

Comparing Quantum Protocols via PRISM

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Motivations

Quantum cryptographic protocols have been proved to be **unconditionally secure**, meaning that every possible attack will fail with probability arbitrarily close to one.

But one would also like to:

- Compare different protocols
- Verify other properties of the protocol
- Compare the behaviour with different parameters

Thats when **model checking** comes in handy!

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Bracket Notation

In linear algebra, we write:

- Vector \mathbf{v}
- Conjugate transpose
 $\mathbf{v}^* = \mathbf{v}^H = \overline{\mathbf{v}^T}$
- Dot product between \mathbf{v} and \mathbf{u}
 $\mathbf{v}^* \cdot \mathbf{u}$
- Norm of \mathbf{v} : $\|\mathbf{v}\| = \sqrt{\mathbf{v}^* \cdot \mathbf{v}}$

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In quantum mechanics, we write:

- Vector $|\psi\rangle$
- conjugate transpose
 $\langle\psi| = |\psi\rangle^H = \overline{|\psi\rangle}^T$
- dot product between $|\psi\rangle$ and $|\phi\rangle$
 $\langle\psi|\phi\rangle$
- Norm of $|\psi\rangle$: $\langle\psi|\psi\rangle$

First Postulate: Quantum State

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$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \begin{pmatrix} i \\ i \end{pmatrix}, \quad \begin{pmatrix} 5 + 5i \\ \frac{1}{3} + \frac{1}{5}i \end{pmatrix}, \quad \dots$$

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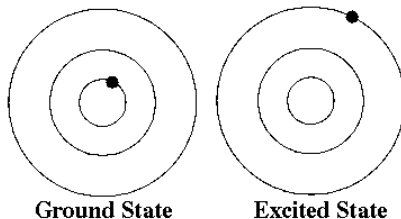
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- **Unit vector:** a vector $|\psi\rangle$ such that $\langle\psi|\psi\rangle = 1$, like:

$$\begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad \begin{pmatrix} 0 \\ i \end{pmatrix}, \quad \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}}i \end{pmatrix}, \quad \dots$$

Qubits

Suppose a physical system with only two states: an particle with spin up or spin down, for example, or an electron that can be in an excited state or in the ground state.



Qubits

A system with only two possible state is described by a **qubit**, a unitary vector in the 2-dimensional Hilbert space $\mathcal{H}_2 \equiv \mathbb{C}^2$.

For example, if an electron is in the ground state, we say that is in state

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

while if it is in the excited state, it's in state

$$|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Computational Basis

$|0\rangle$ and $|1\rangle$ form the so called **computational base** of the (bidimensional) Hilbert space. This means that each vector $|\psi\rangle$ can be expressed as

$$|\psi\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \alpha |0\rangle + \beta |1\rangle, \text{ with } \alpha, \beta \in \mathbb{C}, \langle\psi|\psi\rangle = 1$$

Quantum Superposition

Since $|0\rangle$ and $|1\rangle$ are vectors in \mathcal{H}_2 , so it is any **linear combination** of the two, provided it has unitary length.

$$\begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix}, \quad \begin{pmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{pmatrix}, \quad \begin{pmatrix} \sqrt{0.9} \\ \sqrt{0.1} \end{pmatrix}$$

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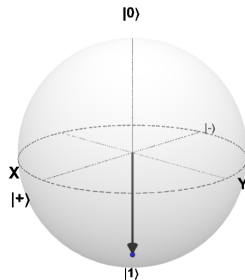
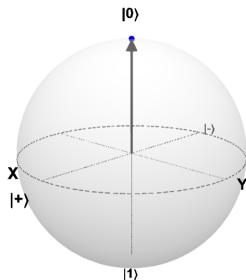
Hadamard basis

The vector $\begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle$ is called the $|+\rangle$ state

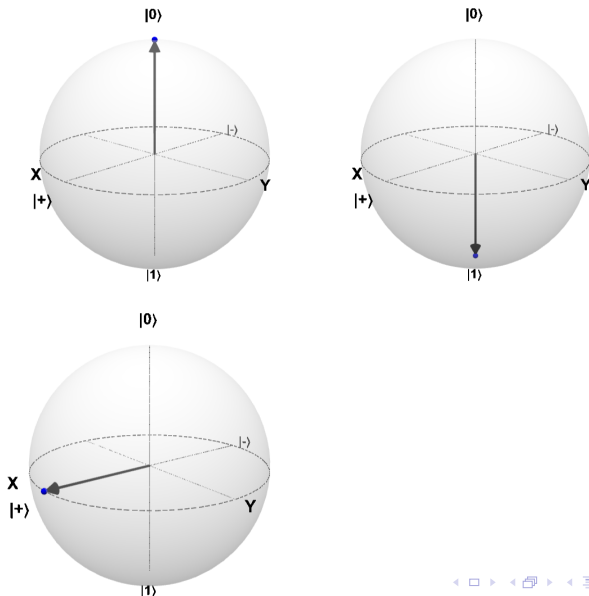
The vector $\begin{pmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \frac{1}{\sqrt{2}} |0\rangle - \frac{1}{\sqrt{2}} |1\rangle$ is called the $|-\rangle$ state

They form the Hadamard Basis of \mathcal{H}_2

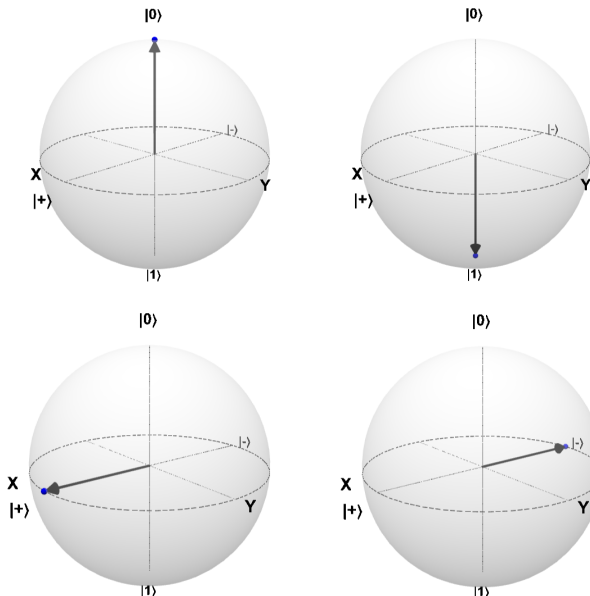
Bloch sphere



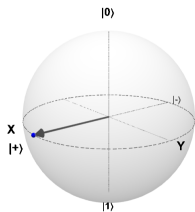
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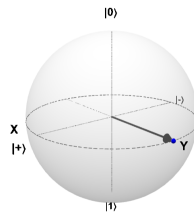
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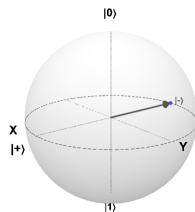
Phase



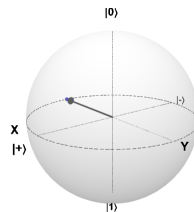
$$|+\rangle = |0\rangle + e^{i0} |1\rangle$$



$$|0\rangle + e^{i\frac{\pi}{2}} |1\rangle$$



$$|-\rangle = |0\rangle + e^{i\pi} |1\rangle$$



$$|0\rangle + e^{i\frac{3\pi}{2}} |1\rangle$$

Compound Systems

How do we describe a system composed of two different qubits? As the **Tensor product** of Hilbert spaces.

$$\text{If } |\psi\rangle, |\phi\rangle \in \mathcal{H} \quad \text{then} \quad |\psi\phi\rangle \in \mathcal{H} \otimes \mathcal{H}$$

$\mathcal{H}_2 \otimes \mathcal{H}_2$ is a **four-dimensional Hilbert space**.

Dot product in a compound space

Given $|\psi\psi'\rangle$ and $|\phi\phi'\rangle$, the dot product is defined as

$$\langle\psi\psi'|\phi\phi'\rangle = \langle\psi|\phi\rangle \langle\psi'|\phi'\rangle$$

Second Postulate: Unitary Transformations

II Postulate

The evolution of a **closed** quantum system is described by a **unitary transformation**. That is, the state $|\psi\rangle$ of the system at time t_1 is related to the state $|\psi'\rangle$ of the system at time t_2 by a unitary operator U which depends only on t_1 and t_2 :

$$|\psi'\rangle = U|\psi\rangle$$

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- **Closed**: without interaction with the environment, i.e. without exchanging with the environment any energy, information, etc..
- **Unitary transformation** on a Hilbert space \mathcal{H}_n : A $n \times n$ matrix such that its *adjoint* it's also its *inverse*:

$$UU^H = U^H U = UU^{-1} = I$$

A useful property

Unitary Matrices preserve dot product

Given $|\psi\rangle, |\phi\rangle$ with dot product x , for a unitary matrix U holds:

$$\langle\psi|\phi\rangle = x$$

$$\langle U\psi|U\phi\rangle = \langle\psi|U^H U|\phi\rangle = \langle\psi|\phi\rangle = x$$

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$$\begin{aligned}\langle\psi|\phi\rangle &= x \\ \langle U\psi|U\phi\rangle &= \langle\psi|U^H U|\phi\rangle = \langle\psi|\phi\rangle = x\end{aligned}$$

Corollary

If $|\psi\rangle$ is a unit vector, $U|\psi\rangle$ will also be a unit vector.

$$\langle U\psi|U\psi\rangle = \langle\psi|\psi\rangle = 1$$

Third Postulate: Measurements

III Postulate

If a state $|\psi\rangle = \alpha|b_0\rangle + \beta|b_1\rangle$ gets **measured** in the base $\{b_0, b_1\}$, it will **collapse**

- In state $|b_0\rangle$ with probability $\|\alpha\|^2$
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If $|+\rangle = \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle$ gets measured in the computational basis $\{|0\rangle, |1\rangle\}$, it will result in $|0\rangle$ or $|1\rangle$ with probability 0.5 each.

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Since $|0\rangle = \frac{1}{\sqrt{2}} |+\rangle + \frac{1}{\sqrt{2}} |-\rangle$, measuring $|0\rangle$ in the hadamard basis $\{|+\rangle, |-\rangle\}$ will result in $|+\rangle$ or $|-\rangle$ with probability 0.5 each.

Takeaway

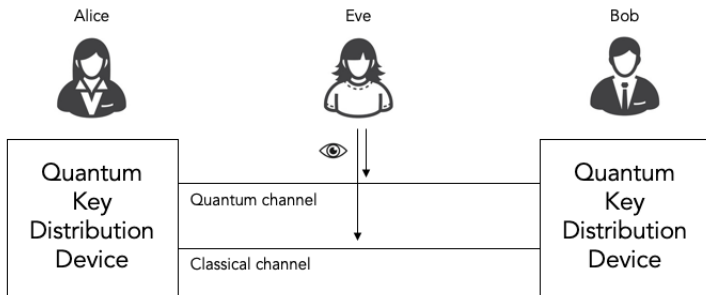
To understand the quantum cryptographic protocols, two properties will be useful:

- The evolution of a quantum system is described by a **unitary matrix**, and unitary matrices preserve the inner product.
- When the $|0\rangle$ and $|1\rangle$ states are **measured** in the computational basis, will give with probability 1 the outcome $|0\rangle$ and $|1\rangle$, respectively. When measured in the Hadamard basis, will give outcome $|+\rangle$ or $|-\rangle$ with probability $\frac{1}{2}$ each. The inverse is true for the $|+\rangle$ and $|-\rangle$ states.

Quantum Key Distribution Protocols

Two communicating parties, Alice and Bob, want to establish a shared **secret** $\mathbf{k} \in \{0, 1\}^N$, $N \geq 0$ for secure communication. They do not share any prior information, but can make use of:

- A **classical**, (public) channel, that can be **passively monitored** but **not tampered** with by an attacker.
- A **quantum** channel, that can be **tampered** with by an attacker, but by its nature can not be **passively monitored**.



Information gains implies disturbance

Theorem

In any attempt to distinguish between two **non-orthogonal** quantum states, information gain is only possible at the expense of **introducing disturbance** to the signal.

Eve, having intercepted a qubit $q \in \{|\psi\rangle, |\phi\rangle\}$ would like to apply a unitary transformation U such that:

$$U(|\psi a\rangle) = |\psi b\rangle$$

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$|b\rangle$ and $|b'\rangle$ must be equal!

The BB84 Protocol

First Phase: Quantum transmission

- 1 Alice creates a random string of bits $\mathbf{d} \in \{0, 1\}^n$, and a random string of bases $\mathbf{b} \in \{\boxplus, \boxtimes\}^n$, where $n > N$.

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$$\begin{aligned} |\phi_{\boxplus 0}\rangle &= |0\rangle & |\phi_{\boxtimes 0}\rangle &= |+\rangle \\ |\phi_{\boxplus 1}\rangle &= |1\rangle & |\phi_{\boxtimes 1}\rangle &= |-\rangle \end{aligned}$$

- 3 Bob creates a random string of bases $\mathbf{b}' \in \{\boxplus, \boxtimes\}^n$, and measures each $|\phi_i\rangle$ in the bases b'_i . When he measures $|0\rangle$ or $|+\rangle$, he stores 0, when he measures $|1\rangle$ or $|-\rangle$, he stores 1. Doing so he obtains a string of bits $\mathbf{d}' \in \{0, 1\}^n$, that will contain some errors with respect to the original \mathbf{d} .

Second Phase: Public discussion

Alice and Bob will keep only the correct bits, $\mathbf{k} \subseteq \mathbf{b}$

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- 1 Alice sends \mathbf{b} over the public channel, Bob confronts it with \mathbf{b}' , and answers to Alice the set $S = \{i \mid b_i = b'_i\}$. Only the corresponding qubits have been correctly measured, and the others are discarded.

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- 2 Alice chooses a subset of the remaining bits in \mathbf{d} and discloses their values to Bob. For each d_i in this subset, since it holds $b_i = b'_i$, it should be $d_i = d'_i$. If Bob notices discrepancies $d_i \neq d'_i$, eavesdropping is detected.

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- ③ The shared secret $\mathbf{k} \in \{0, 1\}^N$ is the string of bits in $\mathbf{d} = \mathbf{b}'$ that have not been disclosed at step 2.

BB84 Observations

- Alice secret, \mathbf{d} , is never completely disclosed. Each bit is sent "encrypted" as $|\phi_{b_id_i}\rangle$ over the quantum channel. Later, the "key" b_i is revealed, but only **after that the quantum information $|\phi_{b_id_i}\rangle$ is destroyed**.
- If Eve guesses the right base each time, her presence is undetectable, but the probability of this event decreases exponentially with n .

Eavesdropper's attacks

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- **Random Substitute attack:** Alice sends qubit $|\psi_{bd}\rangle$, Eve chooses a basis \hat{b} and a bit \hat{d} , and measures Alice's qubit. If $\hat{b} = b$, then Eve has discovered d . Independently from the measurement, Eve will send to Bob $|\psi_{\hat{b}\hat{d}}\rangle$.

PRISM Formalization: Channel

```

module ChannelResend
  ch_state : [0..4];
  ch_bas : [0..1];
  ch_bit : [0..1];
  [aliceput] (ch_state=0) ->
    (ch_state'=1) & (ch_bas'=al_bas) & (ch_bit'=al_bit);
  [evemeasure] (ch_state=1) & (ch_bas=eve_bas) -> (ch_state'=2);
  [evemeasure] (ch_state=1) & (ch_bas!=eve_bas) ->
    0.5 : (ch_state'=2) & (ch_bit'=0) +
    0.5 : (ch_state'=2) & (ch_bit'=1);
  [eveget] (ch_state=2) -> (ch_state'=3);
  [eveput] (ch_state=3) ->
    (ch_state'=4) & (ch_bas'=eve_bas) & (ch_bit'=eve_bit);
  [bobget] (ch_state=4) -> (ch_state'=0);
endmodule

```

PRISM Formalization: Alice

```

const int N;
module Alice
  al_state : [0..5];
  al_index : [1..N];
  al_bas : [0..1];
  al_bit : [0..1];
  [] (al_state=0) & (eve_state>0) ->
    0.5 : (al_state'=1) & (al_bas'=0) +
    0.5 : (al_state'=1) & (al_bas'=1);
  [] (al_state=1) ->
    0.5 : (al_state'=2) & (al_bit'=0) +
    0.5 : (al_state'=2) & (al_bit'=1);
  [aliceput] (al_state=2) -> (al_state'=3);
  [reveal] (al_state=3) -> (al_state'=4);
  [loop] (al_state=4) & (al_index<N) -> (al_state'=0) & (al_index'=al_index+1);
  [loop] (al_state=4) & (al_index=N) -> (al_state'=5);
  [stop] (al_state=4) -> (al_state'=4);
  [stop] (al_state=5) -> (al_state'=5);
endmodule

```

PRISM Formalization: Bob

```

const int CHECKBIT;
module Bob
  bob_state : [0..7];
  bob_bas : [0..1];
  bob_bit : [0..1];
  [] (bob_state=0) & (eve_state>0) ->
    0.5 : (bob_state'=2) & (bob_bas'=0) +
    0.5 : (bob_state'=2) & (bob_bas'=1);
  [bobget] (bob_state=2) & (bob_bas=ch_bas) -> (bob_state'=3) & (bob_bit'=ch_bit);
  [bobget] (bob_state=2) & (bob_bas!=ch_bas) ->
    0.5 : (bob_bit'=ch_bit) & (bob_state'=3) +
    0.5 : (bob_bit'=1-ch_bit) & (bob_state'=3);
  [reveal] (bob_state=3) & (bob_bas!=al_bas) -> (bob_state'=4);
  [reveal] (bob_state=3) & (bob_bas=al_bas) -> (bob_state'=5);
  [] (bob_state=5) & (bob_bit!=al_bit) -> (bob_state'=7);
  [] (bob_state=5) & (bob_bit=al_bit) -> (bob_state'=4);
  [loop] (bob_state=4) & (al_index<N) -> (bob_state'=0);
  [loop] (bob_state=4) & (al_index=N) -> (bob_state'=6);
  [stop] (bob_state=6) -> (bob_state'=6); // no eavesdropping detected
  [stop] (bob_state=7) -> (bob_state'=7); // eavesdropping detected
endmodule

```


PRISM Formalization: Intercept/Resend

```
module EveResend
eve_state : [0..4];
eve_bas : [0..1];
eve_bit : [0..1];
[] (eve_state=0) ->
    0.5 : (eve_state'=1) & (eve_bas'=0) +
    0.5 : (eve_state'=1) & (eve_bas'=1);
[evemeasure] (eve_state=1) -> (eve_state'=2);
[eveget] (eve_state=2) -> (eve_state'=3) & (eve_bit'=ch_bit) ;
[eveput] (eve_state=3) -> (eve_state'=0);
endmodule
```

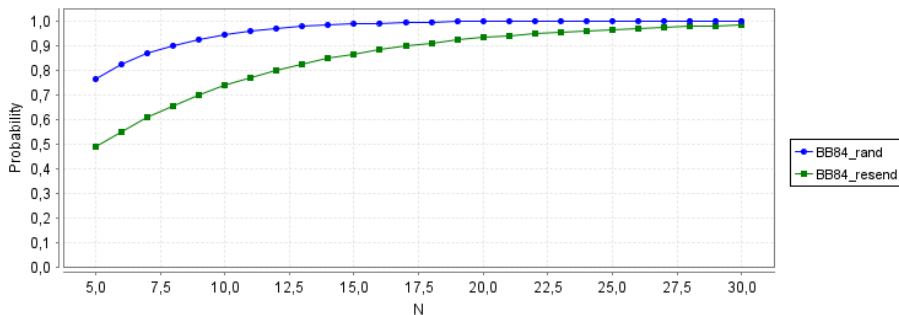
PRISM Formalization: Random Substitution

```

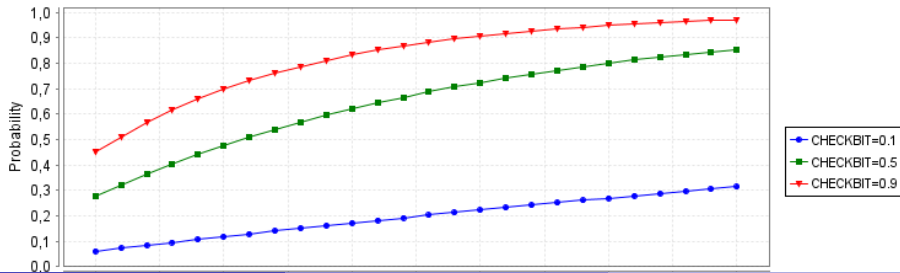
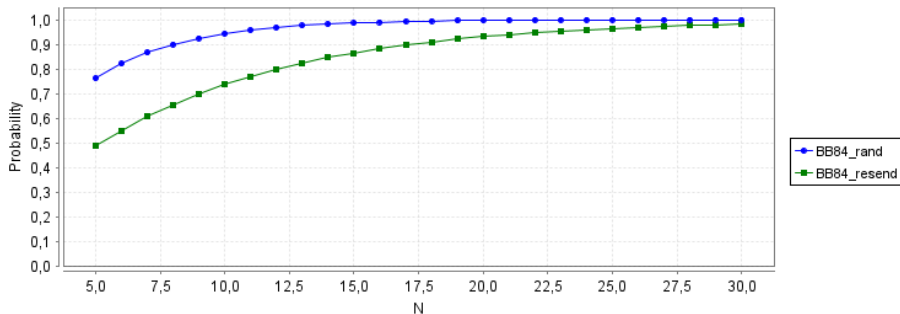
module EveRandom
eve_state : [0..4];
eve_bas : [0..1];
eve_bit : [0..1];
[] (eve_state=0) ->
    0.5 : (eve_state'=1) & (eve_bas'=0) +
    0.5 : (eve_state'=1) & (eve_bas'=1);
[evemeasure] (eve_state=1) -> (eve_state'=2);
[eveget] (eve_state=2) -> (eve_state'=3) & (eve_bit'=ch_bit) ;
[] (eve_state=3) ->
    0.50: (eve_state'=4) & (eve_bit'=0) +
    0.50: (eve_state'=4) & (eve_bit'=1) ;
[eveput] (eve_state=4) -> (eve_state'=0);
endmodule

```

Eve will be detected



Eve will be detected



The B92 Protocol

First Phase: Quantum transmission

- 1 Alice creates a random string of bits $\mathbf{b} \in \{\boxplus, \boxtimes\}^n$, where $n > N$.
- 2 Alice sends to Bob the sequence of qubits $|\phi_i\rangle = |\phi_{b_i}\rangle$, where

$$|\phi_{\boxplus}\rangle = |0\rangle$$

$$|\phi_{\boxtimes}\rangle = |+\rangle$$

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- 3 Bob creates a random string of bases $\mathbf{b}' \in \{\boxplus, \boxtimes\}^n$, and measures each $|\phi_i\rangle$ in the bases b'_i . Doing so he constructs a sequence of bits $\mathbf{t} \in \{0, 1\}^N$, such that

$$t_i = \begin{cases} 0 & \text{if } |\phi_i\rangle, \text{ measured with } b'_i, \text{ yields } |0\rangle \text{ or } |+\rangle \\ 1 & \text{if } |\phi_i\rangle, \text{ measured with } b'_i, \text{ yields } |1\rangle \text{ or } |-\rangle \end{cases}$$

Second Phase: Public discussion

Alice and Bob will keep only the correct bits, $\mathbf{k} \subseteq \mathbf{b}$

The B92 Protocol

First Phase: Quantum transmission

Alice has \mathbf{b} , Bob has \mathbf{b}' and \mathbf{t}

Second Phase: Public discussion

- 1 When $t_i = 1$, it means that Bob has used the wrong basis, and so he knows the original bit b_i , that is $1 - b'_i$. Bob sends \mathbf{t} over the public channel, and both parties keep only the bits for which $t_i = 1$.

The B92 Protocol

First Phase: Quantum transmission

Alice has \mathbf{b} , Bob has \mathbf{b}' and \mathbf{t}

Second Phase: Public discussion

- ① When $t_i = 1$, it means that Bob has used the wrong basis, and so he knows the original bit b_i , that is $1 - b'_i$. Bob sends \mathbf{t} over the public channel, and both parties keep only the bits for which $t_i = 1$.
- ② Alice chooses a subset of the remaining bits in \mathbf{b} and discloses their values to Bob. For each b_i in this subset, it should be $b_i = 1 - b'_i$. If Bob notices discrepancies $b_i = b'_i$, eavesdropping is detected.
- ③ The shared secret $\mathbf{k} \in \{0, 1\}^N$ is the string of bits in $\mathbf{b} = \mathbf{1} - \mathbf{b}'$ that have not been disclosed at step 2.

B92 Observation

- Alice secret, \vec{d} , is never completely disclosed. Each bit is sent codified in a quantum state ($|0\rangle$ or $|+\rangle$), but these two state are not distinguishable without altering them.
- If Bob guesses the right base, he gains no information. If he uses the wrong base, he could measure a state in $\{|0\rangle, |+\rangle\}$, which gives him no information, or a state in $\{|1\rangle, |-\rangle\}$, which is surely an error. He so can deduce the basis that Alice was using.
- With respect to BB84, B92 discards more bits, as the probability of Bob gaining information is $\frac{1}{4}$.

PRISM Formalization: Alice

```

module Alice
  al_state : [0..5];
  al_index : [1..N];
  al_bas : [0..1];
  al_bit : [0..1];
  [] (al_state=0) & (eve_state>0) ->
  0.5 : (al_state'=1) & (al_bit'=0) & (al_bas'=0)+
  0.5 : (al_state'=1) & (al_bit'=0) & (al_bas'=1);
  [aliceput] (al_state=1) -> (al_state'=2);
  [reveal] (al_state=2) -> (al_state'=3);
  [loop] (al_state=3) & (al_index<N) -> (al_state'=0) & (al_index'=al_index+1);
  [loop] (al_state=3) & (al_index=N) -> (al_state'=4);
  [stop] (al_state=3) -> (al_state'=3);
  [stop] (al_state=4) -> (al_state'=4);
endmodule

```

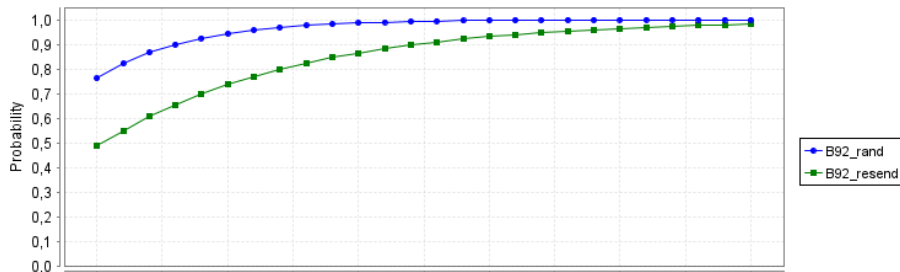
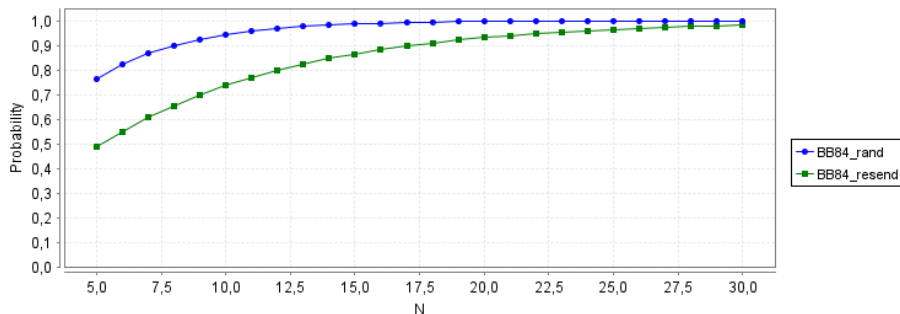
PRISM Formalization: Bob

```

module Bob
  bob_state : [0..7];
  bob_bas : [0..1];
  bob_bit : [0..1];
  [] (bob_state=0) & (eve_state>0) ->
    0.5 : (bob_state'=2) & (bob_bas'=0) +
    0.5 : (bob_state'=2) & (bob_bas'=1);
  [bobget] (bob_state=2) & (bob_bas=ch_bas) -> (bob_state'=3) & (bob_bit'=ch_bit);
  [bobget] (bob_state=2) & (bob_bas!=ch_bas) ->
    0.5 : (bob_bit'=ch_bit) & (bob_state'=3) +
    0.5 : (bob_bit'=1-ch_bit) & (bob_state'=3);
  [reveal] (bob_state=3) & (bob_bit = 0) -> (bob_state'=4);
  [reveal] (bob_state=3) & (bob_bit = 1) -> (bob_state' = 5);
  [] (bob_state=5) & (bob_bas != 1 - al_bas) -> (bob_state'=7);
  [] (bob_state=5) & (bob_bas = 1 - al_bas) -> (bob_state'=4);
  [loop] (bob_state=4) & (al_index<N) -> (bob_state'=0);
  [loop] (bob_state=4) & (al_index=N) -> (bob_state'=6);
  [stop] (bob_state=6) -> (bob_state'=6); // no eavesdropping detected
  [stop] (bob_state=7) -> (bob_state'=7); // eavesdropping detected
endmodule

```

BB84 and B92



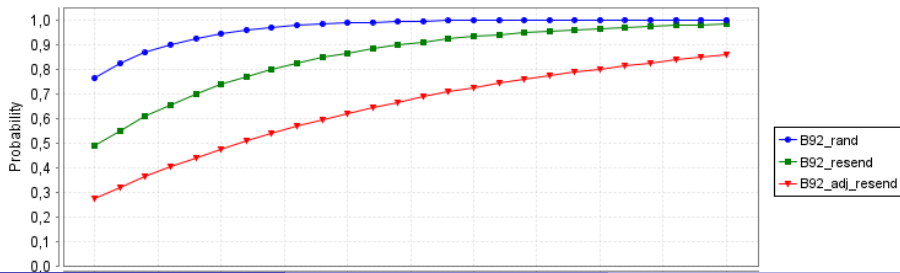
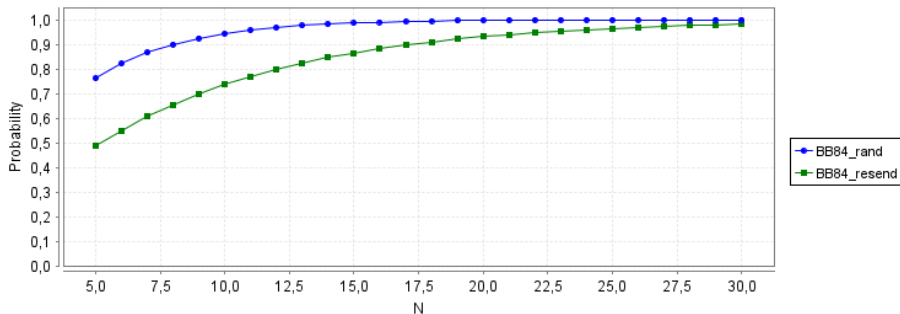
A smarter MIM attack

```

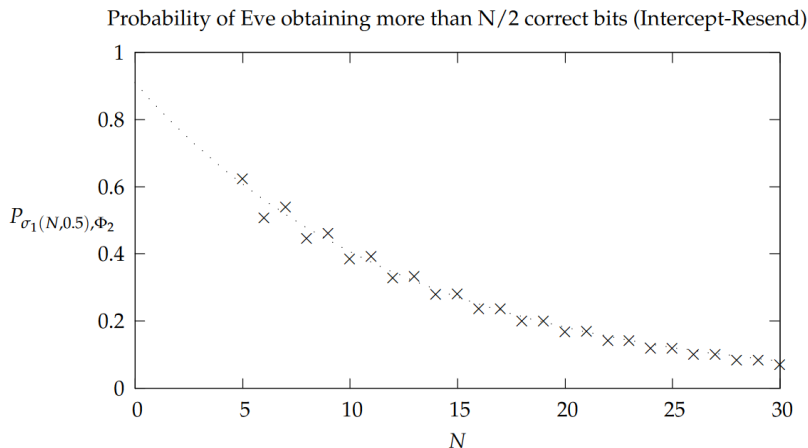
module EveClever
eve_state : [0..4];
eve_bas : [0..1];
eve_bit : [0..1];
[] (eve_state=0) ->
    0.5 : (eve_state'=1) & (eve_bas'=0) +
    0.5 : (eve_state'=1) & (eve_bas'=1);
[evemeasure] (eve_state=1) -> (eve_state'=2);
[eveget] (eve_state=2) -> (eve_state'=3) & (eve_bit'=ch_bit) ;
[] (eve_state=3) & (eve_bit = 1) ->
    (eve_bas'=1-eve_bas) & (eve_bit'=0) & (eve_state' = 4);
[] (eve_state=3) & (eve_bit = 0) ->
    (eve_bas'=eve_bas) & (eve_bit'=0) & (eve_state' = 4);
[eveput] (eve_state=4) -> (eve_state'=0);
endmodule

```

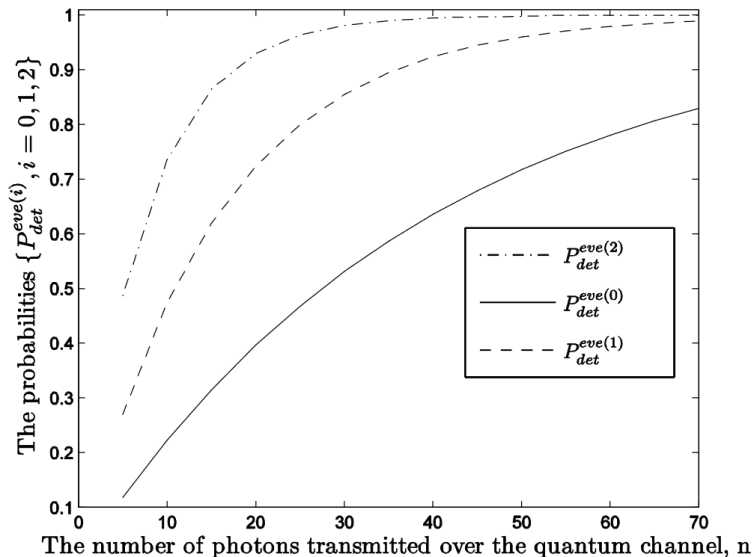
BB84 and B92



Other interesting metrics: bits obtained by Eve



Other interesting metrics: Eve with different power



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