Exploring Quantum Process Calculi

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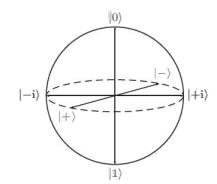
Qubits

The simplest quantum system is a **qubit**. Like a bit, a qubit has two separate states, $|0\rangle$ and $|1\rangle$.

A qubit can also be in a linear combination of states, knows as a **superposition**.

$$|+\rangle = \frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle$$

$$\left|-\right\rangle = \frac{1}{\sqrt{2}}\left|0\right\rangle - \frac{1}{\sqrt{2}}\left|1\right\rangle$$



Measurements

A qubit in superposition cannot be directly observed, because when it get **measured**, it **decays** in one of its basis states.

$$M_{01}\left(\alpha\left|0\right\rangle+\beta\left|1\right\rangle\right)$$
 with probability $|\alpha|^2$
$$|1\rangle$$
 with probability $|\beta|^2$

For example, when measuring the state $|+\rangle=\frac{1}{\sqrt{2}}\,|0\rangle+\frac{1}{\sqrt{2}}\,|1\rangle$, we get either $|0\rangle$ or $|1\rangle$ with the same probability.

No Cloning theorem

No-Cloning

Quantum information cannot be **duplicated**. That is, given a qubit q_1 in state

$$|\psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

it is impossible to prepare a qubit q_2 in the same state, without destroying the information in q_1 .

This means that, contrary to classical bits, qubits cannot be freely copied, stored or broadcasted to multiple receivers.

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In IqCCS, the programmer must **explicitly** describe what happens to each qubit

$$P' = c?q.X(q).c!q$$
 $Q' = c?q.H(q).c!q$
 $P'' = c?q.X(q).disc(q)$ $Q'' = c?q.H(q).disc(q)$

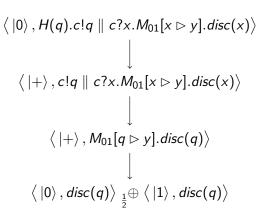
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The two processes seen before

$$P' = c?q.X(q).c!q$$
 $Q' = c?q.H(q).c!q$

are not bisimilar, because there is a context R which tells them apart

$$R = c?q.M_{01}[q > x].(disc(q) || if x = 0 then a!0 else b!0)$$



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$$C = \langle |+\rangle, M_{01}[q \rhd x].c!q \rangle \rightarrow \langle |0\rangle, c!q \rangle_{\frac{1}{2}} \oplus \langle |1\rangle, c!q \rangle$$
$$C' = \langle |0\rangle, M_{\pm}[q \rhd x].c!q \rangle \rightarrow \langle |+\rangle, c!q \rangle_{\frac{1}{2}} \oplus \langle |-\rangle, c!q \rangle$$

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$$\begin{split} \mathcal{C} &= \left\langle \left. \left| + \right\rangle, M_{01}[q \rhd x].c!q \right\rangle \to \left\langle \left. \left| 0 \right\rangle, c!q \right\rangle_{\frac{1}{2}} \oplus \left\langle \left. \left| 1 \right\rangle, c!q \right\rangle \\ \mathcal{C}' &= \left\langle \left. \left| 0 \right\rangle, M_{\pm}[q \rhd x].c!q \right\rangle \to \left\langle \left. \left| + \right\rangle, c!q \right\rangle_{\frac{1}{2}} \oplus \left\langle \left. \left| - \right\rangle, c!q \right\rangle \end{split}$$

According to the usual notion of **probabilistic bisimilarity**, two distributions are bisimilar if they assign the same probability to bisimilar processes. In our example, the two configurations are **not** bisimilar, because they evolve in two non-bisimilar distributions.

Indistinguishable distributions

According to quantum mechanics, it's impossible to distinguish the two sources S and S':

```
S emits a qubit |0\rangle or |1\rangle with the same probability S' emits a qubit |+\rangle or |-\rangle with the same probability
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Suppose we receive a qubit from either S or S', and measure it in the 01 basis. If we measure a qubit from S, it would result in either $|0\rangle$ or $|1\rangle$. If we measure a qubit from S', a $|+\rangle$ qubit would decay in either $|0\rangle$ or $|1\rangle$, and a $|-\rangle$ qubit would decay in either $|0\rangle$ or $|1\rangle$ as well.

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Inadequacy of Probabilistic Bisimilarity

The usual notion of probabilistic bisimilarity is too fine, when comparing distributions of quantum configurations.

Quantum Bisimilarity

We introduce an equivalence relation

$$\equiv \subseteq \mathfrak{D}(\mathit{Conf}) \times \mathfrak{D}(\mathit{Conf})$$

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$$\begin{split} \left\langle \left|+\right\rangle , M_{01}[q\rhd x].c!q\right\rangle &\to \left\langle \left|0\right\rangle , c!q\right\rangle _{\frac{1}{2}}\oplus \left\langle \left|1\right\rangle , c!q\right\rangle \equiv \left\langle \left|+\right\rangle , c!q\right\rangle _{\frac{1}{2}}\oplus \left\langle \left|-\right\rangle , c!q\right\rangle \\ \\ \left\langle \left|0\right\rangle , M_{\pm}[q\rhd x].c!q\right\rangle &\to \left\langle \left|+\right\rangle , c!q\right\rangle _{\frac{1}{2}}\oplus \left\langle \left|-\right\rangle , c!q\right\rangle \equiv \left\langle \left|+\right\rangle , c!q\right\rangle _{\frac{1}{2}}\oplus \left\langle \left|-\right\rangle , c!q\right\rangle \end{split}$$

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Minimal Quantum Process Algebra

Solution II: mQPA

We introduce a new calculus, equipped with a minimal set of features: communication, non-determinism and quantum measurement. In mQPA, the transitions are of the form

$$\rightarrow \subseteq S \times \mathfrak{D}(S)^{\mathcal{H}}$$

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In mQPA, the previous example can be rewritten as

$$egin{aligned} (\mathcal{S}_{\mid 0
angle 0} | \oplus \mathcal{S}') (|+
angle) &= \mathcal{S}_{rac{1}{2}} \oplus \mathcal{S}' \ (\mathcal{S}_{\mid +
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- Even though quantum systems exhibit a probabilistic behaviour,
 probabilistic bisimilarity is not really well suited for the quantum setting.
- mQPA, a minimal calculus, pursuing a foundational approach on which are the dynamics and observable properties of quantum systems.



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- From weak transitions, it is possible to define reachability, temporal logics and model checking.
- Saturated bisimilarity can be cumbersome to prove, it will be interesting to explore how the existing proof techniques adapt to the quantum setting.
- We will investigate the relation between mQPA semantics and the usual, configuration-based semantics, together with their respective bisimilarities.

Thank you for your attention!

