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NOVA SCHOOL OF  
SCIENCE & TECHNOLOGY

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Bsc. in Physics Engineering

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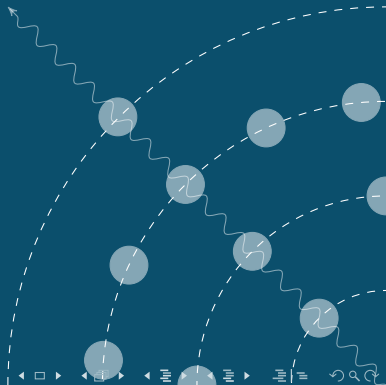
Advisor: Prof. Mauro Guerra

Co-Advisor: Prof. Jorge Machado

*NOVA School of Science and Technology*

# X-ray Resonant Raman Scattering

Spectra simulation from first principles  
for Copper below the ionization threshold  
using high-performance computing



# Overview

## ① Theoretical Introduction

Characteristic x-rays

The Hamiltonian

State-of-the-art

## ② Atomic structure calculations

The system at study

Level calculations

Transition computations

Fundamental atomic  
parameters

## ③ Spectra simulation

Line shape

Photoexcitation

Photoionization

The synthetic spectrum

## ④ Spectral Analysis

The theoretical results

Comparison with  
experimental data

## ⑤ A new parallelization code

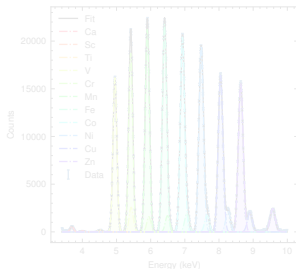
The MPI approach

Speedup comparison

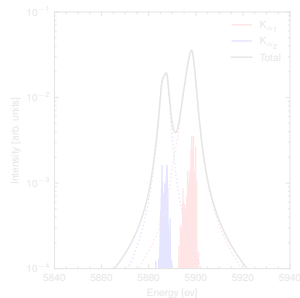
# X-ray applications



(a) Imaging purposes



(b) Sample quantification



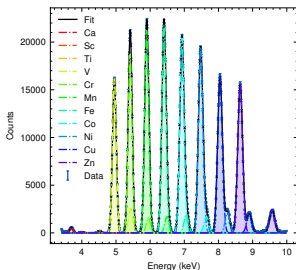
(c) Fundamental parameters

Figure: Application examples of x-ray radiation

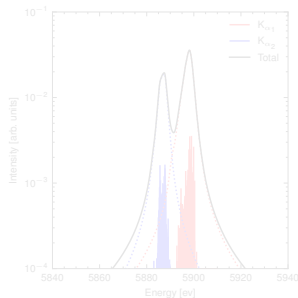
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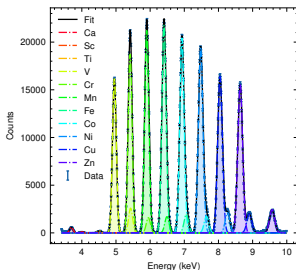
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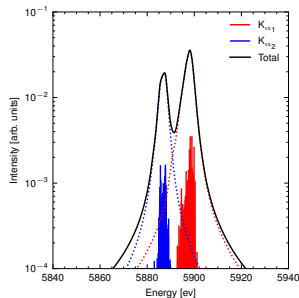
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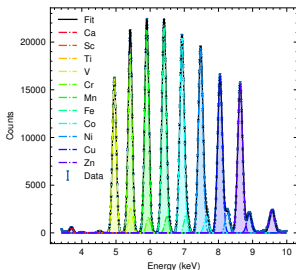
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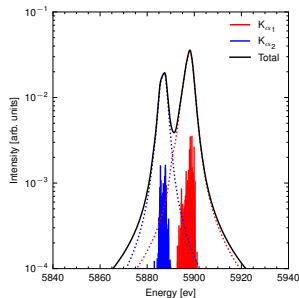
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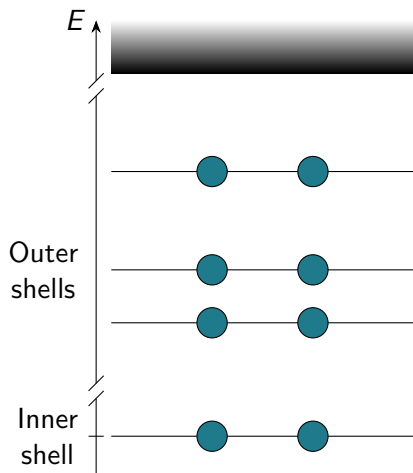
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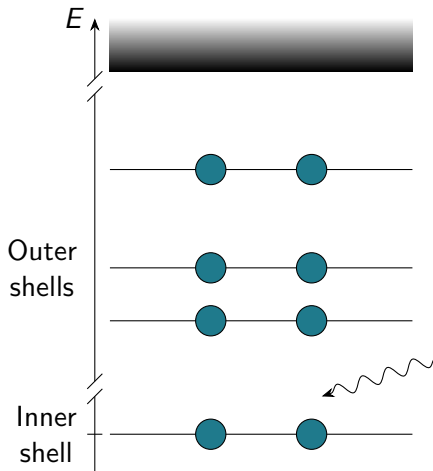
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# Vacancy generation and relaxation processes



- Bound state system
- Energy transfer
- Ionization
- Vacancy generated
- Atomic Relaxation

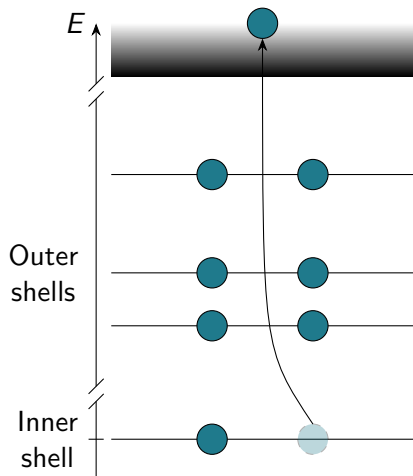
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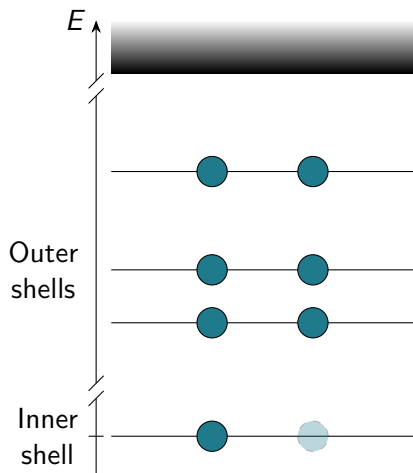


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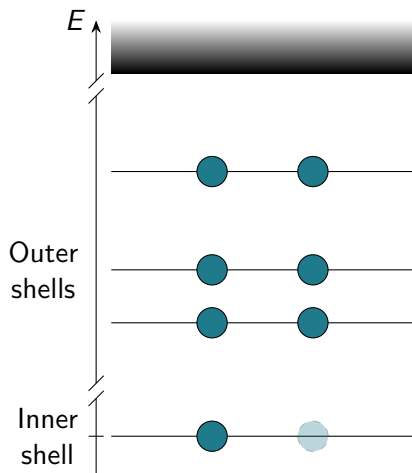
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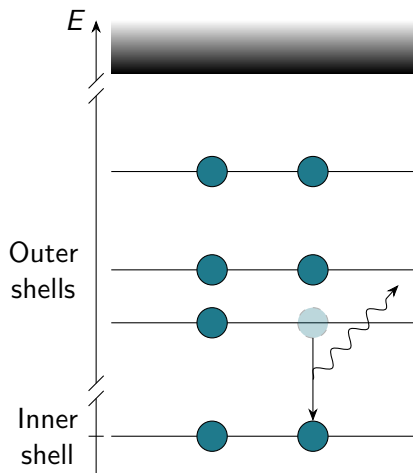
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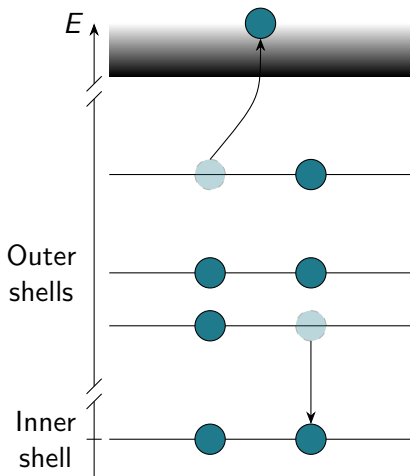
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# Alternative processes

## Shake processes

Post-ionization, the different number of particles in the system leads to a change in the Hamiltonian. This leads to the lack of orthogonality between non-equivalent states and free-wave wavefunctions in the initial and final configurations.

$$\langle 1s_i | 2s_f \rangle \neq 0 \quad \langle 1s_i | e_{\text{free}}^- \rangle \neq 0$$

### Shake-up

Excitation of extra electron(s).

### Shake-off

Ionization of extra electron(s).



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## Shake Processes

The probability of these processes occurring is extremely dependent on the incident energy.

Since for this thesis, the energy ranges studied were below the ionization threshold, no shake processes were accounted for.

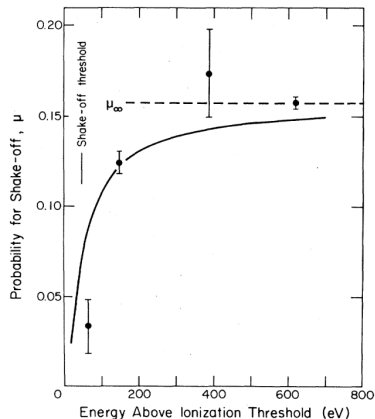
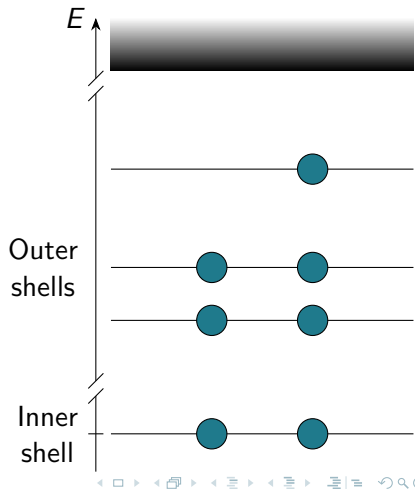


Figure: Thomas model for shake probability. [1]

## Other processes

### Photoexcitation

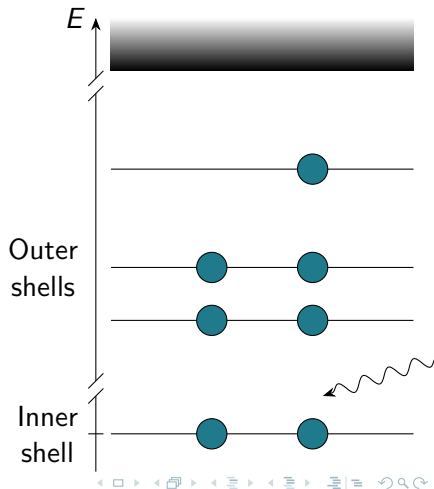
The most relevant process, for the scope of this work, is that of resonant photoexcitation. In it, instead of a vacancy generated due to sending one electron to the continuum, it has, in turn, remained bound but in a higher energy level.



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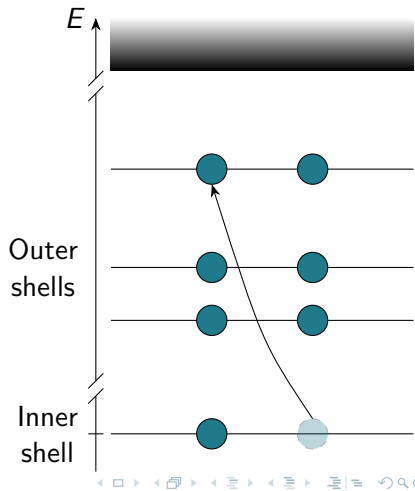
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# Schrödinger's Hamiltonian

In its most basic form, for a "classical" atom (nucleus + electrons), and when relativistic effects are not taken into account, the considered Hamiltonian follows the one use in Schrödinger's equation:

$$\hat{H} = \sum_i^N \frac{1}{2} \nabla_i^2 - \frac{Z}{r_i} + \sum_{j>i} \frac{1}{r_{ij}}$$

It incorporates:

- The kinetic energy of the electron.
- The potential energy of the electron-nucleus attraction.
- The potential energy of the electron-electron repulsion.

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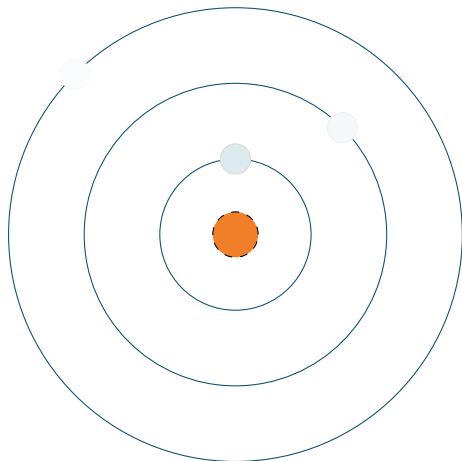
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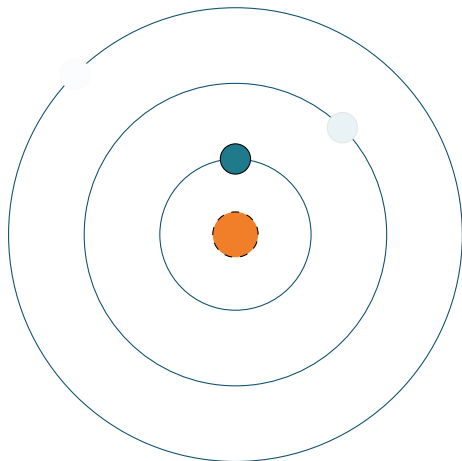
# Solving the non-relativistic many-body problem

Due to the complexity introduced by the many bodies in the system, and their interactions, a numerical method needs to be employed as to obtain eigenfunctions for this Hamiltonian.



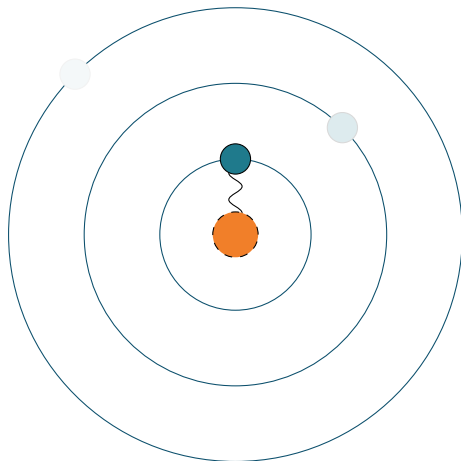
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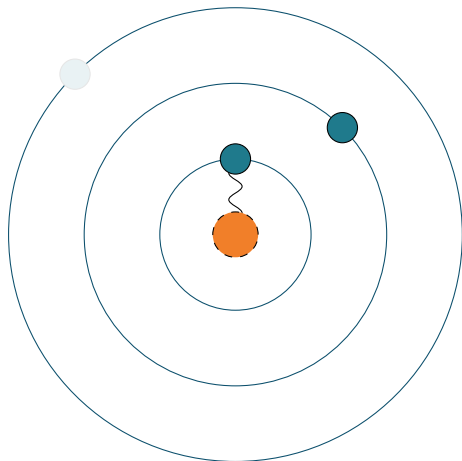
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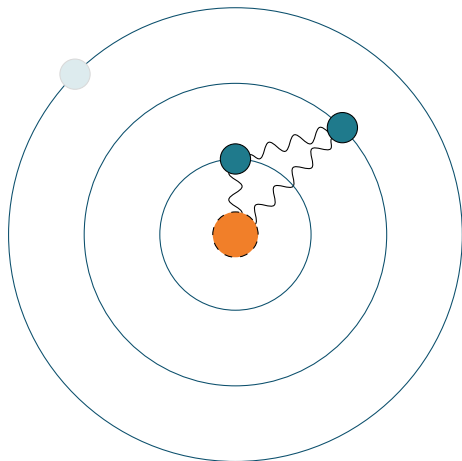
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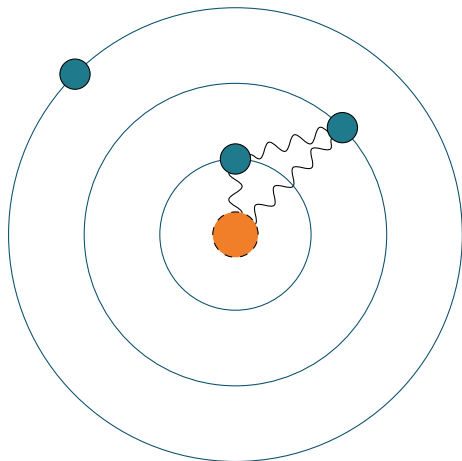
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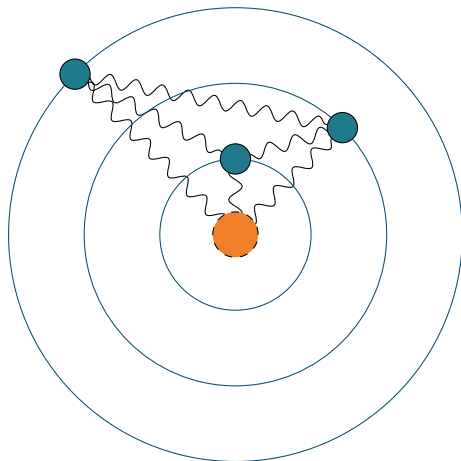
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# Solving the non-relativistic many-body problem

Each of the electrons' wavefunctions,  $u$ , are composed as a product of a spatial part,  $\psi$ , and one related to the electron's spin  $\chi$ .

$$u = \psi\chi$$

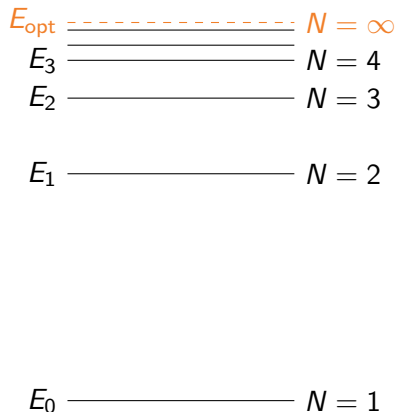
The system's wavefunction should then be written as a Slater determinant as to account for anti-symmetry and the fermionic nature of the electrons.

$$\frac{1}{\sqrt{N!}} \begin{vmatrix} u_1(x_1) & u_2(x_1) & \dots & u_N(x_1) \\ u_1(x_2) & u_2(x_2) & \dots & u_N(x_2) \\ \vdots & \vdots & \ddots & \vdots \\ u_1(x_N) & u_2(x_N) & \dots & u_N(x_N) \end{vmatrix}$$



# Solving the non-relativistic many-body problem

Through a **self-consistent** field approach, the method solves, for each cycle, a set of integro-differential equations as a way to compute new wavefunctions and the new energy for the system.



# Solving the non-relativistic many-body problem

This process is then repeated up until the energy difference in-between two steps is under a pre-defined benchmark value, as to assure the computation has converged.

## METER DIAGRAMA DE BLOCOS

## Limitations of the non-relativistic approach

While Schrodinger's equation may be quite accurate for low energy systems (e.g. Hydrogen), where the speed of the surrounding electrons is not comparable to that of light, such is not the case for more complex and heavier systems.

Speed of 1s electrons in ground state configurations (% of  $c$ )

Hydrogen: 0.516%

Copper: 14.883%

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# Conclusion



Thank you  
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