

## Example 3.2

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Gabe Morris

```
[1]: # npb
import sympy as sp
import numpy as np
import matplotlib.pyplot as plt

plt.style.use('../maroon_ipynb.mplstyle')
```

### Example 3.2

Solve the following ODE with  $f(t) = 200u(t)$  and  $\dot{x}(0) = v_0$ . All other initial conditions are zero.

$$2\ddot{x} + 14\dot{x} + 46.5x + 127\dot{x} + 148x = f(t)$$

### Important Note with Sympy

One of the major caveats of `sympy` is that you should never include a decimal value in a symbolic expression. This may lead to errors depending on what you are trying to accomplish. You can work around this by representing the decimal value as a fraction or using `Rational` from `sympy`.

For example, let's say you want to evaluate  $\frac{0.1}{0.3}x$ .

```
[2]: # do not do this
x = sp.Symbol('x')
expr = 0.1/0.3*x
expr
```

```
[2]: 0.3333333333333333x
```

This might seem ok, but if you expand that floating point value, you will see floating point error. You can adjust the amount of decimals that `sympy` will output with the `n()` method.

```
[3]: expr.n(20)
```

```
[3]: 0.33333333333333337034x
```

Here is the fix using `Rational`.

```
[4]: numerator = sp.Rational(1, 10)
numerator # this will correctly represent 0.1 as a fraction
```

```
[5]: denominator = sp.Rational(3, 10)
      expr_correct = numerator/denominator*x
      expr_correct
```

$$[5] : \frac{x}{3}$$

```
[6]: expr_correct.n(20)
```

[6]:  $0.33333333333333333333x$

## Solving the Easy Way

The best way to solve this with `sympy` would be to use the `dsolve` function.

```
[7]: t, v0 = sp.symbols('t v0', real=True) # sometimes it's necessary to specify
      ↪ that the symbol is real
      x = sp.Function('x')(t)
      eq = sp.Eq(2*x.diff(t, 4) + 14*x.diff(t, 3) + sp.Rational(93, 2)*x.diff(t, 2) +
      ↪ 127*x.diff(t) + 148*x, 200*sp.Heaviside(t))
      eq
```

$$[7]: \quad 148x(t) + 127\frac{d}{dt}x(t) + \frac{93\frac{d^2}{dt^2}x(t)}{2} + 14\frac{d^3}{dt^3}x(t) + 2\frac{d^4}{dt^4}x(t) = 200\theta(t)$$

```
[8]: sol = sp.dsolve(eq, ics={
    x.subs(t, 0): 0,
    x.diff(t, 1).subs(t, 0): v0,
    x.diff(t, 2).subs(t, 0): 0,
    x.diff(t, 3).subs(t, 0): 0
})
sol
```

[8]: 
$$x(t) = \left(-\frac{9v_0}{34} + \frac{10\theta(t)}{17}\right)e^{-4t} + \left(\frac{53v_0}{90} - \frac{20\theta(t)}{9}\right)e^{-2t} + \left(\left(-\frac{248v_0}{765} + \frac{1600\theta(t)}{5661}\right)\cos(3t) + \left(\frac{244v_0}{765} - \frac{3680\theta(t)}{5661}\right)\sin(3t)\right)e^{-t}$$

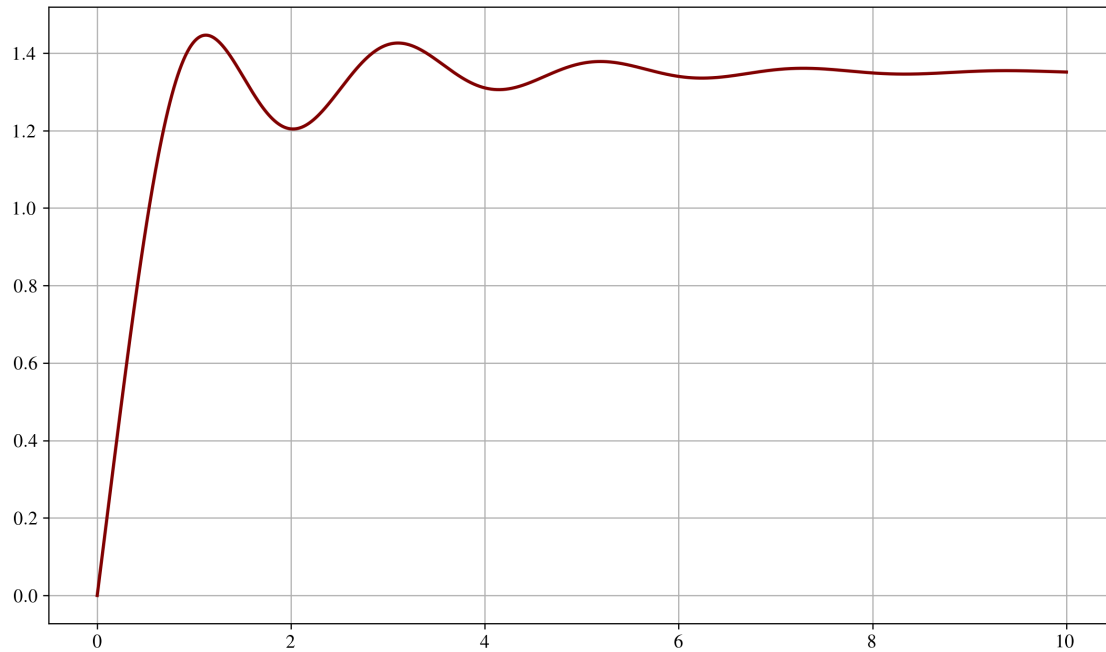
It is very nice that it was able to solve this and retain the  $v_0$  symbol in the solution. We can bring this to the numerical world with `lambdify`.

```
[9]: # Substitute v0 with 2
sol_sub = sol.rhs.subs(v0, 2)

t_array = np.linspace(0, 10, 1000)
x_lamb = sp.lambdify(t, sol_sub, modules='numpy')

fig, ax = plt.subplots()
ax.plot(t_array, x_lamb(t_array))
```

```
plt.show()
```



## Solving with Laplace Transforms

```
[10]: # Putting the original equation into the Laplace domain
s = sp.Symbol('s', real=True)
eq_s = sp.Eq(sp.laplace_transform(eq.lhs, t, s)[0], sp.laplace_transform(eq.
    ↪ rhs, t, s)[0])
eq_s
```

$$\begin{aligned}
 & 2s^4 \mathcal{L}_t[x(t)](s) + 14s^3 \mathcal{L}_t[x(t)](s) - 2s^3 x(0) + \frac{93s^2 \mathcal{L}_t[x(t)](s)}{2} - 14s^2 x(0) - 2s^2 \left. \frac{d}{dt} x(t) \right|_{t=0} + \\
 & 127s \mathcal{L}_t[x(t)](s) - \frac{93sx(0)}{2} - 14s \left. \frac{d}{dt} x(t) \right|_{t=0} - 2s \left. \frac{d^2}{dt^2} x(t) \right|_{t=0} + 148 \mathcal{L}_t[x(t)](s) - 127x(0) - \\
 & \frac{93 \left. \frac{d}{dt} x(t) \right|_{t=0}}{2} - 14 \left. \frac{d^2}{dt^2} x(t) \right|_{t=0} - 2 \left. \frac{d^3}{dt^3} x(t) \right|_{t=0} = \frac{200}{s}
 \end{aligned}$$

In the past, sympy's laplace transforms would not apply the rule with the derivatives, but this changed in 2023 as you can clearly see the  $s$  terms.

```
[11]: # Applying the initial conditions
eq_s = eq_s.subs([
    (x.subs(t, 0), 0),
    (x.diff(t, 1).subs(t, 0), v0),
    (x.diff(t, 2).subs(t, 0), 0),

```

```
(x.diff(t, 3).subs(t, 0), 0)
])
eq_s
```

[11]: 
$$2s^4 \mathcal{L}_t[x(t)](s) + 14s^3 \mathcal{L}_t[x(t)](s) - 2s^2 v_0 + \frac{93s^2 \mathcal{L}_t[x(t)](s)}{2} - 14s v_0 + 127s \mathcal{L}_t[x(t)](s) - \frac{93v_0}{2} + 148 \mathcal{L}_t[x(t)](s) = \frac{200}{s}$$

[12]: 

```
# Solving for X(s)
X_s = sp.solve(eq_s, sp.laplace_transform(x, t, s)[0])[0]
X_s
```

[12]: 
$$\frac{4s^3 v_0 + 28s^2 v_0 + 93s v_0 + 400}{s(4s^4 + 28s^3 + 93s^2 + 254s + 296)}$$

Though I don't think there is a need, you can still use the partial fraction form if you wanted to do this by hand. Sometimes, `sympy` will solve the inverse laplace transform better in this form, but I haven't had any issues with inverse laplace transforms for a while.

[13]: 

```
X_s_partial = sp.apart(X_s, s) # must include the 's' because we have v0 in
↳ the expression
X_s_partial
```

[13]: 
$$-\frac{32 \cdot (1147s v_0 - 1000s - 2812v_0 + 6400)}{28305 \cdot (4s^2 + 4s + 37)} - \frac{9v_0 - 20}{34(s + 4)} + \frac{53v_0 - 200}{90(s + 2)} + \frac{50}{37s}$$

### Pro tip

Let's say we wanted to extract the characteristic polynomial from the above  $X(s)$  expression and find its roots. First off, it would have been better to leave  $f(t)$  in its symbolic form and find the transfer function, but we can get around this with what we have already done by dividing out  $f(t)$ .

[14]: 

```
# Get transfer function
T_s = X_s/(200/s)
T_s
```

[14]: 
$$\frac{4s^3 v_0 + 28s^2 v_0 + 93s v_0 + 400}{200 \cdot (4s^4 + 28s^3 + 93s^2 + 254s + 296)}$$

It may not be in the form that we would perceive as the characteristic polynomial, but its roots will still be the same.

[15]: 

```
# Use sp.fraction to get the numerator and denominator
num, den = sp.fraction(T_s)
den
```

[15]: 
$$800s^4 + 5600s^3 + 18600s^2 + 50800s + 59200$$

[16]: 

```
cp = den/400 # returns original characteristic polynomial
cp
```

[16]: 
$$2s^4 + 14s^3 + \frac{93s^2}{2} + 127s + 148$$

```
[17]: # you can use sp.solve to get the roots
# or you can create a polynomial object and use the nroots method
cp_poly = sp.Poly(cp, s)
roots = cp_poly.nroots()
display(*roots)
```

−4.0

−2.0

−0.5 − 3.0i

−0.5 + 3.0i

## Inverse Laplace Transform

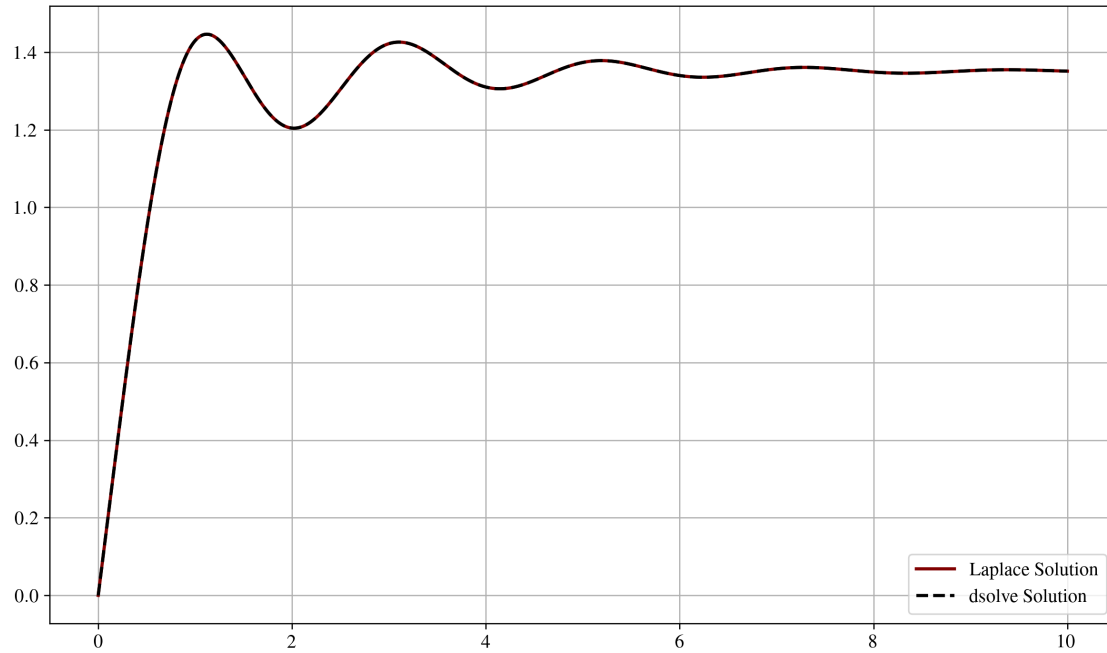
Going back to solving, all we have to do now is take the inverse laplace transform.

```
[18]: sol = sp.inverse_laplace_transform(X_s, s, t)
sol
```

[18]: 
$$\left(\frac{10}{17} - \frac{9v_0}{34}\right) e^{-4t} \theta(t) + \left(\frac{53v_0}{90} - \frac{20}{9}\right) e^{-2t} \theta(t) + \left(\left(\frac{1600}{5661} - \frac{248v_0}{765}\right) e^{-\frac{t}{2}} \cos(3t) + \left(\frac{244v_0}{765} - \frac{3680}{5661}\right) e^{-\frac{t}{2}} \sin(3t)\right) \theta(t) + \frac{50\theta(t)}{37}$$

```
[19]: # Substitute v0 with 2
# Let's see if it's the same as the dsolve solution
x_t = sol.subs(v0, 2)
x_t_lamb = sp.lambdify(t, x_t, modules='numpy')

fig, ax = plt.subplots()
ax.plot(t_array, x_t_lamb(t_array), label='Laplace Solution')
ax.plot(t_array, x_lamb(t_array), ls='--', label='dsolve Solution')
ax.legend()
plt.show()
```



A couple of things to note about this result:

- The dominant root is the root whose real part has the smallest magnitude. This results in the largest time constant. The time constant then for this ODE is  $-\frac{1}{-0.5} = 2$ . Notice that the solution arrives close to steady state at  $4 \cdot 2 = 8$  seconds.
- Since our characteristic polynomial has imaginary roots, we see oscillation. We can recognize that the oscillations occur at  $3 \text{ rad/s}$ . The period in seconds is  $2\pi/3 \approx 2.1 \text{ s}$ , and we can see that the peak to peak distance in the output matches this value.