

04/26/23 - Symbolic Dynamics - Prof. Wolf

A know first ergodic measure: extreme point in the measure space
 On the preimage of a ~~measurable~~ F -invariant set must be 0 or 1.
 e.g. Lebesgue measure not ergodic for a rational rotation of a circle.

Last class: for ergodic measures $h_M(F) = h_{\text{top}}(\mathcal{B}(\mu))$

e.g. let $X \in \Sigma_{\text{inv}}$. $x \in P_{M^d}(f)$

def: $\mu_x = \frac{1}{d} \sum_{k=0}^{d-1} \delta_{f^k(x)}$ period orbit gives you an invariant measure that is ergodic.

compute entropy using: $h_M(x) = \lim_{n \rightarrow \infty} \frac{1}{n} \log \mu(B_n(x, \varepsilon))$

- x is the only point carrying the measure

- so $B_n(x, \varepsilon) = \frac{1}{d}$

- so entropy = 0.

* what is the basin of μ ?
 know that

e.g. Let $X = \sum_d \sum_{p_0, \dots, p_{d-1}} \leq 1 \sum p_i = 1$.

define $\mu(C(x_0 \dots x_k)) = \prod_{i=0}^k p_{x_i}$ def: this is the Bernoulli measure on the cylinders

(\hookrightarrow the Bernoulli measure is shift invariant.)

$\hookrightarrow h_M(F) = \sum p_i \log p_i$

$$p_i = \frac{1}{d}$$

$$\Rightarrow h_M(F) = \log d = h_{\text{top}}(F)$$

($\hookrightarrow \mu$ is a measure of maximal entropy, (mme))
 one can show that this measure is unique

note: an mme always exists for (an SFT?)

e.g. Markov measures.

let $X = X_A \hookrightarrow$ a transitive SFT w/ transition matrix A

finite path between all letters

$B \geq 0$ is a stochastic matrix (rows sum to 1)

def: B is faithfully compatible w/ A if $A_{ij} = 0 \Rightarrow B_{ij} = 0$.

Let $M_{\text{stoch}}(A) = \{B \text{ stochastic, faithfully compatible w/ } A\}$

$P = (p_0, \dots, p_{d-1})$ w/ $p_0 + \dots + p_{d-1} = 1$, $P \in M_{\text{stoch}}(A)$

define $M_{(P,P)}(C(x_0 \dots x_k)) = p_{x_0} \cdot P_{x_0, x_1} \cdot P_{x_1, x_2} \dots P_{x_{k-1}, x_k}$

(\hookrightarrow invariant measure)

little matrix

(\hookrightarrow called a 1-step Markov measure)

$\hookrightarrow h_{M_{(P,P)}}(\sigma) = \sum_{i,j} p_i P_{ij} \log P_{ij}$

e.g. $X \in \Sigma_{\text{inv}}$ $M \in M_X$.

$\Rightarrow h_M(F) = \lim_{n \rightarrow \infty} -\frac{1}{n} \sum_{[\tau] \in \mathcal{L}_n(X)} M([\tau]) \log M([\tau]^{n-1})$ for general invariant measures you need to take the limit (use that the cylinders are generating partitions)

Check: Pesin Dynamical Systems book or Springer lecture notes that has lots of fundamental dynamics.
 (on local entropy by Katok and Brin)

(f is a continuous map on a compact metric space)

philosophically: all a scientist can do is observe measurements. There may be an underlying dynamical systems.

↳ ~~inverse~~ dynamical inverse problem: in a fixed phase space, can you recover the dynamical system from observables?
 → is there an injective map from observables to a dynamical system?
 → is this map constructive

⇒ usually no. Maybe sometimes you can? using topological pressure.

def: Topological pressure $P_{top}: C(X, \mathbb{R}) \xrightarrow{\text{continuous}} \mathbb{R}$

in general $f: X \rightarrow X$ is a continuous dynamical system.
 or work w/ $X \in \Sigma_{inv}$, $\sigma: X \rightarrow X$.

$$Z_n(\phi) = \sum_{T \in \mathcal{Z}_n(X)} \exp\left(\sup_{x \in [T]} S_n \phi(x)\right)$$

this is the n^{th} partition fn

We can show $\log Z_n(\phi) \rightarrow$ subadditive.

⇒ convergent, and the limit is

$$P_{top}(\phi) = \lim_{n \rightarrow \infty} \frac{1}{n} \log Z_n(\phi)$$

↳ we can use this def too on a countable alphabet.

$$= \inf \left\{ \underline{\log} \frac{1}{n} Z_n(\phi) \right\}$$

$$\phi = 0 \text{ gives } Z_n(0) = |\mathcal{X}|^n \text{ so } \Rightarrow P_{top}(0) = h_{top}(f)$$

Note: this is a purely topological construction!

Thm: (Variational principle).

$$P_{top}(\phi) = \sup_{\mu \in M_\phi} \left[h_{\mu}(\phi) + \int \phi d\mu \right]$$

"free energy" $F_U(\phi)$

If M_ϕ satisfies that $P_{top}(\phi) \leq F_{M_\phi}(\phi)$ then we say M_ϕ is an equilibrium state of ϕ .

* Understand Claim: When X is a shift map there is always one equilibrium state.
 how to construct?

$$(M_\phi) \subseteq \mathcal{M}_X \text{ s.t. } h_{\mu_\phi}(\phi) + \int \phi d\mu_\phi \rightarrow P_{top}(\phi)$$

↓
 has a convergent subsequence that must converge to $P_{top}(\phi)$
 upper semi-continuous gives convex + closed

↳ so it has extreme points (the set of equilibrium states)

If ϕ is a transitive SFT, is there a unique equilibrium state?
 ⇒ Yes!

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It's very difficult to study invariant measures but if you have a space of potentials w/ equilibrium states then we can work with observables

We want to talk about coded shifts, and generalize results from SFTs. and so for shifts

Recall $\mathcal{G} = \{\underline{g}_1, \dots\}$ generating words.

$X = X(\mathcal{G}) = \overline{\bigcup_{n \in \mathbb{N}} \text{closure of all concatenations of } \mathcal{G}^n}$

this is not unique for some X , but there's a "minimal" notion and look at \mathcal{G} which uniquely represents X .

If \mathcal{G} not finite the closure includes things not in X_{seq} , which we call $X_{\text{residual}} = X \setminus X_{\text{sequential}}$

Let $X_m = X(\{\underline{g}_1, \dots, \underline{g}_m\})$ w/ finite generators.

↳ this is a sofic shift: produce the graph which is finite, directed (labeled).

We can compute matrices of sofic shifts by finding the adjacency matrix of the graph + computing like an SFT.

One comes is true for pressure. For which coded shifts $X(\mathcal{G})$ and which potentials $\phi: X \rightarrow \mathbb{R}$ do we have

$$P_{\text{top}}(X_m, \phi|_m) \rightarrow P_{\text{top}}(X, \phi)$$

ϕ lives on all of X but we can restrict it to the first m generators

if $\phi \equiv 0$ does $h_{\text{top}}(X_m) \rightarrow h_{\text{top}}(X)$?

↳ this does not hold for all coded shifts

but it does for some!

Recall S-gap shift: $S \subseteq \mathbb{N}_0$, $\mathcal{G} = \{\underline{0^n} \mid n \in S\}$

Recall β -shift: take β shift on $[0, 1]$, do the β expansion of 1, look at all #s that are smaller than the β -expansion of 1.

even w/ full oracle access and full generators you cannot compute the entropy (which is a bit surprising)

indexing (a common technique in dynamical systems research)

let $X = \bigcup_{x \in X} \sigma^{-1}(x)$, μ σ -invariant, fix $E \in \mathcal{B}(X)$ measurable

look at first return map on X :

$$I_E(X) = \min \{n \geq 1 \mid \sigma^n(x) \in E \} \text{ w/ } \min \sigma = +\infty.$$

(Bohr set)
(smallest σ -alg containing upper set)

look at $x \in X$ that return infinitely often:

$$E^\infty = \{x \in E \mid f^n(x) \in E \text{ infinitely many } n \in \mathbb{N}\}.$$

$\mu(E^\infty) = \mu(E)$ by Poincaré recurrence

$$M_E(F) = \frac{M(F \cap E)}{M(E)} \text{ is a prob. measure on } E$$

$\sigma_E : E^\infty \rightarrow E^\infty$ by $\sigma_E(x) = \sigma^{T_E(x)}(x)$
 μ_E is a σ_E invariant measure and if μ is ergodic $\Rightarrow \mu_E$ ergodic.

this is how we push measures nonzero on E onto the first return map.

↪ Sometimes it's easier to study this induced system.

$$\begin{array}{ccc} M_X & \xrightarrow{i} & M_{E^\infty} \\ M & \mapsto & \mu_E \\ & \Downarrow & i(\mu) \\ & & \text{(induced)} \end{array}$$

We have a formula relating entropies for these induced systems.

$$h_{i(\mu)}(\sigma_E) = \frac{h_M(\mu)}{\mu(E)}$$

Now say you have potential ϕ and take the integral?

$$\int_E \phi d\mu \longmapsto \int_E \left(\sum_{k=0}^{T_E(x)-1} \phi(\sigma^k(x)) \right) d\mu$$

↑
point outside
 E get 0 measure this potential is called ϕ_E

so far no μ here
doesn't need to be M_E
(b/c we want to
be able to divide
in the next step)

$$\frac{\int_E \phi d\mu}{\mu(E)} = \int_E \phi_E d\mu_E$$

$$\text{we obtain free energy } \frac{F_{\mu, \phi}(\phi)}{\mu(E)} = F_{\mu_E, \sigma_E}(\phi_E)$$

gives us a relation to study the pressure of the induced system which is often easier

↪ in the induced system the set is often not more compact.

Assume measure $\nu \in M_{\sigma_E}$ w/ integral over first return is finite:

then we can lift the measure: $\int_{E^\infty} \tau_E(x) d\nu(x) < \infty$. (i.e. T_E is L1)

$$\tilde{\nu} \in M_E \text{ by } \tilde{\nu}(B) = \sum_{k=0}^{\infty} \nu(\{x : T_E(x) \geq k\} \cap f^{-k}(B))$$

(↪ finite b/c integral finite)

↪ σ -invariant

↪ not nec. a prob. measure

normalize this measure:

$$l(\nu) = \frac{\tilde{\nu}}{\tilde{\nu}(x)} \text{ is the lift of } \nu, \text{ and it is a } \sigma\text{-inv. prob. measure.}$$

prop: $i(l(\nu)) = \nu$ (lifting and inducing are inverses)

$$M_{\sigma_E} \xrightarrow[i]{e} M_{\sigma_E, E^\infty} \cap \left\{ \int_E \tau_E(x) d\nu(x) < \infty \right\}$$

is a natural bijection.

Why are we doing all of this?

$$X_{\text{seq}} = \{x = g_1 g_2 g_3 \dots g_m g_{m+1} \dots\}$$

induced set where origin lands between generators

for other systems, a different inducing scheme might be better (choose a different return map).

- gives you a full shift on generators (countably infinite)

- then it's good enough to estimate on the first m generators.