

Exploring Quantum Process Calculi

Gabriele Tedeschi

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Qubits

The simplest quantum system is a **qubit**. Like a bit, a qubit has two separate states, $|0\rangle$ and $|1\rangle$.

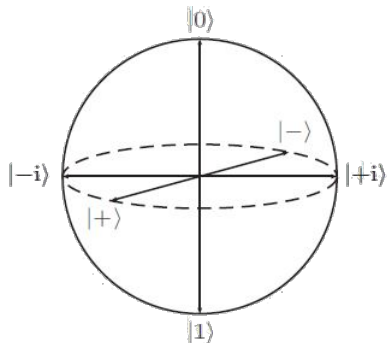
Qubits

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A qubit can also be in a linear combination of states, known as a **superposition**.

$$|+\rangle = \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle$$

$$|-\rangle = \frac{1}{\sqrt{2}} |0\rangle - \frac{1}{\sqrt{2}} |1\rangle$$

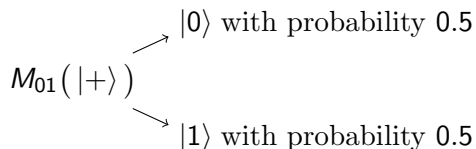


Measurements

A qubit in superposition cannot be directly observed, because when it get **measured**, it **decays** in one of its basis states.

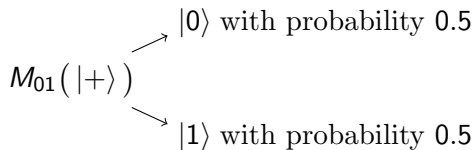
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Interestingly, the same happens when we measure $|0\rangle$ in the $+-$ basis, we get either the outcome $|+\rangle$ either the outcome $|-\rangle$ with the same probability 0.5.

No Cloning theorem

No-Cloning

Quantum information cannot be **duplicated**. That is, given a qubit q_1 in state

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

it is impossible to prepare a qubit q_2 in the same state, without destroying the information in q_1 .

This means that, contrary to classical bits, qubits cannot be freely copied, stored or broadcasted to multiple receivers.

Linear qCCS: Type System

lqCCS

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The type system is needed to prevent **duplication** of quantum information: after receiving a qubit, a process must send it **exactly once**.

$$c?q.H(q).c!q$$

$$a?q.(b!q \parallel c!q)$$

$$c?q.H(q).0$$

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In previous calculi, there were **ambiguous** processes like

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In lqCCS, the programmer must **explicitly** describe what happens to each qubit

$$\begin{aligned} P' &= c?q.X(q).c!q & Q' &= c?q.H(q).c!q \\ P'' &= c?q.X(q).disc(q) & Q'' &= c?q.H(q).disc(q) \end{aligned}$$

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$$\begin{array}{c}
 \langle |0\rangle, H(q).c!q \parallel c?x.M_{01}[x \triangleright y].disc(x) \rangle \\
 \downarrow \\
 \langle |+\rangle, c!q \parallel c?x.M_{01}[x \triangleright y].disc(x) \rangle \\
 \downarrow \\
 \langle |+\rangle, M_{01}[q \triangleright y].disc(q) \rangle \\
 \downarrow \\
 \langle |0\rangle, disc(q) \rangle \frac{1}{2} \oplus \langle |1\rangle, disc(q) \rangle
 \end{array}$$

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- Context: a "program" with a hole, like $[-] \parallel R$. We compare P and Q "inside" this context, i.e. we study $P \parallel R$ and $Q \parallel R$.

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The two processes seen before

$$P' = c?q.X(q).c!q \quad Q' = c?q.H(q).c!q$$

are not bisimilar, because there is a context R which tells them apart

$$R = c?q.M_{01}[q \triangleright x].(\text{disc}(q) \parallel \text{if } x = 0 \text{ then } a!0 \text{ else } b!0)$$

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$$\mathcal{C} = \langle |+\rangle, M_{01}[q \triangleright x].c!q \rangle \rightarrow \langle |0\rangle, c!q \rangle \frac{1}{2} \oplus \langle |1\rangle, c!q \rangle$$

$$\mathcal{C}' = \langle |0\rangle, M_{\pm}[q \triangleright x].c!q \rangle \rightarrow \langle |+\rangle, c!q \rangle \frac{1}{2} \oplus \langle |-\rangle, c!q \rangle$$

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According to the usual notion of **probabilistic bisimilarity**, two distributions are bisimilar if they assign the same probability to bisimilar processes. In our example, the two configurations are **not** bisimilar, because they evolve in two non-bisimilar distributions.

Indistinguishable distributions

According to quantum mechanics, it's impossible to distinguish \mathcal{C} and \mathcal{C}' , if we treat them as two sources of qubits:

\mathcal{C} emits a qubit $|0\rangle$ or $|1\rangle$ with the same probability 0.5

\mathcal{C}' emits a qubit $|+\rangle$ or $|-\rangle$ with the same probability 0.5

Suppose we receive a qubit from either \mathcal{C} or \mathcal{C}' , and measure it in the 01 basis. If we measure a qubit from \mathcal{C} , it would result in either $|0\rangle$ or $|1\rangle$. If we measure a qubit from \mathcal{C}' , a $|+\rangle$ qubit would decay in either $|0\rangle$ or $|1\rangle$, and a $|-\rangle$ qubit would decay in either $|0\rangle$ or $|1\rangle$ as well.

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Inadequacy of Probabilistic Bisimilarity

The usual notion of probabilistic bisimilarity is too fine, when comparing distributions of quantum configurations.

Solution I: Quantum Bisimilarity

Quantum Bisimilarity

We introduce an **equivalence** relation

$$\equiv \subseteq \mathcal{D}(\mathit{Conf}) \times \mathcal{D}(\mathit{Conf})$$

and a new notion of **quantum bisimilarity**. Two processes are quantum bisimilar if they evolve in distributions that are bisimilar **up to equivalence**.

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Solution II: mQPA

We introduce a new calculus, equipped with a minimal set of features: communication, non-determinism and quantum measurement.

In mQPA, the transitions are of the form

$$\rightarrow \subseteq S \times \mathcal{D}(S)^{\mathcal{H}}$$

where $\mathcal{D}(S)^{\mathcal{H}}$ is the set of **quantum distributions**. The probabilistic observable behaviour is **parametric** with respect to an input quantum state.

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In mQPA, the previous example can be rewritten as

$$(S_{|0\rangle\langle 0|} \boxplus S')(|+\rangle) = S_{\frac{1}{2}} \oplus S' \quad (S_{|+\rangle\langle +|} \boxplus S')(|0\rangle) = S_{\frac{1}{2}} \oplus S'$$

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- Even though quantum systems exhibit a probabilistic behaviour, **probabilistic bisimilarity** is not really well suited for the quantum setting.
- **Quantum bisimilarity** relaxes the conditions of probabilistic bisimilarity, better representing quantum systems.
- **mQPA**, a minimal calculus, pursuing a foundational, where the probabilistic classical behaviour is parametric with respect to quantum values.

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- From weak transitions, it is possible to define reachability, temporal logics and **model checking**.
- Saturated bisimilarity can be cumbersome to prove, it will be interesting to explore how the existing **proof techniques** adapt to the quantum setting.
- We will investigate the relation between **mQPA** semantics and the usual, configuration-based semantics.

Thank you for your attention!