r = -\frac{\Delta r}{\Sigma},$$

$$\Gamma_{rr}^r = \frac{Mr^2 - Ma^2 - r\Delta}{\Sigma \Delta},$$

$$\Gamma_{tt}^r = \frac{M\Delta(r^2 - a^2)}{\Sigma^3},$$

$$\Gamma_{r\phi}^t = \frac{Ma(r^2 - 2Mr + a^2)}{\Sigma^2 \Delta},$$

$$\Gamma_{tr}^t = \frac{Mr}{\Sigma^2 \Delta},$$

Geodesic Equations¹⁾ (equatorial):

$$\frac{d^2\phi}{d\tau^2} = -2, \Gamma_{r\phi}^\phi, u^r u^\phi - 2, \Gamma_{tr}^\phi, u^t u^r,$$

$$\frac{d^2r}{d\tau^2} = -\Gamma_{tt}^r(u^t)^2 - \Gamma_{rr}^r(u^r)^2 - \Gamma_{\phi\phi}^r(u^\phi)^2 - 2, \Gamma_{t\phi}^r, u^t u^\phi,$$

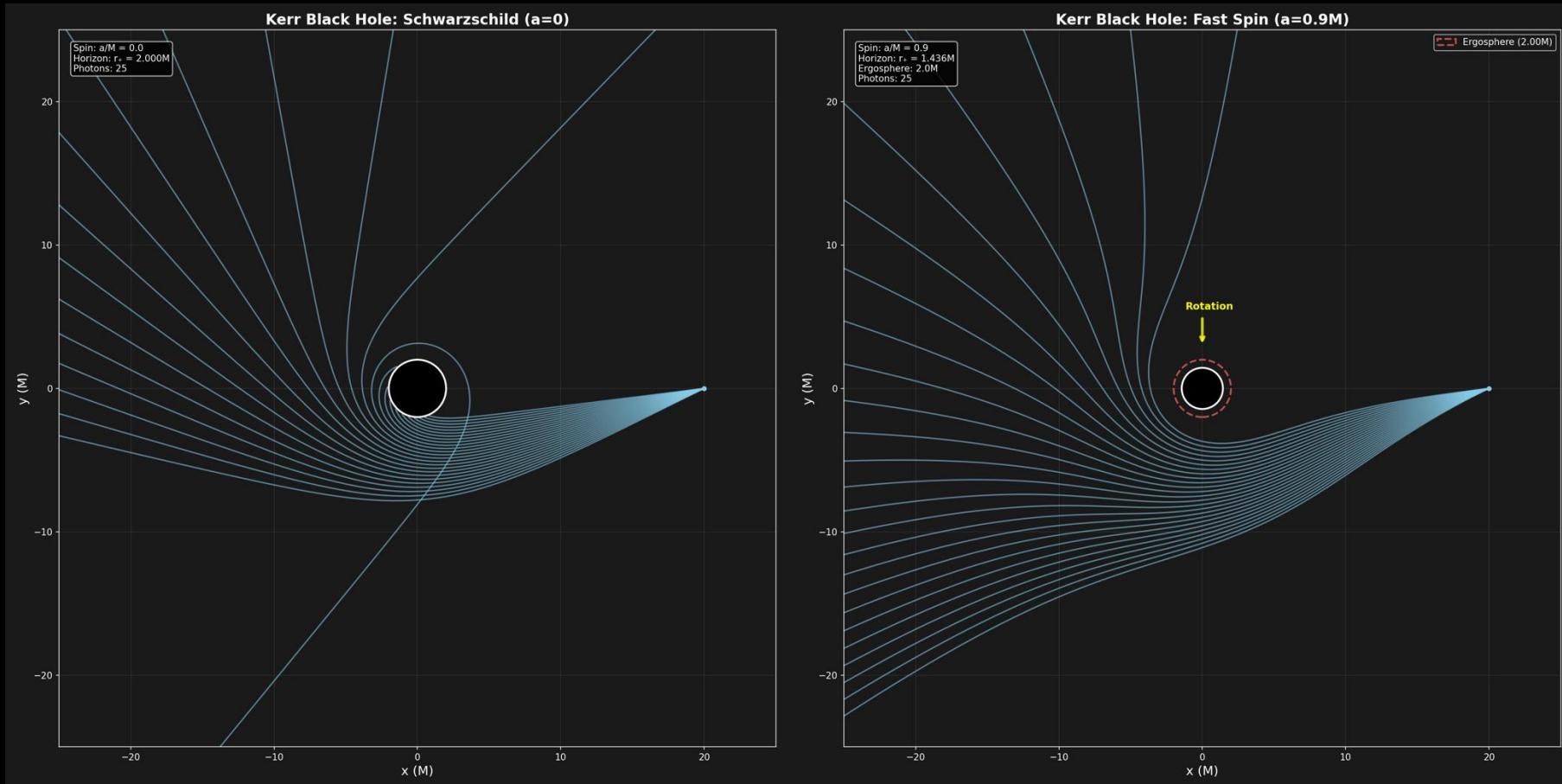
$$\frac{d^2t}{d\tau^2} = -2, \Gamma_{tr}^t, u^t u^r - 2, \Gamma_{r\phi}^t, u^r u^\phi,$$

Conservation (from Killing vectors):

$$L = p_\phi = g_{t\phi} u^t + g_{\phi\phi} u^\phi$$

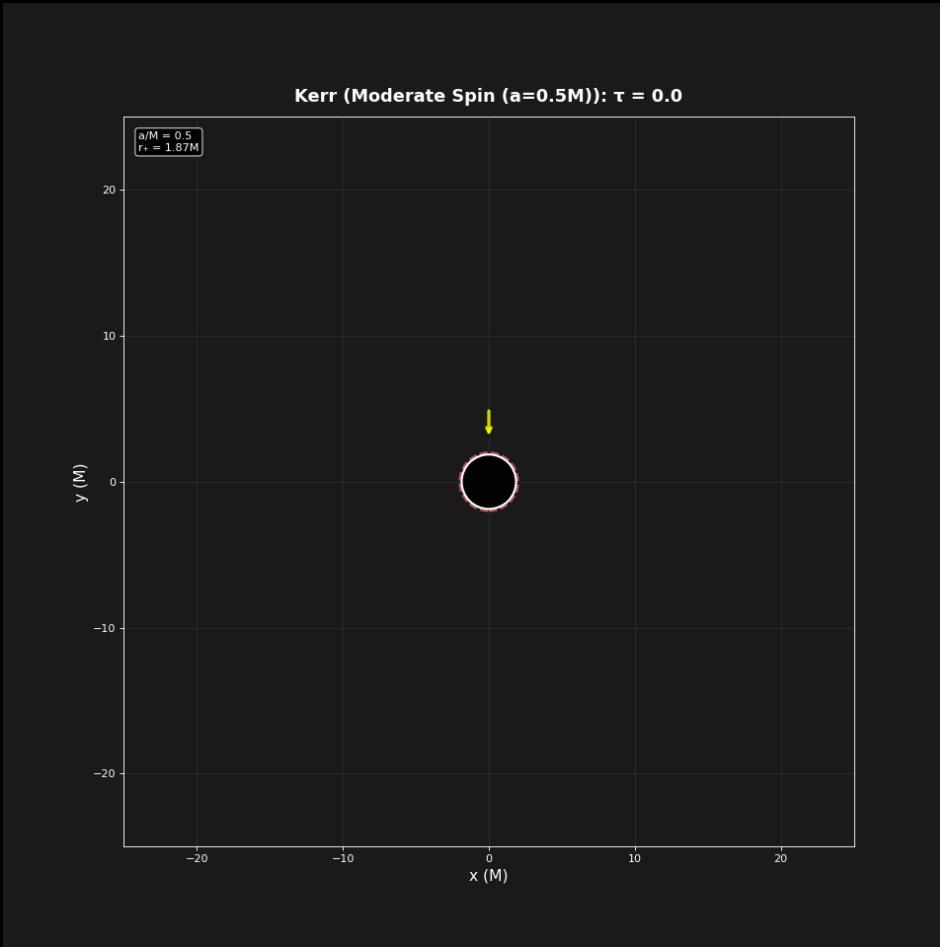
$$E = -p_t = -(g_{tt} u^t + g_{t\phi} u^\phi),$$

Kerr – Results: Null Geodesics



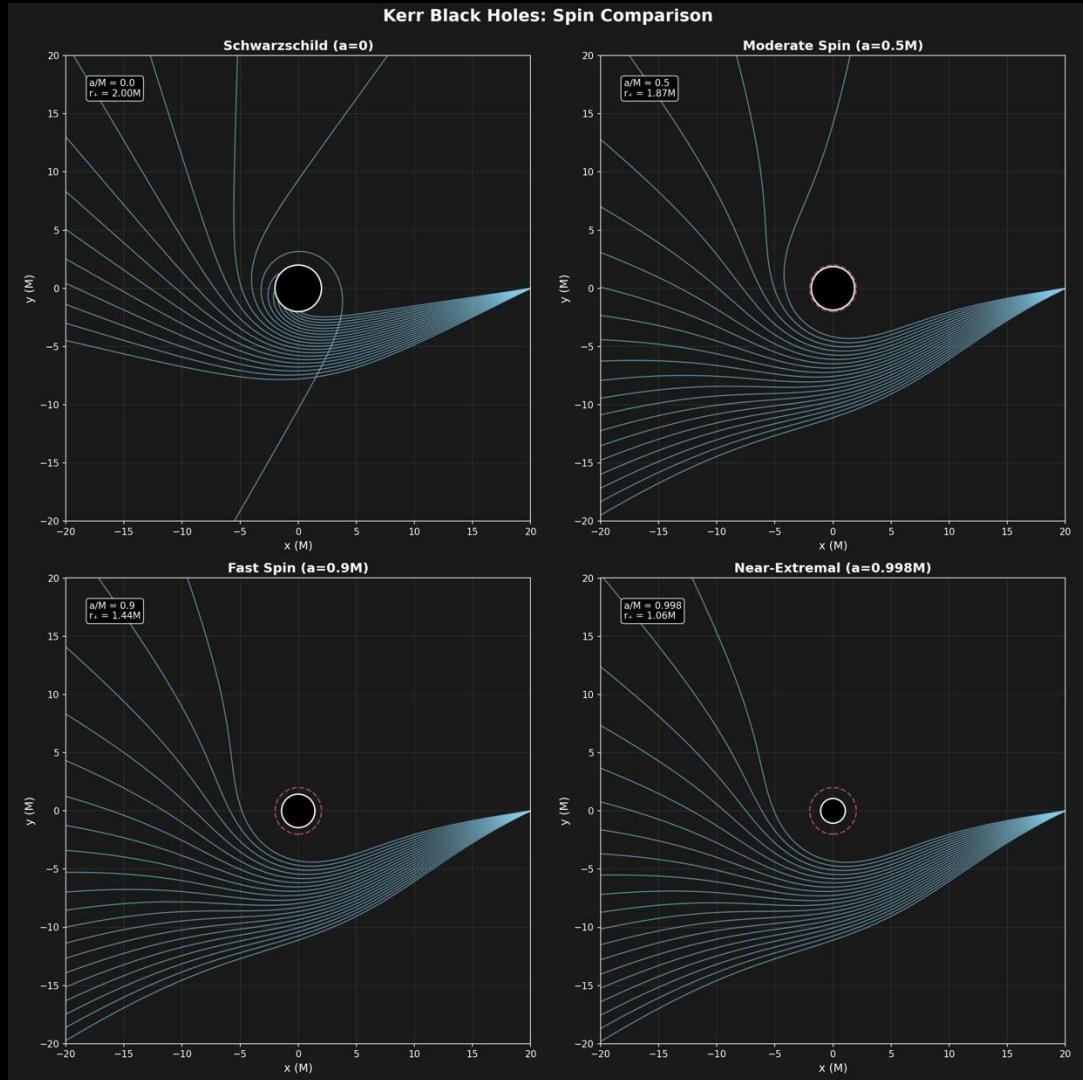
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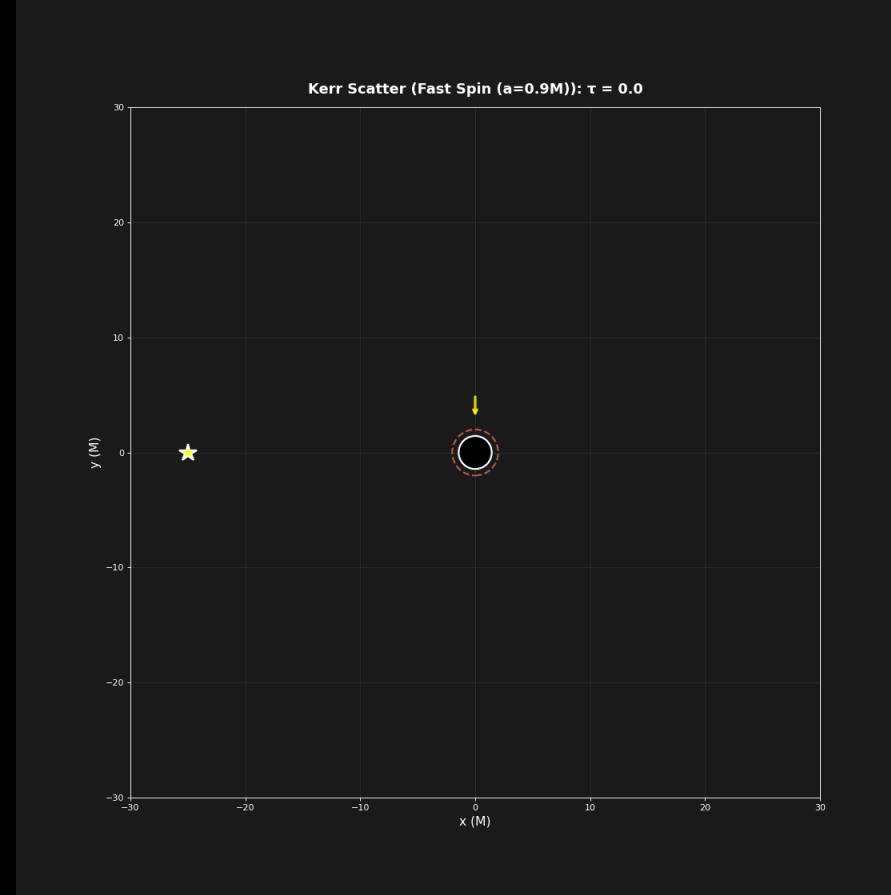
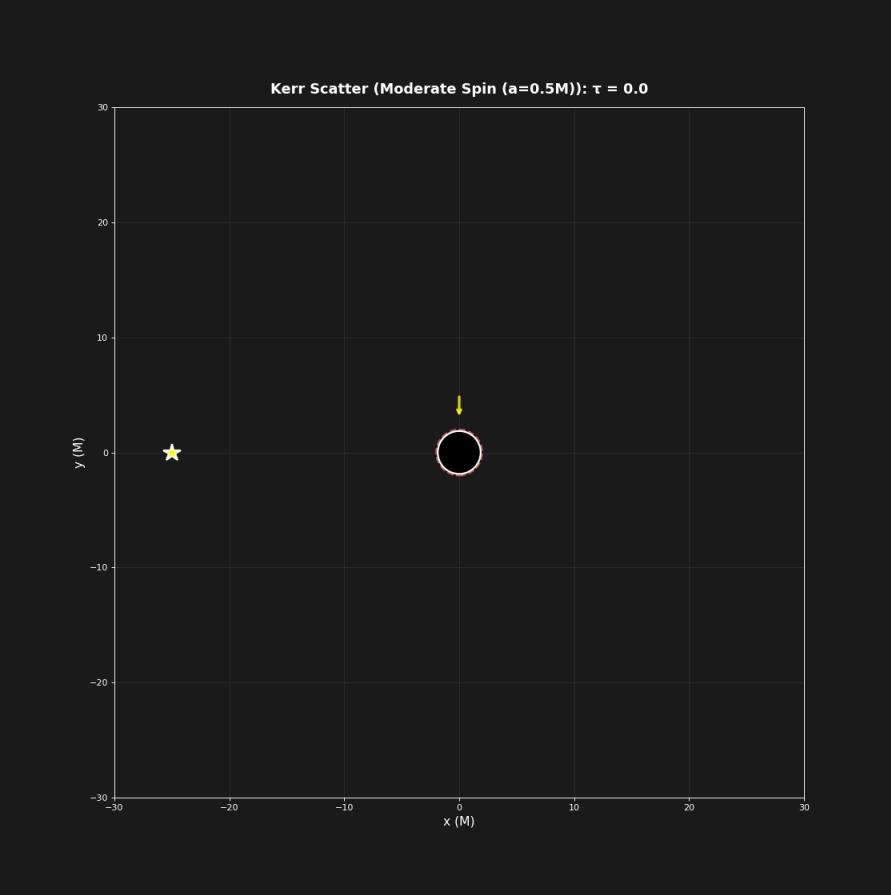


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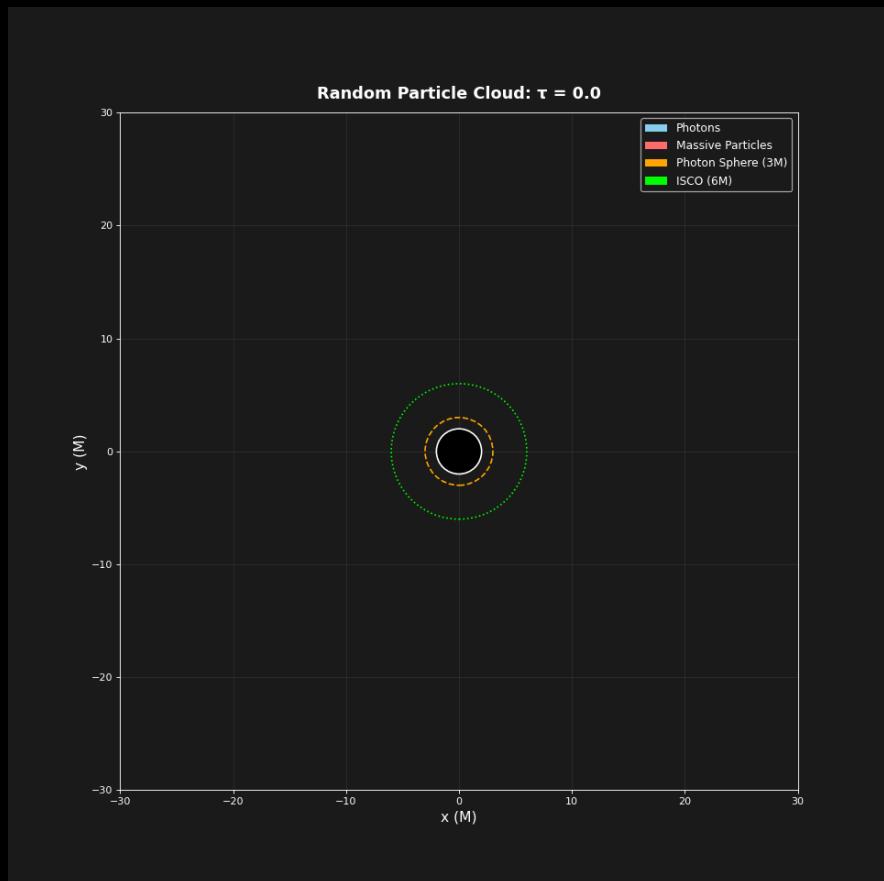


Kerr – Results: Null Geodesics



Conclusion

- Developed a fully general **geodesic simulator** for Schwarzschild and Kerr black holes
- Verified relativistic phenomena:
 - Photon Sphere, ISCO, impact parameters
 - Ergosphere behavior, frame dragging
- ~4000 loc
- Source code, simulations Jupyter Notebook user demo and references can be found on GitHub:
<https://github.com/fynnhufler/General-Relativity-Simulations>



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